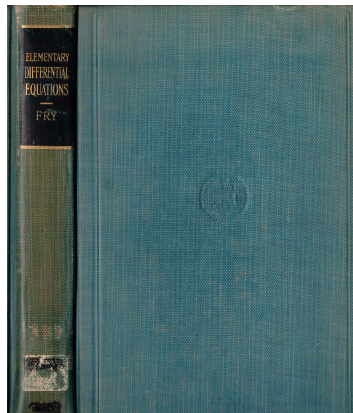


A Solution Manual For

**Elementary Differential Equations. By Thornton
C. Fry. D Van Nostrand. NY. First Edition (1929)**



Nasser M. Abbasi December 31, 2024

Compiled on December 31, 2024 at 3:55am

Contents

1	Lookup tables for all problems in current book	5
2	Book Solved Problems	11

CHAPTER 1

LOOKUP TABLES FOR ALL PROBLEMS IN CURRENT BOOK

1.1	Chapter 1. section 5. Problems at page 19	6
1.2	Chapter IV. Methods of solution: First order equations. section 24. Problems at page 62	6
1.3	Chapter IV. Methods of solution: First order equations. section 29. Problems at page 81	7
1.4	Chapter IV. Methods of solution: First order equations. section 31. Problems at page 85	7
1.5	Chapter IV. Methods of solution: First order equations. section 32. Problems at page 89	8
1.6	Chapter IV. Methods of solution: First order equations. section 33. Problems at page 91	8
1.7	Chapter V. Singular solutions. section 36. Problems at page 99	9
1.8	Chapter VII. Linear equations of order higher than the first. section 56. Problems at page 163	9
1.9	Chapter VII. Linear equations of order higher than the first. section 63. Problems at page 196	10

1.1 Chapter 1. section 5. Problems at page 19

Table 1.1: Lookup table for all problems in current section

ID	problem	ODE
18207	2	$x^2 y'' - \frac{x^2 y'^2}{2y} + 4xy' + 4y = 0$
18208	3	$y' + cy = a$
18209	4	$y'' + \frac{y'}{x} + k^2 y = 0$
18210	5	$\cos(x) y' + \sin(x) y'' + ny \sin(x) = 0$
18211	6	$y' = \frac{\sqrt{1-y^2} \arcsin(y)}{x}$
18212	16 (a)	$v'' = \left(\frac{1}{v} + v'^4\right)^{1/3}$
18213	16 (b)	$v' + u^2 v = \sin(u)$
18214	17 (a)	$\sqrt{y' + y} = (y'' + 2x)^{1/4}$
18215	18	$v' + \frac{2v}{u} = 3$

1.2 Chapter IV. Methods of solution: First order equations. section 24. Problems at page 62

Table 1.2: Lookup table for all problems in current section

ID	problem	ODE
18216	4 (a)	$\sin(x) \cos(y)^2 + \cos(x)^2 y' = 0$
18217	4 (b)	$y' + \sqrt{\frac{1-y^2}{-x^2+1}} = 0$
18218	4 (c)	$y - xy' = b(1 + x^2 y')$
18219	5	$x' = k(A - nx)(M - mx)$
18220	6	$y' = 1 + \frac{1}{x} - \frac{1}{y^2+2} - \frac{1}{x(y^2+2)}$

1.3 Chapter IV. Methods of solution: First order equations. section 29. Problems at page 81

Table 1.3: Lookup table for all problems in current section

ID	problem	ODE
18221	1	$y^2 = x(y - x) y'$
18222	2	$2x^2y + y^3 - x^3y' = 0$
18223	3	$2ax + by + (2cy + bx + e) y' = g$
18224	4	$\sec(x)^2 \tan(y) y' + \sec(y)^2 \tan(x) = 0$
18225	5	$x + yy' = my$
18226	6	$\frac{2x}{y^3} + \left(\frac{1}{y^2} - \frac{3x^2}{y^4}\right) y' = 0$
18227	8	$\left(T + \frac{1}{\sqrt{t^2 - T^2}}\right) T' = \frac{T}{t\sqrt{t^2 - T^2}} - t$

1.4 Chapter IV. Methods of solution: First order equations. section 31. Problems at page 85

Table 1.4: Lookup table for all problems in current section

ID	problem	ODE
18228	1	$y' + yx = x$
18229	2	$y' + \frac{y}{x} = \sin(x)$
18230	3	$y' + \frac{y}{x} = \frac{\sin(x)}{y^3}$
18231	4	$p' = \frac{p+at^3-2pt^2}{t(-t^2+1)}$
18232	5	$(T \ln(t) - 1) T = tT'$
18233	6	$y' + y \cos(x) = \frac{\sin(2x)}{2}$

Continued on next page

Table 1.4 Lookup table
Continued from previous page

ID	problem	ODE
18234	7	$y - \cos(x) y' = y^2 \cos(x) (1 - \sin(x))$

1.5 Chapter IV. Methods of solution: First order equations. section 32. Problems at page 89

Table 1.5: Lookup table for all problems in current section

ID	problem	ODE
18235	2	$xy'^2 - y + 2y' = 0$
18236	3	$2y'^3 + y'^2 - y = 0$
18237	4	$y' = e^{z-y'}$
18238	5	$\sqrt{t^2 + T} = T'$
18239	7	$(x^2 - 1) y'^2 = 1$
18240	8	$y' = (x + y)^2$

1.6 Chapter IV. Methods of solution: First order equations. section 33. Problems at page 91

Table 1.6: Lookup table for all problems in current section

ID	problem	ODE
18241	1	$\theta'' = -p^2\theta$
18242	2 (eq 39)	$\sec(\theta)^2 = \frac{ms'}{k}$
18243	3 (eq 41)	$y'' = \frac{m\sqrt{1+y'^2}}{k}$
18244	4 (eq 50)	$\phi'' = \frac{4\pi nc}{\sqrt{v_0^2 + \frac{2e(\phi - V_0)}{m}}}$

Continued on next page

Table 1.6 Lookup table
Continued from previous page

ID	problem	ODE
18245	8 (eq 68)	$y' = x(ay^2 + b)$
18246	8 (eq 69)	$n' = (n^2 + 1)x$
18247	9 (a)	$v' + \frac{2v}{u} = 3v$
18248	9 (b)	$\sqrt{-u^2 + 1} v' = 2u\sqrt{1 - v^2}$
18249	9 (c)	$\sqrt{1 + v'} = \frac{e^u}{2}$
18250	9 (d)	$\frac{y'}{x} = y \sin(x^2 - 1) - \frac{2y}{\sqrt{x}}$
18251	9 (e)	$y' = 1 + \frac{2y}{x-y}$
18252	10 (a)	$v' + 2vu = 2u$
18253	10 (b)	$1 + v^2 + (u^2 + 1)vv' = 0$
18254	10 (c)	$u \ln(u) v' + \sin(v)^2 = 1$

1.7 Chapter V. Singular solutions. section 36. Problems at page 99

Table 1.7: Lookup table for all problems in current section

ID	problem	ODE
18255	1 (eq 98)	$4yy'^3 - 2x^2y'^2 + 4xyy' + x^3 = 16y^2$

1.8 Chapter VII. Linear equations of order higher than the first. section 56. Problems at page 163

Table 1.8: Lookup table for all problems in current section

ID	problem	ODE
18256	1 (eq 100)	$\theta'' - p^2\theta = 0$
18257	2	$y'' + y = 0$
18258	3	$y'' + 12y = 7y'$
18259	4	$r'' - a^2r = 0$
18260	5	$y'''' - a^4y = 0$
18261	6	$v'' - 6v' + 13v = e^{-2u}$
18262	7	$y'' + 4y' - y = \sin(t)$
18263	8	$y'' + 3y = \sin(x) + \frac{\sin(3x)}{3}$
18264	10	$5x' + x = \sin(3t)$
18265	11	$x'''' - 6x''' + 11x'' - 6x' = e^{-3t}$
18266	14	$x^4y'''' + x^3y''' - 20x^2y'' + 20xy' = 17x^6$
18267	15	$t^4x'''' - 2t^3x''' - 20t^2x'' + 12tx' + 16x = \cos(3 \ln(t))$

1.9 Chapter VII. Linear equations of order higher than the first. section 63. Problems at page 196

Table 1.9: Lookup table for all problems in current section

ID	problem	ODE
18268	1	$y''' - y'' - y' + y = 0$
18269	2	$y'''' - 3y''' + 3y'' - y' = e^{2x}$
18270	3	$y''' - y'' + y' - y = \cos(x)$
18271	8	$x^2y'' + 3xy' + y = \frac{1}{x}$

CHAPTER 2

BOOK SOLVED PROBLEMS

2.1	Chapter 1. section 5. Problems at page 19	12
2.2	Chapter IV. Methods of solution: First order equations. section 24. Problems at page 62	87
2.3	Chapter IV. Methods of solution: First order equations. section 29. Problems at page 81	133
2.4	Chapter IV. Methods of solution: First order equations. section 31. Problems at page 85	271
2.5	Chapter IV. Methods of solution: First order equations. section 32. Problems at page 89	351
2.6	Chapter IV. Methods of solution: First order equations. section 33. Problems at page 91	392
2.7	Chapter V. Singular solutions. section 36. Problems at page 99	567
2.8	Chapter VII. Linear equations of order higher than the first. section 56. Problems at page 163	587
2.9	Chapter VII. Linear equations of order higher than the first. section 63. Problems at page 196	716

2.1 Chapter 1. section 5. Problems at page 19

2.1.1	problem 2	13
2.1.2	problem 3	15
2.1.3	problem 4	26
2.1.4	problem 5	33
2.1.5	problem 6	35
2.1.6	problem 16 (a)	47
2.1.7	problem 16 (b)	50
2.1.8	problem 17 (a)	59
2.1.9	problem 18	62

2.1.1 problem 2

Maple step by step solution	13
Maple trace	13
Maple dsolve solution	14
Mathematica DSolve solution	14

Internal problem ID [18207]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter 1. section 5. Problems at page 19

Problem number : 2

Date solved : Friday, December 20, 2024 at 05:50:33 AM

CAS classification :

[[_2nd_order, _with_linear_symmetries], [_2nd_order, _reducible, _mu_xy]]

Solve

$$x^2 y'' - \frac{x^2 y'^2}{2y} + 4xy' + 4y = 0$$

Maple step by step solution

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying 2nd order Liouville
trying 2nd order WeierstrassP
trying 2nd order JacobiSN
differential order: 2; trying a linearization to 3rd order
trying 2nd order ODE linearizable_by_differentiation
-> Calling odsolve with the ODE`, diff(diff(diff(y(x), x), x), x)+12*(x^2*(diff(diff(y
  Methods for third order ODEs:
  --- Trying classification methods ---
  trying a quadrature
  checking if the LODE has constant coefficients
  checking if the LODE is of Euler type
  <- LODE of Euler type successful
<- 2nd order ODE linearizable_by_differentiation successful`

```

Maple dsolve solution

Solving time : 0.027 (sec)

Leaf size : 21

```
dsolve(x^2*diff(diff(y(x),x),x)-1/2*x^2/y(x)*diff(y(x),x)^2+4*x*diff(y(x),x)+4*y(x) =
      y(x),singsol=all)
```

$$y = \frac{(c_1x + 2c_2)^2}{4c_2x^4}$$

Mathematica DSolve solution

Solving time : 0.493 (sec)

Leaf size : 19

```
DSolve[{x^2*D[y[x],{x,2}]-x^2/(2*y[x])*D[y[x],x]^2+4*x*D[y[x],x]+4*y[x]==0,{}},
      y[x],x,IncludeSingularSolutions->True]
```

$$y(x) \rightarrow \frac{c_2(x + 2c_1)^2}{x^4}$$

2.1.2 problem 3

Solved as first order autonomous ode	15
Solved as first order Exact ode	17
Solved using Lie symmetry for first order ode	20
Maple step by step solution	24
Maple trace	25
Maple dsolve solution	25
Mathematica DSolve solution	25

Internal problem ID [18208]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter 1. section 5. Problems at page 19

Problem number : 3

Date solved : Friday, December 20, 2024 at 05:50:33 AM

CAS classification : [_quadrature]

Solve

$$y' + cy = a$$

Solved as first order autonomous ode

Time used: 0.319 (sec)

Integrating gives

$$\int \frac{1}{-cy + a} dy = dx$$

$$-\frac{\ln(-cy + a)}{c} = x + c_1$$

Singular solutions are found by solving

$$-cy + a = 0$$

for y . This is because we had to divide by this in the above step. This gives the following singular solution(s), which also have to satisfy the given ODE.

$$y = \frac{a}{c}$$

The following diagram is the phase line diagram. It classifies each of the above equilibrium points as stable or not stable or semi-stable.

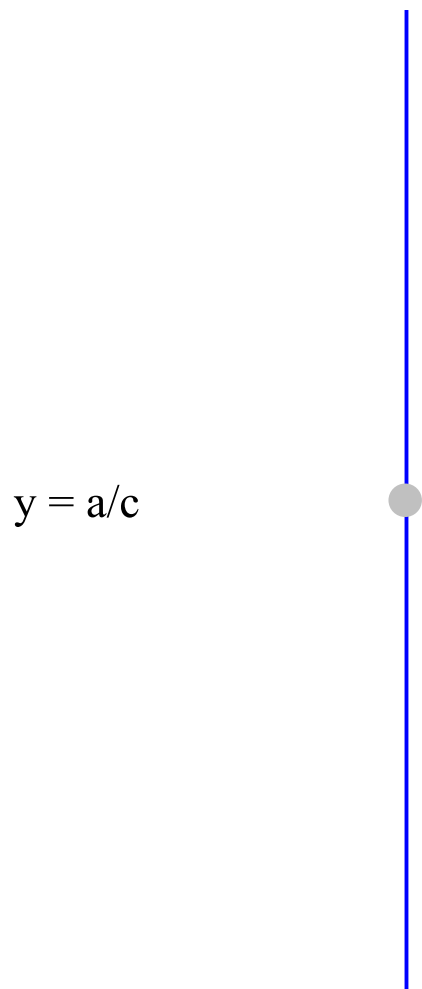


Figure 2.1: Phase line diagram

Solving for y gives

$$y = \frac{a}{c}$$

$$y = -\frac{e^{-cc_1 - xc} - a}{c}$$

Summary of solutions found

$$y = \frac{a}{c}$$

$$y = -\frac{e^{-cc_1 - xc} - a}{c}$$

Solved as first order Exact ode

Time used: 0.138 (sec)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (\text{1A})$$

Therefore

$$\begin{aligned} dy &= (-cy + a) dx \\ (cy - a) dx + dy &= 0 \end{aligned} \quad (\text{2A})$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= cy - a \\ N(x, y) &= 1 \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(cy - a) \\ &= c\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(1) \\ &= 0\end{aligned}$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned}A &= \frac{1}{N} \left(\frac{\partial M}{\partial y} - \frac{\partial N}{\partial x} \right) \\ &= 1((c) - (0)) \\ &= c\end{aligned}$$

Since A does not depend on y , then it can be used to find an integrating factor. The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int A dx} \\ &= e^{\int c dx}\end{aligned}$$

The result of integrating gives

$$\begin{aligned}\mu &= e^{xc} \\ &= e^{xc}\end{aligned}$$

M and N are multiplied by this integrating factor, giving new M and new N which are called \bar{M} and \bar{N} for now so not to confuse them with the original M and N .

$$\begin{aligned}\bar{M} &= \mu M \\ &= e^{xc}(cy - a) \\ &= -(-cy + a)e^{xc}\end{aligned}$$

And

$$\begin{aligned}\bar{N} &= \mu N \\ &= e^{xc}(1) \\ &= e^{xc}\end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned}\bar{M} + \bar{N} \frac{dy}{dx} &= 0 \\ (-(-cy + a)e^{xc}) + (e^{xc}) \frac{dy}{dx} &= 0\end{aligned}$$

The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = \bar{M} \tag{1}$$

$$\frac{\partial \phi}{\partial y} = \bar{N} \tag{2}$$

Integrating (2) w.r.t. y gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial y} dy &= \int \bar{N} dy \\ \int \frac{\partial \phi}{\partial y} dy &= \int e^{xc} dy \\ \phi &= e^{xc}y + f(x)\end{aligned} \tag{3}$$

Where $f(x)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t x gives

$$\frac{\partial \phi}{\partial x} = ce^{xc}y + f'(x) \tag{4}$$

But equation (1) says that $\frac{\partial \phi}{\partial x} = -(-cy + a)e^{xc}$. Therefore equation (4) becomes

$$-(-cy + a)e^{xc} = ce^{xc}y + f'(x) \tag{5}$$

Solving equation (5) for $f'(x)$ gives

$$f'(x) = -e^{xc}a$$

Integrating the above w.r.t x gives

$$\int f'(x) dx = \int (-e^{xc}a) dx$$

$$f(x) = -\frac{e^{xc}a}{c} + c_1$$

Where c_1 is constant of integration. Substituting result found above for $f(x)$ into equation (3) gives ϕ

$$\phi = e^{xc}y - \frac{e^{xc}a}{c} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into the constant c_1 gives the solution as

$$c_1 = e^{xc}y - \frac{e^{xc}a}{c}$$

Solving for y gives

$$y = \frac{(e^{xc}a + cc_1)e^{-xc}}{c}$$

Summary of solutions found

$$y = \frac{(e^{xc}a + cc_1)e^{-xc}}{c}$$

Solved using Lie symmetry for first order ode

Time used: 0.384 (sec)

Writing the ode as

$$y' = -cy + a$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2\xi_y - \omega_x\xi - \omega_y\eta = 0 \quad (\text{A})$$

To determine ξ, η then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$\xi = xa_2 + ya_3 + a_1 \quad (1E)$$

$$\eta = xb_2 + yb_3 + b_1 \quad (2E)$$

Where the unknown coefficients are

$$\{a_1, a_2, a_3, b_1, b_2, b_3\}$$

Substituting equations (1E,2E) and ω into (A) gives

$$b_2 + (-cy + a)(b_3 - a_2) - (-cy + a)^2 a_3 + c(xb_2 + yb_3 + b_1) = 0 \quad (5E)$$

Putting the above in normal form gives

$$-c^2 y^2 a_3 + 2acya_3 - a^2 a_3 + cxb_2 + cya_2 - aa_2 + ab_3 + cb_1 + b_2 = 0$$

Setting the numerator to zero gives

$$-c^2 y^2 a_3 + 2acya_3 - a^2 a_3 + cxb_2 + cya_2 - aa_2 + ab_3 + cb_1 + b_2 = 0 \quad (6E)$$

Looking at the above PDE shows the following are all the terms with $\{x, y\}$ in them.

$$\{x, y\}$$

The following substitution is now made to be able to collect on all terms with $\{x, y\}$ in them

$$\{x = v_1, y = v_2\}$$

The above PDE (6E) now becomes

$$-c^2 a_3 v_2^2 + 2aca_3 v_2 - a^2 a_3 + ca_2 v_2 + cb_2 v_1 - aa_2 + ab_3 + cb_1 + b_2 = 0 \quad (7E)$$

Collecting the above on the terms v_i introduced, and these are

$$\{v_1, v_2\}$$

Equation (7E) now becomes

$$cb_2v_1 - c^2a_3v_2^2 + (2aca_3 + ca_2)v_2 - a^2a_3 - aa_2 + ab_3 + cb_1 + b_2 = 0 \quad (8E)$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$\begin{aligned} cb_2 &= 0 \\ -c^2a_3 &= 0 \\ 2aca_3 + ca_2 &= 0 \\ -a^2a_3 - aa_2 + ab_3 + cb_1 + b_2 &= 0 \end{aligned}$$

Solving the above equations for the unknowns gives

$$\begin{aligned} a_1 &= a_1 \\ a_2 &= 0 \\ a_3 &= 0 \\ b_1 &= -\frac{ab_3}{c} \\ b_2 &= 0 \\ b_3 &= b_3 \end{aligned}$$

Substituting the above solution in the anstaz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$\begin{aligned} \xi &= 0 \\ \eta &= -\frac{-cy + a}{c} \end{aligned}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS \quad (1)$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = x$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{\frac{-cy+a}{c}} dy \end{aligned}$$

Which results in

$$S = \ln(-cy + a)$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = -cy + a$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 1 \\ R_y &= 0 \\ S_x &= 0 \\ S_y &= -\frac{c}{-cy + a} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = -c \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = -c$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S .

Since the ode has the form $\frac{d}{dR}S(R) = f(R)$, then we only need to integrate $f(R)$.

$$\int dS = \int -c dR$$

$$S(R) = -cR + c_2$$

To complete the solution, we just need to transform the above back to x, y coordinates. This results in

$$\ln(-cy + a) = -xc + c_2$$

Which gives

$$y = -\frac{e^{-xc+c_2} - a}{c}$$

Summary of solutions found

$$y = -\frac{e^{-xc+c_2} - a}{c}$$

Maple step by step solution

Let's solve

$$y' + cy = a$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Solve for the highest derivative

$$y' = -cy + a$$

- Separate variables

$$\frac{y'}{-cy+a} = 1$$

- Integrate both sides with respect to x

$$\int \frac{y'}{-cy+a} dx = \int 1 dx + C1$$

- Evaluate integral

$$-\frac{\ln(-cy+a)}{c} = x + C1$$

- Solve for y

$$y = -\frac{e^{-C1c-xc-a}}{c}$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`

```

Maple dsolve solution

Solving time : 0.002 (sec)

Leaf size : 18

```

dsolve(diff(y(x),x)+c*y(x) = a,
        y(x),singsol=all)

```

$$y = \frac{e^{-xc}c_1c + a}{c}$$

Mathematica DSolve solution

Solving time : 0.081 (sec)

Leaf size : 29

```

DSolve[{D[y[x],x]+c*y[x]==a,{}},
        y[x],x,IncludeSingularSolutions->True]

```

$$y(x) \rightarrow \frac{a}{c} + c_1 e^{-cx}$$

$$y(x) \rightarrow \frac{a}{c}$$

2.1.3 problem 4

Solved as second order Bessel ode	26
Solved as second order ode adjoint method	27
Maple step by step solution	30
Maple trace	31
Maple dsolve solution	32
Mathematica DSolve solution	32

Internal problem ID [18209]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter 1. section 5. Problems at page 19

Problem number : 4

Date solved : Friday, December 20, 2024 at 05:50:34 AM

CAS classification : [[_2nd_order, _with_linear_symmetries]]

Solve

$$y'' + \frac{y'}{x} + k^2 y = 0$$

Solved as second order Bessel ode

Time used: 0.057 (sec)

Writing the ode as

$$x^2 y'' + x y' + k^2 x^2 y = 0 \tag{1}$$

Bessel ode has the form

$$x^2 y'' + x y' + (-n^2 + x^2) y = 0 \tag{2}$$

The generalized form of Bessel ode is given by Bowman (1958) as the following

$$x^2 y'' + (1 - 2\alpha) x y' + (\beta^2 \gamma^2 x^{2\gamma} - n^2 \gamma^2 + \alpha^2) y = 0 \tag{3}$$

With the standard solution

$$y = x^\alpha (c_1 \text{BesselJ}(n, \beta x^\gamma) + c_2 \text{BesselY}(n, \beta x^\gamma)) \tag{4}$$

Comparing (3) to (1) and solving for α, β, n, γ gives

$$\alpha = 0$$

$$\beta = k$$

$$n = 0$$

$$\gamma = 1$$

Substituting all the above into (4) gives the solution as

$$y = c_1 \text{BesselJ}(0, kx) + c_2 \text{BesselY}(0, kx)$$

Will add steps showing solving for IC soon.

Summary of solutions found

$$y = c_1 \text{BesselJ}(0, kx) + c_2 \text{BesselY}(0, kx)$$

Solved as second order ode adjoint method

Time used: 0.746 (sec)

In normal form the ode

$$y'' + \frac{y'}{x} + k^2 y = 0 \quad (1)$$

Becomes

$$y'' + p(x)y' + q(x)y = r(x) \quad (2)$$

Where

$$p(x) = \frac{1}{x}$$

$$q(x) = k^2$$

$$r(x) = 0$$

The Lagrange adjoint ode is given by

$$\begin{aligned} \xi'' - (\xi p)' + \xi q &= 0 \\ \xi'' - \left(\frac{\xi(x)}{x}\right)' + (k^2 \xi(x)) &= 0 \\ \frac{\xi(x) k^2 x^2 + \xi''(x) x^2 - \xi'(x) x + \xi(x)}{x^2} &= 0 \end{aligned}$$

Which is solved for $\xi(x)$. Writing the ode as

$$\xi'' x^2 - \xi' x + (k^2 x^2 + 1) \xi = 0 \quad (1)$$

Bessel ode has the form

$$\xi'' x^2 + \xi' x + (-n^2 + x^2) \xi = 0 \quad (2)$$

The generalized form of Bessel ode is given by Bowman (1958) as the following

$$\xi'' x^2 + (1 - 2\alpha) x \xi' + (\beta^2 \gamma^2 x^{2\gamma} - n^2 \gamma^2 + \alpha^2) \xi = 0 \quad (3)$$

With the standard solution

$$\xi = x^\alpha (c_3 \text{BesselJ}(n, \beta x^\gamma) + c_4 \text{BesselY}(n, \beta x^\gamma)) \quad (4)$$

Comparing (3) to (1) and solving for α, β, n, γ gives

$$\alpha = 1$$

$$\beta = k$$

$$n = 0$$

$$\gamma = 1$$

Substituting all the above into (4) gives the solution as

$$\xi = c_3 x \text{BesselJ}(0, kx) + c_4 x \text{BesselY}(0, kx)$$

Will add steps showing solving for IC soon.

The original ode now reduces to first order ode

$$\begin{aligned} \xi(x) y' - y \xi'(x) + \xi(x) p(x) y &= \int \xi(x) r(x) dx \\ y' + y \left(p(x) - \frac{\xi'(x)}{\xi(x)} \right) &= \frac{\int \xi(x) r(x) dx}{\xi(x)} \end{aligned}$$

Or

$$y' + y \left(\frac{1}{x} - \frac{c_3 \text{BesselJ}(0, kx) - c_3 x \text{BesselJ}(1, kx) k + c_4 \text{BesselY}(0, kx) - c_4 x \text{BesselY}(1, kx) k}{c_3 x \text{BesselJ}(0, kx) + c_4 x \text{BesselY}(0, kx)} \right) = 0$$

Which is now a first order ode. This is now solved for y . In canonical form a linear first order is

$$y' + q(x)y = p(x)$$

Comparing the above to the given ode shows that

$$q(x) = \frac{k(\text{BesselJ}(1, kx) c_3 + \text{BesselY}(1, kx) c_4)}{c_3 \text{BesselJ}(0, kx) + c_4 \text{BesselY}(0, kx)}$$

$$p(x) = 0$$

The integrating factor μ is

$$\begin{aligned} \mu &= e^{\int q dx} \\ &= e^{\int \frac{k(\text{BesselJ}(1, kx)c_3 + \text{BesselY}(1, kx)c_4)}{c_3 \text{BesselJ}(0, kx) + c_4 \text{BesselY}(0, kx)} dx} \\ &= \frac{1}{c_3 \text{BesselJ}(0, kx) + c_4 \text{BesselY}(0, kx)} \end{aligned}$$

The ode becomes

$$\frac{d}{dx} \mu y = 0$$

$$\frac{d}{dx} \left(\frac{y}{c_3 \text{BesselJ}(0, kx) + c_4 \text{BesselY}(0, kx)} \right) = 0$$

Integrating gives

$$\begin{aligned} \frac{y}{c_3 \text{BesselJ}(0, kx) + c_4 \text{BesselY}(0, kx)} &= \int 0 dx + c_5 \\ &= c_5 \end{aligned}$$

Dividing throughout by the integrating factor $\frac{1}{c_3 \text{BesselJ}(0, kx) + c_4 \text{BesselY}(0, kx)}$ gives the final solution

$$y = (c_3 \text{BesselJ}(0, kx) + c_4 \text{BesselY}(0, kx)) c_5$$

Hence, the solution found using Lagrange adjoint equation method is

$$y = (c_3 \text{BesselJ}(0, kx) + c_4 \text{BesselY}(0, kx)) c_5$$

The constants can be merged to give

$$y = c_3 \text{BesselJ}(0, kx) + c_4 \text{BesselY}(0, kx)$$

Will add steps showing solving for IC soon.

Summary of solutions found

$$y = c_3 \text{BesselJ}(0, kx) + c_4 \text{BesselY}(0, kx)$$

Maple step by step solution

Let's solve

$$y'' + \frac{y'}{x} + k^2 y = 0$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Check to see if $x_0 = 0$ is a regular singular point

- Define functions

$$[P_2(x) = \frac{1}{x}, P_3(x) = k^2]$$

- $x \cdot P_2(x)$ is analytic at $x = 0$

$$(x \cdot P_2(x)) \Big|_{x=0} = 1$$

- $x^2 \cdot P_3(x)$ is analytic at $x = 0$

$$(x^2 \cdot P_3(x)) \Big|_{x=0} = 0$$

- $x = 0$ is a regular singular point

Check to see if $x_0 = 0$ is a regular singular point

$$x_0 = 0$$

- Multiply by denominators

$$k^2 y x + y'' x + y' = 0$$

- Assume series solution for y

$$y = \sum_{k=0}^{\infty} a_k x^{k+r}$$

- Rewrite ODE with series expansions

- Convert $x \cdot y$ to series expansion

$$x \cdot y = \sum_{k=0}^{\infty} a_k x^{k+r+1}$$

- Shift index using $k \rightarrow k - 1$

$$x \cdot y = \sum_{k=1}^{\infty} a_{k-1} x^{k+r}$$

- Convert y' to series expansion

$$y' = \sum_{k=0}^{\infty} a_k (k+r) x^{k+r-1}$$

- Shift index using $k \rightarrow k + 1$

$$y' = \sum_{k=-1}^{\infty} a_{k+1} (k+r+1) x^{k+r}$$

- Convert $x \cdot y''$ to series expansion

$$x \cdot y'' = \sum_{k=0}^{\infty} a_k (k+r)(k+r-1) x^{k+r-1}$$

- Shift index using $k \rightarrow k+1$

$$x \cdot y'' = \sum_{k=-1}^{\infty} a_{k+1} (k+r+1)(k+r) x^{k+r}$$

Rewrite ODE with series expansions

$$a_0 r^2 x^{-1+r} + a_1 (1+r)^2 x^r + \left(\sum_{k=1}^{\infty} (a_{k+1} (k+r+1)^2 + k^2 a_{k-1}) x^{k+r} \right) = 0$$

- a_0 cannot be 0 by assumption, giving the indicial equation
 $r^2 = 0$
- Values of r that satisfy the indicial equation
 $r = 0$
- Each term must be 0
 $a_1 (1+r)^2 = 0$
- Each term in the series must be 0, giving the recursion relation
 $a_{k+1} (k+1)^2 + k^2 a_{k-1} = 0$
- Shift index using $k \rightarrow k+1$
 $a_{k+2} (k+2)^2 + k^2 a_k = 0$
- Recursion relation that defines series solution to ODE
 $a_{k+2} = -\frac{k^2 a_k}{(k+2)^2}$
- Recursion relation for $r = 0$
 $a_{k+2} = -\frac{k^2 a_k}{(k+2)^2}$
- Solution for $r = 0$
$$\left[y = \sum_{k=0}^{\infty} a_k x^k, a_{k+2} = -\frac{k^2 a_k}{(k+2)^2}, a_1 = 0 \right]$$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
checking if the LODE is of Euler type
trying a symmetry of the form [xi=0, eta=F(x)]
checking if the LODE is missing y

```

```
-> Trying a Liouvillian solution using Kovacic's algorithm
<- No Liouvillian solutions exist
-> Trying a solution in terms of special functions:
    -> Bessel
    <- Bessel successful
<- special function solution successful`
```

Maple dsolve solution

Solving time : 0.002 (sec)

Leaf size : 19

```
dsolve(diff(diff(y(x),x),x)+1/x*diff(y(x),x)+k^2*y(x) = 0,
        y(x),singsol=all)
```

$$y = c_1 \text{BesselJ}(0, kx) + c_2 \text{BesselY}(0, kx)$$

Mathematica DSolve solution

Solving time : 0.029 (sec)

Leaf size : 22

```
DSolve[{D[y[x],{x,2}]+1/x*D[y[x],x]+k^2*y[x]==0,{}},
        y[x],x,IncludeSingularSolutions->True]
```

$$y(x) \rightarrow c_1 \text{BesselJ}(0, kx) + c_2 \text{BesselY}(0, kx)$$

2.1.4 problem 5

Maple step by step solution	33
Maple trace	33
Maple dsolve solution	34
Mathematica DSolve solution	34

Internal problem ID [18210]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter 1. section 5. Problems at page 19

Problem number : 5

Date solved : Friday, December 20, 2024 at 05:50:36 AM

CAS classification : [[_2nd_order, _with_linear_symmetries]]

Solve

$$\cos(x)y' + \sin(x)y'' + ny \sin(x) = 0$$

Maple step by step solution

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a symmetry of the form [xi=0, eta=F(x)]
checking if the LODE is missing y
-> Heun: Equivalence to the GHE or one of its 4 confluent cases under a power @ Moebiu
-> trying a solution of the form r0(x) * Y + r1(x) * Y where Y = exp(int(r(x), dx)) *
-> Trying changes of variables to rationalize or make the ODE simpler
    trying a quadrature
    checking if the LODE has constant coefficients
    checking if the LODE is of Euler type
    trying a symmetry of the form [xi=0, eta=F(x)]
    checking if the LODE is missing y
    -> Trying a Liouvillian solution using Kovacic's algorithm
    <- No Liouvillian solutions exists
    -> Trying a solution in terms of special functions:
        -> Bessel
        -> elliptic
        -> Legendre

```

```

<- Legendre successful
<- special function solution successful
Change of variables used:
[x = arccos(t)]
Linear ODE actually solved:
(-n*t^2+n)*u(t)+(2*t^3-2*t)*diff(u(t),t)+(t^4-2*t^2+1)*diff(diff(u(t),t),t) = 0
<- change of variables successful`

```

Maple dsolve solution

Solving time : 0.040 (sec)

Leaf size : 37

```

dsolve(cos(x)*diff(y(x),x)+sin(x)*diff(diff(y(x),x),x)+n*y(x)*sin(x) = 0,
y(x),singsol=all)

```

$$y = c_1 \text{LegendreP}\left(\frac{\sqrt{4n+1}}{2} - \frac{1}{2}, \cos(x)\right) + c_2 \text{LegendreQ}\left(\frac{\sqrt{4n+1}}{2} - \frac{1}{2}, \cos(x)\right)$$

Mathematica DSolve solution

Solving time : 0.129 (sec)

Leaf size : 48

```

DSolve[{D[Sin[x]*D[y[x],x],x]+n*y[x]*Sin[x]==0,{}},
y[x],x,IncludeSingularSolutions->True]

```

$$y(x) \rightarrow c_1 \text{LegendreP}\left(\frac{1}{2}(\sqrt{4n+1}-1), \cos(x)\right) + c_2 \text{LegendreQ}\left(\frac{1}{2}(\sqrt{4n+1}-1), \cos(x)\right)$$

2.1.5 problem 6

Solved as first order separable ode	35
Solved as first order Exact ode	37
Solved as first order isobaric ode	42
Maple step by step solution	44
Maple trace	45
Maple dsolve solution	45
Mathematica DSolve solution	46

Internal problem ID [18211]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter 1. section 5. Problems at page 19

Problem number : 6

Date solved : Friday, December 20, 2024 at 05:50:38 AM

CAS classification : [_separable]

Solve

$$y' = \frac{\sqrt{1 - y^2} \arcsin(y)}{x}$$

Solved as first order separable ode

Time used: 47.137 (sec)

The ode $y' = \frac{\sqrt{1 - y^2} \arcsin(y)}{x}$ is separable as it can be written as

$$\begin{aligned} y' &= \frac{\sqrt{1 - y^2} \arcsin(y)}{x} \\ &= f(x)g(y) \end{aligned}$$

Where

$$\begin{aligned} f(x) &= \frac{1}{x} \\ g(y) &= \sqrt{-y^2 + 1} \arcsin(y) \end{aligned}$$

Integrating gives

$$\begin{aligned} \int \frac{1}{g(y)} dy &= \int f(x) dx \\ \int \frac{1}{\sqrt{-y^2 + 1} \arcsin(y)} dy &= \int \frac{1}{x} dx \\ \ln(\arcsin(y)) &= \ln(x) + c_1 \end{aligned}$$

We now need to find the singular solutions, these are found by finding for what values $g(y)$ is zero, since we had to divide by this above. Solving $g(y) = 0$ or $\sqrt{-y^2 + 1} \arcsin(y) = 0$ for y gives

$$y = -1$$

$$y = 0$$

$$y = 1$$

Now we go over each such singular solution and check if it verifies the ode itself and any initial conditions given. If it does not then the singular solution will not be used.

Therefore the solutions found are

$$\ln(\arcsin(y)) = \ln(x) + c_1$$

$$y = -1$$

$$y = 0$$

$$y = 1$$

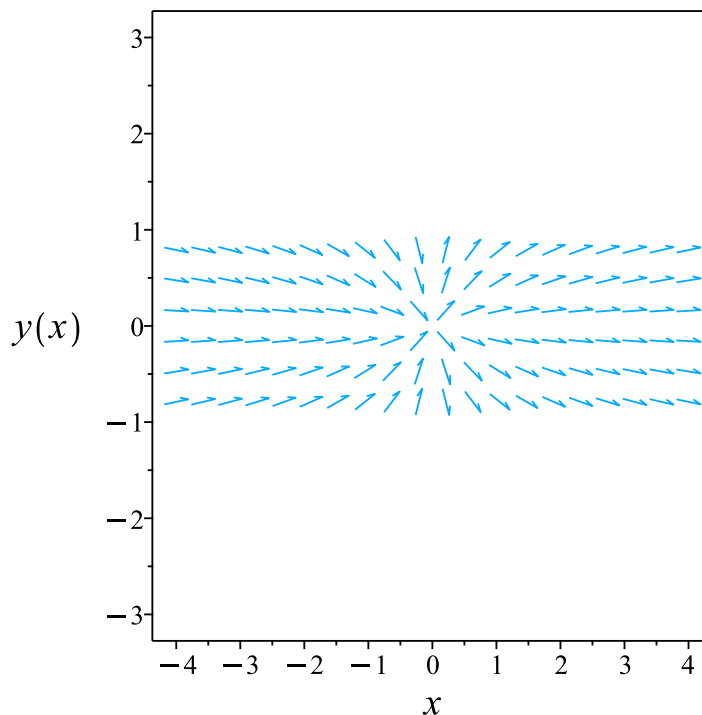


Figure 2.2: Slope field plot

$$y' = \frac{\sqrt{1-y^2} \arcsin(y)}{x}$$

Summary of solutions found

$$\ln(\arcsin(y)) = \ln(x) + c_1$$

$$y = -1$$

$$y = 0$$

$$y = 1$$

Solved as first order Exact ode

Time used: 0.918 (sec)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (\text{1A})$$

Therefore

$$\begin{aligned} dy &= \left(\frac{\sqrt{-y^2 + 1} \arcsin(y)}{x} \right) dx \\ \left(-\frac{\sqrt{-y^2 + 1} \arcsin(y)}{x} \right) dx + dy &= 0 \end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= -\frac{\sqrt{-y^2 + 1} \arcsin(y)}{x} \\ N(x, y) &= 1 \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} \left(-\frac{\sqrt{-y^2 + 1} \arcsin(y)}{x} \right) \\ &= \frac{-1 + \frac{\arcsin(y)y}{\sqrt{-y^2+1}}}{x} \end{aligned}$$

And

$$\begin{aligned} \frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(1) \\ &= 0 \end{aligned}$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned} A &= \frac{1}{N} \left(\frac{\partial M}{\partial y} - \frac{\partial N}{\partial x} \right) \\ &= 1 \left(\left(\frac{\arcsin(y)y}{x\sqrt{-y^2+1}} - \frac{1}{x} \right) - (0) \right) \\ &= \frac{-1 + \frac{\arcsin(y)y}{\sqrt{-y^2+1}}}{x} \end{aligned}$$

Since A depends on y , it can not be used to obtain an integrating factor. We will now try a second method to find an integrating factor. Let

$$\begin{aligned} B &= \frac{1}{M} \left(\frac{\partial N}{\partial x} - \frac{\partial M}{\partial y} \right) \\ &= -\frac{x}{\sqrt{-y^2+1} \arcsin(y)} \left((0) - \left(\frac{\arcsin(y)y}{x\sqrt{-y^2+1}} - \frac{1}{x} \right) \right) \\ &= \frac{-\arcsin(y)y + \sqrt{-y^2+1}}{\arcsin(y)(y^2-1)} \end{aligned}$$

Since B does not depend on x , it can be used to obtain an integrating factor. Let the integrating factor be μ . Then

$$\begin{aligned} \mu &= e^{\int B dy} \\ &= e^{\int \frac{-\arcsin(y)y + \sqrt{-y^2+1}}{\arcsin(y)(y^2-1)} dy} \end{aligned}$$

The result of integrating gives

$$\begin{aligned} \mu &= e^{-\ln(\arcsin(y)) - \frac{\ln(y-1)}{2} - \frac{\ln(y+1)}{2}} \\ &= \frac{1}{\arcsin(y) \sqrt{y-1} \sqrt{y+1}} \end{aligned}$$

M and N are now multiplied by this integrating factor, giving new M and new N which are called \bar{M} and \bar{N} so not to confuse them with the original M and N .

$$\begin{aligned} \bar{M} &= \mu M \\ &= \frac{1}{\arcsin(y) \sqrt{y-1} \sqrt{y+1}} \left(-\frac{\sqrt{-y^2+1} \arcsin(y)}{x} \right) \\ &= -\frac{\sqrt{-y^2+1}}{x\sqrt{y-1} \sqrt{y+1}} \end{aligned}$$

And

$$\begin{aligned} \bar{N} &= \mu N \\ &= \frac{1}{\arcsin(y) \sqrt{y-1} \sqrt{y+1}} (1) \\ &= \frac{1}{\arcsin(y) \sqrt{y-1} \sqrt{y+1}} \end{aligned}$$

So now a modified ODE is obtained from the original ODE which will be exact and can be solved using the standard method. The modified ODE is

$$\begin{aligned} \bar{M} + \bar{N} \frac{dy}{dx} &= 0 \\ \left(-\frac{\sqrt{-y^2+1}}{x\sqrt{y-1} \sqrt{y+1}} \right) + \left(\frac{1}{\arcsin(y) \sqrt{y-1} \sqrt{y+1}} \right) \frac{dy}{dx} &= 0 \end{aligned}$$

The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = \bar{M} \quad (1)$$

$$\frac{\partial \phi}{\partial y} = \bar{N} \quad (2)$$

Integrating (2) w.r.t. y gives

$$\begin{aligned} \int \frac{\partial \phi}{\partial y} dy &= \int \bar{N} dy \\ \int \frac{\partial \phi}{\partial y} dy &= \int \frac{1}{\arcsin(y) \sqrt{y-1} \sqrt{y+1}} dy \\ \phi &= \int \frac{1}{\arcsin(y) \sqrt{y-1} \sqrt{y+1}} dy + f(x) \end{aligned} \quad (3)$$

Where $f(x)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t x gives

$$\frac{\partial \phi}{\partial x} = 0 + f'(x) \quad (4)$$

But equation (1) says that $\frac{\partial \phi}{\partial x} = -\frac{\sqrt{-y^2+1}}{x\sqrt{y-1}\sqrt{y+1}}$. Therefore equation (4) becomes

$$-\frac{\sqrt{-y^2+1}}{x\sqrt{y-1}\sqrt{y+1}} = 0 + f'(x) \quad (5)$$

Solving equation (5) for $f'(x)$ gives

$$f'(x) = -\frac{\sqrt{-y^2+1}}{x\sqrt{y-1}\sqrt{y+1}}$$

Integrating the above w.r.t x gives

$$\begin{aligned} \int f'(x) dx &= \int \left(-\frac{\sqrt{-y^2+1}}{x\sqrt{y-1}\sqrt{y+1}} \right) dx \\ f(x) &= -\frac{\sqrt{-y^2+1} \ln(x)}{\sqrt{y-1}\sqrt{y+1}} + c_1 \end{aligned}$$

Where c_1 is constant of integration. Substituting result found above for $f(x)$ into equation (3) gives ϕ

$$\phi = \int \frac{1}{\arcsin(y) \sqrt{y-1} \sqrt{y+1}} dy - \frac{\sqrt{-y^2+1} \ln(x)}{\sqrt{y-1} \sqrt{y+1}} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into the constant c_1 gives the solution as

$$c_1 = \int \frac{1}{\arcsin(y) \sqrt{y-1} \sqrt{y+1}} dy - \frac{\sqrt{-y^2+1} \ln(x)}{\sqrt{y-1} \sqrt{y+1}}$$

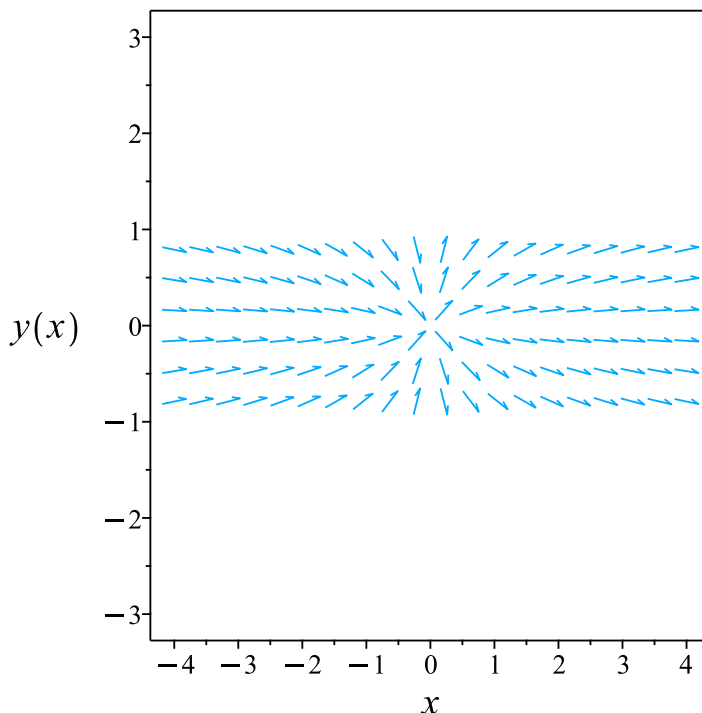


Figure 2.3: Slope field plot

$$y' = \frac{\sqrt{1-y^2} \arcsin(y)}{x}$$

Summary of solutions found

$$\int^y \frac{1}{\arcsin(a) \sqrt{a-1} \sqrt{a+1}} da - \frac{\sqrt{1-y^2} \ln(x)}{\sqrt{y-1} \sqrt{y+1}} = c_1$$

Solved as first order isobaric ode

Time used: 125.872 (sec)

Solving for y' gives

$$y' = \frac{\sqrt{1-y^2} \arcsin(y)}{x} \quad (1)$$

Each of the above ode's is now solved An ode $y' = f(x, y)$ is isobaric if

$$f(tx, t^m y) = t^{m-1} f(x, y) \quad (1)$$

Where here

$$f(x, y) = \frac{\sqrt{1-y^2} \arcsin(y)}{x} \quad (2)$$

 m is the order of isobaric. Substituting (2) into (1) and solving for m gives

$$m = 0$$

Since the ode is isobaric of order $m = 0$, then the substitution

$$\begin{aligned} y &= ux^m \\ &= u \end{aligned}$$

Converts the ODE to a separable in $u(x)$. Performing this substitution gives

$$u'(x) = \frac{\sqrt{1-u(x)^2} \arcsin(u(x))}{x}$$

The ode $u'(x) = \frac{\sqrt{1-u(x)^2} \arcsin(u(x))}{x}$ is separable as it can be written as

$$\begin{aligned} u'(x) &= \frac{\sqrt{1-u(x)^2} \arcsin(u(x))}{x} \\ &= f(x)g(u) \end{aligned}$$

Where

$$\begin{aligned} f(x) &= \frac{1}{x} \\ g(u) &= \sqrt{-u^2 + 1} \arcsin(u) \end{aligned}$$

Integrating gives

$$\int \frac{1}{g(u)} du = \int f(x) dx$$

$$\int \frac{1}{\sqrt{-u^2 + 1} \arcsin(u)} du = \int \frac{1}{x} dx$$

$$\ln(\arcsin(u(x))) = \ln(x) + c_1$$

We now need to find the singular solutions, these are found by finding for what values $g(u)$ is zero, since we had to divide by this above. Solving $g(u) = 0$ or $\sqrt{-u^2 + 1} \arcsin(u) = 0$ for $u(x)$ gives

$$u(x) = -1$$

$$u(x) = 0$$

$$u(x) = 1$$

Now we go over each such singular solution and check if it verifies the ode itself and any initial conditions given. If it does not then the singular solution will not be used.

Therefore the solutions found are

$$\ln(\arcsin(u(x))) = \ln(x) + c_1$$

$$u(x) = -1$$

$$u(x) = 0$$

$$u(x) = 1$$

Converting $\ln(\arcsin(u(x))) = \ln(x) + c_1$ back to y gives

$$\ln(\arcsin(y)) = \ln(x) + c_1$$

Converting $u(x) = -1$ back to y gives

$$y = -1$$

Converting $u(x) = 0$ back to y gives

$$y = 0$$

Converting $u(x) = 1$ back to y gives

$$y = 1$$

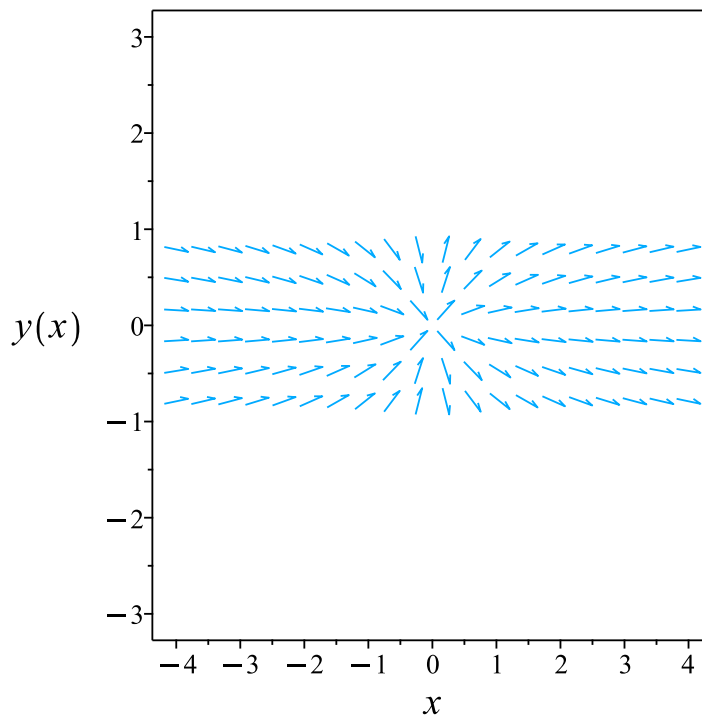


Figure 2.4: Slope field plot

$$y' = \frac{\sqrt{1-y^2} \arcsin(y)}{x}$$

Summary of solutions found

$$\ln(\arcsin(y)) = \ln(x) + c_1$$

$$y = -1$$

$$y = 0$$

$$y = 1$$

Maple step by step solution

Let's solve

$$y' = \frac{\sqrt{1-y^2} \arcsin(y)}{x}$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Solve for the highest derivative

$$y' = \frac{\sqrt{1-y^2} \arcsin(y)}{x}$$

- Separate variables

$$\frac{y'}{\sqrt{1-y^2} \arcsin(y)} = \frac{1}{x}$$

- Integrate both sides with respect to x

$$\int \frac{y'}{\sqrt{1-y^2} \arcsin(y)} dx = \int \frac{1}{x} dx + C1$$

- Evaluate integral

$$\ln(\arcsin(y)) = \ln(x) + C1$$

- Solve for y

$$y = \sin(x e^{C1})$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
<- separable successful`

```

Maple dsolve solution

Solving time : 0.010 (sec)

Leaf size : 8

```

dsolve(diff(y(x),x) = 1/x*(1-y(x)^2)^(1/2)*arcsin(y(x)),
        y(x),singsol=all)

```

$$y = \sin(c_1 x)$$

Mathematica DSolve solution

Solving time : 0.354 (sec)

Leaf size : 27

```
DSolve[{D[y[x],x]==1/x*Sqrt[1-y[x]^2]*ArcSin[y[x]],{}},  
y[x],x,IncludeSingularSolutions->True]
```

$$y(x) \rightarrow \sin(e^{c_1} x)$$

$$y(x) \rightarrow -1$$

$$y(x) \rightarrow 0$$

$$y(x) \rightarrow 1$$

2.1.6 problem 16 (a)

Maple step by step solution	47
Maple trace	47
Maple dsolve solution	49
Mathematica DSolve solution	49

Internal problem ID [18212]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter 1. section 5. Problems at page 19

Problem number : 16 (a)

Date solved : Friday, December 20, 2024 at 05:57:06 AM

CAS classification : [[_2nd_order, _missing_x]]

Solve

$$v'' = \left(\frac{1}{v} + v'^4 \right)^{1/3}$$

Maple step by step solution

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying 2nd order Liouville
trying 2nd order WeierstrassP
trying 2nd order JacobiSN
differential order: 2; trying a linearization to 3rd order
trying 2nd order ODE linearizable_by_differentiation
trying 2nd order, 2 integrating factors of the form mu(x,y)
trying differential order: 2; missing variables
`, `-> Computing symmetries using: way = 3
`, `-> Computing symmetries using: way = exp_sym
-> Calling odsolve with the ODE`, (diff(_b(_a), _a))*_b(_a)-((b(_a)^4*_a+1)/_a)^(1/3)
  Methods for first order ODEs:
  --- Trying classification methods ---
  trying homogeneous types:
  differential order: 1; looking for linear symmetries
  trying exact

```

```

Looking for potential symmetries
trying an equivalence to an Abel ODE
trying 1st order ODE linearizable_by_differentiation
--- Trying Lie symmetry methods, 1st order ---
`, `-> Computing symmetries using: way = 2
`, `-> Computing symmetries using: way = 3
`, `-> Computing symmetries using: way = 4
`, `-> Computing symmetries using: way = 5
trying symmetry patterns for 1st order ODEs
-> trying a symmetry pattern of the form [F(x)*G(y), 0]
-> trying a symmetry pattern of the form [0, F(x)*G(y)]
-> trying symmetry patterns of the forms [F(x),G(y)] and [G(y),F(x)]
-> trying a symmetry pattern of the form [F(x),G(x)]
-> trying a symmetry pattern of the form [F(y),G(y)]
-> trying a symmetry pattern of the form [F(x)+G(y), 0]
-> trying a symmetry pattern of the form [0, F(x)+G(y)]
-> trying a symmetry pattern of the form [F(x),G(x)*y+H(x)]
-> trying a symmetry pattern of conformal type
-> trying 2nd order, dynamical_symmetries, fully reducible to Abel through one integrat
trying 2nd order, integrating factors of the form mu(x,y)/(y)^n, only the singular case
trying differential order: 2; exact nonlinear
trying 2nd order, integrating factor of the form mu(x,y)
-> trying 2nd order, the S-function method
  -> trying a change of variables {x -> y(x), y(x) -> x} and re-entering methods for
  -> trying 2nd order, the S-function method
  -> trying 2nd order, No Point Symmetries Class V
  -> trying 2nd order, No Point Symmetries Class V
  -> trying 2nd order, No Point Symmetries Class V
trying 2nd order, integrating factor of the form mu(x,y)/(y)^n, only the general case
-> trying 2nd order, dynamical_symmetries, only a reduction of order through one integ
--- Trying Lie symmetry methods, 2nd order ---
`, `-> Computing symmetries using: way = 3
`, `-> Computing symmetries using: way = 5
`, `-> Computing symmetries using: way = formal
    *** Sublevel 2 ***
    Methods for first order ODEs:
    --- Trying classification methods ---
    trying a quadrature
    trying 1st order linear

```

```

      <- 1st order linear successful
-> Calling odsolve with the ODE`, diff(y(x), x), y(x)`      *** Sublevel 2 ***
  Methods for first order ODEs:
  --- Trying classification methods ---
  trying a quadrature
  trying 1st order linear
  <- 1st order linear successful`

```

Maple dsolve solution

Solving time : 0.140 (sec)

Leaf size : maple_leaf_size

```

dsolve(diff(diff(v(u),u),u) = (1/v(u)+diff(v(u),u)^4)^(1/3),
        v(u),singsol=all)

```

No solution found

Mathematica DSolve solution

Solving time : 0.0 (sec)

Leaf size : 0

```

DSolve[{D[v[u],{u,2}]==(1/v[u]+D[v[u],u]^4)^(1/3),{}},
        v[u],u,IncludeSingularSolutions->True]

```

Not solved

2.1.7 problem 16 (b)

Solved as first order linear ode	50
Solved as first order Exact ode	52
Maple step by step solution	56
Maple trace	57
Maple dsolve solution	57
Mathematica DSolve solution	58

Internal problem ID [18213]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter 1. section 5. Problems at page 19

Problem number : 16 (b)

Date solved : Friday, December 20, 2024 at 05:57:17 AM

CAS classification : [_linear]

Solve

$$v' + u^2v = \sin(u)$$

Solved as first order linear ode

Time used: 0.414 (sec)

In canonical form a linear first order is

$$v' + q(u)v = p(u)$$

Comparing the above to the given ode shows that

$$q(u) = u^2$$

$$p(u) = \sin(u)$$

The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int q du} \\ &= e^{\int u^2 du} \\ &= e^{\frac{u^3}{3}}\end{aligned}$$

The ode becomes

$$\frac{d}{du}(\mu v) = \mu p$$

$$\frac{d}{du}(\mu v) = (\mu) (\sin(u))$$

$$\frac{d}{du}\left(v e^{\frac{u^3}{3}}\right) = \left(e^{\frac{u^3}{3}}\right) (\sin(u))$$

$$d\left(v e^{\frac{u^3}{3}}\right) = \left(\sin(u) e^{\frac{u^3}{3}}\right) du$$

Integrating gives

$$\begin{aligned} v e^{\frac{u^3}{3}} &= \int \sin(u) e^{\frac{u^3}{3}} du \\ &= \int \sin(u) e^{\frac{u^3}{3}} du + c_1 \end{aligned}$$

Dividing throughout by the integrating factor $e^{\frac{u^3}{3}}$ gives the final solution

$$v = e^{-\frac{u^3}{3}} \left(\int \sin(u) e^{\frac{u^3}{3}} du + c_1 \right)$$

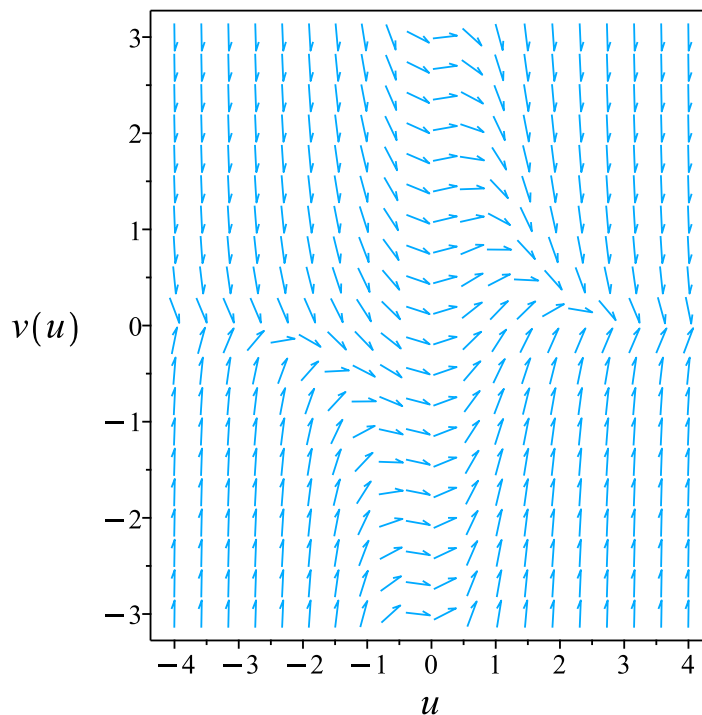


Figure 2.5: Slope field plot
 $v' + u^2v = \sin(u)$

Summary of solutions found

$$v = e^{-\frac{u^3}{3}} \left(\int \sin(u) e^{\frac{u^3}{3}} du + c_1 \right)$$

Solved as first order Exact ode

Time used: 0.348 (sec)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(u, v) du + N(u, v) dv = 0 \quad (\text{1A})$$

Therefore

$$\begin{aligned} dv &= (-v u^2 + \sin(u)) du \\ (v u^2 - \sin(u)) du + dv &= 0 \end{aligned} \quad (\text{2A})$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(u, v) &= v u^2 - \sin(u) \\N(u, v) &= 1\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial v} = \frac{\partial N}{\partial u}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial v} &= \frac{\partial}{\partial v}(v u^2 - \sin(u)) \\&= u^2\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial u} &= \frac{\partial}{\partial u}(1) \\&= 0\end{aligned}$$

Since $\frac{\partial M}{\partial v} \neq \frac{\partial N}{\partial u}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned}A &= \frac{1}{N} \left(\frac{\partial M}{\partial v} - \frac{\partial N}{\partial u} \right) \\&= 1((u^2) - (0)) \\&= u^2\end{aligned}$$

Since A does not depend on v , then it can be used to find an integrating factor. The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int A \, du} \\&= e^{\int u^2 \, du}\end{aligned}$$

The result of integrating gives

$$\begin{aligned}\mu &= e^{\frac{u^3}{3}} \\&= e^{\frac{u^3}{3}}\end{aligned}$$

M and N are multiplied by this integrating factor, giving new M and new N which are called \overline{M} and \overline{N} for now so not to confuse them with the original M and N .

$$\begin{aligned}\overline{M} &= \mu M \\ &= e^{\frac{u^3}{3}} (v u^2 - \sin(u)) \\ &= (v u^2 - \sin(u)) e^{\frac{u^3}{3}}\end{aligned}$$

And

$$\begin{aligned}\overline{N} &= \mu N \\ &= e^{\frac{u^3}{3}} (1) \\ &= e^{\frac{u^3}{3}}\end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned}\overline{M} + \overline{N} \frac{dv}{du} &= 0 \\ \left((v u^2 - \sin(u)) e^{\frac{u^3}{3}} \right) + \left(e^{\frac{u^3}{3}} \right) \frac{dv}{du} &= 0\end{aligned}$$

The following equations are now set up to solve for the function $\phi(u, v)$

$$\frac{\partial \phi}{\partial u} = \overline{M} \tag{1}$$

$$\frac{\partial \phi}{\partial v} = \overline{N} \tag{2}$$

Integrating (2) w.r.t. v gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial v} dv &= \int \overline{N} dv \\ \int \frac{\partial \phi}{\partial v} dv &= \int e^{\frac{u^3}{3}} dv \\ \phi &= v e^{\frac{u^3}{3}} + f(u)\end{aligned} \tag{3}$$

Where $f(u)$ is used for the constant of integration since ϕ is a function of both u and v . Taking derivative of equation (3) w.r.t u gives

$$\frac{\partial \phi}{\partial u} = v u^2 e^{\frac{u^3}{3}} + f'(u) \tag{4}$$

But equation (1) says that $\frac{\partial \phi}{\partial u} = (v u^2 - \sin(u)) e^{\frac{u^3}{3}}$. Therefore equation (4) becomes

$$(v u^2 - \sin(u)) e^{\frac{u^3}{3}} = v u^2 e^{\frac{u^3}{3}} + f'(u) \quad (5)$$

Solving equation (5) for $f'(u)$ gives

$$f'(u) = -\sin(u) e^{\frac{u^3}{3}}$$

Integrating the above w.r.t u gives

$$\begin{aligned} \int f'(u) du &= \int \left(-\sin(u) e^{\frac{u^3}{3}} \right) du \\ f(u) &= \int_0^u -\sin(\tau) e^{\frac{\tau^3}{3}} d\tau + c_1 \end{aligned}$$

Where c_1 is constant of integration. Substituting result found above for $f(u)$ into equation (3) gives ϕ

$$\phi = v e^{\frac{u^3}{3}} + \int_0^u -\sin(\tau) e^{\frac{\tau^3}{3}} d\tau + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into the constant c_1 gives the solution as

$$c_1 = v e^{\frac{u^3}{3}} + \int_0^u -\sin(\tau) e^{\frac{\tau^3}{3}} d\tau$$

Solving for v gives

$$v = -\left(\int_0^u -\sin(\tau) e^{\frac{\tau^3}{3}} d\tau - c_1 \right) e^{-\frac{u^3}{3}}$$

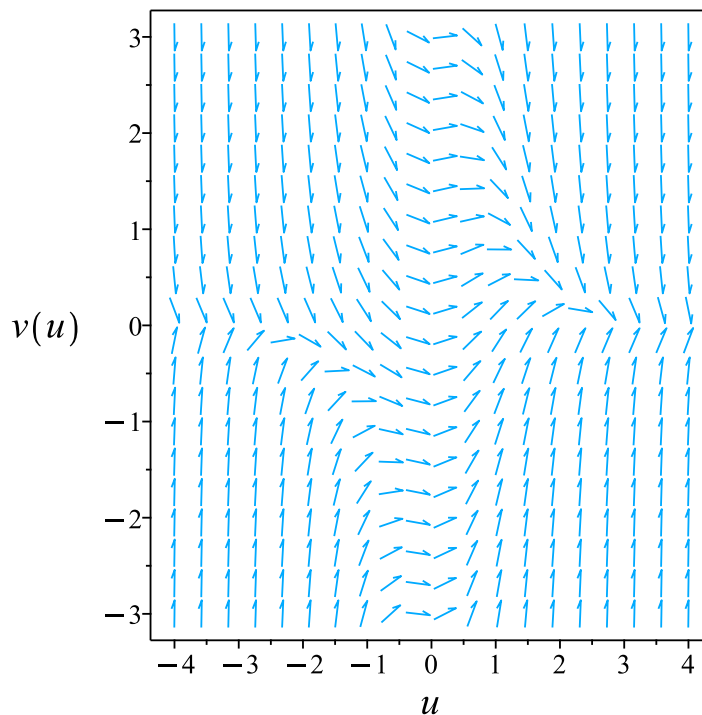


Figure 2.6: Slope field plot
 $v' + u^2v = \sin(u)$

Summary of solutions found

$$v = -\left(\int_0^u -\sin(\tau) e^{\frac{\tau^3}{3}} d\tau - c_1\right) e^{-\frac{u^3}{3}}$$

Maple step by step solution

Let's solve

$$v' + u^2v = \sin(u)$$

- Highest derivative means the order of the ODE is 1
 v'
- Solve for the highest derivative
 $v' = -u^2v + \sin(u)$
- Group terms with v on the lhs of the ODE and the rest on the rhs of the ODE
 $v' + u^2v = \sin(u)$
- The ODE is linear; multiply by an integrating factor $\mu(u)$
 $\mu(u)(v' + u^2v) = \mu(u)\sin(u)$

- Assume the lhs of the ODE is the total derivative $\frac{d}{du}(v\mu(u))$
 $\mu(u)(v' + u^2v) = v'\mu(u) + v\mu'(u)$
- Isolate $\mu'(u)$
 $\mu'(u) = \mu(u)u^2$
- Solve to find the integrating factor
 $\mu(u) = e^{\frac{u^3}{3}}$
- Integrate both sides with respect to u
 $\int \left(\frac{d}{du}(v\mu(u))\right) du = \int \mu(u) \sin(u) du + C1$
- Evaluate the integral on the lhs
 $v\mu(u) = \int \mu(u) \sin(u) du + C1$
- Solve for v
 $v = \frac{\int \mu(u) \sin(u) du + C1}{\mu(u)}$
- Substitute $\mu(u) = e^{\frac{u^3}{3}}$
 $v = \frac{\int \sin(u) e^{\frac{u^3}{3}} du + C1}{e^{\frac{u^3}{3}}}$
- Simplify
 $v = e^{-\frac{u^3}{3}} \left(\int \sin(u) e^{\frac{u^3}{3}} du + C1 \right)$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`

```

Maple dsolve solution

Solving time : 0.002 (sec)
 Leaf size : 24

```

dsolve(diff(v(u),u)+u^2*v(u) = sin(u),
        v(u),singsol=all)

```

$$v = e^{-\frac{u^3}{3}} \left(\int \sin(u) e^{\frac{u^3}{3}} du + c_1 \right)$$

Mathematica DSolve solution

Solving time : 1.978 (sec)

Leaf size : 39

```
DSolve[{D[v[u],u]+u^2*v[u]==Sin[u],{}}],  
v[u],u,IncludeSingularSolutions->True]
```

$$v(u) \rightarrow e^{-\frac{u^3}{3}} \left(\int_1^u e^{\frac{K[1]^3}{3}} \sin(K[1]) dK[1] + c_1 \right)$$

2.1.8 problem 17 (a)

Maple step by step solution	59
Maple trace	59
Maple dsolve solution	60
Mathematica DSolve solution	61

Internal problem ID [18214]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter 1. section 5. Problems at page 19

Problem number : 17 (a)

Date solved : Friday, December 20, 2024 at 05:57:20 AM

CAS classification : [NONE]

Solve

$$\sqrt{y' + y} = (y'' + 2x)^{1/4}$$

Maple step by step solution

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying 2nd order Liouville
trying 2nd order WeierstrassP
trying 2nd order JacobiSN
differential order: 2; trying a linearization to 3rd order
trying 2nd order ODE linearizable_by_differentiation
trying 2nd order, 2 integrating factors of the form mu(x,y)
trying differential order: 2; missing variables
-> trying 2nd order, dynamical_symmetries, fully reducible to Abel through one integrat
    --- trying a change of variables {x -> y(x), y(x) -> x} and re-entering methods for
    -> trying 2nd order, dynamical_symmetries, fully reducible to Abel through one inte
trying 2nd order, integrating factors of the form mu(x,y)/(y)^n, only the singular cas
trying symmetries linear in x and y(x)
trying differential order: 2; exact nonlinear
trying 2nd order, integrating factor of the form mu(y)
trying 2nd order, integrating factor of the form mu(x,y)
trying 2nd order, integrating factor of the form mu(x,y)/(y)^n, only the general case

```

```

trying 2nd order, integrating factor of the form mu(y,y)
trying differential order: 2; mu polynomial in y
trying 2nd order, integrating factor of the form mu(x,y)
differential order: 2; looking for linear symmetries
-> trying 2nd order, the S-function method
  -> trying a change of variables {x -> y(x), y(x) -> x} and re-entering methods for
  -> trying 2nd order, the S-function method
  -> trying 2nd order, No Point Symmetries Class V
    --- trying a change of variables {x -> y(x), y(x) -> x} and re-entering methods
    -> trying 2nd order, No Point Symmetries Class V
  -> trying 2nd order, No Point Symmetries Class V
    --- trying a change of variables {x -> y(x), y(x) -> x} and re-entering methods
    -> trying 2nd order, No Point Symmetries Class V
  -> trying 2nd order, No Point Symmetries Class V
    --- trying a change of variables {x -> y(x), y(x) -> x} and re-entering methods
    -> trying 2nd order, No Point Symmetries Class V
trying 2nd order, integrating factor of the form mu(x,y)/(y)^n, only the general case
-> trying 2nd order, dynamical_symmetries, only a reduction of order through one integ
  --- trying a change of variables {x -> y(x), y(x) -> x} and re-entering methods for
  -> trying 2nd order, dynamical_symmetries, only a reduction of order through one in
solving 2nd order ODE of high degree, Lie methods
`, `-> Computing symmetries using: way = 3
`, `-> Computing symmetries using: way = 5`

```

Maple dsolve solution

Solving time : 0.307 (sec)

Leaf size : maple_leaf_size

```

dsolve((diff(y(x),x)+y(x))^(1/2) = (diff(diff(y(x),x),x)+2*x)^(1/4),
        y(x),singsol=all)

```

No solution found

Mathematica DSolve solution

Solving time : 0.0 (sec)

Leaf size : 0

```
DSolve[{Sqrt[D[y[x],x]+y[x]]== (D[y[x],{x,2}]+2*x)^(1/4),{}}],  
y[x],x,IncludeSingularSolutions->True]
```

Not solved

2.1.9 problem 18

Solved as first order linear ode	62
Solved as first order homogeneous class A ode	64
Solved as first order homogeneous class D2 ode	67
Solved as first order homogeneous class Maple C ode	69
Solved as first order Exact ode	73
Solved as first order isobaric ode	77
Solved using Lie symmetry for first order ode	80
Maple step by step solution	85
Maple trace	86
Maple dsolve solution	86
Mathematica DSolve solution	86

Internal problem ID [18215]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter 1. section 5. Problems at page 19

Problem number : 18

Date solved : Friday, December 20, 2024 at 06:02:20 AM

CAS classification : [_linear]

Solve

$$v' + \frac{2v}{u} = 3$$

Solved as first order linear ode

Time used: 0.183 (sec)

In canonical form a linear first order is

$$v' + q(u)v = p(u)$$

Comparing the above to the given ode shows that

$$q(u) = \frac{2}{u}$$

$$p(u) = 3$$

The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int q \, du} \\ &= e^{\int \frac{2}{u} \, du} \\ &= u^2\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{du}(\mu v) &= \mu p \\ \frac{d}{du}(\mu v) &= (\mu) (3) \\ \frac{d}{du}(v u^2) &= (u^2) (3) \\ d(v u^2) &= (3u^2) \, du\end{aligned}$$

Integrating gives

$$\begin{aligned}v u^2 &= \int 3u^2 \, du \\ &= u^3 + c_1\end{aligned}$$

Dividing throughout by the integrating factor u^2 gives the final solution

$$v = \frac{u^3 + c_1}{u^2}$$

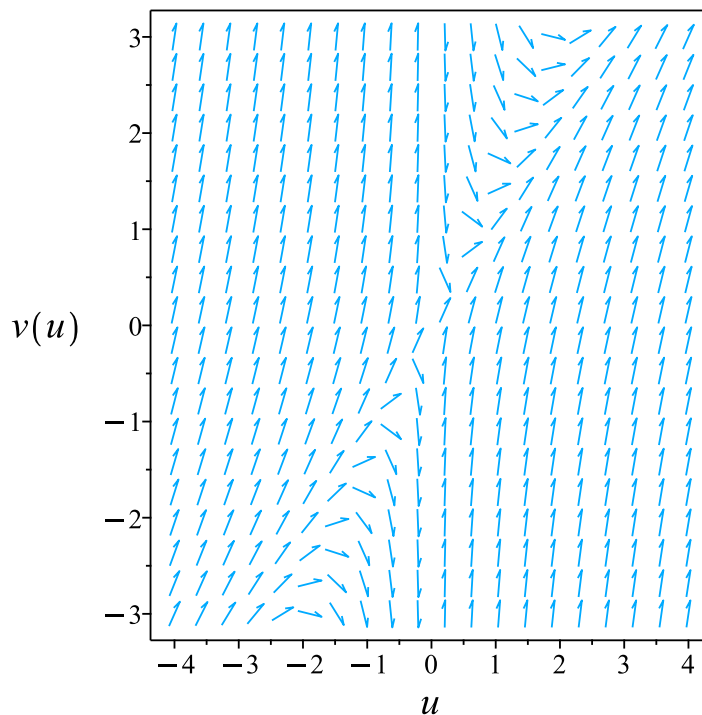


Figure 2.7: Slope field plot
 $v' + \frac{2v}{u} = 3$

Summary of solutions found

$$v = \frac{u^3 + c_1}{u^2}$$

Solved as first order homogeneous class A ode

Time used: 0.300 (sec)

In canonical form, the ODE is

$$\begin{aligned} v' &= F(u, v) \\ &= -\frac{-3u + 2v}{u} \end{aligned} \quad (1)$$

An ode of the form $v' = \frac{M(u, v)}{N(u, v)}$ is called homogeneous if the functions $M(u, v)$ and $N(u, v)$ are both homogeneous functions and of the same order. Recall that a function $f(u, v)$ is homogeneous of order n if

$$f(t^n u, t^n v) = t^n f(u, v)$$

In this case, it can be seen that both $M = 3u - 2v$ and $N = u$ are both homogeneous and of the same order $n = 1$. Therefore this is a homogeneous ode. Since this ode is

homogeneous, it is converted to separable ODE using the substitution $u = \frac{v}{u}$, or $v = uu$. Hence

$$\frac{dv}{du} = \frac{du}{du}u + u$$

Applying the transformation $v = uu$ to the above ODE in (1) gives

$$\begin{aligned}\frac{du}{du}u + u &= 3 - 2u \\ \frac{du}{du} &= \frac{3 - 3u(u)}{u}\end{aligned}$$

Or

$$u'(u) - \frac{3 - 3u(u)}{u} = 0$$

Or

$$u'(u)u + 3u(u) - 3 = 0$$

Which is now solved as separable in $u(u)$.

The ode $u'(u) = -\frac{3(u(u)-1)}{u}$ is separable as it can be written as

$$\begin{aligned}u'(u) &= -\frac{3(u(u) - 1)}{u} \\ &= f(u)g(u)\end{aligned}$$

Where

$$\begin{aligned}f(u) &= \frac{1}{u} \\ g(u) &= -3u + 3\end{aligned}$$

Integrating gives

$$\begin{aligned}\int \frac{1}{g(u)} du &= \int f(u) du \\ \int \frac{1}{-3u + 3} du &= \int \frac{1}{u} du \\ -\frac{\ln(u(u) - 1)}{3} &= \ln(u) + c_1\end{aligned}$$

We now need to find the singular solutions, these are found by finding for what values $g(u)$ is zero, since we had to divide by this above. Solving $g(u) = 0$ or $-3u + 3 = 0$ for $u(u)$ gives

$$u(u) = 1$$

Now we go over each such singular solution and check if it verifies the ode itself and any initial conditions given. If it does not then the singular solution will not be used.

Therefore the solutions found are

$$-\frac{\ln(u(u) - 1)}{3} = \ln(u) + c_1$$

$$u(u) = 1$$

Solving for $u(u)$ gives

$$u(u) = 1$$

$$u(u) = \frac{u^3 + e^{-3c_1}}{u^3}$$

Converting $u(u) = 1$ back to v gives

$$v = u$$

Converting $u(u) = \frac{u^3 + e^{-3c_1}}{u^3}$ back to v gives

$$v = \frac{u^3 + e^{-3c_1}}{u^2}$$

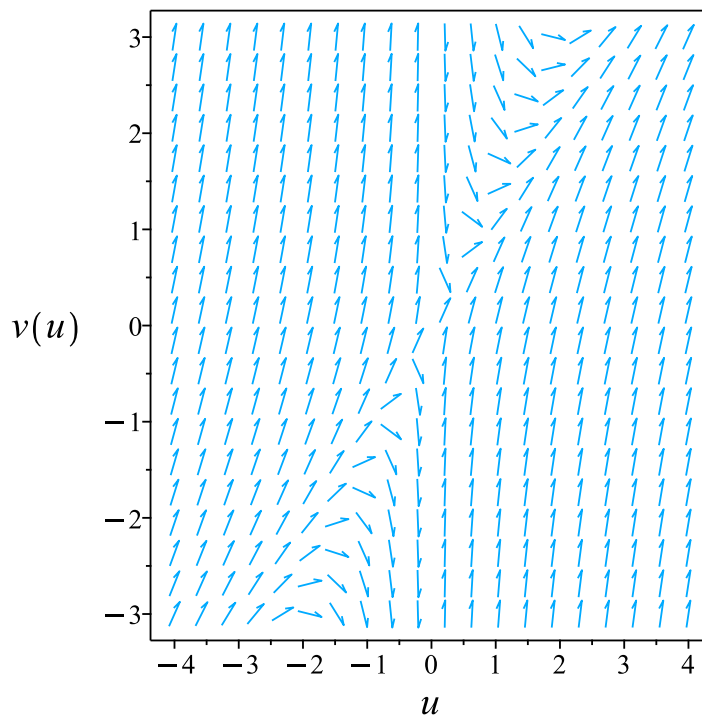


Figure 2.8: Slope field plot
 $v' + \frac{2v}{u} = 3$

Summary of solutions found

$$v = u$$

$$v = \frac{u^3 + e^{-3c_1}}{u^2}$$

Solved as first order homogeneous class D2 ode

Time used: 0.149 (sec)

Applying change of variables $v = u(u)$, then the ode becomes

$$u'(u)u + 3u(u) = 3$$

Which is now solved The ode $u'(u) = -\frac{3(u(u)-1)}{u}$ is separable as it can be written as

$$u'(u) = -\frac{3(u(u) - 1)}{u}$$

$$= f(u)g(u)$$

Where

$$f(u) = \frac{1}{u}$$

$$g(u) = -3u + 3$$

Integrating gives

$$\int \frac{1}{g(u)} du = \int f(u) du$$

$$\int \frac{1}{-3u + 3} du = \int \frac{1}{u} du$$

$$-\frac{\ln(u(u) - 1)}{3} = \ln(u) + c_1$$

We now need to find the singular solutions, these are found by finding for what values $g(u)$ is zero, since we had to divide by this above. Solving $g(u) = 0$ or $-3u + 3 = 0$ for $u(u)$ gives

$$u(u) = 1$$

Now we go over each such singular solution and check if it verifies the ode itself and any initial conditions given. If it does not then the singular solution will not be used.

Therefore the solutions found are

$$-\frac{\ln(u(u) - 1)}{3} = \ln(u) + c_1$$

$$u(u) = 1$$

Solving for $u(u)$ gives

$$u(u) = 1$$

$$u(u) = \frac{u^3 + e^{-3c_1}}{u^3}$$

Converting $u(u) = 1$ back to v gives

$$v = u$$

Converting $u(u) = \frac{u^3 + e^{-3c_1}}{u^3}$ back to v gives

$$v = \frac{u^3 + e^{-3c_1}}{u^2}$$

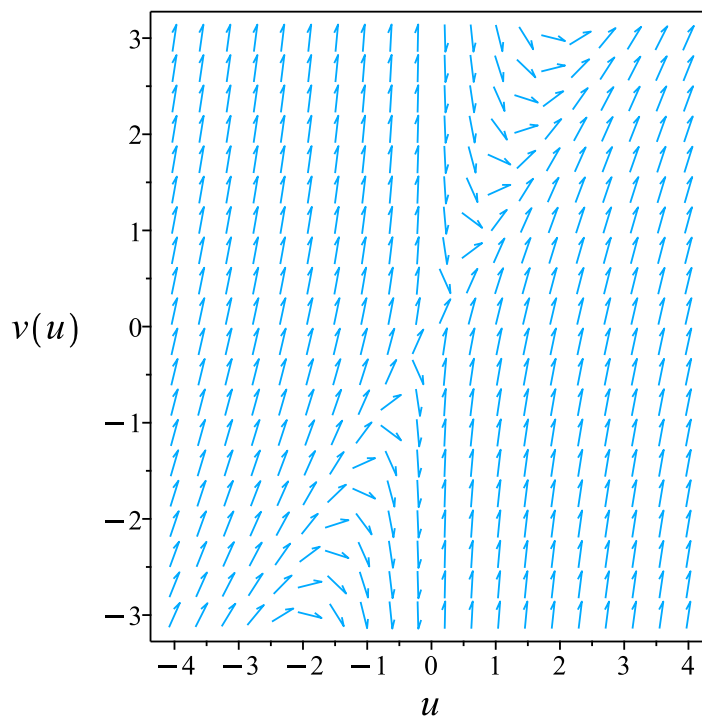


Figure 2.9: Slope field plot

$$v' + \frac{2v}{u} = 3$$

Summary of solutions found

$$v = u$$

$$v = \frac{u^3 + e^{-3c_1}}{u^2}$$

Summary of solutions found

$$v = u$$

$$v = \frac{u^3 + e^{-3c_1}}{u^2}$$

Solved as first order homogeneous class Maple C ode

Time used: 0.302 (sec)

Let $Y = v - y_0$ and $X = u - x_0$ then the above is transformed to new ode in $Y(X)$

$$\frac{d}{dX}Y(X) = -\frac{-3x_0 - 3X + 2Y(X) + 2y_0}{x_0 + X}$$

Solving for possible values of x_0 and y_0 which makes the above ode a homogeneous ode results in

$$x_0 = 0$$

$$y_0 = 0$$

Using these values now it is possible to easily solve for $Y(X)$. The above ode now becomes

$$\frac{d}{dX}Y(X) = -\frac{-3X + 2Y(X)}{X}$$

In canonical form, the ODE is

$$Y' = F(X, Y)$$

$$= -\frac{-3X + 2Y}{X} \tag{1}$$

An ode of the form $Y' = \frac{M(X, Y)}{N(X, Y)}$ is called homogeneous if the functions $M(X, Y)$ and $N(X, Y)$ are both homogeneous functions and of the same order. Recall that a function $f(X, Y)$ is homogeneous of order n if

$$f(t^n X, t^n Y) = t^n f(X, Y)$$

In this case, it can be seen that both $M = 3X - 2Y$ and $N = X$ are both homogeneous and of the same order $n = 1$. Therefore this is a homogeneous ode. Since this ode is homogeneous, it is converted to separable ODE using the substitution $u = \frac{Y}{X}$, or $Y = uX$. Hence

$$\frac{dY}{dX} = \frac{du}{dX}X + u$$

Applying the transformation $Y = uX$ to the above ODE in (1) gives

$$\begin{aligned}\frac{du}{dX}X + u &= 3 - 2u \\ \frac{du}{dX} &= \frac{3 - 3u(X)}{X}\end{aligned}$$

Or

$$\frac{d}{dX}u(X) - \frac{3 - 3u(X)}{X} = 0$$

Or

$$\left(\frac{d}{dX}u(X)\right)X + 3u(X) - 3 = 0$$

Which is now solved as separable in $u(X)$.

The ode $\frac{d}{dX}u(X) = -\frac{3(u(X)-1)}{X}$ is separable as it can be written as

$$\begin{aligned}\frac{d}{dX}u(X) &= -\frac{3(u(X) - 1)}{X} \\ &= f(X)g(u)\end{aligned}$$

Where

$$\begin{aligned}f(X) &= \frac{1}{X} \\ g(u) &= -3u + 3\end{aligned}$$

Integrating gives

$$\begin{aligned}\int \frac{1}{g(u)} du &= \int f(X) dX \\ \int \frac{1}{-3u + 3} du &= \int \frac{1}{X} dX \\ -\frac{\ln(u(X) - 1)}{3} &= \ln(X) + c_1\end{aligned}$$

We now need to find the singular solutions, these are found by finding for what values $g(u)$ is zero, since we had to divide by this above. Solving $g(u) = 0$ or $-3u + 3 = 0$ for $u(X)$ gives

$$u(X) = 1$$

Now we go over each such singular solution and check if it verifies the ode itself and any initial conditions given. If it does not then the singular solution will not be used.

Therefore the solutions found are

$$-\frac{\ln(u(X) - 1)}{3} = \ln(X) + c_1$$

$$u(X) = 1$$

Solving for $u(X)$ gives

$$u(X) = 1$$

$$u(X) = \frac{X^3 + e^{-3c_1}}{X^3}$$

Converting $u(X) = 1$ back to $Y(X)$ gives

$$Y(X) = X$$

Converting $u(X) = \frac{X^3 + e^{-3c_1}}{X^3}$ back to $Y(X)$ gives

$$Y(X) = \frac{X^3 + e^{-3c_1}}{X^2}$$

Using the solution for $Y(X)$

$$Y(X) = X \tag{A}$$

And replacing back terms in the above solution using

$$Y = v + y_0$$

$$X = u + x_0$$

Or

$$Y = v$$

$$X = u$$

Then the solution in v becomes using EQ (A)

$$v = u$$

Using the solution for $Y(X)$

$$Y(X) = \frac{X^3 + e^{-3c_1}}{X^2} \tag{A}$$

And replacing back terms in the above solution using

$$Y = v + y_0$$

$$X = u + x_0$$

Or

$$Y = v$$

$$X = u$$

Then the solution in v becomes using EQ (A)

$$v = \frac{u^3 + e^{-3c_1}}{u^2}$$

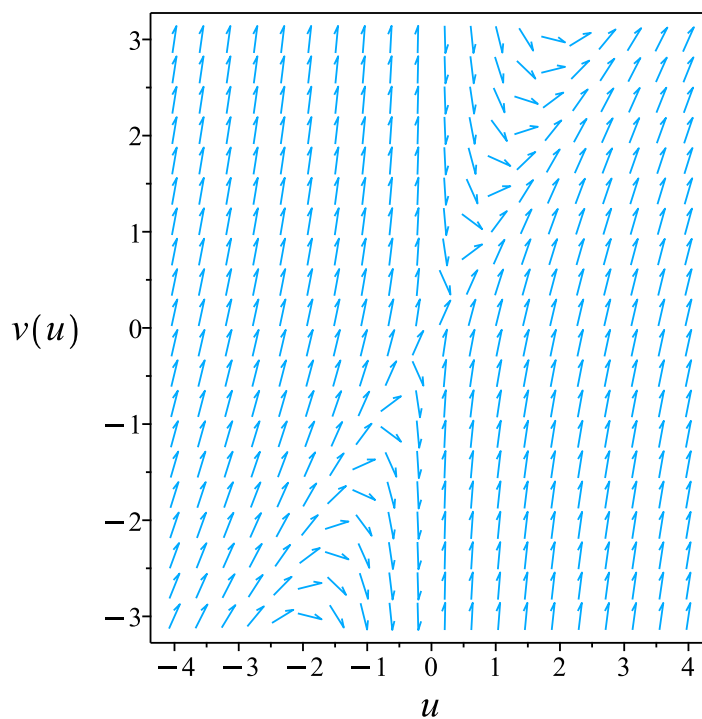


Figure 2.10: Slope field plot

$$v' + \frac{2v}{u} = 3$$

Solved as first order Exact ode

Time used: 0.165 (sec)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx}\phi(x, y) = 0$$

Hence

$$\frac{\partial\phi}{\partial x} + \frac{\partial\phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial\phi}{\partial x} &= M \\ \frac{\partial\phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(u, v) du + N(u, v) dv = 0 \quad (\text{1A})$$

Therefore

$$\begin{aligned} dv &= \left(3 - \frac{2v}{u}\right) du \\ \left(\frac{2v}{u} - 3\right) du + dv &= 0 \end{aligned} \quad (\text{2A})$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(u, v) &= \frac{2v}{u} - 3 \\N(u, v) &= 1\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial v} = \frac{\partial N}{\partial u}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial v} &= \frac{\partial}{\partial v} \left(\frac{2v}{u} - 3 \right) \\&= \frac{2}{u}\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial u} &= \frac{\partial}{\partial u} (1) \\&= 0\end{aligned}$$

Since $\frac{\partial M}{\partial v} \neq \frac{\partial N}{\partial u}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned}A &= \frac{1}{N} \left(\frac{\partial M}{\partial v} - \frac{\partial N}{\partial u} \right) \\&= 1 \left(\left(\frac{2}{u} \right) - (0) \right) \\&= \frac{2}{u}\end{aligned}$$

Since A does not depend on v , then it can be used to find an integrating factor. The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int A \, du} \\&= e^{\int \frac{2}{u} \, du}\end{aligned}$$

The result of integrating gives

$$\begin{aligned}\mu &= e^{2\ln(u)} \\&= u^2\end{aligned}$$

M and N are multiplied by this integrating factor, giving new M and new N which are called \bar{M} and \bar{N} for now so not to confuse them with the original M and N .

$$\begin{aligned}\bar{M} &= \mu M \\ &= u^2 \left(\frac{2v}{u} - 3 \right) \\ &= -3u^2 + 2uv\end{aligned}$$

And

$$\begin{aligned}\bar{N} &= \mu N \\ &= u^2(1) \\ &= u^2\end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned}\bar{M} + \bar{N} \frac{dv}{du} &= 0 \\ (-3u^2 + 2uv) + (u^2) \frac{dv}{du} &= 0\end{aligned}$$

The following equations are now set up to solve for the function $\phi(u, v)$

$$\frac{\partial \phi}{\partial u} = \bar{M} \tag{1}$$

$$\frac{\partial \phi}{\partial v} = \bar{N} \tag{2}$$

Integrating (2) w.r.t. v gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial v} dv &= \int \bar{N} dv \\ \int \frac{\partial \phi}{\partial v} dv &= \int u^2 dv \\ \phi &= v u^2 + f(u)\end{aligned} \tag{3}$$

Where $f(u)$ is used for the constant of integration since ϕ is a function of both u and v . Taking derivative of equation (3) w.r.t u gives

$$\frac{\partial \phi}{\partial u} = 2uv + f'(u) \tag{4}$$

But equation (1) says that $\frac{\partial \phi}{\partial u} = -3u^2 + 2uv$. Therefore equation (4) becomes

$$-3u^2 + 2uv = 2uv + f'(u) \quad (5)$$

Solving equation (5) for $f'(u)$ gives

$$f'(u) = -3u^2$$

Integrating the above w.r.t u gives

$$\begin{aligned} \int f'(u) \, du &= \int (-3u^2) \, du \\ f(u) &= -u^3 + c_1 \end{aligned}$$

Where c_1 is constant of integration. Substituting result found above for $f(u)$ into equation (3) gives ϕ

$$\phi = -u^3 + v u^2 + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into the constant c_1 gives the solution as

$$c_1 = -u^3 + v u^2$$

Solving for v gives

$$v = \frac{u^3 + c_1}{u^2}$$

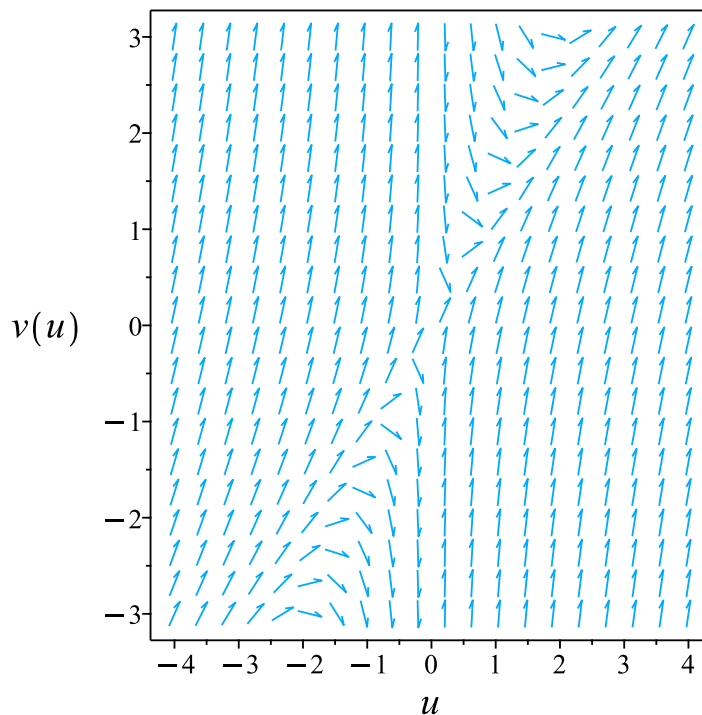


Figure 2.11: Slope field plot
 $v' + \frac{2v}{u} = 3$

Summary of solutions found

$$v = \frac{u^3 + c_1}{u^2}$$

Solved as first order isobaric ode

Time used: 0.195 (sec)

Solving for v' gives

$$v' = -\frac{-3u + 2v}{u} \quad (1)$$

Each of the above ode's is now solved An ode $v' = f(u, v)$ is isobaric if

$$f(tu, t^m v) = t^{m-1} f(u, v) \quad (1)$$

Where here

$$f(u, v) = -\frac{-3u + 2v}{u} \quad (2)$$

m is the order of isobaric. Substituting (2) into (1) and solving for m gives

$$m = 1$$

Since the ode is isobaric of order $m = 1$, then the substitution

$$\begin{aligned} v &= uu^m \\ &= uu \end{aligned}$$

Converts the ODE to a separable in $u(u)$. Performing this substitution gives

$$u(u) + uu'(u) = -\frac{-3u + 2uu(u)}{u}$$

The ode $u'(u) = -\frac{3(u(u)-1)}{u}$ is separable as it can be written as

$$\begin{aligned} u'(u) &= -\frac{3(u(u) - 1)}{u} \\ &= f(u)g(u) \end{aligned}$$

Where

$$\begin{aligned} f(u) &= \frac{1}{u} \\ g(u) &= -3u + 3 \end{aligned}$$

Integrating gives

$$\begin{aligned} \int \frac{1}{g(u)} du &= \int f(u) du \\ \int \frac{1}{-3u + 3} du &= \int \frac{1}{u} du \\ -\frac{\ln(u(u) - 1)}{3} &= \ln(u) + c_1 \end{aligned}$$

We now need to find the singular solutions, these are found by finding for what values $g(u)$ is zero, since we had to divide by this above. Solving $g(u) = 0$ or $-3u + 3 = 0$ for $u(u)$ gives

$$u(u) = 1$$

Now we go over each such singular solution and check if it verifies the ode itself and any initial conditions given. If it does not then the singular solution will not be used.

Therefore the solutions found are

$$\begin{aligned} -\frac{\ln(u(u) - 1)}{3} &= \ln(u) + c_1 \\ u(u) &= 1 \end{aligned}$$

Solving for $u(u)$ gives

$$u(u) = 1$$

$$u(u) = \frac{u^3 + e^{-3c_1}}{u^3}$$

Converting $u(u) = 1$ back to v gives

$$\frac{v}{u} = 1$$

Converting $u(u) = \frac{u^3 + e^{-3c_1}}{u^3}$ back to v gives

$$\frac{v}{u} = \frac{u^3 + e^{-3c_1}}{u^3}$$

Solving for v gives

$$v = u$$

$$v = \frac{u^3 + e^{-3c_1}}{u^2}$$

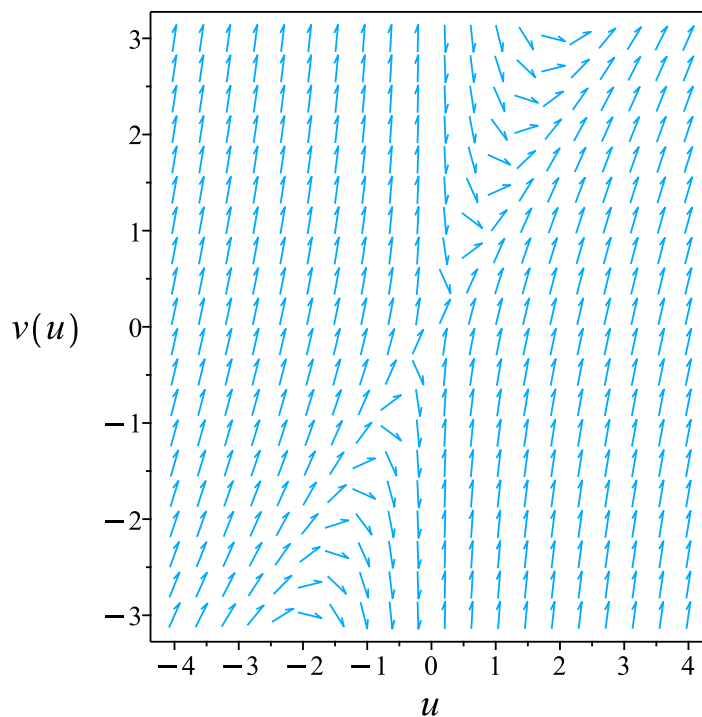


Figure 2.12: Slope field plot

$$v' + \frac{2v}{u} = 3$$

Summary of solutions found

$$v = u$$

$$v = \frac{u^3 + e^{-3c_1}}{u^2}$$

Solved using Lie symmetry for first order ode

Time used: 0.382 (sec)

Writing the ode as

$$v' = -\frac{-3u + 2v}{u}$$

$$v' = \omega(u, v)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_u + \omega(\eta_v - \xi_u) - \omega^2 \xi_v - \omega_u \xi - \omega_v \eta = 0 \quad (\text{A})$$

To determine ξ, η then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$\xi = ua_2 + va_3 + a_1 \quad (\text{1E})$$

$$\eta = ub_2 + vb_3 + b_1 \quad (\text{2E})$$

Where the unknown coefficients are

$$\{a_1, a_2, a_3, b_1, b_2, b_3\}$$

Substituting equations (1E,2E) and ω into (A) gives

$$b_2 - \frac{(-3u + 2v)(b_3 - a_2)}{u} - \frac{(-3u + 2v)^2 a_3}{u^2} \quad (\text{5E})$$

$$- \left(\frac{3}{u} + \frac{-3u + 2v}{u^2} \right) (ua_2 + va_3 + a_1) + \frac{2ub_2 + 2vb_3 + 2b_1}{u} = 0$$

Putting the above in normal form gives

$$-\frac{3u^2 a_2 + 9u^2 a_3 - 3b_2 u^2 - 3u^2 b_3 - 12uva_3 + 6v^2 a_3 - 2ub_1 + 2va_1}{u^2} = 0$$

Setting the numerator to zero gives

$$-3u^2a_2 - 9u^2a_3 + 3b_2u^2 + 3u^2b_3 + 12uva_3 - 6v^2a_3 + 2ub_1 - 2va_1 = 0 \quad (6E)$$

Looking at the above PDE shows the following are all the terms with $\{u, v\}$ in them.

$$\{u, v\}$$

The following substitution is now made to be able to collect on all terms with $\{u, v\}$ in them

$$\{u = v_1, v = v_2\}$$

The above PDE (6E) now becomes

$$-3a_2v_1^2 - 9a_3v_1^2 + 12a_3v_1v_2 - 6a_3v_2^2 + 3b_2v_1^2 + 3b_3v_1^2 - 2a_1v_2 + 2b_1v_1 = 0 \quad (7E)$$

Collecting the above on the terms v_i introduced, and these are

$$\{v_1, v_2\}$$

Equation (7E) now becomes

$$(-3a_2 - 9a_3 + 3b_2 + 3b_3)v_1^2 + 12a_3v_1v_2 + 2b_1v_1 - 6a_3v_2^2 - 2a_1v_2 = 0 \quad (8E)$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$\begin{aligned} -2a_1 &= 0 \\ -6a_3 &= 0 \\ 12a_3 &= 0 \\ 2b_1 &= 0 \\ -3a_2 - 9a_3 + 3b_2 + 3b_3 &= 0 \end{aligned}$$

Solving the above equations for the unknowns gives

$$\begin{aligned} a_1 &= 0 \\ a_2 &= b_2 + b_3 \\ a_3 &= 0 \\ b_1 &= 0 \\ b_2 &= b_2 \\ b_3 &= b_3 \end{aligned}$$

Substituting the above solution in the anstaz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$\begin{aligned}\xi &= u \\ \eta &= u\end{aligned}$$

Shifting is now applied to make $\xi = 0$ in order to simplify the rest of the computation

$$\begin{aligned}\eta &= \eta - \omega(u, v) \xi \\ &= u - \left(-\frac{-3u + 2v}{u} \right) (u) \\ &= -2u + 2v \\ \xi &= 0\end{aligned}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(u, v) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{du}{\xi} = \frac{dv}{\eta} = dS \quad (1)$$

The above comes from the requirements that $(\xi \frac{\partial}{\partial u} + \eta \frac{\partial}{\partial v}) S(u, v) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = u$$

S is found from

$$\begin{aligned}S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{-2u + 2v} dy\end{aligned}$$

Which results in

$$S = \frac{\ln(-u + v)}{2}$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_u + \omega(u, v)S_v}{R_u + \omega(u, v)R_v} \quad (2)$$

Where in the above R_u, R_v, S_u, S_v are all partial derivatives and $\omega(u, v)$ is the right hand side of the original ode given by

$$\omega(u, v) = -\frac{-3u + 2v}{u}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_u &= 1 \\ R_v &= 0 \\ S_u &= \frac{1}{2u - 2v} \\ S_v &= -\frac{1}{2u - 2v} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = -\frac{1}{u} \tag{2A}$$

We now need to express the RHS as function of R only. This is done by solving for u, v in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = -\frac{1}{R}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S .

Since the ode has the form $\frac{d}{dR}S(R) = f(R)$, then we only need to integrate $f(R)$.

$$\begin{aligned} \int dS &= \int -\frac{1}{R} dR \\ S(R) &= -\ln(R) + c_2 \end{aligned}$$

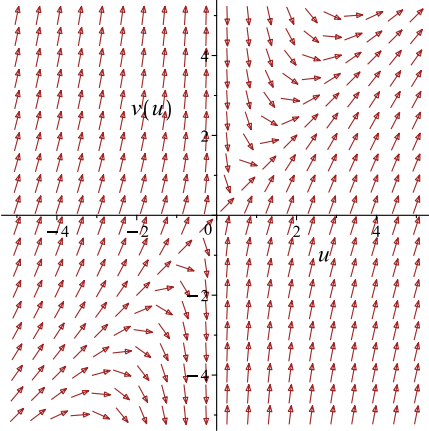
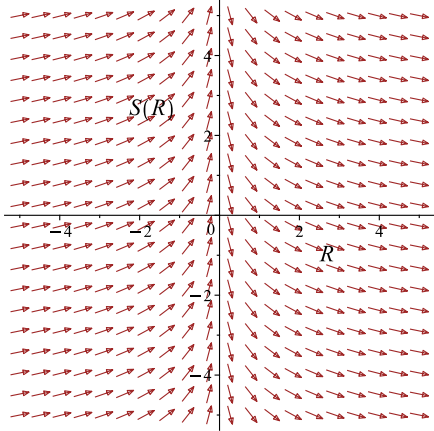
To complete the solution, we just need to transform the above back to u, v coordinates. This results in

$$\frac{\ln(-u + v)}{2} = -\ln(u) + c_2$$

Which gives

$$v = \frac{u^3 + e^{2c_2}}{u^2}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in u, v coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dv}{du} = -\frac{-3u+2v}{u}$ 	$R = u$ $S = \frac{\ln(-u + v)}{2}$	$\frac{dS}{dR} = -\frac{1}{R}$ 

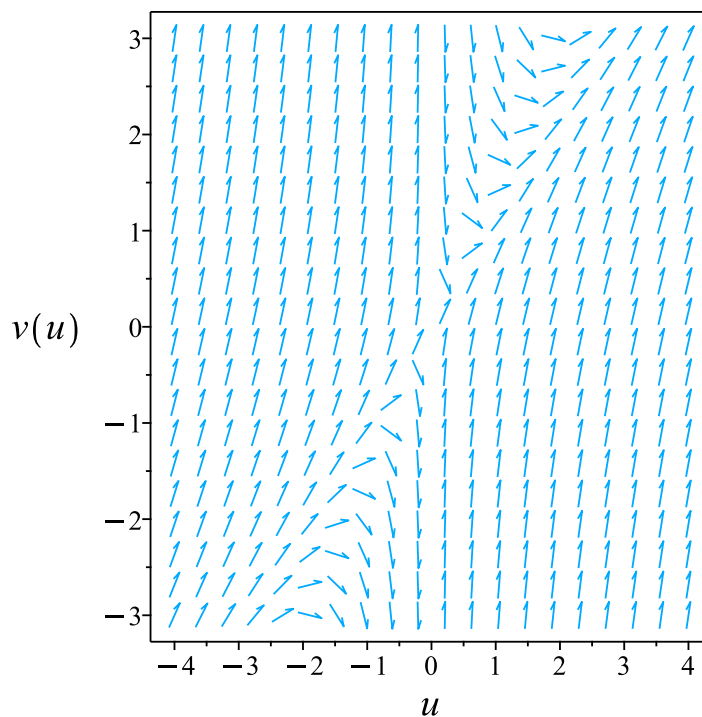


Figure 2.13: Slope field plot

$$v' + \frac{2v}{u} = 3$$

Summary of solutions found

$$v = \frac{u^3 + e^{2c_2}}{u^2}$$

Maple step by step solution

Let's solve

$$v' + \frac{2v}{u} = 3$$

- Highest derivative means the order of the ODE is 1

$$v'$$

- Solve for the highest derivative

$$v' = 3 - \frac{2v}{u}$$

- Group terms with v on the lhs of the ODE and the rest on the rhs of the ODE

$$v' + \frac{2v}{u} = 3$$

- The ODE is linear; multiply by an integrating factor $\mu(u)$

$$\mu(u) \left(v' + \frac{2v}{u} \right) = 3\mu(u)$$

- Assume the lhs of the ODE is the total derivative $\frac{d}{du}(v\mu(u))$

$$\mu(u) \left(v' + \frac{2v}{u} \right) = v'\mu(u) + v\mu'(u)$$

- Isolate $\mu'(u)$

$$\mu'(u) = \frac{2\mu(u)}{u}$$

- Solve to find the integrating factor

$$\mu(u) = u^2$$

- Integrate both sides with respect to u

$$\int \left(\frac{d}{du}(v\mu(u)) \right) du = \int 3\mu(u) du + C1$$

- Evaluate the integral on the lhs

$$v\mu(u) = \int 3\mu(u) du + C1$$

- Solve for v

$$v = \frac{\int 3\mu(u) du + C1}{\mu(u)}$$

- Substitute $\mu(u) = u^2$

$$v = \frac{\int 3u^2 du + C1}{u^2}$$

- Evaluate the integrals on the rhs

$$v = \frac{u^3 + C1}{u^2}$$

Maple trace

```
`Methods for first order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying 1st order linear  
<- 1st order linear successful`
```

Maple dsolve solution

Solving time : 0.002 (sec)

Leaf size : 11

```
dsolve(diff(v(u),u)+2*v(u)/u = 3,  
        v(u),singsol=all)
```

$$v = u + \frac{c_1}{u^2}$$

Mathematica DSolve solution

Solving time : 0.022 (sec)

Leaf size : 13

```
DSolve[{D[v[u],u]+2*v[u]/u==3,{}},  
        v[u],u,IncludeSingularSolutions->True]
```

$$v(u) \rightarrow u + \frac{c_1}{u^2}$$

2.2 Chapter IV. Methods of solution: First order equations. section 24. Problems at page 62

2.2.1	problem 4 (a)	88
2.2.2	problem 4 (b)	92
2.2.3	problem 4 (c)	99
2.2.4	problem 5	112
2.2.5	problem 6	124

2.2.1 problem 4 (a)

Solved as first order separable ode	88
Maple step by step solution	90
Maple trace	90
Maple dsolve solution	91
Mathematica DSolve solution	91

Internal problem ID [18216]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter IV. Methods of solution: First order equations. section 24. Problems at page 62

Problem number : 4 (a)

Date solved : Friday, December 20, 2024 at 06:02:22 AM

CAS classification : [_separable]

Solve

$$\sin(x) \cos(y)^2 + \cos(x)^2 y' = 0$$

Solved as first order separable ode

Time used: 0.204 (sec)

The ode $y' = -\frac{\sin(x) \cos(y)^2}{\cos(x)^2}$ is separable as it can be written as

$$\begin{aligned} y' &= -\frac{\sin(x) \cos(y)^2}{\cos(x)^2} \\ &= f(x)g(y) \end{aligned}$$

Where

$$\begin{aligned} f(x) &= -\frac{\sin(x)}{\cos(x)^2} \\ g(y) &= \cos(y)^2 \end{aligned}$$

Integrating gives

$$\begin{aligned} \int \frac{1}{g(y)} dy &= \int f(x) dx \\ \int \frac{1}{\cos(y)^2} dy &= \int -\frac{\sin(x)}{\cos(x)^2} dx \\ \tan(y) &= -\sec(x) + c_1 \end{aligned}$$

We now need to find the singular solutions, these are found by finding for what values $g(y)$ is zero, since we had to divide by this above. Solving $g(y) = 0$ or $\cos(y)^2 = 0$ for y gives

$$y = \frac{\pi}{2}$$

Now we go over each such singular solution and check if it verifies the ode itself and any initial conditions given. If it does not then the singular solution will not be used.

Therefore the solutions found are

$$\begin{aligned} \tan(y) &= -\sec(x) + c_1 \\ y &= \frac{\pi}{2} \end{aligned}$$

Solving for y gives

$$\begin{aligned} y &= \frac{\pi}{2} \\ y &= \arctan(-\sec(x) + c_1) \end{aligned}$$

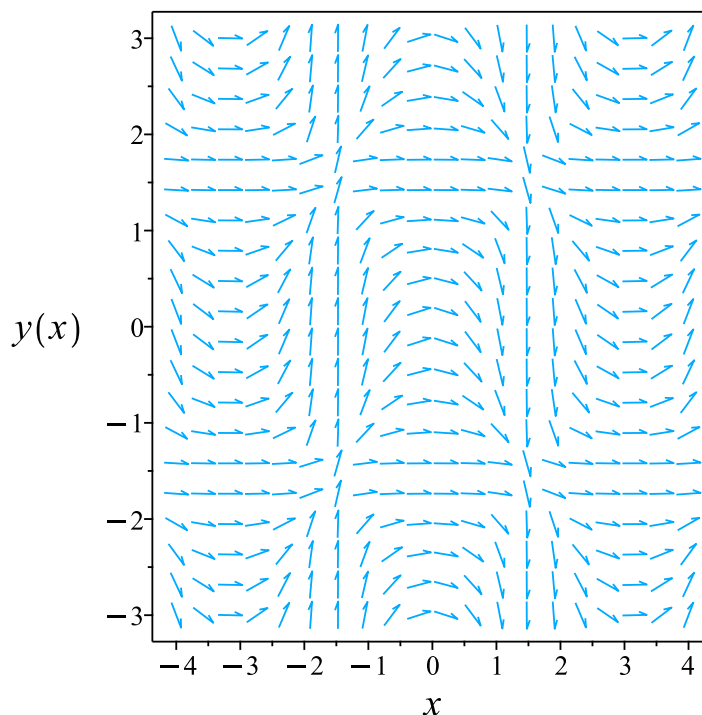


Figure 2.14: Slope field plot
 $\sin(x) \cos(y)^2 + \cos(x)^2 y' = 0$

Summary of solutions found

$$y = \frac{\pi}{2}$$

$$y = \arctan(-\sec(x) + c_1)$$

Maple step by step solution

Let's solve

$$\sin(x) \cos(y)^2 + \cos(x)^2 y' = 0$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Solve for the highest derivative

$$y' = -\frac{\sin(x) \cos(y)^2}{\cos(x)^2}$$

- Separate variables

$$\frac{y'}{\cos(y)^2} = -\frac{\sin(x)}{\cos(x)^2}$$

- Integrate both sides with respect to x

$$\int \frac{y'}{\cos(y)^2} dx = \int -\frac{\sin(x)}{\cos(x)^2} dx + C1$$

- Evaluate integral

$$\tan(y) = -\frac{1}{\cos(x)} + C1$$

- Solve for y

$$y = \arctan\left(\frac{C1 \cos(x) - 1}{\cos(x)}\right)$$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
<- separable successful`
```

Maple dsolve solution

Solving time : 0.005 (sec)

Leaf size : 11

```
dsolve(sin(x)*cos(y(x))^2+cos(x)^2*diff(y(x),x) = 0,  
y(x),singsol=all)
```

$$y = -\arctan(\sec(x) + c_1)$$

Mathematica DSolve solution

Solving time : 1.442 (sec)

Leaf size : 31

```
DSolve[{Sin[x]*Cos[y[x]]^2+ Cos[x]^2*D[y[x],x]==0,{}},  
y[x],x,IncludeSingularSolutions->True]
```

$$y(x) \rightarrow \arctan(-\sec(x) + c_1)$$

$$y(x) \rightarrow -\frac{\pi}{2}$$

$$y(x) \rightarrow \frac{\pi}{2}$$

2.2.2 problem 4 (b)

Solved as first order Exact ode	92
Maple step by step solution	97
Maple dsolve solution	98
Mathematica DSolve solution	98

Internal problem ID [18217]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter IV. Methods of solution: First order equations. section 24. Problems at page 62

Problem number : 4 (b)

Date solved : Friday, December 20, 2024 at 06:02:26 AM

CAS classification : unknown

Solve

$$y' + \sqrt{\frac{1-y^2}{-x^2+1}} = 0$$

Solved as first order Exact ode

Time used: 1025.226 (sec)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned} dy &= \left(-\sqrt{\frac{-y^2 + 1}{-x^2 + 1}} \right) dx \\ \left(\sqrt{\frac{-y^2 + 1}{-x^2 + 1}} \right) dx + dy &= 0 \end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= \sqrt{\frac{-y^2 + 1}{-x^2 + 1}} \\ N(x, y) &= 1 \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} \left(\sqrt{\frac{-y^2 + 1}{-x^2 + 1}} \right) \\ &= \frac{y}{\sqrt{\frac{y^2 - 1}{x^2 - 1}} (x^2 - 1)} \end{aligned}$$

And

$$\begin{aligned} \frac{\partial N}{\partial x} &= \frac{\partial}{\partial x} (1) \\ &= 0 \end{aligned}$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned} A &= \frac{1}{N} \left(\frac{\partial M}{\partial y} - \frac{\partial N}{\partial x} \right) \\ &= 1 \left(\left(-\frac{y}{\sqrt{\frac{-y^2+1}{-x^2+1}} (-x^2+1)} \right) - (0) \right) \\ &= \frac{y}{\sqrt{\frac{y^2-1}{x^2-1}} (x^2-1)} \end{aligned}$$

Since A depends on y , it can not be used to obtain an integrating factor. We will now try a second method to find an integrating factor. Let

$$\begin{aligned} B &= \frac{1}{M} \left(\frac{\partial N}{\partial x} - \frac{\partial M}{\partial y} \right) \\ &= \frac{1}{\sqrt{\frac{y^2-1}{x^2-1}}} \left((0) - \left(-\frac{y}{\sqrt{\frac{-y^2+1}{-x^2+1}} (-x^2+1)} \right) \right) \\ &= -\frac{y}{y^2-1} \end{aligned}$$

Since B does not depend on x , it can be used to obtain an integrating factor. Let the integrating factor be μ . Then

$$\begin{aligned} \mu &= e^{\int B \, dy} \\ &= e^{\int -\frac{y}{y^2-1} \, dy} \end{aligned}$$

The result of integrating gives

$$\begin{aligned} \mu &= e^{-\frac{\ln(y-1)}{2} - \frac{\ln(y+1)}{2}} \\ &= \frac{1}{\sqrt{y-1} \sqrt{y+1}} \end{aligned}$$

M and N are now multiplied by this integrating factor, giving new M and new N which are called \bar{M} and \bar{N} so not to confuse them with the original M and N .

$$\begin{aligned} \bar{M} &= \mu M \\ &= \frac{1}{\sqrt{y-1} \sqrt{y+1}} \left(\sqrt{\frac{-y^2+1}{-x^2+1}} \right) \\ &= \frac{\sqrt{\frac{y^2-1}{x^2-1}}}{\sqrt{y-1} \sqrt{y+1}} \end{aligned}$$

And

$$\begin{aligned}\bar{N} &= \mu N \\ &= \frac{1}{\sqrt{y-1}\sqrt{y+1}} \quad (1) \\ &= \frac{1}{\sqrt{y-1}\sqrt{y+1}}\end{aligned}$$

So now a modified ODE is obtained from the original ODE which will be exact and can be solved using the standard method. The modified ODE is

$$\begin{aligned}\bar{M} + \bar{N} \frac{dy}{dx} &= 0 \\ \left(\frac{\sqrt{\frac{y^2-1}{x^2-1}}}{\sqrt{y-1}\sqrt{y+1}} \right) + \left(\frac{1}{\sqrt{y-1}\sqrt{y+1}} \right) \frac{dy}{dx} &= 0\end{aligned}$$

The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = \bar{M} \quad (1)$$

$$\frac{\partial \phi}{\partial y} = \bar{N} \quad (2)$$

Integrating (2) w.r.t. y gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial y} dy &= \int \bar{N} dy \\ \int \frac{\partial \phi}{\partial y} dy &= \int \frac{1}{\sqrt{y-1}\sqrt{y+1}} dy \\ \phi &= \frac{\sqrt{y^2-1} \ln(y + \sqrt{y^2-1})}{\sqrt{y-1}\sqrt{y+1}} + f(x) \quad (3)\end{aligned}$$

Where $f(x)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t x gives

$$\frac{\partial \phi}{\partial x} = 0 + f'(x) \quad (4)$$

But equation (1) says that $\frac{\partial \phi}{\partial x} = \frac{\sqrt{\frac{y^2-1}{x^2-1}}}{\sqrt{y-1}\sqrt{y+1}}$. Therefore equation (4) becomes

$$\frac{\sqrt{\frac{y^2-1}{x^2-1}}}{\sqrt{y-1}\sqrt{y+1}} = 0 + f'(x) \quad (5)$$

Solving equation (5) for $f'(x)$ gives

$$f'(x) = \frac{\sqrt{\frac{y^2-1}{x^2-1}}}{\sqrt{y-1}\sqrt{y+1}}$$

Integrating the above w.r.t x gives

$$\int f'(x) dx = \int \left(\frac{\sqrt{\frac{y^2-1}{x^2-1}}}{\sqrt{y-1}\sqrt{y+1}} \right) dx$$

$$f(x) = \frac{\sqrt{\frac{y^2-1}{x^2-1}} \sqrt{x^2-1} \ln(x + \sqrt{x^2-1})}{\sqrt{y-1}\sqrt{y+1}} + c_1$$

Where c_1 is constant of integration. Substituting result found above for $f(x)$ into equation (3) gives ϕ

$$\phi = \frac{\sqrt{y^2-1} \ln(y + \sqrt{y^2-1})}{\sqrt{y-1}\sqrt{y+1}} + \frac{\sqrt{\frac{y^2-1}{x^2-1}} \sqrt{x^2-1} \ln(x + \sqrt{x^2-1})}{\sqrt{y-1}\sqrt{y+1}} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into the constant c_1 gives the solution as

$$c_1 = \frac{\sqrt{y^2-1} \ln(y + \sqrt{y^2-1})}{\sqrt{y-1}\sqrt{y+1}} + \frac{\sqrt{\frac{y^2-1}{x^2-1}} \sqrt{x^2-1} \ln(x + \sqrt{x^2-1})}{\sqrt{y-1}\sqrt{y+1}}$$

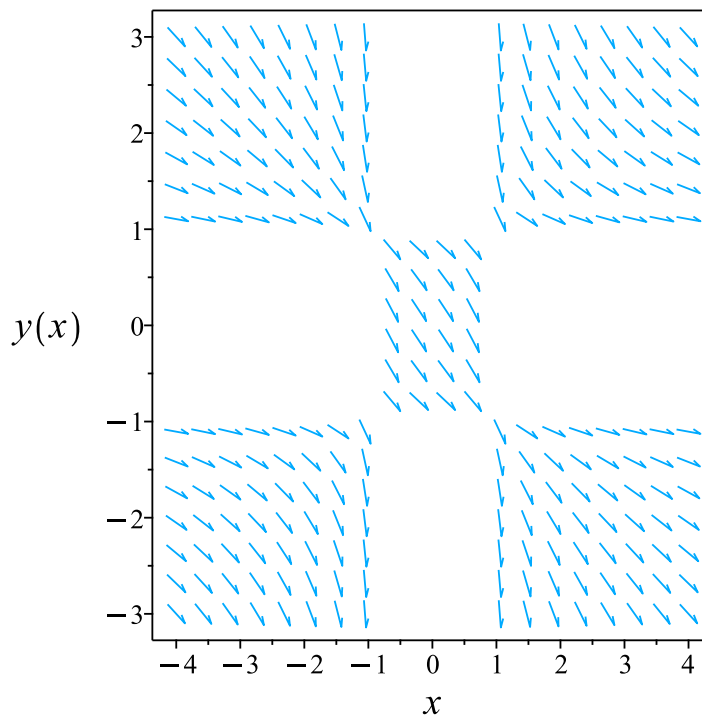


Figure 2.15: Slope field plot

$$y' + \sqrt{\frac{1-y^2}{-x^2+1}} = 0$$

Summary of solutions found

$$\frac{\sqrt{-1+y^2} \ln(y + \sqrt{-1+y^2})}{\sqrt{y-1} \sqrt{y+1}} + \frac{\sqrt{\frac{-1+y^2}{x^2-1}} \sqrt{x^2-1} \ln(x + \sqrt{x^2-1})}{\sqrt{y-1} \sqrt{y+1}} = c_1$$

Maple step by step solution

Let's solve

$$y' + \sqrt{\frac{1-y^2}{-x^2+1}} = 0$$

- Highest derivative means the order of the ODE is 1

y'

- Solve for the highest derivative

$$y' = -\sqrt{\frac{1-y^2}{-x^2+1}}$$

Maple dsolve solution

Solving time : 0.845 (sec)

Leaf size : 84

```
dsolve(diff(y(x),x)+((1-y(x)^2)/(-x^2+1))^(1/2) = 0,
        y(x),singsol=all)
```

$$\frac{\sqrt{\frac{-1+y^2}{x^2-1}} \sqrt{x^2-1} \ln(x + \sqrt{x^2-1})}{\sqrt{y-1} \sqrt{y+1}} + \frac{\sqrt{-1+y^2} \ln(y + \sqrt{-1+y^2})}{\sqrt{y-1} \sqrt{y+1}} + c_1 = 0$$

Mathematica DSolve solution

Solving time : 0.335 (sec)

Leaf size : 39

```
DSolve[{D[y[x],x]+Sqrt[(1-y[x]^2)/(1-x^2)]==0,{}},
        y[x],x,IncludeSingularSolutions->True]
```

$$y(x) \rightarrow -\cosh\left(2\operatorname{arctanh}\left(\frac{1}{\sqrt{\frac{x-1}{x+1}}}\right) - c_1\right)$$

$$y(x) \rightarrow -1$$

$$y(x) \rightarrow 1$$

2.2.3 problem 4 (c)

Solved as first order linear ode	99
Solved as first order separable ode	100
Solved as first order Exact ode	102
Solved using Lie symmetry for first order ode	105
Maple step by step solution	110
Maple trace	110
Maple dsolve solution	110
Mathematica DSolve solution	111

Internal problem ID [18218]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter IV. Methods of solution: First order equations. section 24. Problems at page 62

Problem number : 4 (c)

Date solved : Friday, December 20, 2024 at 06:25:07 AM

CAS classification : [_separable]

Solve

$$y - y'x = b(1 + y'x^2)$$

Solved as first order linear ode

Time used: 0.151 (sec)

In canonical form a linear first order is

$$y' + q(x)y = p(x)$$

Comparing the above to the given ode shows that

$$q(x) = -\frac{1}{x(bx+1)}$$

$$p(x) = -\frac{b}{x(bx+1)}$$

The integrating factor μ is

$$\begin{aligned} \mu &= e^{\int q dx} \\ &= e^{\int -\frac{1}{x(bx+1)} dx} \\ &= \frac{bx+1}{x} \end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dx}(\mu y) &= \mu p \\ \frac{d}{dx}(\mu y) &= (\mu) \left(-\frac{b}{x(bx+1)} \right) \\ \frac{d}{dx} \left(\frac{y(bx+1)}{x} \right) &= \left(\frac{bx+1}{x} \right) \left(-\frac{b}{x(bx+1)} \right) \\ d \left(\frac{y(bx+1)}{x} \right) &= \left(-\frac{b}{x^2} \right) dx\end{aligned}$$

Integrating gives

$$\begin{aligned}\frac{y(bx+1)}{x} &= \int -\frac{b}{x^2} dx \\ &= \frac{b}{x} + c_1\end{aligned}$$

Dividing throughout by the integrating factor $\frac{bx+1}{x}$ gives the final solution

$$y = \frac{c_1 x + b}{bx + 1}$$

Summary of solutions found

$$y = \frac{c_1 x + b}{bx + 1}$$

Solved as first order separable ode

Time used: 2.118 (sec)

The ode $y' = \frac{y-b}{x(bx+1)}$ is separable as it can be written as

$$\begin{aligned}y' &= \frac{y-b}{x(bx+1)} \\ &= f(x)g(y)\end{aligned}$$

Where

$$\begin{aligned}f(x) &= \frac{1}{x(bx+1)} \\ g(y) &= y-b\end{aligned}$$

Integrating gives

$$\int \frac{1}{g(y)} dy = \int f(x) dx$$

$$\int \frac{1}{y-b} dy = \int \frac{1}{x(bx+1)} dx$$

$$\ln(-y+b) = \ln\left(\frac{x}{bx+1}\right) + c_1$$

We now need to find the singular solutions, these are found by finding for what values $g(y)$ is zero, since we had to divide by this above. Solving $g(y) = 0$ or $y - b = 0$ for y gives

$$y = b$$

Now we go over each such singular solution and check if it verifies the ode itself and any initial conditions given. If it does not then the singular solution will not be used.

Therefore the solutions found are

$$\ln(-y+b) = \ln\left(\frac{x}{bx+1}\right) + c_1$$

$$y = b$$

Solving for y gives

$$y = b$$

$$y = -\frac{-b^2x + e^{c_1}x - b}{bx+1}$$

Summary of solutions found

$$y = b$$

$$y = -\frac{-b^2x + e^{c_1}x - b}{bx+1}$$

Solved as first order Exact ode

Time used: 1.023 (sec)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (\text{1A})$$

Therefore

$$\begin{aligned} (-bx^2 - x) dy &= (-y + b) dx \\ (y - b) dx + (-bx^2 - x) dy &= 0 \end{aligned} \quad (\text{2A})$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= y - b \\ N(x, y) &= -bx^2 - x \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(y - b) \\ &= 1\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(-bx^2 - x) \\ &= -2bx - 1\end{aligned}$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned}A &= \frac{1}{N} \left(\frac{\partial M}{\partial y} - \frac{\partial N}{\partial x} \right) \\ &= -\frac{1}{x(bx + 1)}((1) - (-2bx - 1)) \\ &= -\frac{2}{x}\end{aligned}$$

Since A does not depend on y , then it can be used to find an integrating factor. The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int A dx} \\ &= e^{\int -\frac{2}{x} dx}\end{aligned}$$

The result of integrating gives

$$\begin{aligned}\mu &= e^{-2\ln(x)} \\ &= \frac{1}{x^2}\end{aligned}$$

M and N are multiplied by this integrating factor, giving new M and new N which are called \bar{M} and \bar{N} for now so not to confuse them with the original M and N .

$$\begin{aligned}\bar{M} &= \mu M \\ &= \frac{1}{x^2}(y - b) \\ &= \frac{y - b}{x^2}\end{aligned}$$

And

$$\begin{aligned}\bar{N} &= \mu N \\ &= \frac{1}{x^2}(-bx^2 - x) \\ &= \frac{-bx - 1}{x}\end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned}\bar{M} + \bar{N} \frac{dy}{dx} &= 0 \\ \left(\frac{y-b}{x^2}\right) + \left(\frac{-bx-1}{x}\right) \frac{dy}{dx} &= 0\end{aligned}$$

The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = \bar{M} \tag{1}$$

$$\frac{\partial \phi}{\partial y} = \bar{N} \tag{2}$$

Integrating (2) w.r.t. y gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial y} dy &= \int \bar{N} dy \\ \int \frac{\partial \phi}{\partial y} dy &= \int \frac{-bx-1}{x} dy \\ \phi &= -\frac{y(bx+1)}{x} + f(x)\end{aligned} \tag{3}$$

Where $f(x)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t x gives

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= -\frac{yb}{x} + \frac{y(bx+1)}{x^2} + f'(x) \\ &= \frac{y}{x^2} + f'(x)\end{aligned} \tag{4}$$

But equation (1) says that $\frac{\partial \phi}{\partial x} = \frac{y-b}{x^2}$. Therefore equation (4) becomes

$$\frac{y-b}{x^2} = \frac{y}{x^2} + f'(x) \tag{5}$$

Solving equation (5) for $f'(x)$ gives

$$f'(x) = -\frac{b}{x^2}$$

Integrating the above w.r.t x gives

$$\int f'(x) dx = \int \left(-\frac{b}{x^2}\right) dx$$

$$f(x) = \frac{b}{x} + c_1$$

Where c_1 is constant of integration. Substituting result found above for $f(x)$ into equation (3) gives ϕ

$$\phi = -\frac{y(bx+1)}{x} + \frac{b}{x} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into the constant c_1 gives the solution as

$$c_1 = -\frac{y(bx+1)}{x} + \frac{b}{x}$$

Solving for y gives

$$y = -\frac{c_1x - b}{bx + 1}$$

Summary of solutions found

$$y = -\frac{c_1x - b}{bx + 1}$$

Solved using Lie symmetry for first order ode

Time used: 32.554 (sec)

Writing the ode as

$$y' = \frac{y - b}{x(bx + 1)}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

To determine ξ, η then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$\xi = xa_2 + ya_3 + a_1 \quad (\text{1E})$$

$$\eta = xb_2 + yb_3 + b_1 \quad (\text{2E})$$

Where the unknown coefficients are

$$\{a_1, a_2, a_3, b_1, b_2, b_3\}$$

Substituting equations (1E,2E) and ω into (A) gives

$$\begin{aligned} b_2 + \frac{(y-b)(b_3-a_2)}{x(bx+1)} - \frac{(y-b)^2 a_3}{x^2(bx+1)^2} \\ - \left(-\frac{y-b}{x^2(bx+1)} - \frac{(y-b)b}{x(bx+1)^2} \right) (xa_2 + ya_3 + a_1) - \frac{xb_2 + yb_3 + b_1}{x(bx+1)} = 0 \end{aligned} \quad (\text{5E})$$

Putting the above in normal form gives

$$\frac{b^2 x^4 b_2 - b^2 x^2 a_2 - b^2 x^2 b_3 - 2b^2 x y a_3 + b x^3 b_2 + b x^2 y a_2 + 2b x y^2 a_3 - 2b^2 x a_1 - b x^2 b_1 + 2b x y a_1 - b^2 a_3 - b x}{x^2 (bx+1)^2} = 0$$

Setting the numerator to zero gives

$$\begin{aligned} b^2 x^4 b_2 - b^2 x^2 a_2 - b^2 x^2 b_3 - 2b^2 x y a_3 + b x^3 b_2 + b x^2 y a_2 + 2b x y^2 a_3 \\ - 2b^2 x a_1 - b x^2 b_1 + 2b x y a_1 - b^2 a_3 - b x b_3 + b y a_3 - b a_1 - x b_1 + y a_1 = 0 \end{aligned} \quad (\text{6E})$$

Looking at the above PDE shows the following are all the terms with $\{x, y\}$ in them.

$$\{x, y\}$$

The following substitution is now made to be able to collect on all terms with $\{x, y\}$ in them

$$\{x = v_1, y = v_2\}$$

The above PDE (6E) now becomes

$$\begin{aligned} b^2 b_2 v_1^4 - b^2 a_2 v_1^2 - 2b^2 a_3 v_1 v_2 - b^2 b_3 v_1^2 + b a_2 v_1^2 v_2 + 2b a_3 v_1 v_2^2 + b b_2 v_1^3 - 2b^2 a_1 v_1 \\ + 2b a_1 v_1 v_2 - b b_1 v_1^2 - b^2 a_3 + b a_3 v_2 - b b_3 v_1 - b a_1 + a_1 v_2 - b_1 v_1 = 0 \end{aligned} \quad (7E)$$

Collecting the above on the terms v_i introduced, and these are

$$\{v_1, v_2\}$$

Equation (7E) now becomes

$$\begin{aligned} b^2 b_2 v_1^4 + b b_2 v_1^3 + b a_2 v_1^2 v_2 + (-b^2 a_2 - b^2 b_3 - b b_1) v_1^2 + 2b a_3 v_1 v_2^2 \\ + (-2b^2 a_3 + 2b a_1) v_1 v_2 + (-2b^2 a_1 - b b_3 - b_1) v_1 + (b a_3 + a_1) v_2 - b^2 a_3 - b a_1 = 0 \end{aligned} \quad (8E)$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$\begin{aligned} b a_2 &= 0 \\ b b_2 &= 0 \\ b^2 b_3 &= 0 \\ 2b a_3 &= 0 \\ b a_3 + a_1 &= 0 \\ -2b^2 a_3 + 2b a_1 &= 0 \\ -b^2 a_3 - b a_1 &= 0 \\ -2b^2 a_1 - b b_3 - b_1 &= 0 \\ -b^2 a_2 - b^2 b_3 - b b_1 &= 0 \end{aligned}$$

Solving the above equations for the unknowns gives

$$\begin{aligned} a_1 &= 0 \\ a_2 &= 0 \\ a_3 &= 0 \\ b_1 &= -b b_3 \\ b_2 &= 0 \\ b_3 &= b_3 \end{aligned}$$

Substituting the above solution in the anstaz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$\begin{aligned}\xi &= 0 \\ \eta &= y - b\end{aligned}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS \quad (1)$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = x$$

S is found from

$$\begin{aligned}S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{y - b} dy\end{aligned}$$

Which results in

$$S = \ln(y - b)$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = \frac{y - b}{x(bx + 1)}$$

Evaluating all the partial derivatives gives

$$\begin{aligned}R_x &= 1 \\R_y &= 0 \\S_x &= 0 \\S_y &= \frac{1}{y-b}\end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = \frac{1}{x(bx+1)} \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = \frac{1}{R(Rb+1)}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S .

Since the ode has the form $\frac{d}{dR}S(R) = f(R)$, then we only need to integrate $f(R)$.

$$\begin{aligned}\int dS &= \int \frac{1}{R(Rb+1)} dR \\S(R) &= -\ln(Rb+1) + \ln(R) + c_2\end{aligned}$$

This results in

$$\ln(y-b) = -\ln(bx+1) + \ln(x) + c_2$$

Solving for y gives

$$y = \frac{b^2x + xe^{c_2} + b}{bx+1}$$

Summary of solutions found

$$y = \frac{b^2x + xe^{c_2} + b}{bx+1}$$

Maple step by step solution

Let's solve

$$y - y'x = b(1 + y'x^2)$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Solve for the highest derivative

$$y' = \frac{-y+b}{-bx^2-x}$$

- Separate variables

$$\frac{y'}{-y+b} = \frac{1}{-bx^2-x}$$

- Integrate both sides with respect to x

$$\int \frac{y'}{-y+b} dx = \int \frac{1}{-bx^2-x} dx + C1$$

- Evaluate integral

$$-\ln(-y+b) = \ln(bx+1) - \ln(x) + C1$$

- Solve for y

$$y = \frac{e^{C1}b^2x + e^{C1}b - x}{e^{C1}(bx+1)}$$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`
```

Maple dsolve solution

Solving time : 0.002 (sec)

Leaf size : 17

```
dsolve(y(x)-diff(y(x),x)*x = b*(1+diff(y(x),x)*x^2),
       y(x),singsol=all)
```

$$y = \frac{c_1x + b}{bx + 1}$$

Mathematica DSolve solution

Solving time : 0.032 (sec)

Leaf size : 24

```
DSolve[{y[x]-x*D[y[x],x]==b*(1+x^2*D[y[x],x]),{}},  
y[x],x,IncludeSingularSolutions->True]
```

$$y(x) \rightarrow \frac{b + c_1 x}{bx + 1}$$

$$y(x) \rightarrow b$$

2.2.4 problem 5

Solved as first order autonomous ode	112
Solved as first order Exact ode	114
Solved using Lie symmetry for first order ode	118
Maple step by step solution	122
Maple trace	123
Maple dsolve solution	123
Mathematica DSolve solution	123

Internal problem ID [18219]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter IV. Methods of solution: First order equations. section 24. Problems at page 62

Problem number : 5

Date solved : Friday, December 20, 2024 at 06:27:15 AM

CAS classification : [_quadrature]

Solve

$$x' = k(A - nx)(M - mx)$$

Solved as first order autonomous ode

Time used: 5.780 (sec)

Integrating gives

$$\int \frac{1}{k(-nx + A)(-mx + M)} dx = dt$$

$$\frac{-\ln(-mx + M) + \ln(-nx + A)}{k(Am - Mn)} = t + c_1$$

Singular solutions are found by solving

$$k(-nx + A)(-mx + M) = 0$$

for x . This is because we had to divide by this in the above step. This gives the following singular solution(s), which also have to satisfy the given ODE.

$$x = \frac{A}{n}$$

$$x = \frac{M}{m}$$

The following diagram is the phase line diagram. It classifies each of the above equilibrium points as stable or not stable or semi-stable.

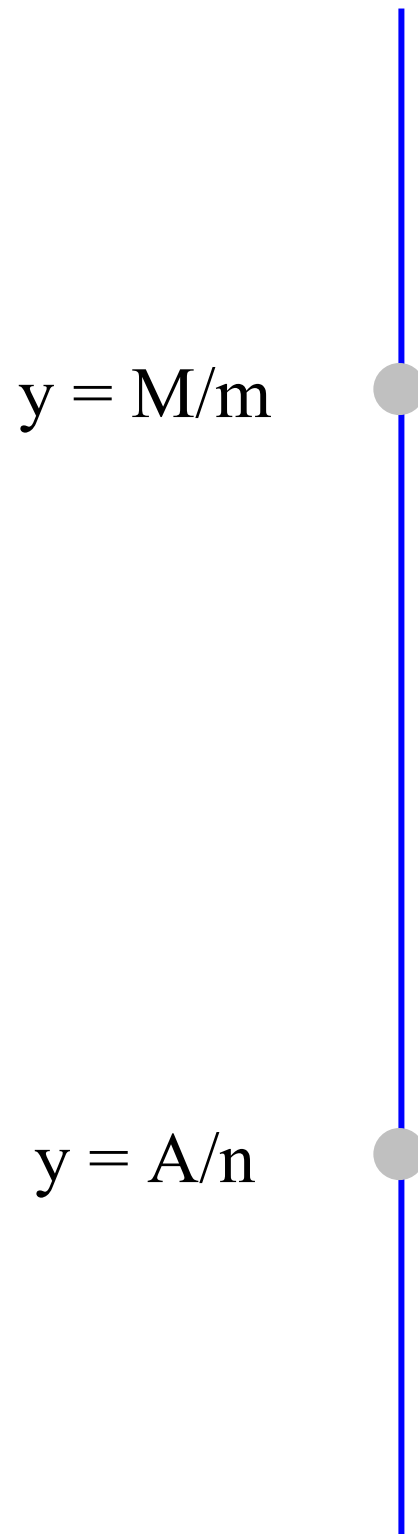


Figure 2.16: Phase line diagram

Solving for x gives

$$x = \frac{A}{n}$$

$$x = \frac{M}{m}$$

$$x = \frac{A e^{-Ac_1 km - Akmt + Mc_1 kn + Mknt} - M}{e^{-Ac_1 km - Akmt + Mc_1 kn + Mknt} n - m}$$

Summary of solutions found

$$x = \frac{A}{n}$$

$$x = \frac{M}{m}$$

$$x = \frac{A e^{-Ac_1 km - Akmt + Mc_1 kn + Mknt} - M}{e^{-Ac_1 km - Akmt + Mc_1 kn + Mknt} n - m}$$

Solved as first order Exact ode

Time used: 2.892 (sec)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\frac{\partial \phi}{\partial x} = M$$

$$\frac{\partial \phi}{\partial y} = N$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(t, x) dt + N(t, x) dx = 0 \quad (1A)$$

Therefore

$$\begin{aligned} dx &= (k(-nx + A)(-mx + M)) dt \\ (-k(-nx + A)(-mx + M)) dt + dx &= 0 \end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(t, x) &= -k(-nx + A)(-mx + M) \\ N(t, x) &= 1 \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial x} = \frac{\partial N}{\partial t}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial x} &= \frac{\partial}{\partial x} (-k(-nx + A)(-mx + M)) \\ &= k((-2nx + A)m + Mn) \end{aligned}$$

And

$$\begin{aligned} \frac{\partial N}{\partial t} &= \frac{\partial}{\partial t} (1) \\ &= 0 \end{aligned}$$

Since $\frac{\partial M}{\partial x} \neq \frac{\partial N}{\partial t}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned} A &= \frac{1}{N} \left(\frac{\partial M}{\partial x} - \frac{\partial N}{\partial t} \right) \\ &= 1((kn(-mx + M) + k(-nx + A)m) - (0)) \\ &= k((-2nx + A)m + Mn) \end{aligned}$$

Since A depends on x , it can not be used to obtain an integrating factor. We will now try a second method to find an integrating factor. Let

$$\begin{aligned} B &= \frac{1}{M} \left(\frac{\partial N}{\partial t} - \frac{\partial M}{\partial x} \right) \\ &= -\frac{1}{k(-nx+A)(-mx+M)} ((0) - (kn(-mx+M) + k(-nx+A)m)) \\ &= \frac{(-2nx+A)m + Mn}{(-nx+A)(-mx+M)} \end{aligned}$$

Since B does not depend on t , it can be used to obtain an integrating factor. Let the integrating factor be μ . Then

$$\begin{aligned} \mu &= e^{\int B dx} \\ &= e^{\int \frac{(-2nx+A)m + Mn}{(-nx+A)(-mx+M)} dx} \end{aligned}$$

The result of integrating gives

$$\begin{aligned} \mu &= e^{-\ln((-nx+A)(-mx+M))} \\ &= \frac{1}{(-nx+A)(-mx+M)} \end{aligned}$$

M and N are now multiplied by this integrating factor, giving new M and new N which are called \bar{M} and \bar{N} so not to confuse them with the original M and N .

$$\begin{aligned} \bar{M} &= \mu M \\ &= \frac{1}{(-nx+A)(-mx+M)} (-k(-nx+A)(-mx+M)) \\ &= -k \end{aligned}$$

And

$$\begin{aligned} \bar{N} &= \mu N \\ &= \frac{1}{(-nx+A)(-mx+M)} (1) \\ &= \frac{1}{(-nx+A)(-mx+M)} \end{aligned}$$

So now a modified ODE is obtained from the original ODE which will be exact and can be solved using the standard method. The modified ODE is

$$\begin{aligned} \bar{M} + \bar{N} \frac{dx}{dt} &= 0 \\ (-k) + \left(\frac{1}{(-nx+A)(-mx+M)} \right) \frac{dx}{dt} &= 0 \end{aligned}$$

The following equations are now set up to solve for the function $\phi(t, x)$

$$\frac{\partial \phi}{\partial t} = \overline{M} \quad (1)$$

$$\frac{\partial \phi}{\partial x} = \overline{N} \quad (2)$$

Integrating (1) w.r.t. t gives

$$\begin{aligned} \int \frac{\partial \phi}{\partial t} dt &= \int \overline{M} dt \\ \int \frac{\partial \phi}{\partial t} dt &= \int -k dt \\ \phi &= -tk + f(x) \end{aligned} \quad (3)$$

Where $f(x)$ is used for the constant of integration since ϕ is a function of both t and x . Taking derivative of equation (3) w.r.t x gives

$$\frac{\partial \phi}{\partial x} = 0 + f'(x) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial x} = \frac{1}{(-nx+A)(-mx+M)}$. Therefore equation (4) becomes

$$\frac{1}{(-nx+A)(-mx+M)} = 0 + f'(x) \quad (5)$$

Solving equation (5) for $f'(x)$ gives

$$f'(x) = \frac{1}{(-nx+A)(-mx+M)}$$

Integrating the above w.r.t x gives

$$\begin{aligned} \int f'(x) dx &= \int \left(\frac{1}{(-nx+A)(-mx+M)} \right) dx \\ f(x) &= -\frac{\ln(-mx+M)}{Am-Mn} + \frac{\ln(-nx+A)}{Am-Mn} + c_1 \end{aligned}$$

Where c_1 is constant of integration. Substituting result found above for $f(x)$ into equation (3) gives ϕ

$$\phi = -tk - \frac{\ln(-mx+M)}{Am-Mn} + \frac{\ln(-nx+A)}{Am-Mn} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into the constant c_1 gives the solution as

$$c_1 = -tk - \frac{\ln(-mx + M)}{Am - Mn} + \frac{\ln(-nx + A)}{Am - Mn}$$

Solving for x gives

$$x = \frac{A e^{-Akm t + Mkn t - mc_1 A + Mc_1 n} - M}{e^{-Akm t + Mkn t - mc_1 A + Mc_1 n} - m}$$

Summary of solutions found

$$x = \frac{A e^{-Akm t + Mkn t - mc_1 A + Mc_1 n} - M}{e^{-Akm t + Mkn t - mc_1 A + Mc_1 n} - m}$$

Solved using Lie symmetry for first order ode

Time used: 7.722 (sec)

Writing the ode as

$$\begin{aligned} x' &= k(-nx + A)(-mx + M) \\ x' &= \omega(t, x) \end{aligned}$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_t + \omega(\eta_x - \xi_t) - \omega^2 \xi_x - \omega_t \xi - \omega_x \eta = 0 \quad (\text{A})$$

To determine ξ, η then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$\xi = ta_2 + xa_3 + a_1 \quad (1\text{E})$$

$$\eta = tb_2 + xb_3 + b_1 \quad (2\text{E})$$

Where the unknown coefficients are

$$\{a_1, a_2, a_3, b_1, b_2, b_3\}$$

Substituting equations (1E,2E) and ω into (A) gives

$$\begin{aligned} b_2 + k(-nx + A)(-mx + M)(b_3 - a_2) - k^2(-nx + A)^2(-mx + M)^2 a_3 \\ - (-kn(-mx + M) - k(-nx + A)m)(tb_2 + xb_3 + b_1) = 0 \end{aligned} \quad (5\text{E})$$

Putting the above in normal form gives

$$\begin{aligned}
& -k^2 m^2 n^2 x^4 a_3 + 2A k^2 m^2 n x^3 a_3 + 2M k^2 m n^2 x^3 a_3 - A^2 k^2 m^2 x^2 a_3 \\
& - 4AM k^2 m n x^2 a_3 - M^2 k^2 n^2 x^2 a_3 + 2A^2 M k^2 m x a_3 + 2A M^2 k^2 n x a_3 \\
& - A^2 M^2 k^2 a_3 - 2kmntxb_2 - kmn x^2 a_2 - kmn x^2 b_3 + Akmtb_2 + Akmxa_2 \\
& + Mkntb_2 + Mknxa_2 - 2kmnxb_1 - AMka_2 + AMkb_3 + Akmb_1 + Mknb_1 + b_2 = 0
\end{aligned}$$

Setting the numerator to zero gives

$$\begin{aligned}
& -k^2 m^2 n^2 x^4 a_3 + 2A k^2 m^2 n x^3 a_3 + 2M k^2 m n^2 x^3 a_3 - A^2 k^2 m^2 x^2 a_3 \\
& - 4AM k^2 m n x^2 a_3 - M^2 k^2 n^2 x^2 a_3 + 2A^2 M k^2 m x a_3 \\
& + 2A M^2 k^2 n x a_3 - A^2 M^2 k^2 a_3 - 2kmntxb_2 - kmn x^2 a_2 \\
& - kmn x^2 b_3 + Akmtb_2 + Akmxa_2 + Mkntb_2 + Mknxa_2 \\
& - 2kmnxb_1 - AMka_2 + AMkb_3 + Akmb_1 + Mknb_1 + b_2 = 0
\end{aligned} \tag{6E}$$

Looking at the above PDE shows the following are all the terms with $\{t, x\}$ in them.

$$\{t, x\}$$

The following substitution is now made to be able to collect on all terms with $\{t, x\}$ in them

$$\{t = v_1, x = v_2\}$$

The above PDE (6E) now becomes

$$\begin{aligned}
& -k^2 m^2 n^2 a_3 v_2^4 + 2A k^2 m^2 n a_3 v_2^3 + 2M k^2 m n^2 a_3 v_2^3 - A^2 k^2 m^2 a_3 v_2^2 \\
& - 4AM k^2 m n a_3 v_2^2 - M^2 k^2 n^2 a_3 v_2^2 + 2A^2 M k^2 m a_3 v_2 \\
& + 2A M^2 k^2 n a_3 v_2 - A^2 M^2 k^2 a_3 - kmna_2 v_2^2 - 2kmnb_2 v_1 v_2 \\
& - kmnb_3 v_2^2 + Akma_2 v_2 + Akmb_2 v_1 + Mkna_2 v_2 + Mknb_2 v_1 \\
& - 2kmnb_1 v_2 - AMka_2 + AMkb_3 + Akmb_1 + Mknb_1 + b_2 = 0
\end{aligned} \tag{7E}$$

Collecting the above on the terms v_i introduced, and these are

$$\{v_1, v_2\}$$

Equation (7E) now becomes

$$\begin{aligned}
& -2kmnb_2 v_1 v_2 + (Akmb_2 + Mknb_2) v_1 \\
& - k^2 m^2 n^2 a_3 v_2^4 + (2A k^2 m^2 n a_3 + 2M k^2 m n^2 a_3) v_2^3 \\
& + (-A^2 k^2 m^2 a_3 - 4AM k^2 m n a_3 - M^2 k^2 n^2 a_3 - kmna_2 - kmnb_3) v_2^2 \\
& + (2A^2 M k^2 m a_3 + 2A M^2 k^2 n a_3 + Akma_2 + Mkna_2 - 2kmnb_1) v_2 \\
& - A^2 M^2 k^2 a_3 - AMka_2 + AMkb_3 + Akmb_1 + Mknb_1 + b_2 = 0
\end{aligned} \tag{8E}$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$\begin{aligned}
 -2kmnb_2 &= 0 \\
 -k^2m^2n^2a_3 &= 0 \\
 2A k^2m^2na_3 + 2M k^2m n^2a_3 &= 0 \\
 Akmb_2 + Mknb_2 &= 0 \\
 2A^2M k^2ma_3 + 2A M^2k^2na_3 + Akma_2 + Mkn a_2 - 2kmnb_1 &= 0 \\
 -A^2k^2m^2a_3 - 4AM k^2mna_3 - M^2k^2n^2a_3 - kmna_2 - kmnb_3 &= 0 \\
 -A^2M^2k^2a_3 - AMka_2 + AMkb_3 + Akmb_1 + Mknb_1 + b_2 &= 0
 \end{aligned}$$

Solving the above equations for the unknowns gives

$$\begin{aligned}
 a_1 &= a_1 \\
 a_2 &= 0 \\
 a_3 &= 0 \\
 b_1 &= 0 \\
 b_2 &= 0 \\
 b_3 &= 0
 \end{aligned}$$

Substituting the above solution in the ansatz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$\begin{aligned}
 \xi &= 1 \\
 \eta &= 0
 \end{aligned}$$

Shifting is now applied to make $\xi = 0$ in order to simplify the rest of the computation

$$\begin{aligned}
 \eta &= \eta - \omega(t, x) \xi \\
 &= 0 - (k(-nx + A) (-mx + M)) (1) \\
 &= -x^2kmn + Axkm + Mxkn - AMk \\
 \xi &= 0
 \end{aligned}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(t, x) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dt}{\xi} = \frac{dx}{\eta} = dS \tag{1}$$

The above comes from the requirements that $(\xi \frac{\partial}{\partial t} + \eta \frac{\partial}{\partial x}) S(t, x) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = t$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{-x^2 kmn + Axkm + Mxkn - AMk} dy \end{aligned}$$

Which results in

$$S = \frac{\ln(-mx + M)}{k(Am - Mn)} - \frac{\ln(-nx + A)}{k(Am - Mn)}$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_t + \omega(t, x)S_x}{R_t + \omega(t, x)R_x} \quad (2)$$

Where in the above R_t, R_x, S_t, S_x are all partial derivatives and $\omega(t, x)$ is the right hand side of the original ode given by

$$\omega(t, x) = k(-nx + A)(-mx + M)$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_t &= 1 \\ R_x &= 0 \\ S_t &= 0 \\ S_x &= -\frac{1}{k(-nx + A)(-mx + M)} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = -1 \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for t, x in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = -1$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S .

Since the ode has the form $\frac{d}{dR}S(R) = f(R)$, then we only need to integrate $f(R)$.

$$\int dS = \int -1 dR$$

$$S(R) = -R + c_2$$

To complete the solution, we just need to transform the above back to t, x coordinates. This results in

$$\frac{\ln(M - mx) - \ln(A - nx)}{k(Am - Mn)} = -t + c_2$$

Which gives

$$x = \frac{A e^{Ac_2 km - Akmt - Mc_2 kn + Mknt} - M}{e^{Ac_2 km - Akmt - Mc_2 kn + Mknt} n - m}$$

Summary of solutions found

$$x = \frac{A e^{Ac_2 km - Akmt - Mc_2 kn + Mknt} - M}{e^{Ac_2 km - Akmt - Mc_2 kn + Mknt} n - m}$$

Maple step by step solution

Let's solve

$$x' = k(A - nx)(M - mx)$$

- Highest derivative means the order of the ODE is 1

$$x'$$

- Solve for the highest derivative

$$x' = k(A - nx)(M - mx)$$

- Separate variables

$$\frac{x'}{(A-nx)(M-mx)} = k$$

- Integrate both sides with respect to t

$$\int \frac{x'}{(A-nx)(M-mx)} dt = \int k dt + C1$$

- Evaluate integral

$$-\frac{\ln(M-mx)}{Am-Mn} + \frac{\ln(A-nx)}{Am-Mn} = tk + C1$$

- Solve for x

$$x = \frac{Ae^{-Akt+Mknt-ACIm+CI Mn} - M}{e^{-Akt+Mknt-ACIm+CI Mn}n - m}$$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
<- separable successful`
```

Maple dsolve solution

Solving time : 0.005 (sec)

Leaf size : 47

```
dsolve(diff(x(t),t) = k*(A-n*x(t))*(M-m*x(t)),
        x(t),singsol=all)
```

$$x = \frac{-Ae^{-k(t+c_1)(Am-Mn)} + M}{-e^{-k(t+c_1)(Am-Mn)}n + m}$$

Mathematica DSolve solution

Solving time : 3.359 (sec)

Leaf size : 82

```
DSolve[{D[x[t],t]==k*(A-n*x[t])*(M-m*x[t]),{}}],
        x[t],t,IncludeSingularSolutions->True]
```

$$x(t) \rightarrow \frac{Ae^{Mn(kt+c_1)} - Me^{Am(kt+c_1)}}{ne^{Mn(kt+c_1)} - me^{Am(kt+c_1)}}$$

$$x(t) \rightarrow \frac{M}{m}$$

$$x(t) \rightarrow \frac{A}{n}$$

2.2.5 problem 6

Solved as first order separable ode	124
Solved as first order Exact ode	126
Maple step by step solution	131
Maple trace	131
Maple dsolve solution	131
Mathematica DSolve solution	132

Internal problem ID [18220]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter IV. Methods of solution: First order equations. section 24. Problems at page 62

Problem number : 6

Date solved : Friday, December 20, 2024 at 06:27:33 AM

CAS classification : [_separable]

Solve

$$y' = 1 + \frac{1}{x} - \frac{1}{y^2 + 2} - \frac{1}{x(y^2 + 2)}$$

Solved as first order separable ode

Time used: 31.273 (sec)

The ode $y' = \frac{y^2x + y^2 + x + 1}{x(y^2 + 2)}$ is separable as it can be written as

$$\begin{aligned} y' &= \frac{y^2x + y^2 + x + 1}{x(y^2 + 2)} \\ &= f(x)g(y) \end{aligned}$$

Where

$$\begin{aligned} f(x) &= \frac{x + 1}{x} \\ g(y) &= \frac{y^2 + 1}{y^2 + 2} \end{aligned}$$

Integrating gives

$$\int \frac{1}{g(y)} dy = \int f(x) dx$$
$$\int \frac{y^2 + 2}{y^2 + 1} dy = \int \frac{x + 1}{x} dx$$
$$y + \arctan(y) = x + \ln(x) + c_1$$

We now need to find the singular solutions, these are found by finding for what values $g(y)$ is zero, since we had to divide by this above. Solving $g(y) = 0$ or $\frac{y^2+1}{y^2+2} = 0$ for y gives

$$y = -i$$
$$y = i$$

Now we go over each such singular solution and check if it verifies the ode itself and any initial conditions given. If it does not then the singular solution will not be used.

Therefore the solutions found are

$$y + \arctan(y) = x + \ln(x) + c_1$$
$$y = -i$$
$$y = i$$

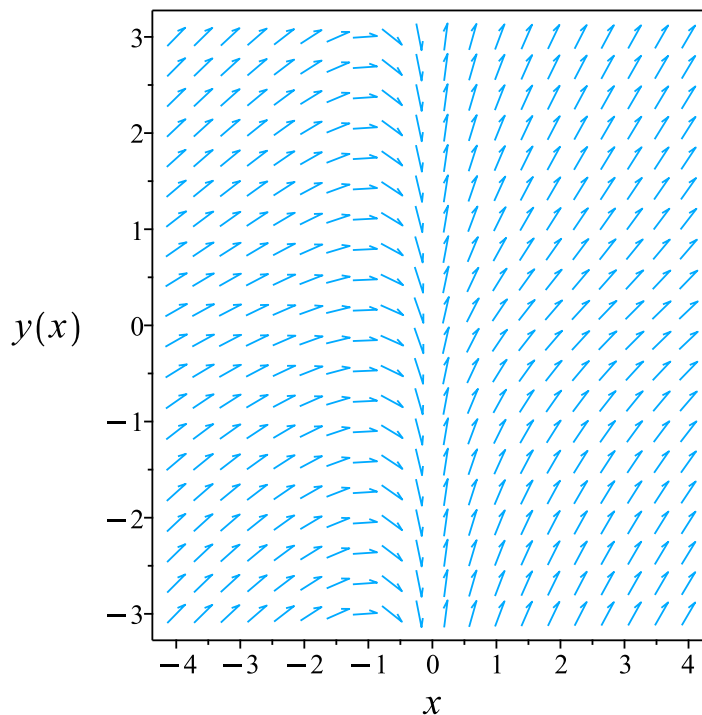


Figure 2.17: Slope field plot
 $y' = 1 + \frac{1}{x} - \frac{1}{y^2+2} - \frac{1}{x(y^2+2)}$

Summary of solutions found

$$y + \arctan(y) = x + \ln(x) + c_1$$

$$y = -i$$

$$y = i$$

Solved as first order Exact ode

Time used: 1.010 (sec)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx}\phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (\text{1A})$$

Therefore

$$\begin{aligned} dy &= \left(1 + \frac{1}{x} - \frac{1}{y^2 + 2} - \frac{1}{x(y^2 + 2)} \right) dx \\ \left(-1 - \frac{1}{x} + \frac{1}{y^2 + 2} + \frac{1}{x(y^2 + 2)} \right) dx + dy &= 0 \end{aligned} \quad (\text{2A})$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= -1 - \frac{1}{x} + \frac{1}{y^2 + 2} + \frac{1}{x(y^2 + 2)} \\ N(x, y) &= 1 \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} \left(-1 - \frac{1}{x} + \frac{1}{y^2 + 2} + \frac{1}{x(y^2 + 2)} \right) \\ &= -\frac{2(x+1)y}{x(y^2 + 2)^2} \end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(1) \\ &= 0\end{aligned}$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned}A &= \frac{1}{N} \left(\frac{\partial M}{\partial y} - \frac{\partial N}{\partial x} \right) \\ &= 1 \left(\left(-\frac{2y}{(y^2+2)^2} - \frac{2y}{x(y^2+2)^2} \right) - (0) \right) \\ &= -\frac{2(x+1)y}{x(y^2+2)^2}\end{aligned}$$

Since A depends on y , it can not be used to obtain an integrating factor. We will now try a second method to find an integrating factor. Let

$$\begin{aligned}B &= \frac{1}{M} \left(\frac{\partial N}{\partial x} - \frac{\partial M}{\partial y} \right) \\ &= -\frac{x(y^2+2)}{(y^2+1)(x+1)} \left((0) - \left(-\frac{2y}{(y^2+2)^2} - \frac{2y}{x(y^2+2)^2} \right) \right) \\ &= -\frac{2y}{(y^2+1)(y^2+2)}\end{aligned}$$

Since B does not depend on x , it can be used to obtain an integrating factor. Let the integrating factor be μ . Then

$$\begin{aligned}\mu &= e^{\int B dy} \\ &= e^{\int -\frac{2y}{(y^2+1)(y^2+2)} dy}\end{aligned}$$

The result of integrating gives

$$\begin{aligned}\mu &= e^{\ln(y^2+2) - \ln(y^2+1)} \\ &= \frac{y^2+2}{y^2+1}\end{aligned}$$

M and N are now multiplied by this integrating factor, giving new M and new N which are called \bar{M} and \bar{N} so not to confuse them with the original M and N .

$$\begin{aligned}\bar{M} &= \mu M \\ &= \frac{y^2+2}{y^2+1} \left(-1 - \frac{1}{x} + \frac{1}{y^2+2} + \frac{1}{x(y^2+2)} \right) \\ &= \frac{-x-1}{x}\end{aligned}$$

And

$$\begin{aligned}\bar{N} &= \mu N \\ &= \frac{y^2 + 2}{y^2 + 1} \quad (1) \\ &= \frac{y^2 + 2}{y^2 + 1}\end{aligned}$$

So now a modified ODE is obtained from the original ODE which will be exact and can be solved using the standard method. The modified ODE is

$$\begin{aligned}\bar{M} + \bar{N} \frac{dy}{dx} &= 0 \\ \left(\frac{-x-1}{x} \right) + \left(\frac{y^2+2}{y^2+1} \right) \frac{dy}{dx} &= 0\end{aligned}$$

The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = \bar{M} \quad (1)$$

$$\frac{\partial \phi}{\partial y} = \bar{N} \quad (2)$$

Integrating (1) w.r.t. x gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int \bar{M} dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int \frac{-x-1}{x} dx \\ \phi &= -x - \ln(x) + f(y)\end{aligned} \quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = 0 + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = \frac{y^2+2}{y^2+1}$. Therefore equation (4) becomes

$$\frac{y^2+2}{y^2+1} = 0 + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = \frac{y^2 + 2}{y^2 + 1}$$

Integrating the above w.r.t y gives

$$\int f'(y) dy = \int \left(\frac{y^2 + 2}{y^2 + 1} \right) dy$$

$$f(y) = y + \arctan(y) + c_1$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = -x - \ln(x) + y + \arctan(y) + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into the constant c_1 gives the solution as

$$c_1 = -x - \ln(x) + y + \arctan(y)$$

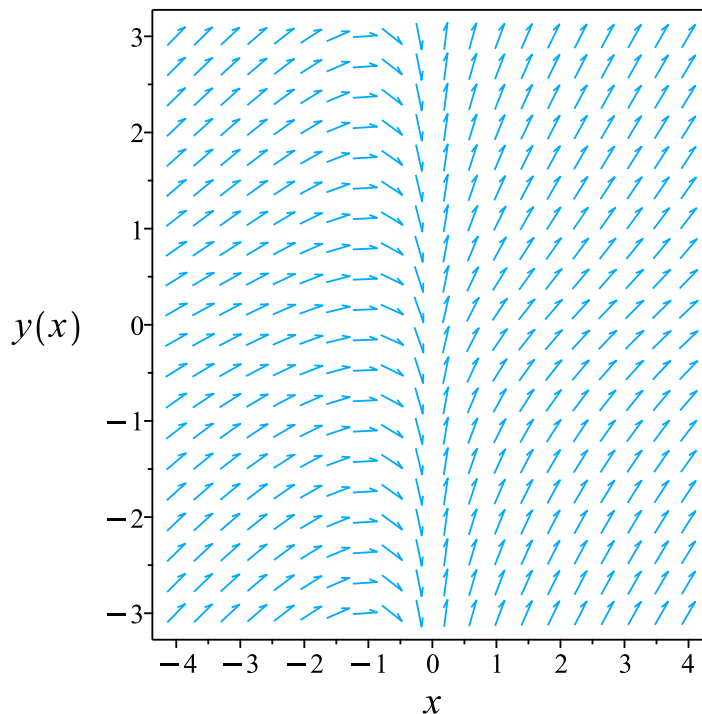


Figure 2.18: Slope field plot

$$y' = 1 + \frac{1}{x} - \frac{1}{y^2+2} - \frac{1}{x(y^2+2)}$$

Summary of solutions found

$$y + \arctan(y) - x - \ln(x) = c_1$$

Maple step by step solution

Let's solve

$$y' = 1 + \frac{1}{x} - \frac{1}{y^2+2} - \frac{1}{x(y^2+2)}$$

- Highest derivative means the order of the ODE is 1

y'

- Solve for the highest derivative

$$y' = 1 + \frac{1}{x} - \frac{1}{y^2+2} - \frac{1}{x(y^2+2)}$$

- Separate variables

$$\frac{y'(y^2+2)}{y^2+1} = \frac{x+1}{x}$$

- Integrate both sides with respect to x

$$\int \frac{y'(y^2+2)}{y^2+1} dx = \int \frac{x+1}{x} dx + C1$$

- Evaluate integral

$$y + \arctan(y) = x + \ln(x) + C1$$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
<- separable successful`
```

Maple dsolve solution

Solving time : 0.013 (sec)

Leaf size : 18

```
dsolve(diff(y(x),x) = 1+1/x-1/(y(x)^2+2)-1/x/(y(x)^2+2),
        y(x),singsol=all)
```

$$y = \tan(\text{RootOf}(-\tan(_Z) - _Z + x + \ln(x) + c_1))$$

Mathematica DSolve solution

Solving time : 0.277 (sec)

Leaf size : 19

```
DSolve[{D[y[x],x]==1+1/x-1/(y[x]^2+2)-1/(x*(y[x]^2+2))},{}},  
y[x],x,IncludeSingularSolutions->True]
```

$$y(x) \rightarrow \text{InverseFunction}[\arctan(\#1) + \#1\&][x + \log(x) + c_1]$$

2.3 Chapter IV. Methods of solution: First order equations. section 29. Problems at page 81

2.3.1	problem 1	134
2.3.2	problem 2	155
2.3.3	problem 3	195
2.3.4	problem 4	215
2.3.5	problem 5	219
2.3.6	problem 6	237
2.3.7	problem 8	264

2.3.1 problem 1

Solved as first order homogeneous class A ode	134
Solved as first order homogeneous class D2 ode	137
Solved as first order Exact ode	139
Solved as first order isobaric ode	144
Solved using Lie symmetry for first order ode	147
Maple step by step solution	153
Maple trace	153
Maple dsolve solution	153
Mathematica DSolve solution	154

Internal problem ID [18221]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter IV. Methods of solution: First order equations. section 29. Problems at page 81

Problem number : 1

Date solved : Friday, December 20, 2024 at 06:31:07 AM

CAS classification :

[[_homogeneous, 'class A'], _rational, [_Abel, '2nd type', 'class B']]

Solve

$$y^2 = x(y - x) y'$$

Solved as first order homogeneous class A ode

Time used: 2.584 (sec)

In canonical form, the ODE is

$$y' = F(x, y) = \frac{y^2}{x(y - x)} \tag{1}$$

An ode of the form $y' = \frac{M(x,y)}{N(x,y)}$ is called homogeneous if the functions $M(x, y)$ and $N(x, y)$ are both homogeneous functions and of the same order. Recall that a function $f(x, y)$ is homogeneous of order n if

$$f(t^n x, t^n y) = t^n f(x, y)$$

In this case, it can be seen that both $M = -y^2$ and $N = x(x - y)$ are both homogeneous and of the same order $n = 2$. Therefore this is a homogeneous ode. Since this ode is homogeneous, it is converted to separable ODE using the substitution $u = \frac{y}{x}$, or $y = ux$. Hence

$$\frac{dy}{dx} = \frac{du}{dx}x + u$$

Applying the transformation $y = ux$ to the above ODE in (1) gives

$$\begin{aligned}\frac{du}{dx}x + u &= \frac{u^2}{u - 1} \\ \frac{du}{dx} &= \frac{\frac{u(x)^2}{u(x)-1} - u(x)}{x}\end{aligned}$$

Or

$$u'(x) - \frac{\frac{u(x)^2}{u(x)-1} - u(x)}{x} = 0$$

Or

$$u'(x) xu(x) - u'(x) x - u(x) = 0$$

Or

$$x(u(x) - 1) u'(x) - u(x) = 0$$

Which is now solved as separable in $u(x)$.

The ode $u'(x) = \frac{u(x)}{x(u(x)-1)}$ is separable as it can be written as

$$\begin{aligned}u'(x) &= \frac{u(x)}{x(u(x) - 1)} \\ &= f(x)g(u)\end{aligned}$$

Where

$$\begin{aligned}f(x) &= \frac{1}{x} \\ g(u) &= \frac{u}{u - 1}\end{aligned}$$

Integrating gives

$$\begin{aligned}\int \frac{1}{g(u)} du &= \int f(x) dx \\ \int \frac{u-1}{u} du &= \int \frac{1}{x} dx \\ u(x) + \ln\left(\frac{1}{u(x)}\right) &= \ln(x) + c_1\end{aligned}$$

We now need to find the singular solutions, these are found by finding for what values $g(u)$ is zero, since we had to divide by this above. Solving $g(u) = 0$ or $\frac{u}{u-1} = 0$ for $u(x)$ gives

$$u(x) = 0$$

Now we go over each such singular solution and check if it verifies the ode itself and any initial conditions given. If it does not then the singular solution will not be used.

Therefore the solutions found are

$$u(x) + \ln\left(\frac{1}{u(x)}\right) = \ln(x) + c_1$$
$$u(x) = 0$$

Solving for $u(x)$ gives

$$u(x) = 0$$
$$u(x) = -\text{LambertW}\left(-\frac{e^{-c_1}}{x}\right)$$

Converting $u(x) = 0$ back to y gives

$$y = 0$$

Converting $u(x) = -\text{LambertW}\left(-\frac{e^{-c_1}}{x}\right)$ back to y gives

$$y = -x \text{LambertW}\left(-\frac{e^{-c_1}}{x}\right)$$

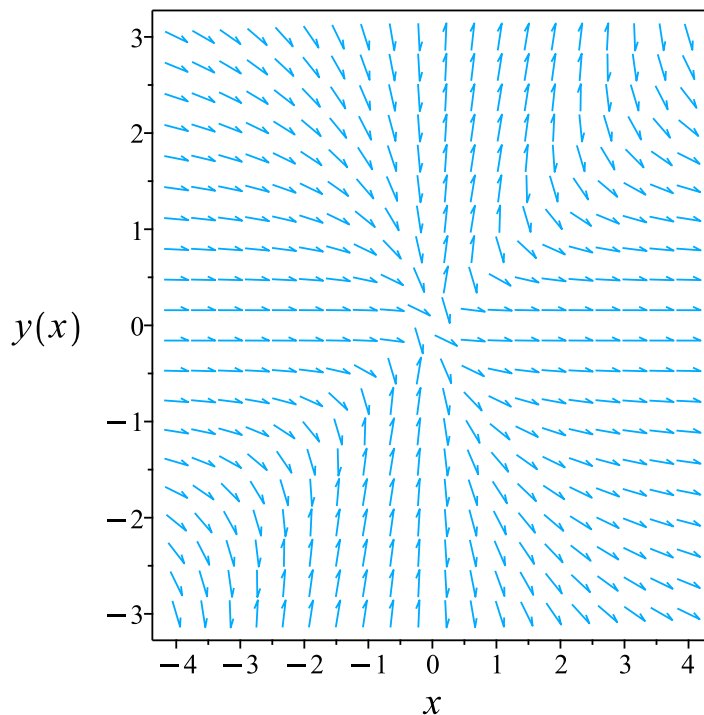


Figure 2.19: Slope field plot
 $y^2 = x(y-x)y'$

Summary of solutions found

$$y = 0$$

$$y = -x \operatorname{LambertW}\left(-\frac{e^{-c_1}}{x}\right)$$

Solved as first order homogeneous class D2 ode

Time used: 0.514 (sec)

Applying change of variables $y = u(x)x$, then the ode becomes

$$u(x)^2 x^2 = x(u(x)x - x)(u'(x)x + u(x))$$

Which is now solved The ode $u'(x) = \frac{u(x)}{(u(x)-1)x}$ is separable as it can be written as

$$\begin{aligned} u'(x) &= \frac{u(x)}{(u(x)-1)x} \\ &= f(x)g(u) \end{aligned}$$

Where

$$f(x) = \frac{1}{x}$$

$$g(u) = \frac{u}{u-1}$$

Integrating gives

$$\int \frac{1}{g(u)} du = \int f(x) dx$$

$$\int \frac{u-1}{u} du = \int \frac{1}{x} dx$$

$$u(x) + \ln\left(\frac{1}{u(x)}\right) = \ln(x) + c_1$$

We now need to find the singular solutions, these are found by finding for what values $g(u)$ is zero, since we had to divide by this above. Solving $g(u) = 0$ or $\frac{u}{u-1} = 0$ for $u(x)$ gives

$$u(x) = 0$$

Now we go over each such singular solution and check if it verifies the ode itself and any initial conditions given. If it does not then the singular solution will not be used.

Therefore the solutions found are

$$u(x) + \ln\left(\frac{1}{u(x)}\right) = \ln(x) + c_1$$

$$u(x) = 0$$

Solving for $u(x)$ gives

$$u(x) = 0$$

$$u(x) = -\text{LambertW}\left(-\frac{e^{-c_1}}{x}\right)$$

Converting $u(x) = 0$ back to y gives

$$y = 0$$

Converting $u(x) = -\text{LambertW}\left(-\frac{e^{-c_1}}{x}\right)$ back to y gives

$$y = -x \text{LambertW}\left(-\frac{e^{-c_1}}{x}\right)$$

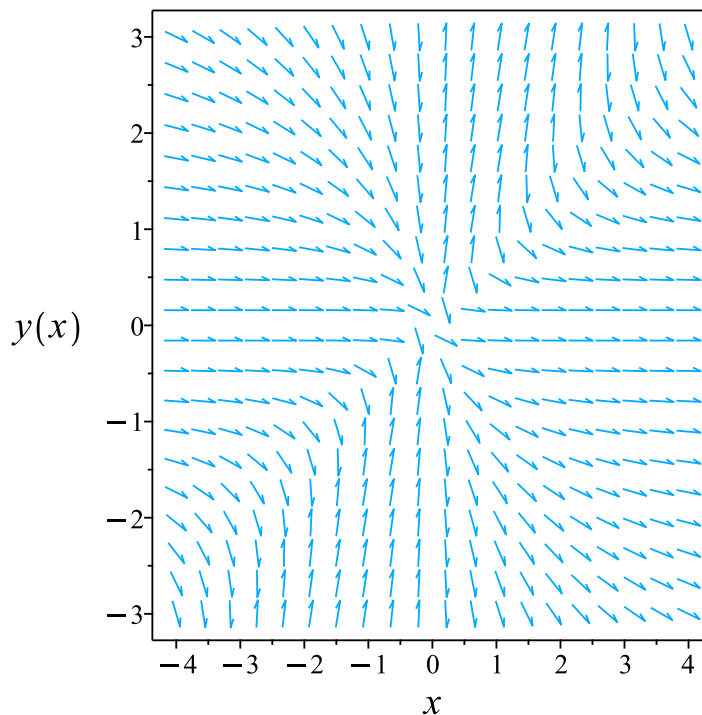


Figure 2.20: Slope field plot
 $y^2 = x(y - x)y'$

Summary of solutions found

$$y = 0$$

$$y = -x \operatorname{LambertW}\left(-\frac{e^{-c_1}}{x}\right)$$

Solved as first order Exact ode

Time used: 2.038 (sec)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx}\phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \tag{1A}$$

Therefore

$$\begin{aligned}(-x(y-x)) dy &= (-y^2) dx \\ (y^2) dx + (-x(y-x)) dy &= 0\end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(x, y) &= y^2 \\ N(x, y) &= -x(y-x)\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(y^2) \\ &= 2y\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(-x(y-x)) \\ &= 2x - y\end{aligned}$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$, then the ODE is not exact. By inspection $\frac{1}{x^2y}$ is an integrating factor. Therefore by multiplying $M = y^2$ and $N = -x(y - x)$ by this integrating factor the ode becomes exact. The new M, N are

$$M = \frac{y}{x^2}$$

$$N = -\frac{y-x}{xy}$$

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx}\phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\frac{\partial \phi}{\partial x} = M$$

$$\frac{\partial \phi}{\partial y} = N$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (\text{1A})$$

Therefore

$$\begin{aligned} \left(-\frac{y-x}{xy}\right) dy &= \left(-\frac{y}{x^2}\right) dx \\ \left(\frac{y}{x^2}\right) dx + \left(-\frac{y-x}{xy}\right) dy &= 0 \end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= \frac{y}{x^2} \\ N(x, y) &= -\frac{y-x}{xy} \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} \left(\frac{y}{x^2}\right) \\ &= \frac{1}{x^2} \end{aligned}$$

And

$$\begin{aligned} \frac{\partial N}{\partial x} &= \frac{\partial}{\partial x} \left(-\frac{y-x}{xy}\right) \\ &= \frac{1}{x^2} \end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \quad (1)$$

$$\frac{\partial \phi}{\partial y} = N \quad (2)$$

Integrating (1) w.r.t. x gives

$$\begin{aligned} \int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int \frac{y}{x^2} dx \\ \phi &= -\frac{y}{x} + f(y) \end{aligned} \quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = -\frac{1}{x} + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = -\frac{y-x}{xy}$. Therefore equation (4) becomes

$$-\frac{y-x}{xy} = -\frac{1}{x} + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = \frac{1}{y}$$

Integrating the above w.r.t y gives

$$\int f'(y) dy = \int \left(\frac{1}{y}\right) dy$$

$$f(y) = \ln(y) + c_1$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = -\frac{y}{x} + \ln(y) + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into the constant c_1 gives the solution as

$$c_1 = -\frac{y}{x} + \ln(y)$$

Solving for y gives

$$y = e^{-\text{LambertW}\left(-\frac{c_1}{x}\right) + c_1}$$

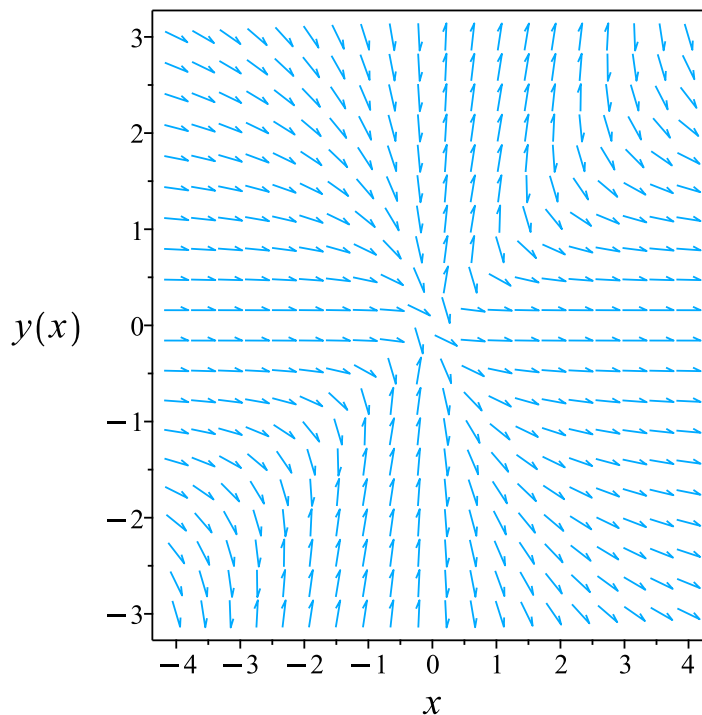


Figure 2.21: Slope field plot
 $y^2 = x(y-x)y'$

Summary of solutions found

$$y = e^{-\text{LambertW}\left(-\frac{e^{c_1}}{x}\right) + c_1}$$

Solved as first order isobaric ode

Time used: 1.270 (sec)

Solving for y' gives

$$y' = \frac{y^2}{x(y-x)} \quad (1)$$

Each of the above ode's is now solved An ode $y' = f(x, y)$ is isobaric if

$$f(tx, t^m y) = t^{m-1} f(x, y) \quad (1)$$

Where here

$$f(x, y) = \frac{y^2}{x(y-x)} \quad (2)$$

m is the order of isobaric. Substituting (2) into (1) and solving for m gives

$$m = 1$$

Since the ode is isobaric of order $m = 1$, then the substitution

$$\begin{aligned} y &= ux^m \\ &= ux \end{aligned}$$

Converts the ODE to a separable in $u(x)$. Performing this substitution gives

$$u(x) + xu'(x) = \frac{xu(x)^2}{xu(x) - x}$$

The ode $u'(x) = \frac{u(x)}{(u(x)-1)x}$ is separable as it can be written as

$$\begin{aligned} u'(x) &= \frac{u(x)}{(u(x) - 1)x} \\ &= f(x)g(u) \end{aligned}$$

Where

$$\begin{aligned} f(x) &= \frac{1}{x} \\ g(u) &= \frac{u}{u - 1} \end{aligned}$$

Integrating gives

$$\begin{aligned} \int \frac{1}{g(u)} du &= \int f(x) dx \\ \int \frac{u-1}{u} du &= \int \frac{1}{x} dx \\ u(x) + \ln\left(\frac{1}{u(x)}\right) &= \ln(x) + c_1 \end{aligned}$$

We now need to find the singular solutions, these are found by finding for what values $g(u)$ is zero, since we had to divide by this above. Solving $g(u) = 0$ or $\frac{u}{u-1} = 0$ for $u(x)$ gives

$$u(x) = 0$$

Now we go over each such singular solution and check if it verifies the ode itself and any initial conditions given. If it does not then the singular solution will not be used.

Therefore the solutions found are

$$u(x) + \ln\left(\frac{1}{u(x)}\right) = \ln(x) + c_1$$
$$u(x) = 0$$

Solving for $u(x)$ gives

$$u(x) = 0$$
$$u(x) = -\text{LambertW}\left(-\frac{e^{-c_1}}{x}\right)$$

Converting $u(x) = 0$ back to y gives

$$\frac{y}{x} = 0$$

Converting $u(x) = -\text{LambertW}\left(-\frac{e^{-c_1}}{x}\right)$ back to y gives

$$\frac{y}{x} = -\text{LambertW}\left(-\frac{e^{-c_1}}{x}\right)$$

Solving for y gives

$$y = 0$$
$$y = -x \text{LambertW}\left(-\frac{e^{-c_1}}{x}\right)$$

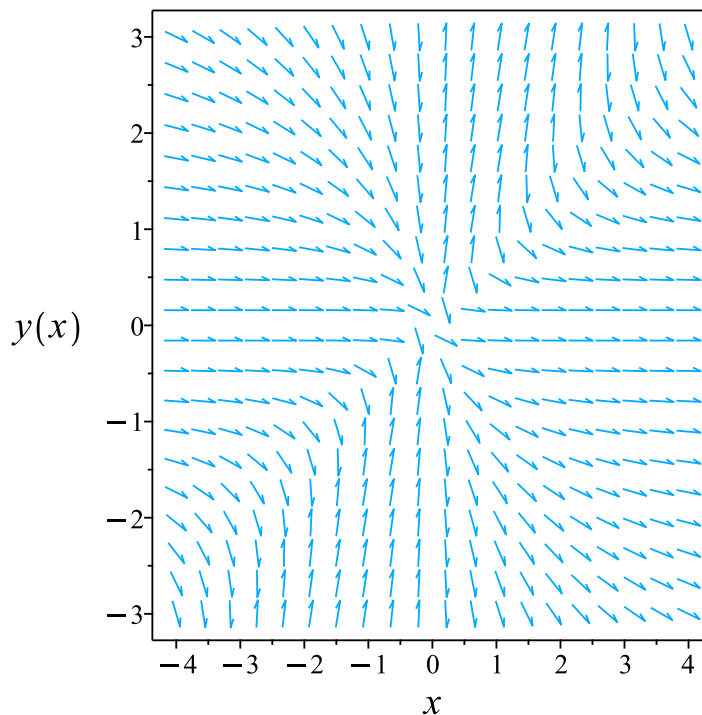


Figure 2.22: Slope field plot
 $y^2 = x(y - x)y'$

Summary of solutions found

$$y = 0$$

$$y = -x \operatorname{LambertW}\left(-\frac{e^{-c_1}}{x}\right)$$

Solved using Lie symmetry for first order ode

Time used: 2.286 (sec)

Writing the ode as

$$y' = \frac{y^2}{x(y-x)}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

To determine ξ, η then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$\xi = xa_2 + ya_3 + a_1 \quad (1E)$$

$$\eta = xb_2 + yb_3 + b_1 \quad (2E)$$

Where the unknown coefficients are

$$\{a_1, a_2, a_3, b_1, b_2, b_3\}$$

Substituting equations (1E,2E) and ω into (A) gives

$$b_2 + \frac{y^2(b_3 - a_2)}{x(y-x)} - \frac{y^4 a_3}{x^2(y-x)^2} - \left(-\frac{y^2}{x^2(y-x)} + \frac{y^2}{x(y-x)^2} \right) (xa_2 + ya_3 + a_1) \quad (5E)$$

$$- \left(\frac{2y}{x(y-x)} - \frac{y^2}{x(y-x)^2} \right) (xb_2 + yb_3 + b_1) = 0$$

Putting the above in normal form gives

$$\frac{x^4 b_2 - x^2 y^2 a_2 + x^2 y^2 b_3 - 2x y^3 a_3 + 2x^2 y b_1 - 2x y^2 a_1 - x y^2 b_1 + y^3 a_1}{x^2 (x-y)^2} = 0$$

Setting the numerator to zero gives

$$x^4 b_2 - x^2 y^2 a_2 + x^2 y^2 b_3 - 2x y^3 a_3 + 2x^2 y b_1 - 2x y^2 a_1 - x y^2 b_1 + y^3 a_1 = 0 \quad (6E)$$

Looking at the above PDE shows the following are all the terms with $\{x, y\}$ in them.

$$\{x, y\}$$

The following substitution is now made to be able to collect on all terms with $\{x, y\}$ in them

$$\{x = v_1, y = v_2\}$$

The above PDE (6E) now becomes

$$-a_2 v_1^2 v_2^2 - 2a_3 v_1 v_2^3 + b_2 v_1^4 + b_3 v_1^2 v_2^2 - 2a_1 v_1 v_2^2 + a_1 v_2^3 + 2b_1 v_1^2 v_2 - b_1 v_1 v_2^2 = 0 \quad (7E)$$

Collecting the above on the terms v_i introduced, and these are

$$\{v_1, v_2\}$$

Equation (7E) now becomes

$$b_2 v_1^4 + (b_3 - a_2) v_1^2 v_2^2 + 2b_1 v_1^2 v_2 - 2a_3 v_1 v_2^3 + (-2a_1 - b_1) v_1 v_2^2 + a_1 v_2^3 = 0 \quad (8E)$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$\begin{aligned} a_1 &= 0 \\ b_2 &= 0 \\ -2a_3 &= 0 \\ 2b_1 &= 0 \\ -2a_1 - b_1 &= 0 \\ b_3 - a_2 &= 0 \end{aligned}$$

Solving the above equations for the unknowns gives

$$\begin{aligned} a_1 &= 0 \\ a_2 &= b_3 \\ a_3 &= 0 \\ b_1 &= 0 \\ b_2 &= 0 \\ b_3 &= b_3 \end{aligned}$$

Substituting the above solution in the anstaz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$\begin{aligned} \xi &= x \\ \eta &= y \end{aligned}$$

Shifting is now applied to make $\xi = 0$ in order to simplify the rest of the computation

$$\begin{aligned} \eta &= \eta - \omega(x, y) \xi \\ &= y - \left(\frac{y^2}{x(y-x)} \right) (x) \\ &= \frac{yx}{x-y} \\ \xi &= 0 \end{aligned}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS \quad (1)$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = x$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{\frac{yx}{x-y}} dy \end{aligned}$$

Which results in

$$S = -\frac{y}{x} + \ln(y)$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = \frac{y^2}{x(y-x)}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 1 \\ R_y &= 0 \\ S_x &= \frac{y}{x^2} \\ S_y &= \frac{x-y}{xy} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = 0 \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = 0$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S .

Since the ode has the form $\frac{d}{dR}S(R) = f(R)$, then we only need to integrate $f(R)$.

$$\int dS = \int 0 dR + c_2$$
$$S(R) = c_2$$

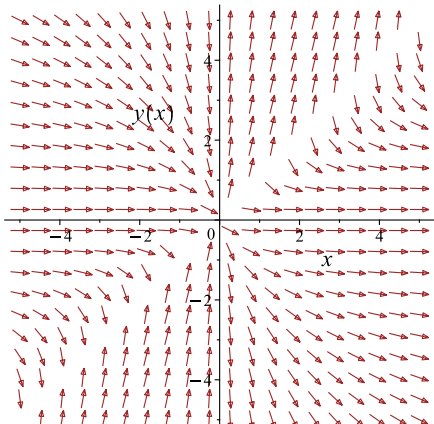
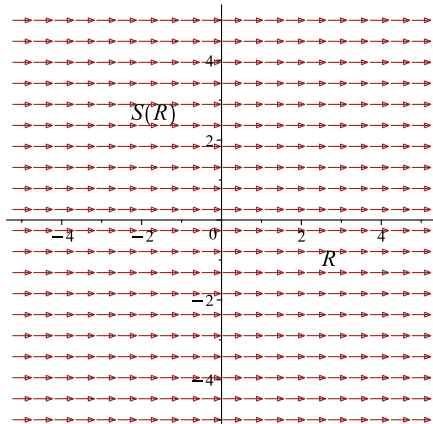
To complete the solution, we just need to transform the above back to x, y coordinates. This results in

$$\frac{\ln(y)x - y}{x} = c_2$$

Which gives

$$y = e^{-\text{LambertW}\left(-\frac{e^{c_2}}{x}\right) + c_2}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = \frac{y^2}{x(y-x)}$ 	$R = x$ $S = \frac{\ln(y) x - y}{x}$	$\frac{dS}{dR} = 0$ 

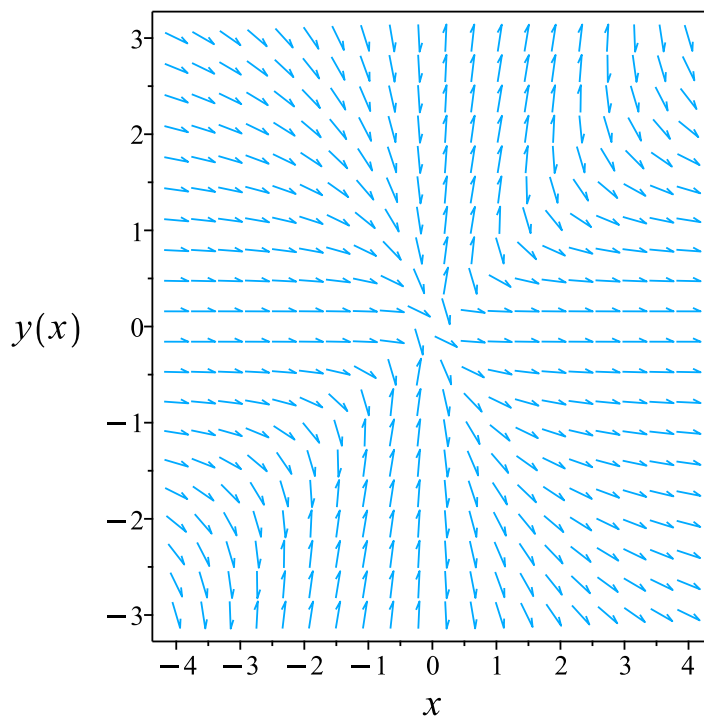


Figure 2.23: Slope field plot
 $y^2 = x(y-x)y'$

Summary of solutions found

$$y = e^{-\text{LambertW}\left(-\frac{e^{-c_2}}{x}\right)+c_2}$$

Maple step by step solution

Let's solve

$$y^2 = x(y-x)y'$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Solve for the highest derivative

$$y' = \frac{y^2}{x(y-x)}$$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
trying inverse linear
trying homogeneous types:
trying homogeneous D
<- homogeneous successful`
```

Maple dsolve solution

Solving time : 0.231 (sec)

Leaf size : 17

```
dsolve(y(x)^2 = x*(y(x)-x)*diff(y(x),x),
       y(x),singsol=all)
```

$$y = -x \text{LambertW}\left(-\frac{e^{-c_1}}{x}\right)$$

Mathematica DSolve solution

Solving time : 1.931 (sec)

Leaf size : 25

```
DSolve[{y[x]^2==x*(y[x]-x)*D[y[x],x],{}},  
y[x],x,IncludeSingularSolutions->True]
```

$$y(x) \rightarrow -xW\left(-\frac{e^{-c_1}}{x}\right)$$

$$y(x) \rightarrow 0$$

2.3.2 problem 2

Solved as first order homogeneous class A ode	155
Solved as first order homogeneous class D2 ode	159
Solved as first order homogeneous class Maple C ode	162
Solved as first order Bernoulli ode	167
Solved as first order isobaric ode	171
Solved using Lie symmetry for first order ode	175
Solved as first order ode of type dAlembert	180
Maple step by step solution	193
Maple trace	193
Maple dsolve solution	194
Mathematica DSolve solution	194

Internal problem ID [18222]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter IV. Methods of solution: First order equations. section 29. Problems at page 81

Problem number : 2

Date solved : Friday, December 20, 2024 at 10:20:09 AM

CAS classification : [[_homogeneous, 'class A'], _rational, _Bernoulli]

Solve

$$2x^2y + y^3 - x^3y' = 0$$

Solved as first order homogeneous class A ode

Time used: 0.669 (sec)

In canonical form, the ODE is

$$\begin{aligned}
 y' &= F(x, y) \\
 &= \frac{y(2x^2 + y^2)}{x^3}
 \end{aligned}
 \tag{1}$$

An ode of the form $y' = \frac{M(x,y)}{N(x,y)}$ is called homogeneous if the functions $M(x, y)$ and $N(x, y)$ are both homogeneous functions and of the same order. Recall that a function $f(x, y)$ is homogeneous of order n if

$$f(t^n x, t^n y) = t^n f(x, y)$$

In this case, it can be seen that both $M = y(2x^2 + y^2)$ and $N = x^3$ are both homogeneous and of the same order $n = 3$. Therefore this is a homogeneous ode. Since this ode is homogeneous, it is converted to separable ODE using the substitution $u = \frac{y}{x}$, or $y = ux$. Hence

$$\frac{dy}{dx} = \frac{du}{dx}x + u$$

Applying the transformation $y = ux$ to the above ODE in (1) gives

$$\begin{aligned} \frac{du}{dx}x + u &= u^3 + 2u \\ \frac{du}{dx} &= \frac{u(x)^3 + u(x)}{x} \end{aligned}$$

Or

$$u'(x) - \frac{u(x)^3 + u(x)}{x} = 0$$

Or

$$-u(x)^3 + u'(x)x - u(x) = 0$$

Which is now solved as separable in $u(x)$.

The ode $u'(x) = \frac{u(x)(u(x)^2+1)}{x}$ is separable as it can be written as

$$\begin{aligned} u'(x) &= \frac{u(x)(u(x)^2 + 1)}{x} \\ &= f(x)g(u) \end{aligned}$$

Where

$$\begin{aligned} f(x) &= \frac{1}{x} \\ g(u) &= u(u^2 + 1) \end{aligned}$$

Integrating gives

$$\begin{aligned} \int \frac{1}{g(u)} du &= \int f(x) dx \\ \int \frac{1}{u(u^2 + 1)} du &= \int \frac{1}{x} dx \\ \ln \left(\frac{u(x)}{\sqrt{u(x)^2 + 1}} \right) &= \ln(x) + c_1 \end{aligned}$$

We now need to find the singular solutions, these are found by finding for what values $g(u)$ is zero, since we had to divide by this above. Solving $g(u) = 0$ or $u(u^2 + 1) = 0$ for $u(x)$ gives

$$u(x) = 0$$

$$u(x) = -i$$

$$u(x) = i$$

Now we go over each such singular solution and check if it verifies the ode itself and any initial conditions given. If it does not then the singular solution will not be used.

Therefore the solutions found are

$$\ln \left(\frac{u(x)}{\sqrt{u(x)^2 + 1}} \right) = \ln(x) + c_1$$

$$u(x) = 0$$

$$u(x) = -i$$

$$u(x) = i$$

Solving for $u(x)$ gives

$$u(x) = 0$$

$$u(x) = -i$$

$$u(x) = i$$

$$u(x) = \frac{e^{c_1} x}{\sqrt{1 - x^2 e^{2c_1}}}$$

$$u(x) = -\frac{e^{c_1} x}{\sqrt{1 - x^2 e^{2c_1}}}$$

Converting $u(x) = 0$ back to y gives

$$y = 0$$

Converting $u(x) = -i$ back to y gives

$$y = -ix$$

Converting $u(x) = i$ back to y gives

$$y = ix$$

Converting $u(x) = \frac{e^{c_1}x}{\sqrt{1-x^2e^{2c_1}}}$ back to y gives

$$y = \frac{x^2e^{c_1}}{\sqrt{1-x^2e^{2c_1}}}$$

Converting $u(x) = -\frac{e^{c_1}x}{\sqrt{1-x^2e^{2c_1}}}$ back to y gives

$$y = -\frac{x^2e^{c_1}}{\sqrt{1-x^2e^{2c_1}}}$$

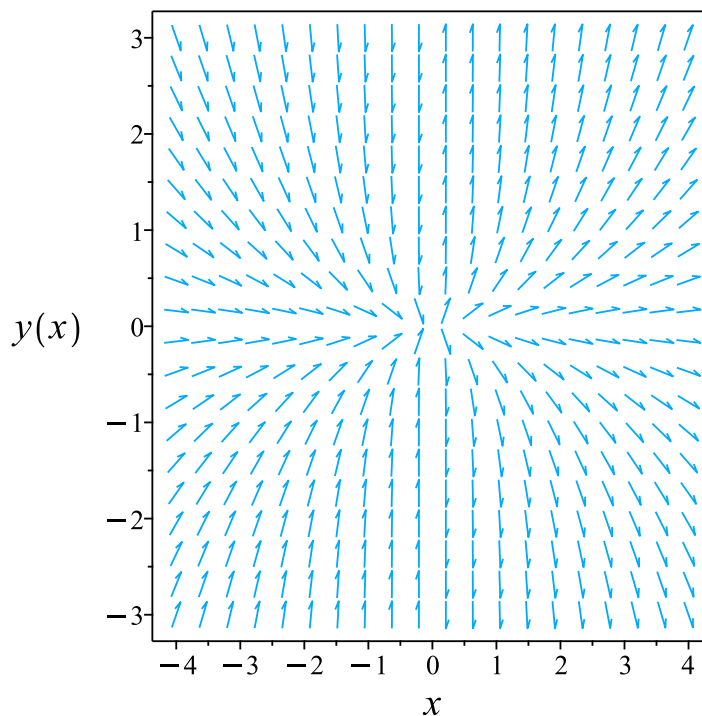


Figure 2.24: Slope field plot
 $2x^2y + y^3 - x^3y' = 0$

Summary of solutions found

$$y = 0$$

$$y = -ix$$

$$y = ix$$

$$y = \frac{x^2e^{c_1}}{\sqrt{1-x^2e^{2c_1}}}$$

$$y = -\frac{x^2e^{c_1}}{\sqrt{1-x^2e^{2c_1}}}$$

Solved as first order homogeneous class D2 ode

Time used: 0.170 (sec)

Applying change of variables $y = u(x)x$, then the ode becomes

$$2x^3u(x) + u(x)^3x^3 - x^3(u'(x)x + u(x)) = 0$$

Which is now solved The ode $u'(x) = \frac{u(x)(u(x)^2+1)}{x}$ is separable as it can be written as

$$\begin{aligned} u'(x) &= \frac{u(x)(u(x)^2+1)}{x} \\ &= f(x)g(u) \end{aligned}$$

Where

$$\begin{aligned} f(x) &= \frac{1}{x} \\ g(u) &= u(u^2+1) \end{aligned}$$

Integrating gives

$$\begin{aligned} \int \frac{1}{g(u)} du &= \int f(x) dx \\ \int \frac{1}{u(u^2+1)} du &= \int \frac{1}{x} dx \\ \ln \left(\frac{u(x)}{\sqrt{u(x)^2+1}} \right) &= \ln(x) + c_1 \end{aligned}$$

We now need to find the singular solutions, these are found by finding for what values $g(u)$ is zero, since we had to divide by this above. Solving $g(u) = 0$ or $u(u^2+1) = 0$ for $u(x)$ gives

$$\begin{aligned} u(x) &= 0 \\ u(x) &= -i \\ u(x) &= i \end{aligned}$$

Now we go over each such singular solution and check if it verifies the ode itself and any initial conditions given. If it does not then the singular solution will not be used.

Therefore the solutions found are

$$\ln \left(\frac{u(x)}{\sqrt{u(x)^2 + 1}} \right) = \ln(x) + c_1$$

$$u(x) = 0$$

$$u(x) = -i$$

$$u(x) = i$$

Solving for $u(x)$ gives

$$u(x) = 0$$

$$u(x) = -i$$

$$u(x) = i$$

$$u(x) = \frac{e^{c_1} x}{\sqrt{1 - x^2 e^{2c_1}}}$$

$$u(x) = -\frac{e^{c_1} x}{\sqrt{1 - x^2 e^{2c_1}}}$$

Converting $u(x) = 0$ back to y gives

$$y = 0$$

Converting $u(x) = -i$ back to y gives

$$y = -ix$$

Converting $u(x) = i$ back to y gives

$$y = ix$$

Converting $u(x) = \frac{e^{c_1} x}{\sqrt{1 - x^2 e^{2c_1}}}$ back to y gives

$$y = \frac{x^2 e^{c_1}}{\sqrt{1 - x^2 e^{2c_1}}}$$

Converting $u(x) = -\frac{e^{c_1} x}{\sqrt{1 - x^2 e^{2c_1}}}$ back to y gives

$$y = -\frac{x^2 e^{c_1}}{\sqrt{1 - x^2 e^{2c_1}}}$$

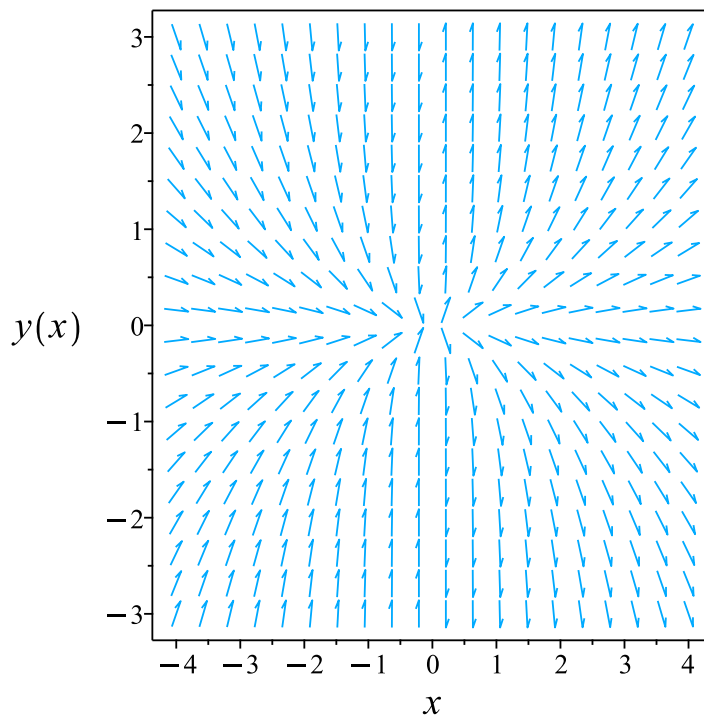


Figure 2.25: Slope field plot
 $2x^2y + y^3 - x^3y' = 0$

Summary of solutions found

$$y = 0$$

$$y = -ix$$

$$y = ix$$

$$y = \frac{x^2 e^{c_1}}{\sqrt{1 - x^2 e^{2c_1}}}$$

$$y = -\frac{x^2 e^{c_1}}{\sqrt{1 - x^2 e^{2c_1}}}$$

Summary of solutions found

$$y = 0$$

$$y = -ix$$

$$y = ix$$

$$y = \frac{x^2 e^{c_1}}{\sqrt{1 - x^2 e^{2c_1}}}$$

$$y = -\frac{x^2 e^{c_1}}{\sqrt{1 - x^2 e^{2c_1}}}$$

Solved as first order homogeneous class Maple C ode

Time used: 0.674 (sec)

Let $Y = y - y_0$ and $X = x - x_0$ then the above is transformed to new ode in $Y(X)$

$$\frac{d}{dX}Y(X) = \frac{(Y(X) + y_0) ((Y(X) + y_0)^2 + 2(x_0 + X)^2)}{(x_0 + X)^3}$$

Solving for possible values of x_0 and y_0 which makes the above ode a homogeneous ode results in

$$x_0 = 0$$

$$y_0 = 0$$

Using these values now it is possible to easily solve for $Y(X)$. The above ode now becomes

$$\frac{d}{dX}Y(X) = \frac{2Y(X)X^2 + Y(X)^3}{X^3}$$

In canonical form, the ODE is

$$\begin{aligned} Y' &= F(X, Y) \\ &= \frac{Y(2X^2 + Y^2)}{X^3} \end{aligned} \tag{1}$$

An ode of the form $Y' = \frac{M(X, Y)}{N(X, Y)}$ is called homogeneous if the functions $M(X, Y)$ and $N(X, Y)$ are both homogeneous functions and of the same order. Recall that a function $f(X, Y)$ is homogeneous of order n if

$$f(t^n X, t^n Y) = t^n f(X, Y)$$

In this case, it can be seen that both $M = Y(2X^2 + Y^2)$ and $N = X^3$ are both homogeneous and of the same order $n = 3$. Therefore this is a homogeneous ode. Since this ode is homogeneous, it is converted to separable ODE using the substitution $u = \frac{Y}{X}$, or $Y = uX$. Hence

$$\frac{dY}{dX} = \frac{du}{dX}X + u$$

Applying the transformation $Y = uX$ to the above ODE in (1) gives

$$\begin{aligned}\frac{du}{dX}X + u &= u^3 + 2u \\ \frac{du}{dX} &= \frac{u(X)^3 + u(X)}{X}\end{aligned}$$

Or

$$\frac{d}{dX}u(X) - \frac{u(X)^3 + u(X)}{X} = 0$$

Or

$$-u(X)^3 + \left(\frac{d}{dX}u(X)\right)X - u(X) = 0$$

Which is now solved as separable in $u(X)$.

The ode $\frac{d}{dX}u(X) = \frac{u(X)(u(X)^2+1)}{X}$ is separable as it can be written as

$$\begin{aligned}\frac{d}{dX}u(X) &= \frac{u(X)(u(X)^2+1)}{X} \\ &= f(X)g(u)\end{aligned}$$

Where

$$\begin{aligned}f(X) &= \frac{1}{X} \\ g(u) &= u(u^2 + 1)\end{aligned}$$

Integrating gives

$$\begin{aligned}\int \frac{1}{g(u)} du &= \int f(X) dX \\ \int \frac{1}{u(u^2+1)} du &= \int \frac{1}{X} dX \\ \ln\left(\frac{u(X)}{\sqrt{u(X)^2+1}}\right) &= \ln(X) + c_1\end{aligned}$$

We now need to find the singular solutions, these are found by finding for what values $g(u)$ is zero, since we had to divide by this above. Solving $g(u) = 0$ or $u(u^2 + 1) = 0$ for $u(X)$ gives

$$\begin{aligned}u(X) &= 0 \\ u(X) &= -i \\ u(X) &= i\end{aligned}$$

Now we go over each such singular solution and check if it verifies the ode itself and any initial conditions given. If it does not then the singular solution will not be used.

Therefore the solutions found are

$$\ln \left(\frac{u(X)}{\sqrt{u(X)^2 + 1}} \right) = \ln(X) + c_1$$

$$u(X) = 0$$

$$u(X) = -i$$

$$u(X) = i$$

Solving for $u(X)$ gives

$$u(X) = 0$$

$$u(X) = -i$$

$$u(X) = i$$

$$u(X) = \frac{e^{c_1} X}{\sqrt{1 - X^2 e^{2c_1}}}$$

$$u(X) = -\frac{e^{c_1} X}{\sqrt{1 - X^2 e^{2c_1}}}$$

Converting $u(X) = 0$ back to $Y(X)$ gives

$$Y(X) = 0$$

Converting $u(X) = -i$ back to $Y(X)$ gives

$$Y(X) = -iX$$

Converting $u(X) = i$ back to $Y(X)$ gives

$$Y(X) = iX$$

Converting $u(X) = \frac{e^{c_1} X}{\sqrt{1 - X^2 e^{2c_1}}}$ back to $Y(X)$ gives

$$Y(X) = \frac{X^2 e^{c_1}}{\sqrt{1 - X^2 e^{2c_1}}}$$

Converting $u(X) = -\frac{e^{c_1} X}{\sqrt{1 - X^2 e^{2c_1}}}$ back to $Y(X)$ gives

$$Y(X) = -\frac{X^2 e^{c_1}}{\sqrt{1 - X^2 e^{2c_1}}}$$

Using the solution for $Y(X)$

$$Y(X) = 0 \tag{A}$$

And replacing back terms in the above solution using

$$\begin{aligned} Y &= y + y_0 \\ X &= x + x_0 \end{aligned}$$

Or

$$\begin{aligned} Y &= y \\ X &= x \end{aligned}$$

Then the solution in y becomes using EQ (A)

$$y = 0$$

Using the solution for $Y(X)$

$$Y(X) = -iX \tag{A}$$

And replacing back terms in the above solution using

$$\begin{aligned} Y &= y + y_0 \\ X &= x + x_0 \end{aligned}$$

Or

$$\begin{aligned} Y &= y \\ X &= x \end{aligned}$$

Then the solution in y becomes using EQ (A)

$$y = -ix$$

Using the solution for $Y(X)$

$$Y(X) = iX \tag{A}$$

And replacing back terms in the above solution using

$$\begin{aligned} Y &= y + y_0 \\ X &= x + x_0 \end{aligned}$$

Or

$$Y = y$$

$$X = x$$

Then the solution in y becomes using EQ (A)

$$y = ix$$

Using the solution for $Y(X)$

$$Y(X) = \frac{X^2 e^{c_1}}{\sqrt{1 - X^2 e^{2c_1}}} \quad (\text{A})$$

And replacing back terms in the above solution using

$$Y = y + y_0$$

$$X = x + x_0$$

Or

$$Y = y$$

$$X = x$$

Then the solution in y becomes using EQ (A)

$$y = \frac{x^2 e^{c_1}}{\sqrt{1 - x^2 e^{2c_1}}}$$

Using the solution for $Y(X)$

$$Y(X) = -\frac{X^2 e^{c_1}}{\sqrt{1 - X^2 e^{2c_1}}} \quad (\text{A})$$

And replacing back terms in the above solution using

$$Y = y + y_0$$

$$X = x + x_0$$

Or

$$Y = y$$

$$X = x$$

Then the solution in y becomes using EQ (A)

$$y = -\frac{x^2 e^{c_1}}{\sqrt{1 - x^2 e^{2c_1}}}$$

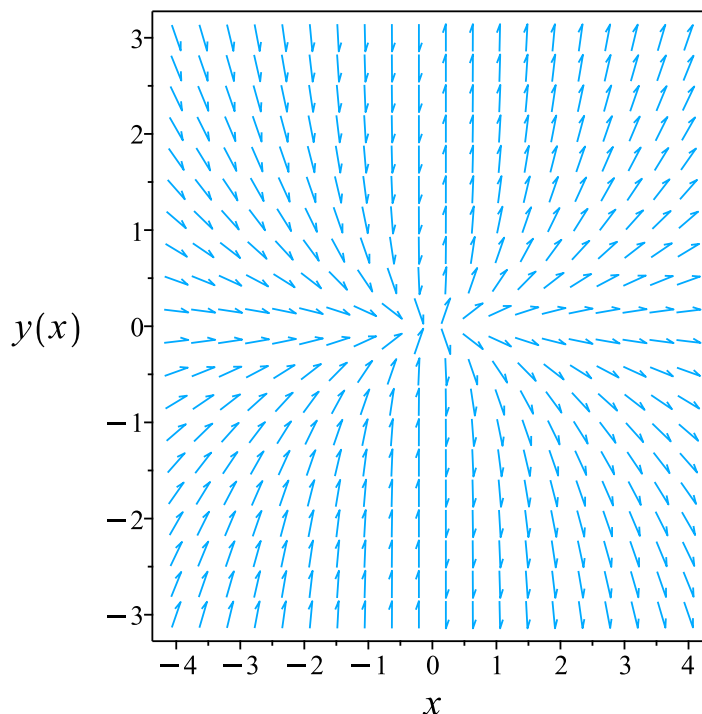


Figure 2.26: Slope field plot
 $2x^2y + y^3 - x^3y' = 0$

Solved as first order Bernoulli ode

Time used: 0.218 (sec)

In canonical form, the ODE is

$$\begin{aligned} y' &= F(x, y) \\ &= \frac{y(2x^2 + y^2)}{x^3} \end{aligned}$$

This is a Bernoulli ODE.

$$y' = \left(\frac{2}{x}\right)y + \left(\frac{1}{x^3}\right)y^3 \quad (1)$$

The standard Bernoulli ODE has the form

$$y' = f_0(x)y + f_1(x)y^n \quad (2)$$

Comparing this to (1) shows that

$$\begin{aligned} f_0 &= \frac{2}{x} \\ f_1 &= \frac{1}{x^3} \end{aligned}$$

The first step is to divide the above equation by y^n which gives

$$\frac{y'}{y^n} = f_0(x)y^{1-n} + f_1(x) \quad (3)$$

The next step is use the substitution $v = y^{1-n}$ in equation (3) which generates a new ODE in $v(x)$ which will be linear and can be easily solved using an integrating factor. Backsubstitution then gives the solution $y(x)$ which is what we want.

This method is now applied to the ODE at hand. Comparing the ODE (1) With (2) Shows that

$$\begin{aligned} f_0(x) &= \frac{2}{x} \\ f_1(x) &= \frac{1}{x^3} \\ n &= 3 \end{aligned}$$

Dividing both sides of ODE (1) by $y^n = y^3$ gives

$$y' \frac{1}{y^3} = \frac{2}{x y^2} + \frac{1}{x^3} \quad (4)$$

Let

$$\begin{aligned} v &= y^{1-n} \\ &= \frac{1}{y^2} \end{aligned} \quad (5)$$

Taking derivative of equation (5) w.r.t x gives

$$v' = -\frac{2}{y^3} y' \quad (6)$$

Substituting equations (5) and (6) into equation (4) gives

$$\begin{aligned} -\frac{v'(x)}{2} &= \frac{2v(x)}{x} + \frac{1}{x^3} \\ v' &= -\frac{4v}{x} - \frac{2}{x^3} \end{aligned} \quad (7)$$

The above now is a linear ODE in $v(x)$ which is now solved.

In canonical form a linear first order is

$$v'(x) + q(x)v(x) = p(x)$$

Comparing the above to the given ode shows that

$$q(x) = \frac{4}{x}$$
$$p(x) = -\frac{2}{x^3}$$

The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int q dx} \\ &= e^{\int \frac{4}{x} dx} \\ &= x^4\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dx}(\mu v) &= \mu p \\ \frac{d}{dx}(\mu v) &= (\mu) \left(-\frac{2}{x^3}\right) \\ \frac{d}{dx}(v x^4) &= (x^4) \left(-\frac{2}{x^3}\right) \\ d(v x^4) &= (-2x) dx\end{aligned}$$

Integrating gives

$$\begin{aligned}v x^4 &= \int -2x dx \\ &= -x^2 + c_1\end{aligned}$$

Dividing throughout by the integrating factor x^4 gives the final solution

$$v(x) = \frac{-x^2 + c_1}{x^4}$$

The substitution $v = y^{1-n}$ is now used to convert the above solution back to y which results in

$$\frac{1}{y^2} = \frac{-x^2 + c_1}{x^4}$$

Solving for y gives

$$y = \frac{x^2}{\sqrt{-x^2 + c_1}}$$

$$y = -\frac{x^2}{\sqrt{-x^2 + c_1}}$$

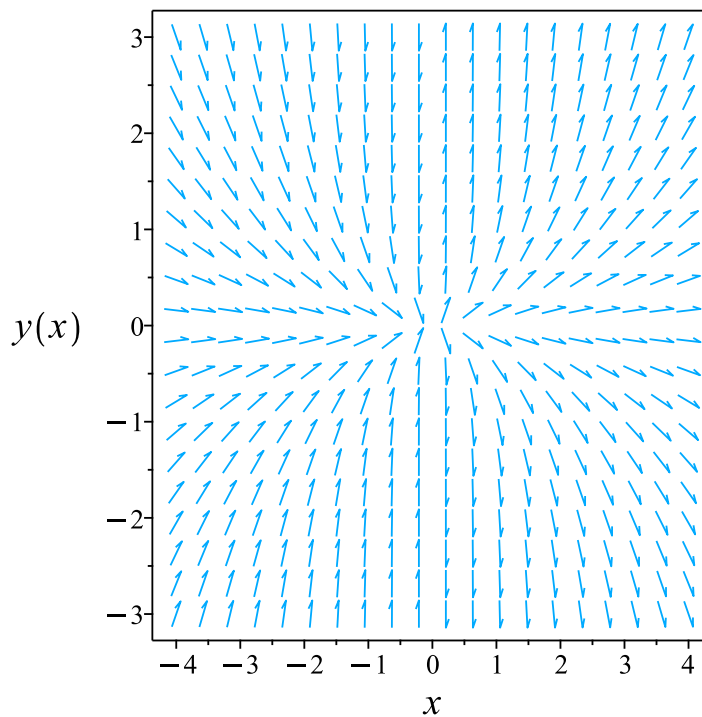


Figure 2.27: Slope field plot
 $2x^2y + y^3 - x^3y' = 0$

Summary of solutions found

$$y = \frac{x^2}{\sqrt{-x^2 + c_1}}$$

$$y = -\frac{x^2}{\sqrt{-x^2 + c_1}}$$

Solved as first order isobaric ode

Time used: 0.273 (sec)

Solving for y' gives

$$y' = \frac{y(y^2 + 2x^2)}{x^3} \quad (1)$$

Each of the above ode's is now solved An ode $y' = f(x, y)$ is isobaric if

$$f(tx, t^m y) = t^{m-1} f(x, y) \quad (1)$$

Where here

$$f(x, y) = \frac{y(y^2 + 2x^2)}{x^3} \quad (2)$$

 m is the order of isobaric. Substituting (2) into (1) and solving for m gives

$$m = 1$$

Since the ode is isobaric of order $m = 1$, then the substitution

$$\begin{aligned} y &= ux^m \\ &= ux \end{aligned}$$

Converts the ODE to a separable in $u(x)$. Performing this substitution gives

$$u(x) + xu'(x) = \frac{u(x)(x^2u(x)^2 + 2x^2)}{x^2}$$

The ode $u'(x) = \frac{u(x)(u(x)^2+1)}{x}$ is separable as it can be written as

$$\begin{aligned} u'(x) &= \frac{u(x)(u(x)^2 + 1)}{x} \\ &= f(x)g(u) \end{aligned}$$

Where

$$\begin{aligned} f(x) &= \frac{1}{x} \\ g(u) &= u(u^2 + 1) \end{aligned}$$

Integrating gives

$$\int \frac{1}{g(u)} du = \int f(x) dx$$

$$\int \frac{1}{u(u^2 + 1)} du = \int \frac{1}{x} dx$$

$$\ln \left(\frac{u(x)}{\sqrt{u(x)^2 + 1}} \right) = \ln(x) + c_1$$

We now need to find the singular solutions, these are found by finding for what values $g(u)$ is zero, since we had to divide by this above. Solving $g(u) = 0$ or $u(u^2 + 1) = 0$ for $u(x)$ gives

$$u(x) = 0$$

$$u(x) = -i$$

$$u(x) = i$$

Now we go over each such singular solution and check if it verifies the ode itself and any initial conditions given. If it does not then the singular solution will not be used.

Therefore the solutions found are

$$\ln \left(\frac{u(x)}{\sqrt{u(x)^2 + 1}} \right) = \ln(x) + c_1$$

$$u(x) = 0$$

$$u(x) = -i$$

$$u(x) = i$$

Solving for $u(x)$ gives

$$u(x) = 0$$

$$u(x) = -i$$

$$u(x) = i$$

$$u(x) = \frac{e^{c_1} x}{\sqrt{1 - x^2 e^{2c_1}}}$$

$$u(x) = -\frac{e^{c_1} x}{\sqrt{1 - x^2 e^{2c_1}}}$$

Converting $u(x) = 0$ back to y gives

$$\frac{y}{x} = 0$$

Converting $u(x) = -i$ back to y gives

$$\frac{y}{x} = -i$$

Converting $u(x) = i$ back to y gives

$$\frac{y}{x} = i$$

Converting $u(x) = \frac{e^{c_1 x}}{\sqrt{1-x^2 e^{2c_1}}}$ back to y gives

$$\frac{y}{x} = \frac{e^{c_1 x}}{\sqrt{1-x^2 e^{2c_1}}}$$

Converting $u(x) = -\frac{e^{c_1 x}}{\sqrt{1-x^2 e^{2c_1}}}$ back to y gives

$$\frac{y}{x} = -\frac{e^{c_1 x}}{\sqrt{1-x^2 e^{2c_1}}}$$

Solving for y gives

$$y = 0$$

$$y = -ix$$

$$y = ix$$

$$y = \frac{x^2 e^{c_1}}{\sqrt{1-x^2 e^{2c_1}}}$$

$$y = -\frac{x^2 e^{c_1}}{\sqrt{1-x^2 e^{2c_1}}}$$

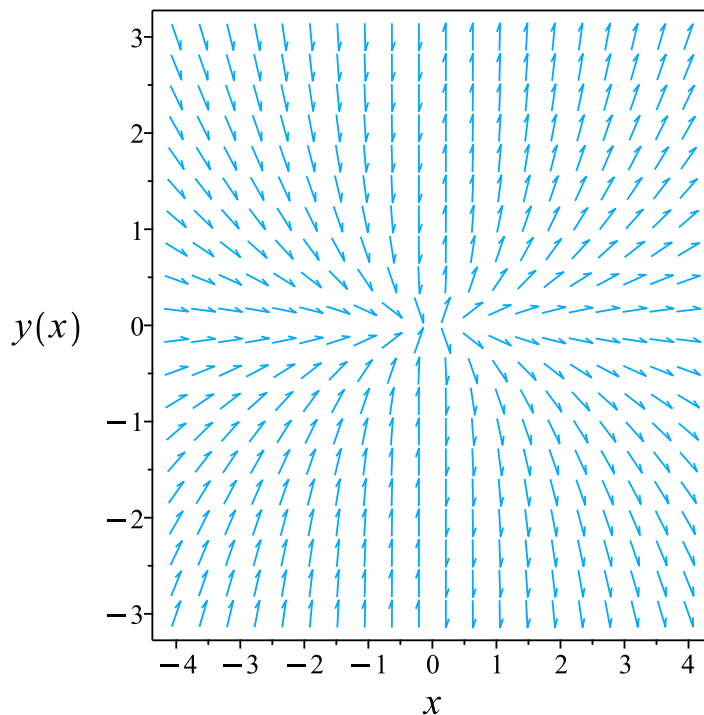


Figure 2.28: Slope field plot
 $2x^2y + y^3 - x^3y' = 0$

Summary of solutions found

$$y = 0$$

$$y = -ix$$

$$y = ix$$

$$y = \frac{x^2 e^{c_1}}{\sqrt{1 - x^2 e^{2c_1}}}$$

$$y = -\frac{x^2 e^{c_1}}{\sqrt{1 - x^2 e^{2c_1}}}$$

Solved using Lie symmetry for first order ode

Time used: 0.959 (sec)

Writing the ode as

$$y' = \frac{y(2x^2 + y^2)}{x^3}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

To determine ξ, η then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$\xi = xa_2 + ya_3 + a_1 \quad (\text{1E})$$

$$\eta = xb_2 + yb_3 + b_1 \quad (\text{2E})$$

Where the unknown coefficients are

$$\{a_1, a_2, a_3, b_1, b_2, b_3\}$$

Substituting equations (1E,2E) and ω into (A) gives

$$b_2 + \frac{y(2x^2 + y^2)(b_3 - a_2)}{x^3} - \frac{y^2(2x^2 + y^2)^2 a_3}{x^6}$$

$$- \left(\frac{4y}{x^2} - \frac{3y(2x^2 + y^2)}{x^4} \right) (xa_2 + ya_3 + a_1) \quad (\text{5E})$$

$$- \left(\frac{2x^2 + y^2}{x^3} + \frac{2y^2}{x^3} \right) (xb_2 + yb_3 + b_1) = 0$$

Putting the above in normal form gives

$$\frac{b_2x^6 + 2x^4y^2a_3 + 3x^4y^2b_2 - 2x^3y^3a_2 + 2x^3y^3b_3 + x^2y^4a_3 + y^6a_3 + 2x^5b_1 - 2x^4ya_1 + 3x^3y^2b_1 - 3x^2y^3a_1}{x^6}$$

$$= 0$$

Setting the numerator to zero gives

$$-b_2x^6 - 2x^4y^2a_3 - 3x^4y^2b_2 + 2x^3y^3a_2 - 2x^3y^3b_3 - x^2y^4a_3 \quad (\text{6E})$$

$$- y^6a_3 - 2x^5b_1 + 2x^4ya_1 - 3x^3y^2b_1 + 3x^2y^3a_1 = 0$$

Looking at the above PDE shows the following are all the terms with $\{x, y\}$ in them.

$$\{x, y\}$$

The following substitution is now made to be able to collect on all terms with $\{x, y\}$ in them

$$\{x = v_1, y = v_2\}$$

The above PDE (6E) now becomes

$$\begin{aligned} 2a_2v_1^3v_2^3 - 2a_3v_1^4v_2^2 - a_3v_1^2v_2^4 - a_3v_2^6 - b_2v_1^6 - 3b_2v_1^4v_2^2 \\ - 2b_3v_1^3v_2^3 + 2a_1v_1^4v_2 + 3a_1v_1^2v_2^3 - 2b_1v_1^5 - 3b_1v_1^3v_2^2 = 0 \end{aligned} \quad (7E)$$

Collecting the above on the terms v_i introduced, and these are

$$\{v_1, v_2\}$$

Equation (7E) now becomes

$$\begin{aligned} -b_2v_1^6 - 2b_1v_1^5 + (-2a_3 - 3b_2)v_1^4v_2^2 + 2a_1v_1^4v_2 \\ + (2a_2 - 2b_3)v_1^3v_2^3 - 3b_1v_1^3v_2^2 - a_3v_1^2v_2^4 + 3a_1v_1^2v_2^3 - a_3v_2^6 = 0 \end{aligned} \quad (8E)$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$\begin{aligned} 2a_1 &= 0 \\ 3a_1 &= 0 \\ -a_3 &= 0 \\ -3b_1 &= 0 \\ -2b_1 &= 0 \\ -b_2 &= 0 \\ 2a_2 - 2b_3 &= 0 \\ -2a_3 - 3b_2 &= 0 \end{aligned}$$

Solving the above equations for the unknowns gives

$$a_1 = 0$$

$$a_2 = b_3$$

$$a_3 = 0$$

$$b_1 = 0$$

$$b_2 = 0$$

$$b_3 = b_3$$

Substituting the above solution in the anstaz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$\xi = x$$

$$\eta = y$$

Shifting is now applied to make $\xi = 0$ in order to simplify the rest of the computation

$$\begin{aligned} \eta &= \eta - \omega(x, y) \xi \\ &= y - \left(\frac{y(2x^2 + y^2)}{x^3} \right) (x) \\ &= \frac{-x^2y - y^3}{x^2} \\ \xi &= 0 \end{aligned}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS \quad (1)$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y} \right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = x$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{\frac{-x^2y - y^3}{x^2}} dy \end{aligned}$$

Which results in

$$S = -\ln(y) + \frac{\ln(x^2 + y^2)}{2}$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = \frac{y(2x^2 + y^2)}{x^3}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 1 \\ R_y &= 0 \\ S_x &= \frac{x}{x^2 + y^2} \\ S_y &= -\frac{x^2}{y(x^2 + y^2)} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = -\frac{1}{x} \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = -\frac{1}{R}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S .

Since the ode has the form $\frac{d}{dR}S(R) = f(R)$, then we only need to integrate $f(R)$.

$$\begin{aligned} \int dS &= \int -\frac{1}{R} dR \\ S(R) &= -\ln(R) + c_2 \end{aligned}$$

To complete the solution, we just need to transform the above back to x, y coordinates. This results in

$$-\ln(y) + \frac{\ln(x^2 + y^2)}{2} = -\ln(x) + c_2$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = \frac{y(2x^2 + y^2)}{x^3}$	$R = x$ $S = -\ln(y) + \frac{\ln(x^2 + y^2)}{2}$	$\frac{dS}{dR} = -\frac{1}{R}$

Solving for y gives

$$y = \frac{x^2}{\sqrt{e^{2c_2} - x^2}}$$

$$y = -\frac{x^2}{\sqrt{e^{2c_2} - x^2}}$$

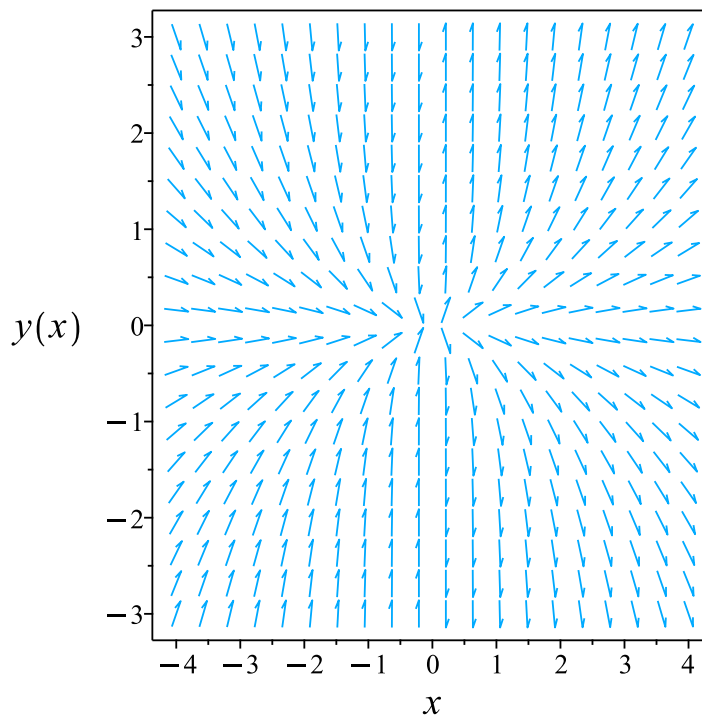


Figure 2.29: Slope field plot
 $2x^2y + y^3 - x^3y' = 0$

Summary of solutions found

$$y = \frac{x^2}{\sqrt{e^{2c_2} - x^2}}$$

$$y = -\frac{x^2}{\sqrt{e^{2c_2} - x^2}}$$

Solved as first order ode of type dAlembert

Time used: 758.165 (sec)

Let $p = y'$ the ode becomes

$$-x^3p + 2x^2y + y^3 = 0$$

Solving for y from the above results in

$$y = \left(\frac{(108p + 12\sqrt{81p^2 + 96})^{1/3}}{6} - \frac{4}{(108p + 12\sqrt{81p^2 + 96})^{1/3}} \right) x \quad (1)$$

$$y = \left(-\frac{(108p + 12\sqrt{81p^2 + 96})^{1/3}}{12} + \frac{2}{(108p + 12\sqrt{81p^2 + 96})^{1/3}} + \frac{i\sqrt{3} \left(\frac{(108p + 12\sqrt{81p^2 + 96})^{1/3}}{6} + \frac{4}{(108p + 12\sqrt{81p^2 + 96})^{1/3}} \right)}{2} \right) x \quad (2)$$

$$y = \left(-\frac{(108p + 12\sqrt{81p^2 + 96})^{1/3}}{12} + \frac{2}{(108p + 12\sqrt{81p^2 + 96})^{1/3}} - \frac{i\sqrt{3} \left(\frac{(108p + 12\sqrt{81p^2 + 96})^{1/3}}{6} + \frac{4}{(108p + 12\sqrt{81p^2 + 96})^{1/3}} \right)}{2} \right) x \quad (3)$$

This has the form

$$y = xf(p) + g(p) \quad (*)$$

Where f, g are functions of $p = y'(x)$. Each of the above ode's is dAlembert ode which is now solved.

Solving ode 1A

Taking derivative of (*) w.r.t. x gives

$$\begin{aligned} p &= f + (xf' + g') \frac{dp}{dx} \\ p - f &= (xf' + g') \frac{dp}{dx} \end{aligned} \quad (2)$$

Comparing the form $y = xf + g$ to (1A) shows that

$$f = \frac{(108p + 12\sqrt{81p^2 + 96})^{2/3} - 24}{6(108p + 12\sqrt{81p^2 + 96})^{1/3}}$$

$$g = 0$$

Hence (2) becomes

$$p - \frac{(108p + 12\sqrt{81p^2 + 96})^{2/3} - 24}{6(108p + 12\sqrt{81p^2 + 96})^{1/3}} = \left(\frac{6x}{(108p + 12\sqrt{81p^2 + 96})^{2/3}} + \frac{54xp}{(108p + 12\sqrt{81p^2 + 96})^{2/3} \sqrt{81p^2 + 96}} \right) \quad (2A)$$

The singular solution is found by setting $\frac{dp}{dx} = 0$ in the above which gives

$$p - \frac{(108p + 12\sqrt{81p^2 + 96})^{2/3} - 24}{6(108p + 12\sqrt{81p^2 + 96})^{1/3}} = 0$$

Solving the above for p results in

$$p_1 = 0$$

Substituting these in (1A) and keeping singular solution that verifies the ode gives

$$y = 0$$

The general solution is found when $\frac{dp}{dx} \neq 0$. From eq. (2A). This results in

$$p'(x) = \frac{p(x) - \frac{(108p(x) + 12\sqrt{81p(x)^2 + 96})^{2/3} - 24}{6(108p(x) + 12\sqrt{81p(x)^2 + 96})^{1/3}}}{\frac{6x}{(108p(x) + 12\sqrt{81p(x)^2 + 96})^{2/3}} + \frac{54xp(x)}{(108p(x) + 12\sqrt{81p(x)^2 + 96})^{2/3} \sqrt{81p(x)^2 + 96}} + \frac{144x}{(108p(x) + 12\sqrt{81p(x)^2 + 96})^{4/3}} + \frac{108p(x)}{(108p(x) + 12\sqrt{81p(x)^2 + 96})^{2/3}}} \quad (3)$$

This ODE is now solved for $p(x)$. No inversion is needed. The ode $p'(x) = -\frac{\sqrt{81p(x)^2 + 96} \left((108p(x) + 12\sqrt{81p(x)^2 + 96})^{2/3} - 6p(x) \left((108p(x) + 12\sqrt{81p(x)^2 + 96})^{2/3} + 24 \right) \right)}{3 \left((108p(x) + 12\sqrt{81p(x)^2 + 96})^{2/3} + 24 \right) x}$

is separable as it can be written as

$$p'(x) = -\frac{\sqrt{81p(x)^2 + 96} \left((108p(x) + 12\sqrt{81p(x)^2 + 96})^{2/3} - 6p(x) \left((108p(x) + 12\sqrt{81p(x)^2 + 96})^{2/3} + 24 \right) \right)}{3 \left((108p(x) + 12\sqrt{81p(x)^2 + 96})^{2/3} + 24 \right) x}$$

$$= f(x)g(p)$$

Where

$$f(x) = -\frac{1}{3x}$$

$$g(p) = \frac{\sqrt{81p^2 + 96} \left((108p + 12\sqrt{81p^2 + 96})^{2/3} - 6p(108p + 12\sqrt{81p^2 + 96})^{1/3} - 24 \right)}{(108p + 12\sqrt{81p^2 + 96})^{2/3} + 24}$$

Integrating gives

$$\int \frac{1}{g(p)} \dots$$

$$\int \frac{(108p + 12\sqrt{81p^2 + 96})^{2/3} + 24}{\sqrt{81p^2 + 96} \left((108p + 12\sqrt{81p^2 + 96})^{2/3} - 6p(108p + 12\sqrt{81p^2 + 96})^{1/3} - 24 \right)} \dots$$

$$\int^{p(x)} \frac{(108\tau + 12\sqrt{81\tau^2 + 96})^{2/3} + 24}{\sqrt{81\tau^2 + 96} \left((108\tau + 12\sqrt{81\tau^2 + 96})^{2/3} - 6\tau(108\tau + 12\sqrt{81\tau^2 + 96})^{1/3} - 24 \right)} d\tau = \ln \left(\frac{1}{x^{1/3}} \right) + \dots$$

We now need to find the singular solutions, these are found by finding for what values

$g(p)$ is zero, since we had to divide by this above. Solving $g(p) = 0$ or $\frac{\sqrt{81p^2 + 96} \left((108p + 12\sqrt{81p^2 + 96})^{2/3} - 6p(108p + 12\sqrt{81p^2 + 96})^{1/3} - 24 \right)}{(108p + 12\sqrt{81p^2 + 96})^{2/3} + 24} = 0$ for $p(x)$ gives

$$p(x) = 0$$

$$p(x) = -\frac{4i\sqrt{6}}{9}$$

$$p(x) = \frac{4i\sqrt{6}}{9}$$

Now we go over each such singular solution and check if it verifies the ode itself and any initial conditions given. If it does not then the singular solution will not be used.

Therefore the solutions found are

$$\int^{p(x)} \frac{(108\tau + 12\sqrt{81\tau^2 + 96})^{2/3} + 24}{\sqrt{81\tau^2 + 96} \left((108\tau + 12\sqrt{81\tau^2 + 96})^{2/3} - 6\tau(108\tau + 12\sqrt{81\tau^2 + 96})^{1/3} - 24 \right)} d\tau = \ln \left(\frac{1}{x^{1/3}} \right) + \dots$$

$$p(x) = \dots$$

$$p(x) = -\frac{4i\sqrt{6}}{9}$$

$$p(x) = \frac{4i\sqrt{6}}{9}$$

Substituting the above solution for p in (2A) gives

$$y = \frac{x \left(\left(108 \operatorname{RootOf} \left(-3 \left(f^{-Z} - \frac{(108\tau + 12\sqrt{81\tau^2 + 96})^{2/3} + 24}{\sqrt{81\tau^2 + 96} \left(-(108\tau + 12\sqrt{81\tau^2 + 96})^{2/3} + 6\tau(108\tau + 12\sqrt{81\tau^2 + 96})^{1/3} + 24 \right)} \right) d\tau \right) - \ln(x) \right)}{6 \left(108 \operatorname{RootOf} \left(-3 \left(f^{-Z} - \frac{(108\tau + 12\sqrt{81\tau^2 + 96})^{2/3} + 24}{\sqrt{81\tau^2 + 96} \left(-(108\tau + 12\sqrt{81\tau^2 + 96})^{2/3} + 6\tau(108\tau + 12\sqrt{81\tau^2 + 96})^{1/3} + 24 \right)} \right) d\tau \right) - \ln(x) \right)}$$

Solving ode 2A

Taking derivative of (*) w.r.t. x gives

$$\begin{aligned} p &= f + (xf' + g') \frac{dp}{dx} \\ p - f &= (xf' + g') \frac{dp}{dx} \end{aligned} \quad (2)$$

Comparing the form $y = xf + g$ to (1A) shows that

$$\begin{aligned} f &= \frac{(i\sqrt{3} - 1) (108p + 12\sqrt{81p^2 + 96})^{2/3} + 24i\sqrt{3} + 24}{12 (108p + 12\sqrt{81p^2 + 96})^{1/3}} \\ g &= 0 \end{aligned}$$

Hence (2) becomes

$$p - \frac{(i\sqrt{3} - 1) (108p + 12\sqrt{81p^2 + 96})^{2/3} + 24i\sqrt{3} + 24}{12 (108p + 12\sqrt{81p^2 + 96})^{1/3}} = \left(\frac{3ix\sqrt{3}}{(108p + 12\sqrt{81p^2 + 96})^{2/3}} + \frac{1}{(108p + 12\sqrt{81p^2 + 96})^{1/3}} \right) p \quad (2A)$$

The singular solution is found by setting $\frac{dp}{dx} = 0$ in the above which gives

$$p - \frac{(i\sqrt{3} - 1) (108p + 12\sqrt{81p^2 + 96})^{2/3} + 24i\sqrt{3} + 24}{12 (108p + 12\sqrt{81p^2 + 96})^{1/3}} = 0$$

Solving the above for p results in

$$p_1 = \frac{i(3i\sqrt{3} - 3 + 3\sqrt{30 + 30i\sqrt{3}}) \sqrt{3} + 21i\sqrt{3} + 27 - 3\sqrt{30 + 30i\sqrt{3}}}{12\sqrt{3i\sqrt{3} - 3 + 3\sqrt{30 + 30i\sqrt{3}}}}$$

Substituting these in (1A) and keeping singular solution that verifies the ode gives

$$y = ix$$

The general solution is found when $\frac{dp}{dx} \neq 0$. From eq. (2A). This results in

$$p'(x) = \frac{\frac{3ix\sqrt{3}}{(108p(x)+12\sqrt{81p(x)^2+96})^{2/3}} + \frac{27ix\sqrt{3}p(x)}{(108p(x)+12\sqrt{81p(x)^2+96})^{2/3} \sqrt{81p(x)^2+96}} - \frac{3x}{(108p(x)+12\sqrt{81p(x)^2+96})^{2/3}} - \frac{3x}{(108p(x)+12\sqrt{81p(x)^2+96})^{2/3}}}{(3)}$$

This ODE is now solved for $p(x)$. No inversion is needed. The ode $p'(x) = -\frac{\sqrt{81p(x)^2+96} \left(\sqrt{3} (108p(x)+12\sqrt{81p(x)^2+96})^{2/3} + 24\sqrt{3} + 12ip(x) (108p(x)+12\sqrt{81p(x)^2+96})^{1/3} + i(108p(x)+12\sqrt{81p(x)^2+96})^{1/3} \right)}{3 \left(\sqrt{3} (108p(x)+12\sqrt{81p(x)^2+96})^{2/3} - 24\sqrt{3} + i(108p(x)+12\sqrt{81p(x)^2+96})^{1/3} \right)}$

is separable as it can be written as

$$p'(x) = -\frac{\sqrt{81p(x)^2+96} \left(\sqrt{3} (108p(x)+12\sqrt{81p(x)^2+96})^{2/3} + 24\sqrt{3} + 12ip(x) (108p(x)+12\sqrt{81p(x)^2+96})^{1/3} + i(108p(x)+12\sqrt{81p(x)^2+96})^{1/3} \right)}{3 \left(\sqrt{3} (108p(x)+12\sqrt{81p(x)^2+96})^{2/3} - 24\sqrt{3} + i(108p(x)+12\sqrt{81p(x)^2+96})^{1/3} \right)}$$

$$= f(x)g(p)$$

Where

$$f(x) = -\frac{1}{3x}$$

$$g(p) = \frac{\sqrt{81p^2+96} \left(\sqrt{3} (108p+12\sqrt{81p^2+96})^{2/3} + 24\sqrt{3} + 12ip(108p+12\sqrt{81p^2+96})^{1/3} + i(108p+12\sqrt{81p^2+96})^{1/3} \right)}{\sqrt{3} (108p+12\sqrt{81p^2+96})^{2/3} - 24\sqrt{3} + i(108p+12\sqrt{81p^2+96})^{1/3} + 24\sqrt{3} + i(108p+12\sqrt{81p^2+96})^{1/3}}$$

Integrating gives

$$\int \frac{\sqrt{3} (108p + 12\sqrt{81p^2 + 96})^{2/3} - 24\sqrt{3} + i(108p + 12\sqrt{81p^2 + 96})^{2/3}}{\sqrt{81p^2 + 96} \left(\sqrt{3} (108p + 12\sqrt{81p^2 + 96})^{2/3} + 24\sqrt{3} + 12ip (108p + 12\sqrt{81p^2 + 96})^{2/3} + 24\sqrt{3} + 12ip (108p + 12\sqrt{81p^2 + 96})^{2/3} \right)} dx$$

$$\int^{p(x)} \frac{\sqrt{3} (108\tau + 12\sqrt{81\tau^2 + 96})^{2/3} - 24\sqrt{3} + i(108\tau + 12\sqrt{81\tau^2 + 96})^{2/3}}{\sqrt{81\tau^2 + 96} \left(\sqrt{3} (108\tau + 12\sqrt{81\tau^2 + 96})^{2/3} + 24\sqrt{3} + 12i\tau (108\tau + 12\sqrt{81\tau^2 + 96})^{1/3} + i (108\tau + 12\sqrt{81\tau^2 + 96})^{2/3} \right)} d\tau$$

We now need to find the singular solutions, these are found by finding for what values

$g(p)$ is zero, since we had to divide by this above. Solving $g(p) = 0$ or $\frac{\sqrt{81p^2+96} \left(\sqrt{3} (108p+12\sqrt{81p^2+96})^{2/3} + 24\sqrt{3} + 12ip (108p+12\sqrt{81p^2+96})^{2/3} + 24\sqrt{3} + 12ip (108p+12\sqrt{81p^2+96})^{2/3} \right)}{\sqrt{3} (108p+12\sqrt{81p^2+96})^{2/3}}$

0 for $p(x)$ gives

$$p(x) = -\frac{4i\sqrt{6}}{9}$$

$$p(x) = \frac{4i\sqrt{6}}{9}$$

$$p(x) = \frac{i(3i\sqrt{3} - 3 + 3\sqrt{30 + 30i\sqrt{3}}) \sqrt{3} + 21i\sqrt{3} + 27 - 3\sqrt{30 + 30i\sqrt{3}}}{12\sqrt{3i\sqrt{3} - 3 + 3\sqrt{30 + 30i\sqrt{3}}}}$$

Now we go over each such singular solution and check if it verifies the ode itself and any initial conditions given. If it does not then the singular solution will not be used.

Therefore the solutions found are

$$\int^{p(x)} \frac{\sqrt{3} (108\tau + 12\sqrt{81\tau^2 + 96})^{2/3} - 24\sqrt{3} + i(108\tau + 12\sqrt{81\tau^2 + 96})^{2/3}}{\sqrt{81\tau^2 + 96} \left(\sqrt{3} (108\tau + 12\sqrt{81\tau^2 + 96})^{2/3} + 24\sqrt{3} + 12i\tau (108\tau + 12\sqrt{81\tau^2 + 96})^{1/3} + i (108\tau + 12\sqrt{81\tau^2 + 96})^{2/3} \right)} d\tau$$

$$p(x) = \frac{i(3i\sqrt{3} - 3 + 3\sqrt{30 + 30i\sqrt{3}})}{12\sqrt{3i\sqrt{3} - 3 + 3\sqrt{30 + 30i\sqrt{3}}}}$$

Substituting the above solution for p in (2A) gives

$$y = \frac{x \left((i\sqrt{3} - 1) \left(108 \operatorname{RootOf} \left(-3 \left(\int^{-Z} \frac{\sqrt{3} (108\tau + 12\sqrt{81\tau^2 + 96})^{2/3} - 24\sqrt{3} + i(108\tau + 12\sqrt{81\tau^2 + 96})}{\sqrt{81\tau^2 + 96} (\sqrt{3} (108\tau + 12\sqrt{81\tau^2 + 96})^{2/3} + 24\sqrt{3} + 12i\tau (108\tau + 12\sqrt{81\tau^2 + 96})^{1/3} + 24)} \right) \right) \right)}{12 \left(108 \operatorname{RootOf} \left(-3 \left(\int^{-Z} \frac{\sqrt{3} (108\tau + 12\sqrt{81\tau^2 + 96})^{2/3} - 24\sqrt{3} + i(108\tau + 12\sqrt{81\tau^2 + 96})}{\sqrt{81\tau^2 + 96} (\sqrt{3} (108\tau + 12\sqrt{81\tau^2 + 96})^{2/3} + 24\sqrt{3} + 12i\tau (108\tau + 12\sqrt{81\tau^2 + 96})^{1/3} + 24)} \right) \right) \right)}$$

Solving ode 3A

Taking derivative of (*) w.r.t. x gives

$$\begin{aligned} p &= f + (xf' + g') \frac{dp}{dx} \\ p - f &= (xf' + g') \frac{dp}{dx} \end{aligned} \quad (2)$$

Comparing the form $y = xf + g$ to (1A) shows that

$$\begin{aligned} f &= -\frac{(1 + i\sqrt{3}) (108p + 12\sqrt{81p^2 + 96})^{2/3} + 24i\sqrt{3} - 24}{12 (108p + 12\sqrt{81p^2 + 96})^{1/3}} \\ g &= 0 \end{aligned}$$

Hence (2) becomes

$$p + \frac{(1 + i\sqrt{3}) (108p + 12\sqrt{81p^2 + 96})^{2/3} + 24i\sqrt{3} - 24}{12 (108p + 12\sqrt{81p^2 + 96})^{1/3}} = \left(-\frac{3x}{(108p + 12\sqrt{81p^2 + 96})^{2/3}} - \frac{1}{(108p + 12\sqrt{81p^2 + 96})^{1/3}} \right) \quad (2A)$$

The singular solution is found by setting $\frac{dp}{dx} = 0$ in the above which gives

$$p + \frac{(1 + i\sqrt{3})(108p + 12\sqrt{81p^2 + 96})^{2/3} + 24i\sqrt{3} - 24}{12(108p + 12\sqrt{81p^2 + 96})^{1/3}} = 0$$

Solving the above for p results in

$$p_1 = -\frac{i(-3i\sqrt{3} - 3 + 3\sqrt{30 - 30i\sqrt{3}})\sqrt{3} + 21i\sqrt{3} - 27 + 3\sqrt{30 - 30i\sqrt{3}}}{12\sqrt{-3i\sqrt{3} - 3 + 3\sqrt{30 - 30i\sqrt{3}}}}$$

Substituting these in (1A) and keeping singular solution that verifies the ode gives

$$y = -ix$$

The general solution is found when $\frac{dp}{dx} \neq 0$. From eq. (2A). This results in

$$p'(x) = \frac{3x}{(108p(x) + 12\sqrt{81p(x)^2 + 96})^{2/3}} - \frac{27xp(x)}{(108p(x) + 12\sqrt{81p(x)^2 + 96})^{2/3}\sqrt{81p(x)^2 + 96}} - \frac{3ix\sqrt{3}}{(108p(x) + 12\sqrt{81p(x)^2 + 96})^{2/3}} - \frac{10}{(108p(x) + 12\sqrt{81p(x)^2 + 96})^{1/3}} \quad (3)$$

This ODE is now solved for $p(x)$. No inversion is needed. The ode $p'(x) = -\frac{(\sqrt{3}(108p(x) + 12\sqrt{81p(x)^2 + 96})^{2/3} + 24\sqrt{3} - 12ip(x))}{3x(\sqrt{3}(108p(x) + 12\sqrt{81p(x)^2 + 96})^{2/3} - 24\sqrt{3} - i(108p(x) + 12\sqrt{81p(x)^2 + 96})^{1/3})}$

is separable as it can be written as

$$p'(x) = -\frac{(\sqrt{3}(108p(x) + 12\sqrt{81p(x)^2 + 96})^{2/3} + 24\sqrt{3} - 12ip(x))(108p(x) + 12\sqrt{81p(x)^2 + 96})^{1/3}}{3x(\sqrt{3}(108p(x) + 12\sqrt{81p(x)^2 + 96})^{2/3} - 24\sqrt{3} - i(108p(x) + 12\sqrt{81p(x)^2 + 96})^{1/3})}$$

$$= f(x)g(p)$$

Where

$$f(x) = -\frac{1}{3x}$$

$$g(p) = \frac{(\sqrt{3}(108p + 12\sqrt{81p^2 + 96})^{2/3} + 24\sqrt{3} - 12ip)(108p + 12\sqrt{81p^2 + 96})^{1/3} - i(108p + 12\sqrt{81p^2 + 96})^{1/3}}{\sqrt{3}(108p + 12\sqrt{81p^2 + 96})^{2/3} - 24\sqrt{3} - i(108p + 12\sqrt{81p^2 + 96})^{1/3} - 2}$$

Integrating gives

$$\int \frac{\sqrt{3} (108p + 12\sqrt{81p^2 + 96})^{2/3} - 24\sqrt{3} - i(108p + 12\sqrt{81p^2 + 96})^{1/3} - i(108p + 12\sqrt{81p^2 + 96})^{2/3}}{(\sqrt{3} (108p + 12\sqrt{81p^2 + 96})^{2/3} + 24\sqrt{3} - 12ip (108p + 12\sqrt{81p^2 + 96})^{1/3} - i(108p + 12\sqrt{81p^2 + 96})^{2/3} - 24\sqrt{3})^{2/3}} dp$$

$$\int^{p(x)} \frac{\sqrt{3} (108\tau + 12\sqrt{81\tau^2 + 96})^{2/3} - 24\sqrt{3} - i(108\tau + 12\sqrt{81\tau^2 + 96})^{1/3} - i(108\tau + 12\sqrt{81\tau^2 + 96})^{2/3}}{(\sqrt{3} (108\tau + 12\sqrt{81\tau^2 + 96})^{2/3} + 24\sqrt{3} - 12i\tau (108\tau + 12\sqrt{81\tau^2 + 96})^{1/3} - i(108\tau + 12\sqrt{81\tau^2 + 96})^{2/3} - 24\sqrt{3})^{2/3}} d\tau$$

We now need to find the singular solutions, these are found by finding for what values

$g(p)$ is zero, since we had to divide by this above. Solving $g(p) = 0$ or $\frac{(\sqrt{3} (108p + 12\sqrt{81p^2 + 96})^{2/3} + 24\sqrt{3} - 12ip (108p + 12\sqrt{81p^2 + 96})^{1/3} - i(108p + 12\sqrt{81p^2 + 96})^{2/3} - 24\sqrt{3})^{2/3}}{\sqrt{3} (108p + 12\sqrt{81p^2 + 96})^{2/3}}$

0 for $p(x)$ gives

$$p(x) = -\frac{4i\sqrt{6}}{9}$$

$$p(x) = \frac{4i\sqrt{6}}{9}$$

$$p(x) = -\frac{i(-3i\sqrt{3} - 3 + 3\sqrt{30 - 30i\sqrt{3}})\sqrt{3} + 21i\sqrt{3} - 27 + 3\sqrt{30 - 30i\sqrt{3}}}{12\sqrt{-3i\sqrt{3} - 3 + 3\sqrt{30 - 30i\sqrt{3}}}}$$

Now we go over each such singular solution and check if it verifies the ode itself and any initial conditions given. If it does not then the singular solution will not be used.

Therefore the solutions found are

$$\int^{p(x)} \frac{\sqrt{3} (108\tau + 12\sqrt{81\tau^2 + 96})^{2/3} - 24\sqrt{3} - i(108\tau + 12\sqrt{81\tau^2 + 96})^{1/3} - i(108\tau + 12\sqrt{81\tau^2 + 96})^{2/3}}{(\sqrt{3} (108\tau + 12\sqrt{81\tau^2 + 96})^{2/3} + 24\sqrt{3} - 12i\tau (108\tau + 12\sqrt{81\tau^2 + 96})^{1/3} - i(108\tau + 12\sqrt{81\tau^2 + 96})^{2/3} - 24\sqrt{3})^{2/3}} d\tau$$

$$p(x) = -\frac{i(-3i\sqrt{3} - 3 + 3\sqrt{30 - 30i\sqrt{3}})\sqrt{3} + 21i\sqrt{3} - 27 + 3\sqrt{30 - 30i\sqrt{3}}}{12\sqrt{-3i\sqrt{3} - 3 + 3\sqrt{30 - 30i\sqrt{3}}}}$$

Substituting the above solution for p in (2A) gives

$$y = -\frac{x \left((1 + i\sqrt{3}) \left(108 \operatorname{RootOf} \left(-3 \left(f^{-Z} \frac{\sqrt{3} (108\tau + 12\sqrt{81\tau^2 + 96})^{2/3} - 24\sqrt{3} - i(108\tau + 12\sqrt{81\tau^2 + 96})}{(\sqrt{3} (108\tau + 12\sqrt{81\tau^2 + 96})^{2/3} + 24\sqrt{3} - 12i\tau (108\tau + 12\sqrt{81\tau^2 + 96})^{1/3} - i(108\tau + 12\sqrt{81\tau^2 + 96})) \right)} \right) \right)}{12 \left(108 \operatorname{RootOf} \left(-3 \left(f^{-Z} \frac{\sqrt{3} (108\tau + 12\sqrt{81\tau^2 + 96})^{2/3} - 24\sqrt{3} - i(108\tau + 12\sqrt{81\tau^2 + 96})}{(\sqrt{3} (108\tau + 12\sqrt{81\tau^2 + 96})^{2/3} + 24\sqrt{3} - 12i\tau (108\tau + 12\sqrt{81\tau^2 + 96})^{1/3} - i(108\tau + 12\sqrt{81\tau^2 + 96})) \right)} \right) \right)}$$

The solution

$$y = \frac{x \left((-48i\sqrt{6})^{2/3} - 24 \right)}{6 (-48i\sqrt{6})^{1/3}}$$

was found not to satisfy the ode or the IC. Hence it is removed. The solution

$$y = \frac{x \left(48^{2/3} (i\sqrt{6})^{2/3} - 24 \right) 48^{2/3}}{288 (i\sqrt{6})^{1/3}}$$

was found not to satisfy the ode or the IC. Hence it is removed. The solution

$$y = -\frac{x \left((1 + i\sqrt{3}) (-48i\sqrt{6})^{2/3} + 24i\sqrt{3} - 24 \right)}{12 (-48i\sqrt{6})^{1/3}}$$

was found not to satisfy the ode or the IC. Hence it is removed. The solution

$$y = -\frac{x\left((1+i\sqrt{3})48^{2/3}(i\sqrt{6})^{2/3}+24i\sqrt{3}-24\right)48^{2/3}}{576(i\sqrt{6})^{1/3}}$$

was found not to satisfy the ode or the IC. Hence it is removed. The solution

$$y = \frac{x\left((1+i\sqrt{3})\left(108\text{RootOf}\left(-3\left(\int^{-Z}\frac{\sqrt{3}(108\tau+12\sqrt{81\tau^2+96})^{2/3}-24\sqrt{3}-i(108\tau+12\sqrt{81\tau^2+96})}{(\sqrt{3}(108\tau+12\sqrt{81\tau^2+96})^{2/3}+24\sqrt{3}-12i\tau(108\tau+12\sqrt{81\tau^2+96})^{1/3}-i(108\tau+12\sqrt{81\tau^2+96}))}\right)\right)\right)}{12\left(108\text{RootOf}\left(-3\left(\int^{-Z}\frac{\sqrt{3}(108\tau+12\sqrt{81\tau^2+96})^{2/3}-24\sqrt{3}-i(108\tau+12\sqrt{81\tau^2+96})}{(\sqrt{3}(108\tau+12\sqrt{81\tau^2+96})^{2/3}+24\sqrt{3}-12i\tau(108\tau+12\sqrt{81\tau^2+96})^{1/3}-i(108\tau+12\sqrt{81\tau^2+96}))}\right)\right)\right)}$$

was found not to satisfy the ode or the IC. Hence it is removed. The solution

$$y = \frac{x\left((i\sqrt{3}-1)(-48i\sqrt{6})^{2/3}+24i\sqrt{3}+24\right)}{12(-48i\sqrt{6})^{1/3}}$$

was found not to satisfy the ode or the IC. Hence it is removed. The solution

$$y = \frac{x\left((i\sqrt{3}-1)48^{2/3}(i\sqrt{6})^{2/3}+24i\sqrt{3}+24\right)48^{2/3}}{576(i\sqrt{6})^{1/3}}$$

was found not to satisfy the ode or the IC. Hence it is removed. The solution

$$y = \frac{x\left((i\sqrt{3}-1)\left(108\text{RootOf}\left(-3\left(\int^{-Z}\frac{\sqrt{3}(108\tau+12\sqrt{81\tau^2+96})^{2/3}-24\sqrt{3}+i(108\tau+12\sqrt{81\tau^2+96})}{\sqrt{81\tau^2+96}\left(\sqrt{3}(108\tau+12\sqrt{81\tau^2+96})^{2/3}+24\sqrt{3}+12i\tau(108\tau+12\sqrt{81\tau^2+96})^{1/3}+i(108\tau+12\sqrt{81\tau^2+96})\right)}\right)\right)\right)}{12\left(108\text{RootOf}\left(-3\left(\int^{-Z}\frac{\sqrt{3}(108\tau+12\sqrt{81\tau^2+96})^{2/3}-24\sqrt{3}+i(108\tau+12\sqrt{81\tau^2+96})}{\sqrt{81\tau^2+96}\left(\sqrt{3}(108\tau+12\sqrt{81\tau^2+96})^{2/3}+24\sqrt{3}+12i\tau(108\tau+12\sqrt{81\tau^2+96})^{1/3}+i(108\tau+12\sqrt{81\tau^2+96})\right)}\right)\right)\right)}$$

was found not to satisfy the ode or the IC. Hence it is removed.

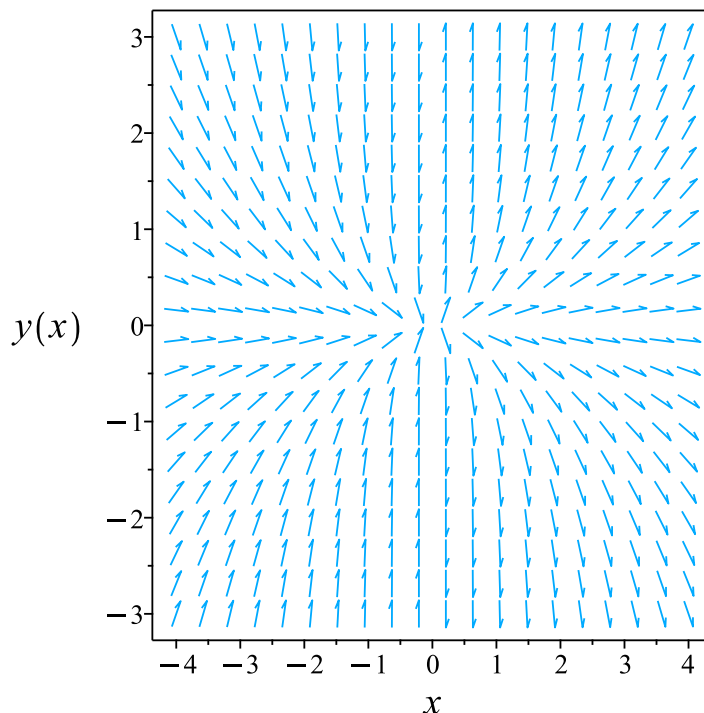


Figure 2.30: Slope field plot
 $2x^2y + y^3 - x^3y' = 0$

Summary of solutions found

$$y = 0$$

$$y = -ix$$

$$y = ix$$

$$y$$

$$x \left(\left(108 \operatorname{RootOf} \left(-3 \left(\int^{-Z} - \frac{(108\tau + 12\sqrt{81\tau^2 + 96})^{2/3} + 24}{\sqrt{81\tau^2 + 96} \left(-(108\tau + 12\sqrt{81\tau^2 + 96})^{2/3} + 6\tau(108\tau + 12\sqrt{81\tau^2 + 96})^{1/3} + 24 \right)} d\tau \right) - \ln(x) + \dots \right) \right)$$

$$6 \left(108 \operatorname{RootOf} \left(-3 \left(\int^{-Z} - \frac{(108\tau + 12\sqrt{81\tau^2 + 96})^{2/3} + 24}{\sqrt{81\tau^2 + 96} \left(-(108\tau + 12\sqrt{81\tau^2 + 96})^{2/3} + 6\tau(108\tau + 12\sqrt{81\tau^2 + 96})^{1/3} + 24 \right)} d\tau \right) - \ln(x) \right)$$

$$y = \frac{x \left((1 + i\sqrt{3}) \left(-\frac{9(i(-3i\sqrt{3}-3+3\sqrt{30-30i\sqrt{3}})\sqrt{3}+21i\sqrt{3}-27+3\sqrt{30-30i\sqrt{3}})}{\sqrt{-3i\sqrt{3}-3+3\sqrt{30-30i\sqrt{3}}}} \right) + 12\sqrt{\frac{9(i(-3i\sqrt{3}-3+3\sqrt{30-30i\sqrt{3}})\sqrt{3}+21i\sqrt{3}-27+3\sqrt{30-30i\sqrt{3}})}{16(-3i\sqrt{3}-3+3\sqrt{30-30i\sqrt{3}})}} \right)}{12 \left(-\frac{9(i(-3i\sqrt{3}-3+3\sqrt{30-30i\sqrt{3}})\sqrt{3}+21i\sqrt{3}-27+3\sqrt{30-30i\sqrt{3}})}{\sqrt{-3i\sqrt{3}-3+3\sqrt{30-30i\sqrt{3}}}} \right) + 12\sqrt{\frac{9(i(-3i\sqrt{3}-3+3\sqrt{30-30i\sqrt{3}})\sqrt{3}+21i\sqrt{3}-27+3\sqrt{30-30i\sqrt{3}})}{16(-3i\sqrt{3}-3+3\sqrt{30-30i\sqrt{3}})}}}$$

$$y = \frac{x \left((i\sqrt{3} - 1) \left(\frac{9i(3i\sqrt{3}-3+3\sqrt{30+30i\sqrt{3}})\sqrt{3}+189i\sqrt{3}+243-27\sqrt{30+30i\sqrt{3}}}{\sqrt{3i\sqrt{3}-3+3\sqrt{30+30i\sqrt{3}}}} \right) + 12\sqrt{\frac{9(i(3i\sqrt{3}-3+3\sqrt{30+30i\sqrt{3}})\sqrt{3}+21i\sqrt{3}-27+3\sqrt{30+30i\sqrt{3}})}{16(3i\sqrt{3}-3+3\sqrt{30+30i\sqrt{3}})}} \right)}{12 \left(\frac{9i(3i\sqrt{3}-3+3\sqrt{30+30i\sqrt{3}})\sqrt{3}+189i\sqrt{3}+243-27\sqrt{30+30i\sqrt{3}}}{\sqrt{3i\sqrt{3}-3+3\sqrt{30+30i\sqrt{3}}}} \right) + 12\sqrt{\frac{9(i(3i\sqrt{3}-3+3\sqrt{30+30i\sqrt{3}})\sqrt{3}+21i\sqrt{3}-27+3\sqrt{30+30i\sqrt{3}})}{16(3i\sqrt{3}-3+3\sqrt{30+30i\sqrt{3}})}}}$$

$$y = \frac{x(12^{2/3}96^{1/3} - 24) 12^{2/3}96^{5/6}}{6912}$$

Maple step by step solution

Let's solve

$$2x^2y + y^3 - x^3y' = 0$$

- Highest derivative means the order of the ODE is 1

y'

- Solve for the highest derivative

$$y' = -\frac{-y^3-2x^2y}{x^3}$$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
<- Bernoulli successful`
```

Maple dsolve solution

Solving time : 0.003 (sec)

Leaf size : 34

```
dsolve(2*x^2*y(x)+y(x)^3-x^3*diff(y(x),x) = 0,
      y(x),singsol=all)
```

$$y = \frac{x^2}{\sqrt{-x^2 + c_1}}$$

$$y = -\frac{x^2}{\sqrt{-x^2 + c_1}}$$

Mathematica DSolve solution

Solving time : 0.193 (sec)

Leaf size : 47

```
DSolve[{2*x^2*y[x]+y[x]^3-x^3*D[y[x],x]==0,{}},
      y[x],x,IncludeSingularSolutions->True]
```

$$y(x) \rightarrow -\frac{x^2}{\sqrt{-x^2 + c_1}}$$

$$y(x) \rightarrow \frac{x^2}{\sqrt{-x^2 + c_1}}$$

$$y(x) \rightarrow 0$$

2.3.3 problem 3

Solved as first order homogeneous class Maple C ode	196
Solved as first order Exact ode	201
Solved using Lie symmetry for first order ode	204
Solved as first order ode of type dAlembert	208
Maple step by step solution	212
Maple trace	213
Maple dsolve solution	213
Mathematica DSolve solution	214

Internal problem ID [18223]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter IV. Methods of solution: First order equations. section 29. Problems at page 81

Problem number : 3

Date solved : Friday, December 20, 2024 at 10:45:23 AM

CAS classification :

[[_homogeneous, 'class C'], _exact, _rational, [_Abel, '2nd type', 'class A']]

Solve

$$2ax + by + (2cy + bx + e)y' = g$$

Summary of solutions found

$$y = \frac{-bx + \sqrt{-4 \left(x - \frac{be+2cg}{4ac-b^2}\right)^2 ac + b^2 \left(x - \frac{be+2cg}{4ac-b^2}\right)^2 + 4 e^{c_1} c - e}}{2c}$$

$$y = -\frac{bx + \sqrt{-4 \left(x - \frac{be+2cg}{4ac-b^2}\right)^2 ac + b^2 \left(x - \frac{be+2cg}{4ac-b^2}\right)^2 + 4 e^{c_1} c + e}}{2c}$$

$$y = \frac{4\sqrt{-4ac + b^2} acx - \sqrt{-4ac + b^2} b^2 x - 4abcx + b^3 x - \sqrt{-4ac + b^2} be - 2\sqrt{-4ac + b^2} cg - 4ace + b^2 e}{2c(4ac - b^2)}$$

$$y = -\frac{4\sqrt{-4ac + b^2} acx - \sqrt{-4ac + b^2} b^2 x + 4abcx - b^3 x - \sqrt{-4ac + b^2} be - 2\sqrt{-4ac + b^2} cg + 4ace - b^2 e}{2c(4ac - b^2)}$$

Solved as first order homogeneous class Maple C ode

Time used: 0.974 (sec)

Let $Y = y - y_0$ and $X = x - x_0$ then the above is transformed to new ode in $Y(X)$

$$\frac{d}{dX}Y(X) = -\frac{b(Y(X) + y_0) + 2a(x_0 + X) - g}{2c(Y(X) + y_0) + b(x_0 + X) + e}$$

Solving for possible values of x_0 and y_0 which makes the above ode a homogeneous ode results in

$$x_0 = \frac{be + 2cg}{4ac - b^2}$$

$$y_0 = \frac{-2ae - bg}{4ac - b^2}$$

Using these values now it is possible to easily solve for $Y(X)$. The above ode now becomes

$$\frac{d}{dX}Y(X) = -\frac{2aX + bY(X) + \frac{2a(be+2cg)}{4ac-b^2} + \frac{b(-2ae-bg)}{4ac-b^2} - g}{bX + 2cY(X) + \frac{b(be+2cg)}{4ac-b^2} + \frac{2c(-2ae-bg)}{4ac-b^2} + e}$$

In canonical form, the ODE is

$$Y' = F(X, Y)$$

$$= -\frac{2aX + bY}{bX + 2cY} \quad (1)$$

An ode of the form $Y' = \frac{M(X, Y)}{N(X, Y)}$ is called homogeneous if the functions $M(X, Y)$ and $N(X, Y)$ are both homogeneous functions and of the same order. Recall that a function $f(X, Y)$ is homogeneous of order n if

$$f(t^n X, t^n Y) = t^n f(X, Y)$$

In this case, it can be seen that both $M = -2aX - bY$ and $N = bX + 2cY$ are both homogeneous and of the same order $n = 1$. Therefore this is a homogeneous ode. Since this ode is homogeneous, it is converted to separable ODE using the substitution $u = \frac{Y}{X}$, or $Y = uX$. Hence

$$\frac{dY}{dX} = \frac{du}{dX}X + u$$

Applying the transformation $Y = uX$ to the above ODE in (1) gives

$$\frac{du}{dX}X + u = \frac{-bu - 2a}{2cu + b}$$

$$\frac{du}{dX} = \frac{\frac{-bu(X)-2a}{2cu(X)+b} - u(X)}{X}$$

Or

$$\frac{d}{dX}u(X) - \frac{\frac{-bu(X)-2a}{2cu(X)+b} - u(X)}{X} = 0$$

Or

$$2\left(\frac{d}{dX}u(X)\right)u(X)Xc + \left(\frac{d}{dX}u(X)\right)Xb + 2u(X)^2c + 2bu(X) + 2a = 0$$

Or

$$X(2cu(X) + b)\left(\frac{d}{dX}u(X)\right) + 2u(X)^2c + 2bu(X) + 2a = 0$$

Which is now solved as separable in $u(X)$.

The ode $\frac{d}{dX}u(X) = -\frac{2(u(X)^2c+bu(X)+a)}{X(2cu(X)+b)}$ is separable as it can be written as

$$\begin{aligned}\frac{d}{dX}u(X) &= -\frac{2(u(X)^2c + bu(X) + a)}{X(2cu(X) + b)} \\ &= f(X)g(u)\end{aligned}$$

Where

$$\begin{aligned}f(X) &= -\frac{2}{X} \\ g(u) &= \frac{u^2c + bu + a}{2cu + b}\end{aligned}$$

Integrating gives

$$\begin{aligned}\int \frac{1}{g(u)} du &= \int f(X) dX \\ \int \frac{2cu + b}{u^2c + bu + a} du &= \int -\frac{2}{X} dX \\ \ln(u(X)^2c + bu(X) + a) &= \ln\left(\frac{1}{X^2}\right) + c_1\end{aligned}$$

We now need to find the singular solutions, these are found by finding for what values $g(u)$ is zero, since we had to divide by this above. Solving $g(u) = 0$ or $\frac{u^2c+bu+a}{2cu+b} = 0$ for $u(X)$ gives

$$\begin{aligned}u(X) &= \frac{-b + \sqrt{-4ac + b^2}}{2c} \\ u(X) &= -\frac{b + \sqrt{-4ac + b^2}}{2c}\end{aligned}$$

Now we go over each such singular solution and check if it verifies the ode itself and any initial conditions given. If it does not then the singular solution will not be used.

Therefore the solutions found are

$$\ln(u(X)^2 c + bu(X) + a) = \ln\left(\frac{1}{X^2}\right) + c_1$$

$$u(X) = \frac{-b + \sqrt{-4ac + b^2}}{2c}$$

$$u(X) = -\frac{b + \sqrt{-4ac + b^2}}{2c}$$

Solving for $u(X)$ gives

$$u(X) = \frac{-b + \sqrt{-4ac + b^2}}{2c}$$

$$u(X) = -\frac{b + \sqrt{-4ac + b^2}}{2c}$$

$$u(X) = \frac{-bX + \sqrt{-4X^2ac + b^2X^2 + 4e^{c_1}c}}{2cX}$$

$$u(X) = -\frac{bX + \sqrt{-4X^2ac + b^2X^2 + 4e^{c_1}c}}{2cX}$$

Converting $u(X) = \frac{-b + \sqrt{-4ac + b^2}}{2c}$ back to $Y(X)$ gives

$$Y(X) = \frac{X(-b + \sqrt{-4ac + b^2})}{2c}$$

Converting $u(X) = -\frac{b + \sqrt{-4ac + b^2}}{2c}$ back to $Y(X)$ gives

$$Y(X) = -\frac{X(b + \sqrt{-4ac + b^2})}{2c}$$

Converting $u(X) = \frac{-bX + \sqrt{-4X^2ac + b^2X^2 + 4e^{c_1}c}}{2cX}$ back to $Y(X)$ gives

$$Y(X) = \frac{-bX + \sqrt{-4X^2ac + b^2X^2 + 4e^{c_1}c}}{2c}$$

Converting $u(X) = -\frac{bX + \sqrt{-4X^2ac + b^2X^2 + 4e^{c_1}c}}{2cX}$ back to $Y(X)$ gives

$$Y(X) = -\frac{bX + \sqrt{-4X^2ac + b^2X^2 + 4e^{c_1}c}}{2c}$$

Using the solution for $Y(X)$

$$Y(X) = \frac{-bX + \sqrt{-4X^2ac + b^2X^2 + 4e^{c_1}c}}{2c} \quad (\text{A})$$

And replacing back terms in the above solution using

$$\begin{aligned} Y &= y + y_0 \\ X &= x + x_0 \end{aligned}$$

Or

$$\begin{aligned} Y &= y + \frac{-2ae - bg}{4ac - b^2} \\ X &= x + \frac{be + 2cg}{4ac - b^2} \end{aligned}$$

Then the solution in y becomes using EQ (A)

$$y - \frac{-2ae - bg}{4ac - b^2} = \frac{-b\left(x - \frac{be+2cg}{4ac-b^2}\right) + \sqrt{-4\left(x - \frac{be+2cg}{4ac-b^2}\right)^2 ac + b^2\left(x - \frac{be+2cg}{4ac-b^2}\right)^2 + 4e^{c_1}c}}{2c}$$

Using the solution for $Y(X)$

$$Y(X) = -\frac{bX + \sqrt{-4X^2ac + b^2X^2 + 4e^{c_1}c}}{2c} \quad (\text{A})$$

And replacing back terms in the above solution using

$$\begin{aligned} Y &= y + y_0 \\ X &= x + x_0 \end{aligned}$$

Or

$$\begin{aligned} Y &= y + \frac{-2ae - bg}{4ac - b^2} \\ X &= x + \frac{be + 2cg}{4ac - b^2} \end{aligned}$$

Then the solution in y becomes using EQ (A)

$$y - \frac{-2ae - bg}{4ac - b^2} = -\frac{b\left(x - \frac{be+2cg}{4ac-b^2}\right) + \sqrt{-4\left(x - \frac{be+2cg}{4ac-b^2}\right)^2 ac + b^2\left(x - \frac{be+2cg}{4ac-b^2}\right)^2 + 4e^{c_1}c}}{2c}$$

Using the solution for $Y(X)$

$$Y(X) = \frac{X(-b + \sqrt{-4ac + b^2})}{2c} \quad (\text{A})$$

And replacing back terms in the above solution using

$$Y = y + y_0$$

$$X = x + x_0$$

Or

$$Y = y + \frac{-2ae - bg}{4ac - b^2}$$

$$X = x + \frac{be + 2cg}{4ac - b^2}$$

Then the solution in y becomes using EQ (A)

$$y - \frac{-2ae - bg}{4ac - b^2} = \frac{\left(x - \frac{be+2cg}{4ac-b^2}\right) (-b + \sqrt{-4ac + b^2})}{2c}$$

Using the solution for $Y(X)$

$$Y(X) = -\frac{X(b + \sqrt{-4ac + b^2})}{2c} \quad (\text{A})$$

And replacing back terms in the above solution using

$$Y = y + y_0$$

$$X = x + x_0$$

Or

$$Y = y + \frac{-2ae - bg}{4ac - b^2}$$

$$X = x + \frac{be + 2cg}{4ac - b^2}$$

Then the solution in y becomes using EQ (A)

$$y - \frac{-2ae - bg}{4ac - b^2} = -\frac{\left(x - \frac{be+2cg}{4ac-b^2}\right) (b + \sqrt{-4ac + b^2})}{2c}$$

Solving for y gives

$$y = \frac{-bx + \sqrt{-4\left(x - \frac{be+2cg}{4ac-b^2}\right)^2 ac + b^2\left(x - \frac{be+2cg}{4ac-b^2}\right)^2 + 4e^{c_1}c} - e}{2c}$$

$$y = -\frac{bx + \sqrt{-4\left(x - \frac{be+2cg}{4ac-b^2}\right)^2 ac + b^2\left(x - \frac{be+2cg}{4ac-b^2}\right)^2 + 4e^{c_1}c} + e}{2c}$$

$$y = \frac{4\sqrt{-4ac + b^2} acx - \sqrt{-4ac + b^2} b^2x - 4abcx + b^3x - \sqrt{-4ac + b^2} be - 2\sqrt{-4ac + b^2} cg - 4ace + b^2e}{2c(4ac - b^2)}$$

$$y = \frac{4\sqrt{-4ac + b^2} acx - \sqrt{-4ac + b^2} b^2x + 4abcx - b^3x - \sqrt{-4ac + b^2} be - 2\sqrt{-4ac + b^2} cg + 4ace - b^2e}{2c(4ac - b^2)}$$

Solved as first order Exact ode

Time used: 0.306 (sec)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\frac{\partial \phi}{\partial x} = M$$

$$\frac{\partial \phi}{\partial y} = N$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (\text{1A})$$

Therefore

$$\begin{aligned}(bx + 2cy + e) dy &= (-2ax - by + g) dx \\ (2ax + by - g) dx + (bx + 2cy + e) dy &= 0\end{aligned}\tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(x, y) &= 2ax + by - g \\ N(x, y) &= bx + 2cy + e\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(2ax + by - g) \\ &= b\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(bx + 2cy + e) \\ &= b\end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M\tag{1}$$

$$\frac{\partial \phi}{\partial y} = N\tag{2}$$

Integrating (1) w.r.t. x gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int M dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int 2ax + by - g dx \\ \phi &= x(ax + by - g) + f(y)\end{aligned}\tag{3}$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = bx + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = bx + 2cy + e$. Therefore equation (4) becomes

$$bx + 2cy + e = bx + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 2cy + e$$

Integrating the above w.r.t y gives

$$\int f'(y) dy = \int (2cy + e) dy$$

$$f(y) = y^2c + ey + c_1$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = x(ax + by - g) + y^2c + ey + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into the constant c_1 gives the solution as

$$c_1 = x(ax + by - g) + y^2c + ey$$

Solving for y gives

$$y = \frac{-bx - e + \sqrt{-4acx^2 + b^2x^2 + 2bex + 4cgx + 4c_1c + e^2}}{2c}$$

$$y = -\frac{bx + \sqrt{-4acx^2 + b^2x^2 + 2bex + 4cgx + 4c_1c + e^2} + e}{2c}$$

Summary of solutions found

$$y = \frac{-bx - e + \sqrt{-4acx^2 + b^2x^2 + 2bex + 4cgx + 4c_1c + e^2}}{2c}$$

$$y = -\frac{bx + \sqrt{-4acx^2 + b^2x^2 + 2bex + 4cgx + 4c_1c + e^2} + e}{2c}$$

Solved using Lie symmetry for first order ode

Time used: 0.747 (sec)

Writing the ode as

$$y' = -\frac{2ax + by - g}{bx + 2cy + e}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

To determine ξ, η then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$\xi = xa_2 + ya_3 + a_1 \quad (\text{1E})$$

$$\eta = xb_2 + yb_3 + b_1 \quad (\text{2E})$$

Where the unknown coefficients are

$$\{a_1, a_2, a_3, b_1, b_2, b_3\}$$

Substituting equations (1E,2E) and ω into (A) gives

$$b_2 - \frac{(2ax + by - g)(b_3 - a_2)}{bx + 2cy + e} - \frac{(2ax + by - g)^2 a_3}{(bx + 2cy + e)^2}$$

$$- \left(-\frac{2a}{bx + 2cy + e} + \frac{(2ax + by - g)b}{(bx + 2cy + e)^2} \right) (xa_2 + ya_3 + a_1) \quad (\text{5E})$$

$$- \left(-\frac{b}{bx + 2cy + e} + \frac{2(2ax + by - g)c}{(bx + 2cy + e)^2} \right) (xb_2 + yb_3 + b_1) = 0$$

Putting the above in normal form gives

$$\frac{4a^2x^2a_3 - 2abx^2a_2 + 2abx^2b_3 + 4abxya_3 + 4acx^2b_2 - 8acxya_2 + 8acxyb_3 - 4acy^2a_3 - 2b^2x^2b_2 + 2b^2y^2}{= 0}$$

Setting the numerator to zero gives

$$\begin{aligned}
& -4a^2x^2a_3 + 2abx^2a_2 - 2abx^2b_3 - 4abxya_3 - 4acx^2b_2 + 8acxya_2 - 8acxyb_3 \\
& + 4acy^2a_3 + 2b^2x^2b_2 - 2b^2y^2a_3 + 4bcxyb_2 + 2bcy^2a_2 - 2bcy^2b_3 + 4c^2y^2b_2 \quad (6E) \\
& - 4acxb_1 + 4acya_1 + 4aexa_2 - 2aexb_3 + 2aeya_3 + 4agxa_3 + b^2xb_1 \\
& - b^2ya_1 + 3bexb_2 + beya_2 + bgxb_3 + 3bgya_3 + 4ceyb_2 + 2cgxb_2 - 2cgya_2 \\
& + 4cgyb_3 + 2aea_1 + beb_1 + bga_1 + 2cgb_1 + e^2b_2 - ega_2 + egb_3 - g^2a_3 = 0
\end{aligned}$$

Looking at the above PDE shows the following are all the terms with $\{x, y\}$ in them.

$$\{x, y\}$$

The following substitution is now made to be able to collect on all terms with $\{x, y\}$ in them

$$\{x = v_1, y = v_2\}$$

The above PDE (6E) now becomes

$$\begin{aligned}
& -4a^2a_3v_1^2 + 2aba_2v_1^2 - 4aba_3v_1v_2 - 2abb_3v_1^2 + 8aca_2v_1v_2 + 4aca_3v_2^2 \\
& - 4acb_2v_1^2 - 8acb_3v_1v_2 - 2b^2a_3v_2^2 + 2b^2b_2v_1^2 + 2bca_2v_2^2 + 4bcb_2v_1v_2 - 2bcb_3v_2^2 \quad (7E) \\
& + 4c^2b_2v_2^2 + 4aca_1v_2 - 4acb_1v_1 + 4aea_2v_1 + 2aea_3v_2 - 2aeb_3v_1 + 4aga_3v_1 \\
& - b^2a_1v_2 + b^2b_1v_1 + bea_2v_2 + 3beb_2v_1 + 3bga_3v_2 + bgb_3v_1 + 4ceb_2v_2 - 2cga_2v_2 \\
& + 2cgb_2v_1 + 4cgb_3v_2 + 2aea_1 + beb_1 + bga_1 + 2cgb_1 + e^2b_2 - ega_2 + egb_3 - g^2a_3 \\
& = 0
\end{aligned}$$

Collecting the above on the terms v_i introduced, and these are

$$\{v_1, v_2\}$$

Equation (7E) now becomes

$$\begin{aligned}
& (-4a^2a_3 + 2aba_2 - 2abb_3 - 4acb_2 + 2b^2b_2) v_1^2 \\
& + (-4aba_3 + 8aca_2 - 8acb_3 + 4bcb_2) v_1v_2 \\
& + (-4acb_1 + 4aea_2 - 2aeb_3 + 4aga_3 + b^2b_1 + 3beb_2 + bgb_3 + 2cgb_2) v_1 \quad (8E) \\
& + (4aca_3 - 2b^2a_3 + 2bca_2 - 2bcb_3 + 4c^2b_2) v_2^2 \\
& + (4aca_1 + 2aea_3 - b^2a_1 + bea_2 + 3bga_3 + 4ceb_2 - 2cga_2 + 4cgb_3) v_2 \\
& + 2aea_1 + beb_1 + bga_1 + 2cgb_1 + e^2b_2 - ega_2 + egb_3 - g^2a_3 = 0
\end{aligned}$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$\begin{aligned}
 -4aba_3 + 8aca_2 - 8acb_3 + 4cb_2 &= 0 \\
 4aca_3 - 2b^2a_3 + 2bca_2 - 2bcb_3 + 4c^2b_2 &= 0 \\
 -4a^2a_3 + 2aba_2 - 2abb_3 - 4acb_2 + 2b^2b_2 &= 0 \\
 4aca_1 + 2aea_3 - b^2a_1 + bea_2 + 3bga_3 + 4ceb_2 - 2cga_2 + 4cgb_3 &= 0 \\
 -4acb_1 + 4aea_2 - 2aeb_3 + 4aga_3 + b^2b_1 + 3beb_2 + bgb_3 + 2cgb_2 &= 0 \\
 2aea_1 + beb_1 + bga_1 + 2cgb_1 + e^2b_2 - ega_2 + egb_3 - g^2a_3 &= 0
 \end{aligned}$$

Solving the above equations for the unknowns gives

$$\begin{aligned}
 a_1 &= \frac{2acea_3 - b^2ea_3 - bceb_3 - bcga_3 - 2c^2gb_3}{c(4ac - b^2)} \\
 a_2 &= \frac{ba_3 + cb_3}{c} \\
 a_3 &= a_3 \\
 b_1 &= \frac{abea_3 + 2aceb_3 + 2acga_3 + bcgb_3}{c(4ac - b^2)} \\
 b_2 &= -\frac{aa_3}{c} \\
 b_3 &= b_3
 \end{aligned}$$

Substituting the above solution in the anstaz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$\begin{aligned}
 \xi &= \frac{4acx - b^2x - be - 2cg}{4ac - b^2} \\
 \eta &= \frac{4acy - b^2y + 2ae + bg}{4ac - b^2}
 \end{aligned}$$

Shifting is now applied to make $\xi = 0$ in order to simplify the rest of the computation

$$\begin{aligned}
 \eta &= \eta - \omega(x, y) \xi \\
 &= \frac{4acy - b^2y + 2ae + bg}{4ac - b^2} - \left(-\frac{2ax + by - g}{bx + 2cy + e} \right) \left(\frac{4acx - b^2x - be - 2cg}{4ac - b^2} \right) \\
 &= \frac{8a^2cx^2 - 2ab^2x^2 + 8abcxy + 8ac^2y^2 - 2b^3xy - 2b^2cy^2 + 8acey - 8acgx - 2b^2ey + 2b^2gx + 2ae^2 + 2bgx + 2cgy + 2e^2}{4abcx + 8ac^2y - b^3x - 2b^2cy + 4ace - b^2e}
 \end{aligned}$$

$$\xi = 0$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS \quad (1)$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = x$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{\frac{8a^2c x^2 - 2ab^2x^2 + 8abcxy + 8ac^2y^2 - 2b^3xy - 2b^2cy^2 + 8acey - 8acgx - 2b^2ey + 2b^2gx + 2ae^2 + 2beg + 2cg^2}{4abcx + 8ac^2y - b^3x - 2b^2cy + 4ace - b^2e}} dy \end{aligned}$$

Which results in

$$S = \frac{\ln(4a^2cx^2 - ab^2x^2 + 4abcxy + 4ac^2y^2 - b^3xy - b^2cy^2 + 4acey - 4acgx - b^2ey + b^2gx + ae^2 + beg)}{2}$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = -\frac{2ax + by - g}{bx + 2cy + e}$$

Evaluating all the partial derivatives gives

$$R_x = 1$$

$$R_y = 0$$

$$S_x = \frac{(2ax + by - g)(4ac - b^2)}{8a^2cx^2 + (-2b^2x^2 + 8cxyb + 8c^2y^2 + (8ey - 8gx)c + 2e^2)a - 2(by - g)(b^2x + (cy + e)b + cg)}$$

$$S_y = \frac{(bx + 2cy + e)(4ac - b^2)}{8a^2cx^2 + (-2b^2x^2 + 8cxyb + 8c^2y^2 + (8ey - 8gx)c + 2e^2)a - 2(by - g)(b^2x + (cy + e)b + cg)}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = 0 \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = 0$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S .

Since the ode has the form $\frac{d}{dR}S(R) = f(R)$, then we only need to integrate $f(R)$.

$$\int dS = \int 0 dR + c_2$$

$$S(R) = c_2$$

To complete the solution, we just need to transform the above back to x, y coordinates. This results in

$$\frac{\ln \left(4a^2c x^2 + \left(-b^2x^2 + 4ybcx - 4cgx + 4\left(cy + \frac{e}{2} \right)^2 \right) a - (by - g)(b^2x + (cy + e)b + cg) \right)}{2} = c_2$$

Summary of solutions found

$$\frac{\ln \left(4a^2c x^2 + \left(-b^2x^2 + 4ybcx - 4cgx + 4\left(cy + \frac{e}{2} \right)^2 \right) a - (by - g)(b^2x + (cy + e)b + cg) \right)}{2} = c_2$$

Solved as first order ode of type dAlembert

Time used: 1.000 (sec)

Let $p = y'$ the ode becomes

$$2ax + by + (bx + 2cy + e)p = g$$

Solving for y from the above results in

$$y = -\frac{(bp + 2a)x}{2cp + b} - \frac{ep - g}{2cp + b} \quad (1)$$

This has the form

$$y = xf(p) + g(p) \quad (*)$$

Where f, g are functions of $p = y'(x)$. The above ode is d'Alembert ode which is now solved.

Taking derivative of (*) w.r.t. x gives

$$\begin{aligned} p &= f + (xf' + g') \frac{dp}{dx} \\ p - f &= (xf' + g') \frac{dp}{dx} \end{aligned} \quad (2)$$

Comparing the form $y = xf + g$ to (1A) shows that

$$\begin{aligned} f &= \frac{-bp - 2a}{2cp + b} \\ g &= \frac{-ep + g}{2cp + b} \end{aligned}$$

Hence (2) becomes

$$p - \frac{-bp - 2a}{2cp + b} = \left(-\frac{xb}{2cp + b} + \frac{2xcbp}{(2cp + b)^2} + \frac{4xca}{(2cp + b)^2} - \frac{e}{2cp + b} + \frac{2cep}{(2cp + b)^2} - \frac{2cg}{(2cp + b)^2} \right) p'(x) \quad (2A)$$

The singular solution is found by setting $\frac{dp}{dx} = 0$ in the above which gives

$$p - \frac{-bp - 2a}{2cp + b} = 0$$

Solving the above for p results in

$$\begin{aligned} p_1 &= \frac{-b + \sqrt{-4ac + b^2}}{2c} \\ p_2 &= -\frac{b + \sqrt{-4ac + b^2}}{2c} \end{aligned}$$

Substituting these in (1A) and keeping singular solution that verifies the ode gives

$$\begin{aligned} y &= \frac{-bx\sqrt{-4ac + b^2} - 4acx + b^2x - e\sqrt{-4ac + b^2} + be + 2cg}{2\sqrt{-4ac + b^2}c} \\ y &= \frac{-bx\sqrt{-4ac + b^2} + 4acx - b^2x - e\sqrt{-4ac + b^2} - be - 2cg}{2\sqrt{-4ac + b^2}c} \end{aligned}$$

The general solution is found when $\frac{dp}{dx} \neq 0$. From eq. (2A). This results in

$$p'(x) = \frac{p(x) - \frac{-bp(x)-2a}{2cp(x)+b}}{-\frac{xb}{2cp(x)+b} + \frac{2xcbp(x)}{(2cp(x)+b)^2} + \frac{4xca}{(2cp(x)+b)^2} - \frac{e}{2cp(x)+b} + \frac{2cep(x)}{(2cp(x)+b)^2} - \frac{2cg}{(2cp(x)+b)^2}} \quad (3)$$

This ODE is now solved for $p(x)$. No inversion is needed. The ode $p'(x) = \frac{2(2cp(x)+b)(p(x)^2c+bp(x)+a)}{4acx-b^2x-be-2cg}$ is separable as it can be written as

$$\begin{aligned} p'(x) &= \frac{2(2cp(x)+b)(p(x)^2c+bp(x)+a)}{4acx-b^2x-be-2cg} \\ &= f(x)g(p) \end{aligned}$$

Where

$$\begin{aligned} f(x) &= \frac{2}{4acx-b^2x-be-2cg} \\ g(p) &= (2cp+b)(cp^2+bp+a) \end{aligned}$$

Integrating gives

$$\begin{aligned} \int \frac{1}{g(p)} dp &= \int f(x) dx \\ \int \frac{1}{(2cp+b)(cp^2+bp+a)} dp &= \int \frac{2}{4acx-b^2x-be-2cg} dx \\ \frac{\ln\left(\frac{(2cp(x)+b)^2}{p(x)^2c+bp(x)+a}\right)}{4ac-b^2} &= \frac{2 \ln((4ac-b^2)x-be-2cg)}{4ac-b^2} + c_1 \end{aligned}$$

We now need to find the singular solutions, these are found by finding for what values $g(p)$ is zero, since we had to divide by this above. Solving $g(p) = 0$ or $(2cp+b)(cp^2+bp+a) = 0$ for $p(x)$ gives

$$\begin{aligned} p(x) &= -\frac{b}{2c} \\ p(x) &= \frac{-b + \sqrt{-4ac + b^2}}{2c} \\ p(x) &= -\frac{b + \sqrt{-4ac + b^2}}{2c} \end{aligned}$$

Now we go over each such singular solution and check if it verifies the ode itself and any initial conditions given. If it does not then the singular solution will not be used.

Therefore the solutions found are

$$\ln \left(\frac{(2cp(x)+b)^2}{p(x)^2c+bp(x)+a} \right) = \frac{2 \ln ((4ac - b^2)x - be - 2cg)}{4ac - b^2} + c_1$$

$$p(x) = -\frac{b}{2c}$$

$$p(x) = \frac{-b + \sqrt{-4ac + b^2}}{2c}$$

$$p(x) = -\frac{b + \sqrt{-4ac + b^2}}{2c}$$

Substituting the above solution for p in (2A) gives

Expression too large to display

$$y = \frac{x \left(-\frac{b(-b+\sqrt{-4ac+b^2})}{2c} - 2a \right)}{\sqrt{-4ac+b^2}} + \frac{-\frac{e(-b+\sqrt{-4ac+b^2})}{2c} + g}{\sqrt{-4ac+b^2}}$$

$$y = -\frac{x \left(\frac{b(b+\sqrt{-4ac+b^2})}{2c} - 2a \right)}{\sqrt{-4ac+b^2}} - \frac{\frac{e(b+\sqrt{-4ac+b^2})}{2c} + g}{\sqrt{-4ac+b^2}}$$

Summary of solutions found

$$y = \frac{-bx\sqrt{-4ac+b^2} - 4acx + b^2x - e\sqrt{-4ac+b^2} + be + 2cg}{2\sqrt{-4ac+b^2}c}$$

$$y = \frac{-bx\sqrt{-4ac+b^2} + 4acx - b^2x - e\sqrt{-4ac+b^2} - be - 2cg}{2\sqrt{-4ac+b^2}c}$$

$$y = \frac{x \left(-\frac{b(-b+\sqrt{-4ac+b^2})}{2c} - 2a \right)}{\sqrt{-4ac+b^2}} + \frac{-\frac{e(-b+\sqrt{-4ac+b^2})}{2c} + g}{\sqrt{-4ac+b^2}}$$

$$y = -\frac{x \left(\frac{b(b+\sqrt{-4ac+b^2})}{2c} - 2a \right)}{\sqrt{-4ac+b^2}} - \frac{\frac{e(b+\sqrt{-4ac+b^2})}{2c} + g}{\sqrt{-4ac+b^2}}$$

Expression too large to display

Maple step by step solution

Let's solve

$$2ax + by + (2cy + bx + e) y' = g$$

- Highest derivative means the order of the ODE is 1
- Check if ODE is exact
 - ODE is exact if the lhs is the total derivative of a C^2 function

$$F'(x, y) = 0$$
 - Compute derivative of lhs

$$F'(x, y) + \left(\frac{\partial}{\partial y} F(x, y)\right) y' = 0$$
 - Evaluate derivatives

$$b = b$$
 - Condition met, ODE is exact
- Exact ODE implies solution will be of this form

$$\left[F(x, y) = C1, M(x, y) = F'(x, y), N(x, y) = \frac{\partial}{\partial y} F(x, y) \right]$$
- Solve for $F(x, y)$ by integrating $M(x, y)$ with respect to x

$$F(x, y) = \int (2ax + by - g) dx + _F1(y)$$
- Evaluate integral

$$F(x, y) = a x^2 + bxy - gx + _F1(y)$$
- Take derivative of $F(x, y)$ with respect to y

$$N(x, y) = \frac{\partial}{\partial y} F(x, y)$$
- Compute derivative

$$bx + 2cy + e = bx + \frac{d}{dy} _F1(y)$$
- Isolate for $\frac{d}{dy} _F1(y)$

$$\frac{d}{dy} _F1(y) = 2cy + e$$
- Solve for $_F1(y)$

$$_F1(y) = y^2 c + ey$$
- Substitute $_F1(y)$ into equation for $F(x, y)$

$$F(x, y) = a x^2 + bxy + y^2 c + ey - gx$$
- Substitute $F(x, y)$ into the solution of the ODE

$$a x^2 + bxy + y^2 c + ey - gx = C1$$
- Solve for y

$$\left\{ y = \frac{-bx-e+\sqrt{-4acx^2+b^2x^2+2bex+4cgx+4C1c+e^2}}{2c}, y = -\frac{bx+\sqrt{-4acx^2+b^2x^2+2bex+4cgx+4C1c+e^2+e}}{2c} \right\}$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
trying inverse linear
trying homogeneous types:
trying homogeneous C
trying homogeneous types:
trying homogeneous D
<- homogeneous successful
<- homogeneous successful`

```

Maple dsolve solution

Solving time : 0.096 (sec)

Leaf size : 88

```

dsolve(2*a*x+b*y(x)+(2*c*y(x)+b*x+e)*diff(y(x),x) = g,
       y(x),singsol=all)

```

$$y = \frac{\sqrt{-64 \left(ac - \frac{b^2}{4} \right) \left(c \left(ax - \frac{g}{2} \right) - \frac{b(bx+e)}{4} \right)^2 c_1^2 + 4c + (-4abcx + b^3x - 4ace + b^2e) c_1}}{8cc_1 \left(ac - \frac{b^2}{4} \right)}$$

Mathematica DSolve solution

Solving time : 17.055 (sec)

Leaf size : 132

```
DSolve[{(2*a*x+b*y[x])+(2*c*y[x]+b*x+e)*D[y[x],x]==g,{}},
y[x],x,IncludeSingularSolutions->True]
```

$$y(x) \rightarrow -\frac{\frac{\sqrt{4cx(g-ax)+b^2x^2+2bex+4c^2c_1+e^2}}{\sqrt{\frac{1}{c}}} + bx + e}{2c}$$

$$y(x) \rightarrow -\frac{-\frac{\sqrt{4cx(g-ax)+b^2x^2+2bex+4c^2c_1+e^2}}{\sqrt{\frac{1}{c}}} + bx + e}{2c}$$

2.3.4 problem 4

Solved as first order separable ode	215
Maple step by step solution	217
Maple trace	217
Maple dsolve solution	218
Mathematica DSolve solution	218

Internal problem ID [18224]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter IV. Methods of solution: First order equations. section 29. Problems at page 81

Problem number : 4

Date solved : Friday, December 20, 2024 at 10:45:27 AM

CAS classification : [_separable]

Solve

$$\sec(x)^2 \tan(y) y' + \sec(y)^2 \tan(x) = 0$$

Solved as first order separable ode

Time used: 0.308 (sec)

The ode $y' = -\frac{\sec(y)^2 \tan(x)}{\sec(x)^2 \tan(y)}$ is separable as it can be written as

$$\begin{aligned} y' &= -\frac{\sec(y)^2 \tan(x)}{\sec(x)^2 \tan(y)} \\ &= f(x)g(y) \end{aligned}$$

Where

$$\begin{aligned} f(x) &= -\frac{\tan(x)}{\sec(x)^2} \\ g(y) &= \frac{\sec(y)^2}{\tan(y)} \end{aligned}$$

Integrating gives

$$\int \frac{1}{g(y)} dy = \int f(x) dx$$

$$\int \frac{\tan(y)}{\sec(y)^2} dy = \int -\frac{\tan(x)}{\sec(x)^2} dx$$

$$-\frac{\cos(y)^2}{2} = \frac{\cos(x)^2}{2} + c_1$$

Solving for y gives

$$y = \pi - \arccos\left(\sqrt{-\cos(x)^2 - 2c_1}\right)$$

$$y = \arccos\left(\sqrt{-\cos(x)^2 - 2c_1}\right)$$

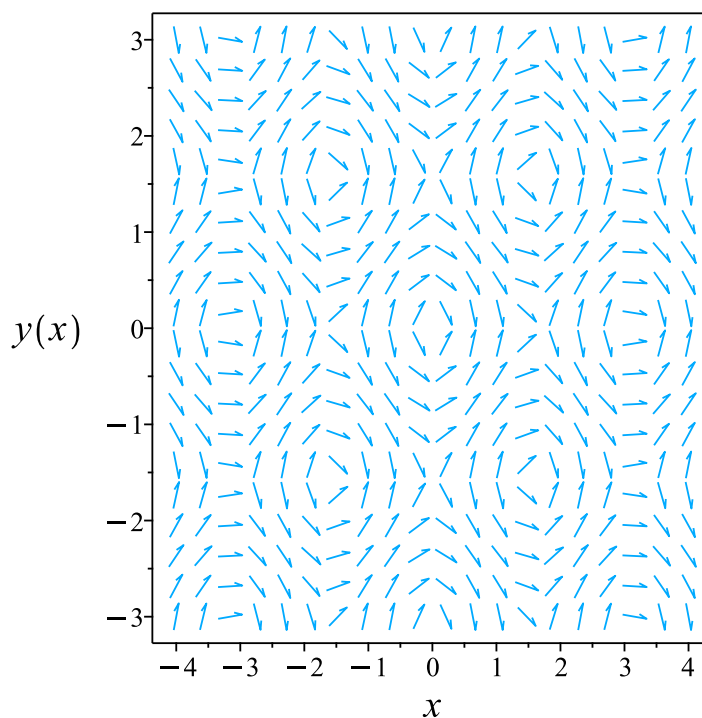


Figure 2.31: Slope field plot
 $\sec(x)^2 \tan(y) y' + \sec(y)^2 \tan(x) = 0$

Summary of solutions found

$$y = \pi - \arccos\left(\sqrt{-\cos(x)^2 - 2c_1}\right)$$

$$y = \arccos \left(\sqrt{-\cos(x)^2 - 2c_1} \right)$$

Maple step by step solution

Let's solve

$$\sec(x)^2 \tan(y) y' + \sec(y)^2 \tan(x) = 0$$

- Highest derivative means the order of the ODE is 1

y'

- Solve for the highest derivative

$$y' = -\frac{\sec(y)^2 \tan(x)}{\sec(x)^2 \tan(y)}$$

- Separate variables

$$\frac{y' \tan(y)}{\sec(y)^2} = -\frac{\tan(x)}{\sec(x)^2}$$

- Integrate both sides with respect to x

$$\int \frac{y' \tan(y)}{\sec(y)^2} dx = \int -\frac{\tan(x)}{\sec(x)^2} dx + C1$$

- Evaluate integral

$$-\frac{1}{2\sec(y)^2} = \frac{1}{2\sec(x)^2} + C1$$

- Solve for y

$$\left\{ y = \pi - \operatorname{arcsec} \left(\frac{1}{\sqrt{-\cos(x)^2 - 2C1}} \right), y = \operatorname{arcsec} \left(\frac{1}{\sqrt{-\cos(x)^2 - 2C1}} \right) \right\}$$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
<- separable successful`
```

Maple dsolve solution

Solving time : 0.006 (sec)

Leaf size : 41

```
dsolve(sec(x)^2*tan(y(x))*diff(y(x),x)+sec(y(x))^2*tan(x) = 0,
        y(x),singsol=all)
```

$$y = \operatorname{arcsec} \left(\frac{2}{\sqrt{-2 \cos(2x) + 8c_1}} \right)$$

$$y = \frac{\pi}{2} + \operatorname{arccsc} \left(\frac{2}{\sqrt{-2 \cos(2x) + 8c_1}} \right)$$

Mathematica DSolve solution

Solving time : 0.522 (sec)

Leaf size : 41

```
DSolve[{Sec[x]^2*Tan[y[x]]*D[y[x],x]+Sec[y[x]]^2*Tan[x]==0,{}},
        y[x],x,IncludeSingularSolutions->True]
```

$$y(x) \rightarrow -\frac{1}{2} \arccos(-\cos(2x) - 2c_1)$$

$$y(x) \rightarrow \frac{1}{2} \arccos(-\cos(2x) - 2c_1)$$

2.3.5 problem 5

Solved as first order homogeneous class A ode	219
Solved as first order homogeneous class D2 ode	222
Solved as first order homogeneous class Maple C ode	224
Solved as first order isobaric ode	228
Solved using Lie symmetry for first order ode	230
Maple step by step solution	235
Maple trace	235
Maple dsolve solution	235
Mathematica DSolve solution	236

Internal problem ID [18225]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter IV. Methods of solution: First order equations. section 29. Problems at page 81

Problem number : 5

Date solved : Friday, December 20, 2024 at 10:46:05 AM

CAS classification :

[[_homogeneous, 'class A'], _rational, [_Abel, '2nd type', 'class A']]

Solve

$$x + yy' = my$$

Solved as first order homogeneous class A ode

Time used: 0.606 (sec)

In canonical form, the ODE is

$$\begin{aligned} y' &= F(x, y) \\ &= \frac{my - x}{y} \end{aligned} \tag{1}$$

An ode of the form $y' = \frac{M(x,y)}{N(x,y)}$ is called homogeneous if the functions $M(x, y)$ and $N(x, y)$ are both homogeneous functions and of the same order. Recall that a function $f(x, y)$ is homogeneous of order n if

$$f(t^n x, t^n y) = t^n f(x, y)$$

In this case, it can be seen that both $M = my - x$ and $N = y$ are both homogeneous and of the same order $n = 1$. Therefore this is a homogeneous ode. Since this ode is homogeneous, it is converted to separable ODE using the substitution $u = \frac{y}{x}$, or $y = ux$. Hence

$$\frac{dy}{dx} = \frac{du}{dx}x + u$$

Applying the transformation $y = ux$ to the above ODE in (1) gives

$$\begin{aligned} \frac{du}{dx}x + u &= m - \frac{1}{u} \\ \frac{du}{dx} &= \frac{m - \frac{1}{u(x)} - u(x)}{x} \end{aligned}$$

Or

$$u'(x) - \frac{m - \frac{1}{u(x)} - u(x)}{x} = 0$$

Or

$$u'(x)u(x)x + u(x)^2 - mu(x) + 1 = 0$$

Which is now solved as separable in $u(x)$.

The ode $u'(x) = -\frac{u(x)^2 - mu(x) + 1}{u(x)x}$ is separable as it can be written as

$$\begin{aligned} u'(x) &= -\frac{u(x)^2 - mu(x) + 1}{u(x)x} \\ &= f(x)g(u) \end{aligned}$$

Where

$$\begin{aligned} f(x) &= -\frac{1}{x} \\ g(u) &= \frac{-mu + u^2 + 1}{u} \end{aligned}$$

Integrating gives

$$\begin{aligned} \int \frac{1}{g(u)} du &= \int f(x) dx \\ \int \frac{u}{-mu + u^2 + 1} du &= \int -\frac{1}{x} dx \\ \frac{\ln(u(x)^2 - mu(x) + 1)}{2} + \frac{m \operatorname{arctanh}\left(\frac{m-2u(x)}{\sqrt{m^2-4}}\right)}{\sqrt{m^2-4}} &= \ln\left(\frac{1}{x}\right) + c_1 \end{aligned}$$

We now need to find the singular solutions, these are found by finding for what values $g(u)$ is zero, since we had to divide by this above. Solving $g(u) = 0$ or $\frac{-mu+u^2+1}{u} = 0$ for $u(x)$ gives

$$u(x) = \frac{m}{2} - \frac{\sqrt{m^2 - 4}}{2}$$

$$u(x) = \frac{m}{2} + \frac{\sqrt{m^2 - 4}}{2}$$

Now we go over each such singular solution and check if it verifies the ode itself and any initial conditions given. If it does not then the singular solution will not be used.

Therefore the solutions found are

$$\frac{\ln(u(x)^2 - mu(x) + 1)}{2} + \frac{m \operatorname{arctanh}\left(\frac{m-2u(x)}{\sqrt{m^2-4}}\right)}{\sqrt{m^2-4}} = \ln\left(\frac{1}{x}\right) + c_1$$

$$u(x) = \frac{m}{2} - \frac{\sqrt{m^2-4}}{2}$$

$$u(x) = \frac{m}{2} + \frac{\sqrt{m^2-4}}{2}$$

Converting $\frac{\ln(u(x)^2 - mu(x) + 1)}{2} + \frac{m \operatorname{arctanh}\left(\frac{m-2u(x)}{\sqrt{m^2-4}}\right)}{\sqrt{m^2-4}} = \ln\left(\frac{1}{x}\right) + c_1$ back to y gives

$$\frac{\operatorname{arctanh}\left(\frac{mx-2y}{x\sqrt{(m-2)(m+2)}}\right) m}{\sqrt{(m-2)(m+2)}} + \frac{\ln\left(\frac{-myx+y^2+x^2}{x^2}\right)}{2} = \ln\left(\frac{1}{x}\right) + c_1$$

Converting $u(x) = \frac{m}{2} - \frac{\sqrt{m^2-4}}{2}$ back to y gives

$$y = x \left(\frac{m}{2} - \frac{\sqrt{m^2-4}}{2} \right)$$

Converting $u(x) = \frac{m}{2} + \frac{\sqrt{m^2-4}}{2}$ back to y gives

$$y = x \left(\frac{m}{2} + \frac{\sqrt{m^2-4}}{2} \right)$$

Summary of solutions found

$$\frac{\operatorname{arctanh}\left(\frac{mx-2y}{x\sqrt{(m-2)(m+2)}}\right) m}{\sqrt{(m-2)(m+2)}} + \frac{\ln\left(\frac{-myx+y^2+x^2}{x^2}\right)}{2} = \ln\left(\frac{1}{x}\right) + c_1$$

$$y = x \left(\frac{m}{2} - \frac{\sqrt{m^2 - 4}}{2} \right)$$

$$y = x \left(\frac{m}{2} + \frac{\sqrt{m^2 - 4}}{2} \right)$$

Solved as first order homogeneous class D2 ode

Time used: 0.245 (sec)

Applying change of variables $y = u(x)x$, then the ode becomes

$$x + u(x)x(u'(x)x + u(x)) = mu(x)x$$

Which is now solved The ode $u'(x) = -\frac{u(x)^2 - mu(x) + 1}{u(x)x}$ is separable as it can be written as

$$\begin{aligned} u'(x) &= -\frac{u(x)^2 - mu(x) + 1}{u(x)x} \\ &= f(x)g(u) \end{aligned}$$

Where

$$\begin{aligned} f(x) &= -\frac{1}{x} \\ g(u) &= \frac{-mu + u^2 + 1}{u} \end{aligned}$$

Integrating gives

$$\begin{aligned} \int \frac{1}{g(u)} du &= \int f(x) dx \\ \int \frac{u}{-mu + u^2 + 1} du &= \int -\frac{1}{x} dx \\ \frac{\ln(u(x)^2 - mu(x) + 1)}{2} + \frac{m \operatorname{arctanh}\left(\frac{m-2u(x)}{\sqrt{m^2-4}}\right)}{\sqrt{m^2-4}} &= \ln\left(\frac{1}{x}\right) + c_1 \end{aligned}$$

We now need to find the singular solutions, these are found by finding for what values $g(u)$ is zero, since we had to divide by this above. Solving $g(u) = 0$ or $\frac{-mu+u^2+1}{u} = 0$ for $u(x)$ gives

$$\begin{aligned} u(x) &= \frac{m}{2} - \frac{\sqrt{m^2 - 4}}{2} \\ u(x) &= \frac{m}{2} + \frac{\sqrt{m^2 - 4}}{2} \end{aligned}$$

Now we go over each such singular solution and check if it verifies the ode itself and any initial conditions given. If it does not then the singular solution will not be used.

Therefore the solutions found are

$$\frac{\ln(u(x)^2 - mu(x) + 1)}{2} + \frac{m \operatorname{arctanh}\left(\frac{m-2u(x)}{\sqrt{m^2-4}}\right)}{\sqrt{m^2-4}} = \ln\left(\frac{1}{x}\right) + c_1$$

$$u(x) = \frac{m}{2} - \frac{\sqrt{m^2-4}}{2}$$

$$u(x) = \frac{m}{2} + \frac{\sqrt{m^2-4}}{2}$$

Converting $\frac{\ln(u(x)^2 - mu(x) + 1)}{2} + \frac{m \operatorname{arctanh}\left(\frac{m-2u(x)}{\sqrt{m^2-4}}\right)}{\sqrt{m^2-4}} = \ln\left(\frac{1}{x}\right) + c_1$ back to y gives

$$\frac{\ln\left(\frac{y^2}{x^2} - \frac{my}{x} + 1\right)}{2} + \frac{m \operatorname{arctanh}\left(\frac{m-\frac{2y}{x}}{\sqrt{m^2-4}}\right)}{\sqrt{m^2-4}} = \ln\left(\frac{1}{x}\right) + c_1$$

Converting $u(x) = \frac{m}{2} - \frac{\sqrt{m^2-4}}{2}$ back to y gives

$$y = x \left(\frac{m}{2} - \frac{\sqrt{m^2-4}}{2} \right)$$

Converting $u(x) = \frac{m}{2} + \frac{\sqrt{m^2-4}}{2}$ back to y gives

$$y = x \left(\frac{m}{2} + \frac{\sqrt{m^2-4}}{2} \right)$$

Summary of solutions found

$$\frac{\ln\left(\frac{y^2}{x^2} - \frac{my}{x} + 1\right)}{2} + \frac{m \operatorname{arctanh}\left(\frac{m-\frac{2y}{x}}{\sqrt{m^2-4}}\right)}{\sqrt{m^2-4}} = \ln\left(\frac{1}{x}\right) + c_1$$

$$y = x \left(\frac{m}{2} - \frac{\sqrt{m^2-4}}{2} \right)$$

$$y = x \left(\frac{m}{2} + \frac{\sqrt{m^2-4}}{2} \right)$$

Summary of solutions found

$$\frac{\operatorname{arctanh}\left(\frac{mx-2y}{x\sqrt{(m-2)(m+2)}}\right) m}{\sqrt{(m-2)(m+2)}} + \frac{\ln\left(\frac{-myx+y^2+x^2}{x^2}\right)}{2} = \ln\left(\frac{1}{x}\right) + c_1$$

$$y = x\left(\frac{m}{2} - \frac{\sqrt{m^2-4}}{2}\right)$$

$$y = x\left(\frac{m}{2} + \frac{\sqrt{m^2-4}}{2}\right)$$

Solved as first order homogeneous class Maple C ode

Time used: 0.757 (sec)

Let $Y = y - y_0$ and $X = x - x_0$ then the above is transformed to new ode in $Y(X)$

$$\frac{d}{dX}Y(X) = \frac{m(Y(X) + y_0) - x_0 - X}{Y(X) + y_0}$$

Solving for possible values of x_0 and y_0 which makes the above ode a homogeneous ode results in

$$x_0 = 0$$

$$y_0 = 0$$

Using these values now it is possible to easily solve for $Y(X)$. The above ode now becomes

$$\frac{d}{dX}Y(X) = \frac{mY(X) - X}{Y(X)}$$

In canonical form, the ODE is

$$\begin{aligned} Y' &= F(X, Y) \\ &= \frac{mY - X}{Y} \end{aligned} \tag{1}$$

An ode of the form $Y' = \frac{M(X, Y)}{N(X, Y)}$ is called homogeneous if the functions $M(X, Y)$ and $N(X, Y)$ are both homogeneous functions and of the same order. Recall that a function $f(X, Y)$ is homogeneous of order n if

$$f(t^n X, t^n Y) = t^n f(X, Y)$$

In this case, it can be seen that both $M = mY - X$ and $N = Y$ are both homogeneous and of the same order $n = 1$. Therefore this is a homogeneous ode. Since this ode is homogeneous, it is converted to separable ODE using the substitution $u = \frac{Y}{X}$, or $Y = uX$. Hence

$$\frac{dY}{dX} = \frac{du}{dX}X + u$$

Applying the transformation $Y = uX$ to the above ODE in (1) gives

$$\begin{aligned} \frac{du}{dX}X + u &= m - \frac{1}{u} \\ \frac{du}{dX} &= \frac{m - \frac{1}{u(X)} - u(X)}{X} \end{aligned}$$

Or

$$\frac{d}{dX}u(X) - \frac{m - \frac{1}{u(X)} - u(X)}{X} = 0$$

Or

$$\left(\frac{d}{dX}u(X) \right) u(X) X + u(X)^2 - mu(X) + 1 = 0$$

Which is now solved as separable in $u(X)$.

The ode $\frac{d}{dX}u(X) = -\frac{u(X)^2 - mu(X) + 1}{u(X)X}$ is separable as it can be written as

$$\begin{aligned} \frac{d}{dX}u(X) &= -\frac{u(X)^2 - mu(X) + 1}{u(X)X} \\ &= f(X)g(u) \end{aligned}$$

Where

$$\begin{aligned} f(X) &= -\frac{1}{X} \\ g(u) &= \frac{-mu + u^2 + 1}{u} \end{aligned}$$

Integrating gives

$$\begin{aligned} \int \frac{1}{g(u)} du &= \int f(X) dX \\ \int \frac{u}{-mu + u^2 + 1} du &= \int -\frac{1}{X} dX \\ \frac{\ln(u(X)^2 - mu(X) + 1)}{2} + \frac{m \operatorname{arctanh}\left(\frac{m-2u(X)}{\sqrt{m^2-4}}\right)}{\sqrt{m^2-4}} &= \ln\left(\frac{1}{X}\right) + c_1 \end{aligned}$$

We now need to find the singular solutions, these are found by finding for what values $g(u)$ is zero, since we had to divide by this above. Solving $g(u) = 0$ or $\frac{-mu+u^2+1}{u} = 0$ for $u(X)$ gives

$$u(X) = \frac{m}{2} - \frac{\sqrt{m^2 - 4}}{2}$$

$$u(X) = \frac{m}{2} + \frac{\sqrt{m^2 - 4}}{2}$$

Now we go over each such singular solution and check if it verifies the ode itself and any initial conditions given. If it does not then the singular solution will not be used.

Therefore the solutions found are

$$\frac{\ln(u(X)^2 - mu(X) + 1)}{2} + \frac{m \operatorname{arctanh}\left(\frac{m-2u(X)}{\sqrt{m^2-4}}\right)}{\sqrt{m^2-4}} = \ln\left(\frac{1}{X}\right) + c_1$$

$$u(X) = \frac{m}{2} - \frac{\sqrt{m^2-4}}{2}$$

$$u(X) = \frac{m}{2} + \frac{\sqrt{m^2-4}}{2}$$

Converting $\frac{\ln(u(X)^2 - mu(X) + 1)}{2} + \frac{m \operatorname{arctanh}\left(\frac{m-2u(X)}{\sqrt{m^2-4}}\right)}{\sqrt{m^2-4}} = \ln\left(\frac{1}{X}\right) + c_1$ back to $Y(X)$ gives

$$\frac{\operatorname{arctanh}\left(\frac{mX-2Y(X)}{X\sqrt{(m-2)(m+2)}}\right) m}{\sqrt{(m-2)(m+2)}} + \frac{\ln\left(\frac{-mY(X)X+Y(X)^2+X^2}{X^2}\right)}{2} = \ln\left(\frac{1}{X}\right) + c_1$$

Converting $u(X) = \frac{m}{2} - \frac{\sqrt{m^2-4}}{2}$ back to $Y(X)$ gives

$$Y(X) = X\left(\frac{m}{2} - \frac{\sqrt{m^2-4}}{2}\right)$$

Converting $u(X) = \frac{m}{2} + \frac{\sqrt{m^2-4}}{2}$ back to $Y(X)$ gives

$$Y(X) = X\left(\frac{m}{2} + \frac{\sqrt{m^2-4}}{2}\right)$$

Using the solution for $Y(X)$

$$\frac{\operatorname{arctanh}\left(\frac{mX-2Y(X)}{X\sqrt{(m-2)(m+2)}}\right) m}{\sqrt{(m-2)(m+2)}} + \frac{\ln\left(\frac{-mY(X)X+Y(X)^2+X^2}{X^2}\right)}{2} = \ln\left(\frac{1}{X}\right) + c_1 \quad (\text{A})$$

And replacing back terms in the above solution using

$$\begin{aligned} Y &= y + y_0 \\ X &= x + x_0 \end{aligned}$$

Or

$$\begin{aligned} Y &= y \\ X &= x \end{aligned}$$

Then the solution in y becomes using EQ (A)

$$\frac{\operatorname{arctanh}\left(\frac{mx-2y}{x\sqrt{(m-2)(m+2)}}\right) m}{\sqrt{(m-2)(m+2)}} + \frac{\ln\left(\frac{-myx+y^2+x^2}{x^2}\right)}{2} = \ln\left(\frac{1}{x}\right) + c_1$$

Using the solution for $Y(X)$

$$Y(X) = X \left(\frac{m}{2} - \frac{\sqrt{m^2 - 4}}{2} \right) \quad (\text{A})$$

And replacing back terms in the above solution using

$$\begin{aligned} Y &= y + y_0 \\ X &= x + x_0 \end{aligned}$$

Or

$$\begin{aligned} Y &= y \\ X &= x \end{aligned}$$

Then the solution in y becomes using EQ (A)

$$y = x \left(\frac{m}{2} - \frac{\sqrt{m^2 - 4}}{2} \right)$$

Using the solution for $Y(X)$

$$Y(X) = X \left(\frac{m}{2} + \frac{\sqrt{m^2 - 4}}{2} \right) \quad (\text{A})$$

And replacing back terms in the above solution using

$$\begin{aligned} Y &= y + y_0 \\ X &= x + x_0 \end{aligned}$$

Or

$$Y = y$$

$$X = x$$

Then the solution in y becomes using EQ (A)

$$y = x \left(\frac{m}{2} + \frac{\sqrt{m^2 - 4}}{2} \right)$$

Solved as first order isobaric ode

Time used: 0.287 (sec)

Solving for y' gives

$$y' = \frac{my - x}{y} \quad (1)$$

Each of the above ode's is now solved An ode $y' = f(x, y)$ is isobaric if

$$f(tx, t^m y) = t^{m-1} f(x, y) \quad (1)$$

Where here

$$f(x, y) = \frac{my - x}{y} \quad (2)$$

m is the order of isobaric. Substituting (2) into (1) and solving for m gives

$$m = 1$$

Since the ode is isobaric of order $m = 1$, then the substitution

$$\begin{aligned} y &= ux^m \\ &= ux \end{aligned}$$

Converts the ODE to a separable in $u(x)$. Performing this substitution gives

$$u(x) + xu'(x) = \frac{mxu(x) - x}{xu(x)}$$

The ode $u'(x) = -\frac{u(x)^2 - u(x)m + 1}{u(x)x}$ is separable as it can be written as

$$\begin{aligned} u'(x) &= -\frac{u(x)^2 - u(x)m + 1}{u(x)x} \\ &= f(x)g(u) \end{aligned}$$

Where

$$f(x) = -\frac{1}{x}$$

$$g(u) = \frac{-um + u^2 + 1}{u}$$

Integrating gives

$$\int \frac{1}{g(u)} du = \int f(x) dx$$

$$\int \frac{u}{-um + u^2 + 1} du = \int -\frac{1}{x} dx$$

$$\frac{\ln(u(x)^2 - u(x)m + 1)}{2} + \frac{m \operatorname{arctanh}\left(\frac{m-2u(x)}{\sqrt{m^2-4}}\right)}{\sqrt{m^2-4}} = \ln\left(\frac{1}{x}\right) + c_1$$

We now need to find the singular solutions, these are found by finding for what values $g(u)$ is zero, since we had to divide by this above. Solving $g(u) = 0$ or $\frac{-um+u^2+1}{u} = 0$ for $u(x)$ gives

$$u(x) = \frac{m}{2} - \frac{\sqrt{m^2-4}}{2}$$

$$u(x) = \frac{m}{2} + \frac{\sqrt{m^2-4}}{2}$$

Now we go over each such singular solution and check if it verifies the ode itself and any initial conditions given. If it does not then the singular solution will not be used.

Therefore the solutions found are

$$\frac{\ln(u(x)^2 - u(x)m + 1)}{2} + \frac{m \operatorname{arctanh}\left(\frac{m-2u(x)}{\sqrt{m^2-4}}\right)}{\sqrt{m^2-4}} = \ln\left(\frac{1}{x}\right) + c_1$$

$$u(x) = \frac{m}{2} - \frac{\sqrt{m^2-4}}{2}$$

$$u(x) = \frac{m}{2} + \frac{\sqrt{m^2-4}}{2}$$

Converting $\frac{\ln(u(x)^2 - u(x)m + 1)}{2} + \frac{m \operatorname{arctanh}\left(\frac{m-2u(x)}{\sqrt{m^2-4}}\right)}{\sqrt{m^2-4}} = \ln\left(\frac{1}{x}\right) + c_1$ back to y gives

$$\frac{\ln\left(\frac{y^2}{x^2} - \frac{my}{x} + 1\right)}{2} + \frac{m \operatorname{arctanh}\left(\frac{m-\frac{2y}{x}}{\sqrt{m^2-4}}\right)}{\sqrt{m^2-4}} = \ln\left(\frac{1}{x}\right) + c_1$$

Converting $u(x) = \frac{m}{2} - \frac{\sqrt{m^2-4}}{2}$ back to y gives

$$\frac{y}{x} = \frac{m}{2} - \frac{\sqrt{m^2-4}}{2}$$

Converting $u(x) = \frac{m}{2} + \frac{\sqrt{m^2-4}}{2}$ back to y gives

$$\frac{y}{x} = \frac{m}{2} + \frac{\sqrt{m^2-4}}{2}$$

Solving for y gives

$$\frac{\ln\left(\frac{y^2}{x^2} - \frac{my}{x} + 1\right)}{2} + \frac{m \operatorname{arctanh}\left(\frac{m-\frac{2y}{x}}{\sqrt{m^2-4}}\right)}{\sqrt{m^2-4}} = \ln\left(\frac{1}{x}\right) + c_1$$

$$y = -\frac{x(-m + \sqrt{m^2-4})}{2}$$

$$y = \frac{x(m + \sqrt{m^2-4})}{2}$$

Summary of solutions found

$$\frac{\ln\left(\frac{y^2}{x^2} - \frac{my}{x} + 1\right)}{2} + \frac{m \operatorname{arctanh}\left(\frac{m-\frac{2y}{x}}{\sqrt{m^2-4}}\right)}{\sqrt{m^2-4}} = \ln\left(\frac{1}{x}\right) + c_1$$

$$y = -\frac{x(-m + \sqrt{m^2-4})}{2}$$

$$y = \frac{x(m + \sqrt{m^2-4})}{2}$$

Solved using Lie symmetry for first order ode

Time used: 7.844 (sec)

Writing the ode as

$$y' = \frac{my - x}{y}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

To determine ξ, η then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$\xi = xa_2 + ya_3 + a_1 \quad (1E)$$

$$\eta = xb_2 + yb_3 + b_1 \quad (2E)$$

Where the unknown coefficients are

$$\{a_1, a_2, a_3, b_1, b_2, b_3\}$$

Substituting equations (1E,2E) and ω into (A) gives

$$b_2 + \frac{(my-x)(b_3-a_2)}{y} - \frac{(my-x)^2 a_3}{y^2} + \frac{xa_2 + ya_3 + a_1}{y} \quad (5E)$$

$$- \left(\frac{m}{y} - \frac{my-x}{y^2} \right) (xb_2 + yb_3 + b_1) = 0$$

Putting the above in normal form gives

$$\frac{m^2 y^2 a_3 - 2mxy a_3 + m y^2 a_2 - m y^2 b_3 + x^2 a_3 + x^2 b_2 - 2xy a_2 + 2xy b_3 - y^2 a_3 - b_2 y^2 + x b_1 - y a_1}{y^2} = 0$$

Setting the numerator to zero gives

$$-m^2 y^2 a_3 + 2mxy a_3 - m y^2 a_2 + m y^2 b_3 - x^2 a_3 - x^2 b_2 \quad (6E)$$

$$+ 2xy a_2 - 2xy b_3 + y^2 a_3 + b_2 y^2 - x b_1 + y a_1 = 0$$

Looking at the above PDE shows the following are all the terms with $\{x, y\}$ in them.

$$\{x, y\}$$

The following substitution is now made to be able to collect on all terms with $\{x, y\}$ in them

$$\{x = v_1, y = v_2\}$$

The above PDE (6E) now becomes

$$\begin{aligned} & -m^2 a_3 v_2^2 - m a_2 v_2^2 + 2m a_3 v_1 v_2 + m b_3 v_2^2 + 2a_2 v_1 v_2 - a_3 v_1^2 \\ & + a_3 v_2^2 - b_2 v_1^2 + b_2 v_2^2 - 2b_3 v_1 v_2 + a_1 v_2 - b_1 v_1 = 0 \end{aligned} \quad (7E)$$

Collecting the above on the terms v_i introduced, and these are

$$\{v_1, v_2\}$$

Equation (7E) now becomes

$$\begin{aligned} & (-a_3 - b_2) v_1^2 + (2m a_3 + 2a_2 - 2b_3) v_1 v_2 - b_1 v_1 \\ & + (-m^2 a_3 - m a_2 + m b_3 + a_3 + b_2) v_2^2 + a_1 v_2 = 0 \end{aligned} \quad (8E)$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$\begin{aligned} a_1 &= 0 \\ -b_1 &= 0 \\ -a_3 - b_2 &= 0 \\ 2m a_3 + 2a_2 - 2b_3 &= 0 \\ -m^2 a_3 - m a_2 + m b_3 + a_3 + b_2 &= 0 \end{aligned}$$

Solving the above equations for the unknowns gives

$$\begin{aligned} a_1 &= 0 \\ a_2 &= m b_2 + b_3 \\ a_3 &= -b_2 \\ b_1 &= 0 \\ b_2 &= b_2 \\ b_3 &= b_3 \end{aligned}$$

Substituting the above solution in the anstaz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$\begin{aligned} \xi &= x \\ \eta &= y \end{aligned}$$

Shifting is now applied to make $\xi = 0$ in order to simplify the rest of the computation

$$\begin{aligned}\eta &= \eta - \omega(x, y) \xi \\ &= y - \left(\frac{my - x}{y} \right) (x) \\ &= \frac{-myx + x^2 + y^2}{y} \\ \xi &= 0\end{aligned}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS \quad (1)$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y} \right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = x$$

S is found from

$$\begin{aligned}S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{\frac{-myx + x^2 + y^2}{y}} dy\end{aligned}$$

Which results in

$$S = \frac{\ln(-myx + x^2 + y^2)}{2} - \frac{mx \operatorname{arctanh}\left(\frac{-mx+2y}{\sqrt{m^2x^2-4x^2}}\right)}{\sqrt{m^2x^2-4x^2}}$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = \frac{my - x}{y}$$

Evaluating all the partial derivatives gives

$$\begin{aligned}R_x &= 1 \\R_y &= 0 \\S_x &= \frac{my - x}{myx - x^2 - y^2} \\S_y &= -\frac{y}{myx - x^2 - y^2}\end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = 0 \tag{2A}$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = 0$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S .

Since the ode has the form $\frac{d}{dR}S(R) = f(R)$, then we only need to integrate $f(R)$.

$$\begin{aligned}\int dS &= \int 0 dR + c_2 \\S(R) &= c_2\end{aligned}$$

To complete the solution, we just need to transform the above back to x, y coordinates. This results in

$$\frac{\ln(-myx + y^2 + x^2)}{2} + \frac{m \operatorname{arctanh}\left(\frac{mx-2y}{x\sqrt{m^2-4}}\right)}{\sqrt{m^2-4}} = c_2$$

Summary of solutions found

$$\frac{\ln(-myx + y^2 + x^2)}{2} + \frac{m \operatorname{arctanh}\left(\frac{mx-2y}{x\sqrt{m^2-4}}\right)}{\sqrt{m^2-4}} = c_2$$

Maple step by step solution

Let's solve

$$x + yy' = my$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Solve for the highest derivative

$$y' = \frac{my-x}{y}$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
trying inverse linear
trying homogeneous types:
trying homogeneous D
<- homogeneous successful`

```

Maple dsolve solution

Solving time : 0.076 (sec)

Leaf size : 57

```

dsolve(x+y(x)*diff(y(x),x) = m*y(x),
       y(x),singsol=all)

```

$$y = \text{RootOf} \left(_Z^2 - e^{\text{RootOf} \left(x^2 \left(4 e^{-Z \cosh \left(\frac{\sqrt{m^2-4} (2c_1 + _Z + 2 \ln(x))}{2m} \right)^2 + m^2 - 4} \right)} + 1 - m_Z \right)} x \right)$$

Mathematica DSolve solution

Solving time : 0.093 (sec)

Leaf size : 72

```
DSolve[{x+y[x]*D[y[x],x]==m*y[x],{}},  
y[x],x,IncludeSingularSolutions->True]
```

$$\text{Solve} \left[\frac{m \arctan \left(\frac{\frac{2y(x)}{x} - m}{\sqrt{4 - m^2}} \right)}{\sqrt{4 - m^2}} + \frac{1}{2} \log \left(-\frac{my(x)}{x} + \frac{y(x)^2}{x^2} + 1 \right) = -\log(x) + c_1, y(x) \right]$$

2.3.6 problem 6

Solved as first order homogeneous class A ode	237
Solved as first order homogeneous class D2 ode	240
Solved as first order homogeneous class Maple C ode	243
Solved as first order Exact ode	248
Solved as first order isobaric ode	252
Solved using Lie symmetry for first order ode	255
Maple step by step solution	260
Maple trace	261
Maple dsolve solution	262
Mathematica DSolve solution	263

Internal problem ID [18226]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter IV. Methods of solution: First order equations. section 29. Problems at page 81

Problem number : 6

Date solved : Friday, December 20, 2024 at 10:46:16 AM

CAS classification :

[[_homogeneous, 'class A'], _exact, _rational, _dAlembert]

Solve

$$\frac{2x}{y^3} + \left(\frac{1}{y^2} - \frac{3x^2}{y^4} \right) y' = 0$$

Solved as first order homogeneous class A ode

Time used: 0.372 (sec)

In canonical form, the ODE is

$$\begin{aligned} y' &= F(x, y) \\ &= -\frac{2yx}{-3x^2 + y^2} \end{aligned} \tag{1}$$

An ode of the form $y' = \frac{M(x,y)}{N(x,y)}$ is called homogeneous if the functions $M(x, y)$ and $N(x, y)$ are both homogeneous functions and of the same order. Recall that a function $f(x, y)$ is homogeneous of order n if

$$f(t^n x, t^n y) = t^n f(x, y)$$

In this case, it can be seen that both $M = 2xy$ and $N = 3x^2 - y^2$ are both homogeneous and of the same order $n = 2$. Therefore this is a homogeneous ode. Since this ode is homogeneous, it is converted to separable ODE using the substitution $u = \frac{y}{x}$, or $y = ux$. Hence

$$\frac{dy}{dx} = \frac{du}{dx}x + u$$

Applying the transformation $y = ux$ to the above ODE in (1) gives

$$\begin{aligned} \frac{du}{dx}x + u &= -\frac{2u}{u^2 - 3} \\ \frac{du}{dx} &= \frac{-\frac{2u(x)}{u(x)^2 - 3} - u(x)}{x} \end{aligned}$$

Or

$$u'(x) - \frac{-\frac{2u(x)}{u(x)^2 - 3} - u(x)}{x} = 0$$

Or

$$u'(x)u(x)^2x + u(x)^3 - 3u'(x)x - u(x) = 0$$

Or

$$x(u(x)^2 - 3)u'(x) + u(x)^3 - u(x) = 0$$

Which is now solved as separable in $u(x)$.

The ode $u'(x) = -\frac{u(x)(u(x)^2 - 1)}{x(u(x)^2 - 3)}$ is separable as it can be written as

$$\begin{aligned} u'(x) &= -\frac{u(x)(u(x)^2 - 1)}{x(u(x)^2 - 3)} \\ &= f(x)g(u) \end{aligned}$$

Where

$$\begin{aligned} f(x) &= -\frac{1}{x} \\ g(u) &= \frac{u(u^2 - 1)}{u^2 - 3} \end{aligned}$$

Integrating gives

$$\begin{aligned} \int \frac{1}{g(u)} du &= \int f(x) dx \\ \int \frac{u^2 - 3}{u(u^2 - 1)} du &= \int -\frac{1}{x} dx \\ -\ln \left(\frac{(u(x) - 1)(u(x) + 1)}{u(x)^3} \right) &= \ln \left(\frac{1}{x} \right) + c_1 \end{aligned}$$

We now need to find the singular solutions, these are found by finding for what values $g(u)$ is zero, since we had to divide by this above. Solving $g(u) = 0$ or $\frac{u(u^2-1)}{u^2-3} = 0$ for $u(x)$ gives

$$u(x) = -1$$

$$u(x) = 0$$

$$u(x) = 1$$

Now we go over each such singular solution and check if it verifies the ode itself and any initial conditions given. If it does not then the singular solution will not be used.

Therefore the solutions found are

$$-\ln\left(\frac{(u(x)-1)(u(x)+1)}{u(x)^3}\right) = \ln\left(\frac{1}{x}\right) + c_1$$

$$u(x) = -1$$

$$u(x) = 0$$

$$u(x) = 1$$

Converting $-\ln\left(\frac{(u(x)-1)(u(x)+1)}{u(x)^3}\right) = \ln\left(\frac{1}{x}\right) + c_1$ back to y gives

$$-\ln\left(-\frac{x(-y+x)(y+x)}{y^3}\right) = \ln\left(\frac{1}{x}\right) + c_1$$

Converting $u(x) = -1$ back to y gives

$$y = -x$$

Converting $u(x) = 0$ back to y gives

$$y = 0$$

Converting $u(x) = 1$ back to y gives

$$y = x$$

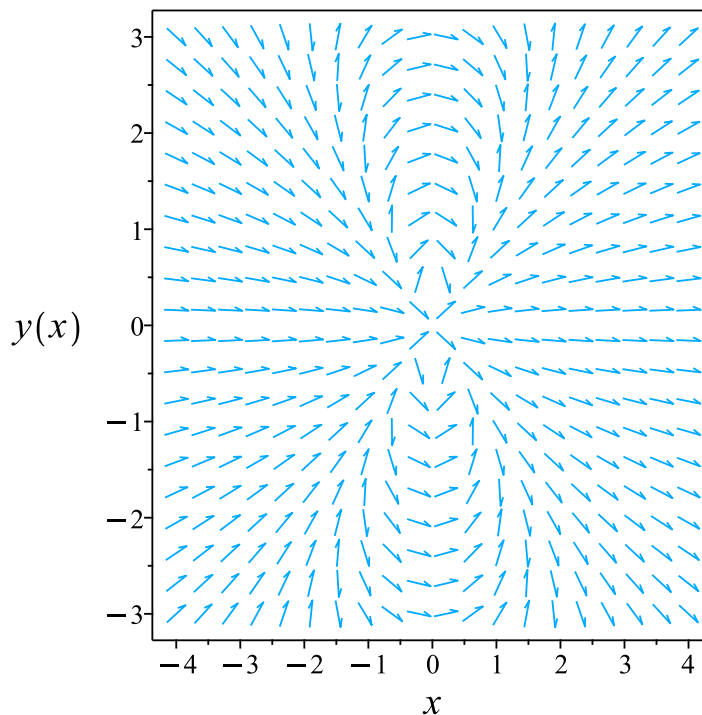


Figure 2.32: Slope field plot

$$\frac{2x}{y^3} + \left(\frac{1}{y^2} - \frac{3x^2}{y^4} \right) y' = 0$$

Summary of solutions found

$$-\ln \left(-\frac{x(-y+x)(y+x)}{y^3} \right) = \ln \left(\frac{1}{x} \right) + c_1$$

$$y = 0$$

$$y = x$$

$$y = -x$$

Solved as first order homogeneous class D2 ode

Time used: 0.259 (sec)

Applying change of variables $y = u(x)x$, then the ode becomes

$$\frac{2}{x^2 u(x)^3} + \left(\frac{1}{u(x)^2 x^2} - \frac{3}{x^2 u(x)^4} \right) (u'(x)x + u(x)) = 0$$

Which is now solved The ode $u'(x) = -\frac{u(x)(u(x)^2-1)}{x(u(x)^2-3)}$ is separable as it can be written as

$$\begin{aligned} u'(x) &= -\frac{u(x)(u(x)^2-1)}{x(u(x)^2-3)} \\ &= f(x)g(u) \end{aligned}$$

Where

$$\begin{aligned} f(x) &= -\frac{1}{x} \\ g(u) &= \frac{u(u^2-1)}{u^2-3} \end{aligned}$$

Integrating gives

$$\begin{aligned} \int \frac{1}{g(u)} du &= \int f(x) dx \\ \int \frac{u^2-3}{u(u^2-1)} du &= \int -\frac{1}{x} dx \\ -\ln\left(\frac{(u(x)-1)(u(x)+1)}{u(x)^3}\right) &= \ln\left(\frac{1}{x}\right) + c_1 \end{aligned}$$

We now need to find the singular solutions, these are found by finding for what values $g(u)$ is zero, since we had to divide by this above. Solving $g(u) = 0$ or $\frac{u(u^2-1)}{u^2-3} = 0$ for $u(x)$ gives

$$\begin{aligned} u(x) &= -1 \\ u(x) &= 0 \\ u(x) &= 1 \end{aligned}$$

Now we go over each such singular solution and check if it verifies the ode itself and any initial conditions given. If it does not then the singular solution will not be used.

Therefore the solutions found are

$$\begin{aligned} -\ln\left(\frac{(u(x)-1)(u(x)+1)}{u(x)^3}\right) &= \ln\left(\frac{1}{x}\right) + c_1 \\ u(x) &= -1 \\ u(x) &= 0 \\ u(x) &= 1 \end{aligned}$$

Converting $-\ln\left(\frac{(u(x)-1)(u(x)+1)}{u(x)^3}\right) = \ln\left(\frac{1}{x}\right) + c_1$ back to y gives

$$-\ln\left(\frac{\left(\frac{y}{x}-1\right)\left(\frac{y}{x}+1\right)x^3}{y^3}\right) = \ln\left(\frac{1}{x}\right) + c_1$$

Converting $u(x) = -1$ back to y gives

$$y = -x$$

Converting $u(x) = 0$ back to y gives

$$y = 0$$

Converting $u(x) = 1$ back to y gives

$$y = x$$

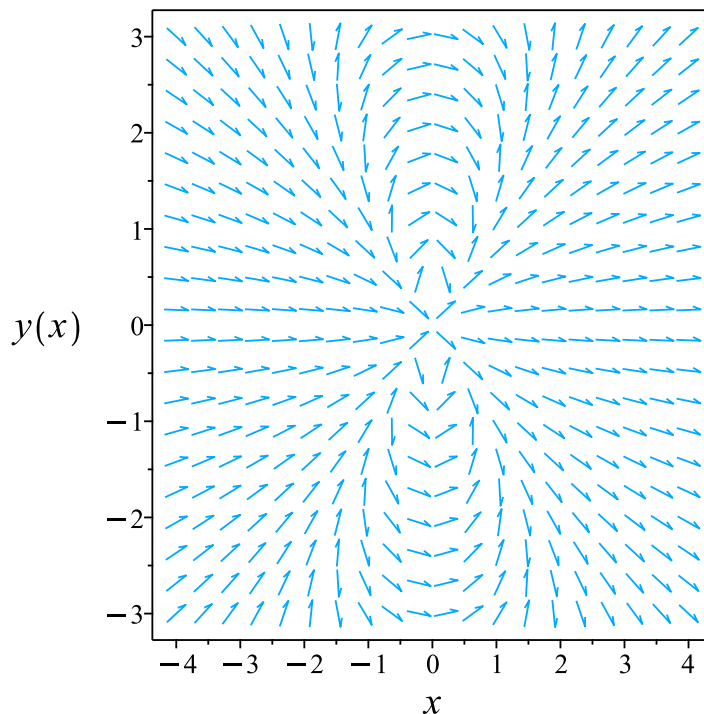


Figure 2.33: Slope field plot

$$\frac{2x}{y^3} + \left(\frac{1}{y^2} - \frac{3x^2}{y^4}\right) y' = 0$$

Summary of solutions found

$$-\ln\left(\frac{\left(\frac{y}{x}-1\right)\left(\frac{y}{x}+1\right)x^3}{y^3}\right) = \ln\left(\frac{1}{x}\right) + c_1$$

$$y = 0$$

$$y = x$$

$$y = -x$$

Summary of solutions found

$$-\ln\left(-\frac{x(-y+x)(y+x)}{y^3}\right) = \ln\left(\frac{1}{x}\right) + c_1$$

$$y = 0$$

$$y = x$$

$$y = -x$$

Solved as first order homogeneous class Maple C ode

Time used: 0.715 (sec)

Let $Y = y - y_0$ and $X = x - x_0$ then the above is transformed to new ode in $Y(X)$

$$\frac{d}{dX}Y(X) = -\frac{2(Y(X) + y_0)(x_0 + X)}{(Y(X) + y_0)^2 - 3(x_0 + X)^2}$$

Solving for possible values of x_0 and y_0 which makes the above ode a homogeneous ode results in

$$x_0 = 0$$

$$y_0 = 0$$

Using these values now it is possible to easily solve for $Y(X)$. The above ode now becomes

$$\frac{d}{dX}Y(X) = -\frac{2Y(X)X}{-3X^2 + Y(X)^2}$$

In canonical form, the ODE is

$$\begin{aligned} Y' &= F(X, Y) \\ &= -\frac{2YX}{-3X^2 + Y^2} \end{aligned} \tag{1}$$

An ode of the form $Y' = \frac{M(X,Y)}{N(X,Y)}$ is called homogeneous if the functions $M(X, Y)$ and $N(X, Y)$ are both homogeneous functions and of the same order. Recall that a function $f(X, Y)$ is homogeneous of order n if

$$f(t^n X, t^n Y) = t^n f(X, Y)$$

In this case, it can be seen that both $M = 2YX$ and $N = 3X^2 - Y^2$ are both homogeneous and of the same order $n = 2$. Therefore this is a homogeneous ode. Since this ode is homogeneous, it is converted to separable ODE using the substitution $u = \frac{Y}{X}$, or $Y = uX$. Hence

$$\frac{dY}{dX} = \frac{du}{dX}X + u$$

Applying the transformation $Y = uX$ to the above ODE in (1) gives

$$\begin{aligned} \frac{du}{dX}X + u &= -\frac{2u}{u^2 - 3} \\ \frac{du}{dX} &= \frac{-\frac{2u(X)}{u(X)^2 - 3} - u(X)}{X} \end{aligned}$$

Or

$$\frac{d}{dX}u(X) - \frac{-\frac{2u(X)}{u(X)^2 - 3} - u(X)}{X} = 0$$

Or

$$\left(\frac{d}{dX}u(X)\right)u(X)^2 X + u(X)^3 - 3\left(\frac{d}{dX}u(X)\right)X - u(X) = 0$$

Or

$$X(u(X)^2 - 3)\left(\frac{d}{dX}u(X)\right) + u(X)^3 - u(X) = 0$$

Which is now solved as separable in $u(X)$.

The ode $\frac{d}{dX}u(X) = -\frac{u(X)(u(X)^2 - 1)}{X(u(X)^2 - 3)}$ is separable as it can be written as

$$\begin{aligned} \frac{d}{dX}u(X) &= -\frac{u(X)(u(X)^2 - 1)}{X(u(X)^2 - 3)} \\ &= f(X)g(u) \end{aligned}$$

Where

$$\begin{aligned} f(X) &= -\frac{1}{X} \\ g(u) &= \frac{u(u^2 - 1)}{u^2 - 3} \end{aligned}$$

Integrating gives

$$\int \frac{1}{g(u)} du = \int f(X) dX$$

$$\int \frac{u^2 - 3}{u(u^2 - 1)} du = \int -\frac{1}{X} dX$$

$$-\ln \left(\frac{(u(X) - 1)(u(X) + 1)}{u(X)^3} \right) = \ln \left(\frac{1}{X} \right) + c_1$$

We now need to find the singular solutions, these are found by finding for what values $g(u)$ is zero, since we had to divide by this above. Solving $g(u) = 0$ or $\frac{u(u^2-1)}{u^2-3} = 0$ for $u(X)$ gives

$$u(X) = -1$$

$$u(X) = 0$$

$$u(X) = 1$$

Now we go over each such singular solution and check if it verifies the ode itself and any initial conditions given. If it does not then the singular solution will not be used.

Therefore the solutions found are

$$-\ln \left(\frac{(u(X) - 1)(u(X) + 1)}{u(X)^3} \right) = \ln \left(\frac{1}{X} \right) + c_1$$

$$u(X) = -1$$

$$u(X) = 0$$

$$u(X) = 1$$

Converting $-\ln \left(\frac{(u(X)-1)(u(X)+1)}{u(X)^3} \right) = \ln \left(\frac{1}{X} \right) + c_1$ back to $Y(X)$ gives

$$-\ln \left(-\frac{X(-Y(X) + X)(Y(X) + X)}{Y(X)^3} \right) = \ln \left(\frac{1}{X} \right) + c_1$$

Converting $u(X) = -1$ back to $Y(X)$ gives

$$Y(X) = -X$$

Converting $u(X) = 0$ back to $Y(X)$ gives

$$Y(X) = 0$$

Converting $u(X) = 1$ back to $Y(X)$ gives

$$Y(X) = X$$

Using the solution for $Y(X)$

$$-\ln\left(-\frac{X(-Y(X)+X)(Y(X)+X)}{Y(X)^3}\right) = \ln\left(\frac{1}{X}\right) + c_1 \quad (\text{A})$$

And replacing back terms in the above solution using

$$\begin{aligned} Y &= y + y_0 \\ X &= x + x_0 \end{aligned}$$

Or

$$\begin{aligned} Y &= y \\ X &= x \end{aligned}$$

Then the solution in y becomes using EQ (A)

$$-\ln\left(-\frac{x(-y+x)(y+x)}{y^3}\right) = \ln\left(\frac{1}{x}\right) + c_1$$

Using the solution for $Y(X)$

$$Y(X) = 0 \quad (\text{A})$$

And replacing back terms in the above solution using

$$\begin{aligned} Y &= y + y_0 \\ X &= x + x_0 \end{aligned}$$

Or

$$\begin{aligned} Y &= y \\ X &= x \end{aligned}$$

Then the solution in y becomes using EQ (A)

$$y = 0$$

Using the solution for $Y(X)$

$$Y(X) = X \quad (\text{A})$$

And replacing back terms in the above solution using

$$\begin{aligned} Y &= y + y_0 \\ X &= x + x_0 \end{aligned}$$

Or

$$Y = y$$

$$X = x$$

Then the solution in y becomes using EQ (A)

$$y = x$$

Using the solution for $Y(X)$

$$Y(X) = -X \tag{A}$$

And replacing back terms in the above solution using

$$Y = y + y_0$$

$$X = x + x_0$$

Or

$$Y = y$$

$$X = x$$

Then the solution in y becomes using EQ (A)

$$y = -x$$

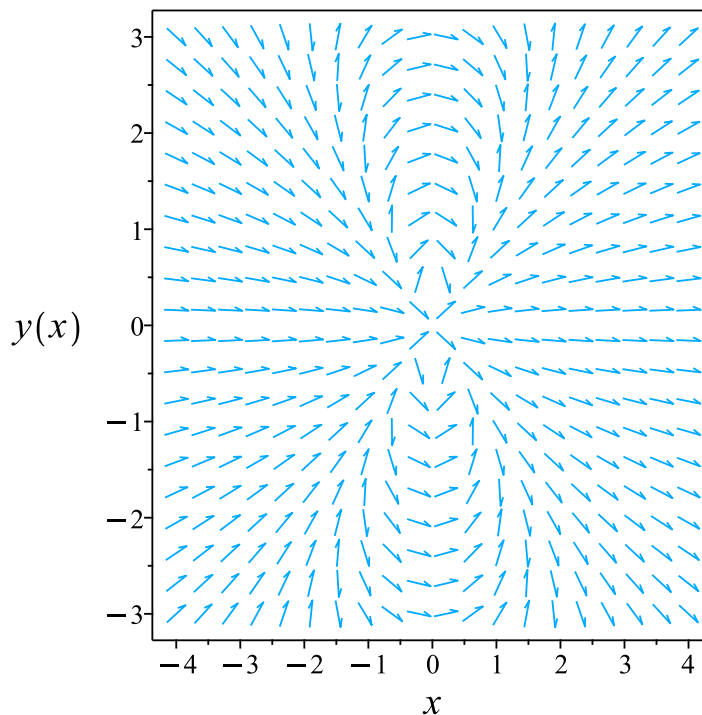


Figure 2.34: Slope field plot

$$\frac{2x}{y^3} + \left(\frac{1}{y^2} - \frac{3x^2}{y^4} \right) y' = 0$$

Solved as first order Exact ode

Time used: 0.115 (sec)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned} \left(\frac{1}{y^2} - \frac{3x^2}{y^4}\right) dy &= \left(-\frac{2x}{y^3}\right) dx \\ \left(\frac{2x}{y^3}\right) dx + \left(\frac{1}{y^2} - \frac{3x^2}{y^4}\right) dy &= 0 \end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= \frac{2x}{y^3} \\ N(x, y) &= \frac{1}{y^2} - \frac{3x^2}{y^4} \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} \left(\frac{2x}{y^3}\right) \\ &= -\frac{6x}{y^4} \end{aligned}$$

And

$$\begin{aligned} \frac{\partial N}{\partial x} &= \frac{\partial}{\partial x} \left(\frac{1}{y^2} - \frac{3x^2}{y^4}\right) \\ &= -\frac{6x}{y^4} \end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \quad (1)$$

$$\frac{\partial \phi}{\partial y} = N \quad (2)$$

Integrating (1) w.r.t. x gives

$$\int \frac{\partial \phi}{\partial x} dx = \int M dx$$

$$\int \frac{\partial \phi}{\partial x} dx = \int \frac{2x}{y^3} dx$$

$$\phi = \frac{x^2}{y^3} + f(y) \quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = -\frac{3x^2}{y^4} + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = \frac{1}{y^2} - \frac{3x^2}{y^4}$. Therefore equation (4) becomes

$$\frac{1}{y^2} - \frac{3x^2}{y^4} = -\frac{3x^2}{y^4} + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = \frac{1}{y^2}$$

Integrating the above w.r.t y gives

$$\int f'(y) dy = \int \left(\frac{1}{y^2} \right) dy$$

$$f(y) = -\frac{1}{y} + c_1$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = \frac{x^2}{y^3} - \frac{1}{y} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into the constant c_1 gives the solution as

$$c_1 = \frac{x^2}{y^3} - \frac{1}{y}$$

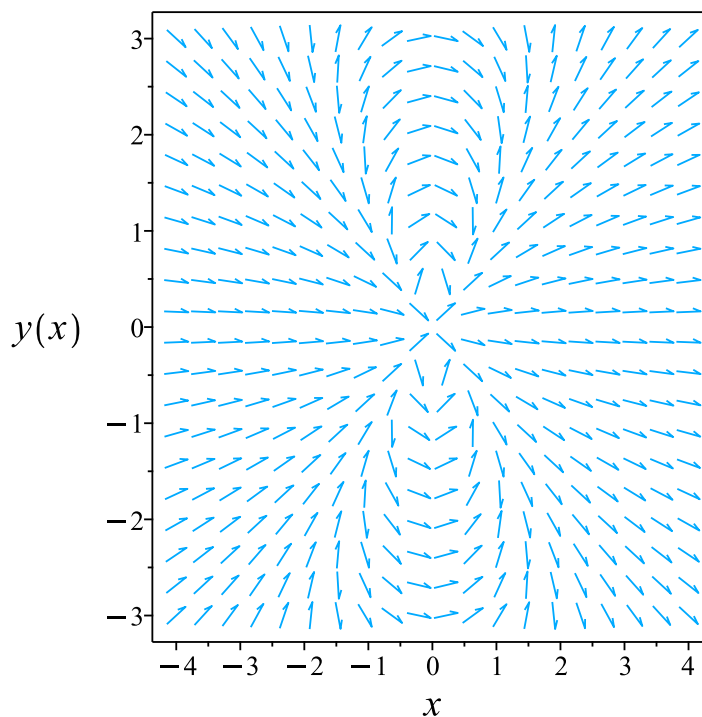


Figure 2.35: Slope field plot

$$\frac{2x}{y^3} + \left(\frac{1}{y^2} - \frac{3x^2}{y^4} \right) y' = 0$$

Summary of solutions found

$$\frac{x^2}{y^3} - \frac{1}{y} = c_1$$

Solved as first order isobaric ode

Time used: 0.209 (sec)

Solving for y' gives

$$y' = -\frac{2yx}{y^2 - 3x^2} \quad (1)$$

Each of the above ode's is now solved An ode $y' = f(x, y)$ is isobaric if

$$f(tx, t^m y) = t^{m-1} f(x, y) \quad (1)$$

Where here

$$f(x, y) = -\frac{2yx}{y^2 - 3x^2} \quad (2)$$

 m is the order of isobaric. Substituting (2) into (1) and solving for m gives

$$m = 1$$

Since the ode is isobaric of order $m = 1$, then the substitution

$$\begin{aligned} y &= ux^m \\ &= ux \end{aligned}$$

Converts the ODE to a separable in $u(x)$. Performing this substitution gives

$$u(x) + xu'(x) = -\frac{2x^2u(x)}{x^2u(x)^2 - 3x^2}$$

The ode $u'(x) = -\frac{u(x)(u(x)^2-1)}{(u(x)^2-3)x}$ is separable as it can be written as

$$\begin{aligned} u'(x) &= -\frac{u(x)(u(x)^2-1)}{(u(x)^2-3)x} \\ &= f(x)g(u) \end{aligned}$$

Where

$$\begin{aligned} f(x) &= -\frac{1}{x} \\ g(u) &= \frac{u(u^2-1)}{u^2-3} \end{aligned}$$

Integrating gives

$$\int \frac{1}{g(u)} du = \int f(x) dx$$

$$\int \frac{u^2 - 3}{u(u^2 - 1)} du = \int -\frac{1}{x} dx$$

$$-\ln \left(\frac{(u(x) - 1)(u(x) + 1)}{u(x)^3} \right) = \ln \left(\frac{1}{x} \right) + c_1$$

We now need to find the singular solutions, these are found by finding for what values $g(u)$ is zero, since we had to divide by this above. Solving $g(u) = 0$ or $\frac{u(u^2-1)}{u^2-3} = 0$ for $u(x)$ gives

$$u(x) = -1$$

$$u(x) = 0$$

$$u(x) = 1$$

Now we go over each such singular solution and check if it verifies the ode itself and any initial conditions given. If it does not then the singular solution will not be used.

Therefore the solutions found are

$$-\ln \left(\frac{(u(x) - 1)(u(x) + 1)}{u(x)^3} \right) = \ln \left(\frac{1}{x} \right) + c_1$$

$$u(x) = -1$$

$$u(x) = 0$$

$$u(x) = 1$$

Converting $-\ln \left(\frac{(u(x)-1)(u(x)+1)}{u(x)^3} \right) = \ln \left(\frac{1}{x} \right) + c_1$ back to y gives

$$-\ln \left(\frac{\left(\frac{y}{x} - 1\right) \left(\frac{y}{x} + 1\right) x^3}{y^3} \right) = \ln \left(\frac{1}{x} \right) + c_1$$

Converting $u(x) = -1$ back to y gives

$$\frac{y}{x} = -1$$

Converting $u(x) = 0$ back to y gives

$$\frac{y}{x} = 0$$

Converting $u(x) = 1$ back to y gives

$$\frac{y}{x} = 1$$

Solving for y gives

$$-\ln\left(\frac{\left(\frac{y}{x} - 1\right)\left(\frac{y}{x} + 1\right)x^3}{y^3}\right) = \ln\left(\frac{1}{x}\right) + c_1$$

$$y = 0$$

$$y = x$$

$$y = -x$$

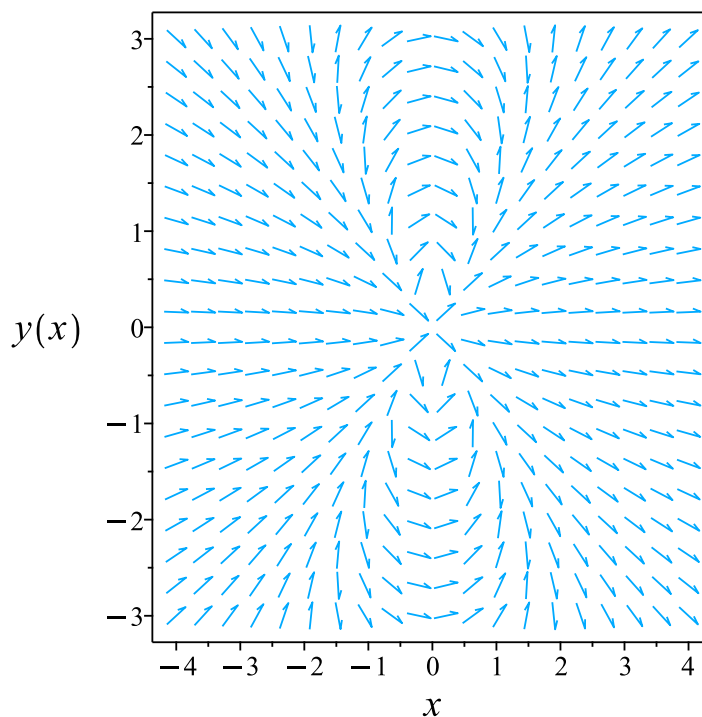


Figure 2.36: Slope field plot

$$\frac{2x}{y^3} + \left(\frac{1}{y^2} - \frac{3x^2}{y^4}\right)y' = 0$$

Summary of solutions found

$$-\ln\left(\frac{\left(\frac{y}{x} - 1\right)\left(\frac{y}{x} + 1\right)x^3}{y^3}\right) = \ln\left(\frac{1}{x}\right) + c_1$$

$$y = 0$$

$$y = x$$

$$y = -x$$

Solved using Lie symmetry for first order ode

Time used: 0.955 (sec)

Writing the ode as

$$y' = -\frac{2yx}{-3x^2 + y^2}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

To determine ξ, η then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$\xi = xa_2 + ya_3 + a_1 \quad (\text{1E})$$

$$\eta = xb_2 + yb_3 + b_1 \quad (\text{2E})$$

Where the unknown coefficients are

$$\{a_1, a_2, a_3, b_1, b_2, b_3\}$$

Substituting equations (1E,2E) and ω into (A) gives

$$b_2 - \frac{2yx(b_3 - a_2)}{-3x^2 + y^2} - \frac{4y^2x^2a_3}{(-3x^2 + y^2)^2}$$

$$- \left(-\frac{2y}{-3x^2 + y^2} - \frac{12yx^2}{(-3x^2 + y^2)^2} \right) (xa_2 + ya_3 + a_1) \quad (\text{5E})$$

$$- \left(-\frac{2x}{-3x^2 + y^2} + \frac{4y^2x}{(-3x^2 + y^2)^2} \right) (xb_2 + yb_3 + b_1) = 0$$

Putting the above in normal form gives

$$\frac{3x^4b_2 + 2y^2x^2a_3 - 8x^2y^2b_2 + 4xy^3a_2 - 4xy^3b_3 + 2y^4a_3 + y^4b_2 - 6x^3b_1 + 6x^2ya_1 - 2xy^2b_1 + 2y^3a_1}{(3x^2 - y^2)^2} = 0$$

Setting the numerator to zero gives

$$\begin{aligned} 3x^4b_2 + 2y^2x^2a_3 - 8x^2y^2b_2 + 4xy^3a_2 - 4xy^3b_3 + 2y^4a_3 \\ + y^4b_2 - 6x^3b_1 + 6x^2ya_1 - 2xy^2b_1 + 2y^3a_1 = 0 \end{aligned} \quad (6E)$$

Looking at the above PDE shows the following are all the terms with $\{x, y\}$ in them.

$$\{x, y\}$$

The following substitution is now made to be able to collect on all terms with $\{x, y\}$ in them

$$\{x = v_1, y = v_2\}$$

The above PDE (6E) now becomes

$$\begin{aligned} 4a_2v_1v_2^3 + 2a_3v_1^2v_2^2 + 2a_3v_2^4 + 3b_2v_1^4 - 8b_2v_1^2v_2^2 + b_2v_2^4 \\ - 4b_3v_1v_2^3 + 6a_1v_1^2v_2 + 2a_1v_2^3 - 6b_1v_1^3 - 2b_1v_1v_2^2 = 0 \end{aligned} \quad (7E)$$

Collecting the above on the terms v_i introduced, and these are

$$\{v_1, v_2\}$$

Equation (7E) now becomes

$$\begin{aligned} 3b_2v_1^4 - 6b_1v_1^3 + (2a_3 - 8b_2)v_1^2v_2^2 + 6a_1v_1^2v_2 \\ + (4a_2 - 4b_3)v_1v_2^3 - 2b_1v_1v_2^2 + (2a_3 + b_2)v_2^4 + 2a_1v_2^3 = 0 \end{aligned} \quad (8E)$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$2a_1 = 0$$

$$6a_1 = 0$$

$$-6b_1 = 0$$

$$-2b_1 = 0$$

$$3b_2 = 0$$

$$4a_2 - 4b_3 = 0$$

$$2a_3 - 8b_2 = 0$$

$$2a_3 + b_2 = 0$$

Solving the above equations for the unknowns gives

$$a_1 = 0$$

$$a_2 = b_3$$

$$a_3 = 0$$

$$b_1 = 0$$

$$b_2 = 0$$

$$b_3 = b_3$$

Substituting the above solution in the anstaz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$\xi = x$$

$$\eta = y$$

Shifting is now applied to make $\xi = 0$ in order to simplify the rest of the computation

$$\begin{aligned} \eta &= \eta - \omega(x, y) \xi \\ &= y - \left(-\frac{2yx}{-3x^2 + y^2} \right) (x) \\ &= \frac{x^2y - y^3}{3x^2 - y^2} \\ \xi &= 0 \end{aligned}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS \quad (1)$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y} \right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = x$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{\frac{x^2y - y^3}{3x^2 - y^2}} dy \end{aligned}$$

Which results in

$$S = -\ln(x + y) + 3\ln(y) - \ln(-x + y)$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = -\frac{2yx}{-3x^2 + y^2}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 1 \\ R_y &= 0 \\ S_x &= -\frac{2x}{x^2 - y^2} \\ S_y &= -\frac{1}{x + y} + \frac{3}{y} + \frac{1}{x - y} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = 0 \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = 0$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S .

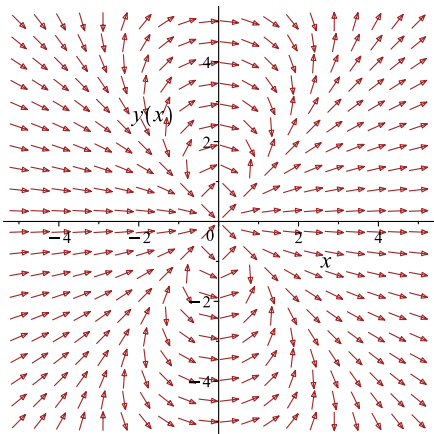
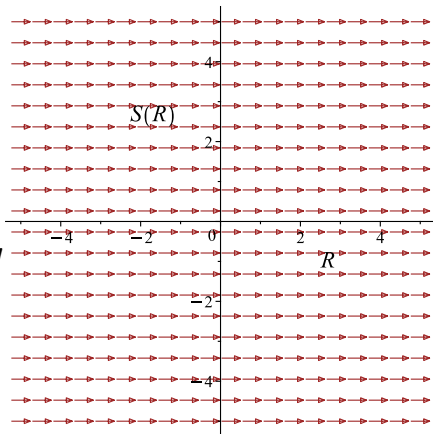
Since the ode has the form $\frac{d}{dR}S(R) = f(R)$, then we only need to integrate $f(R)$.

$$\begin{aligned} \int dS &= \int 0 dR + c_2 \\ S(R) &= c_2 \end{aligned}$$

To complete the solution, we just need to transform the above back to x, y coordinates. This results in

$$-\ln(y + x) + 3 \ln(y) - \ln(y - x) = c_2$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = -\frac{2yx}{-3x^2+y^2}$ 	$R = x$ $S = -\ln(x + y) + 3 \ln(y)$	$\frac{dS}{dR} = 0$ 

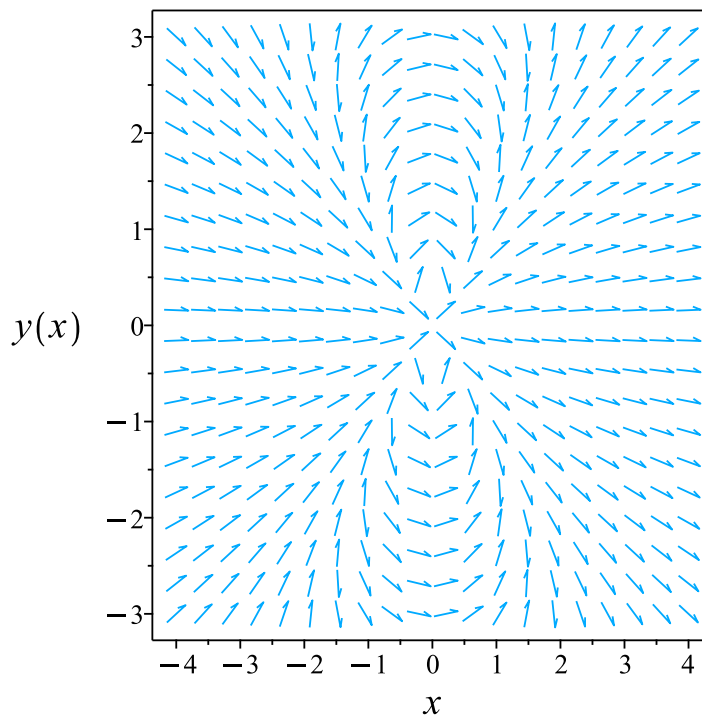


Figure 2.37: Slope field plot

$$\frac{2x}{y^3} + \left(\frac{1}{y^2} - \frac{3x^2}{y^4}\right) y' = 0$$

Summary of solutions found

$$-\ln(y+x) + 3\ln(y) - \ln(y-x) = c_2$$

Maple step by step solution

Let's solve

$$\frac{2x}{y^3} + \left(\frac{1}{y^2} - \frac{3x^2}{y^4}\right) y' = 0$$

- Highest derivative means the order of the ODE is 1
 y'
- Check if ODE is exact
 - ODE is exact if the lhs is the total derivative of a C^2 function
 $F'(x, y) = 0$
 - Compute derivative of lhs
 $F'(x, y) + \left(\frac{\partial}{\partial y} F(x, y)\right) y' = 0$
 - Evaluate derivatives

$$-\frac{6x}{y^4} = -\frac{6x}{y^4}$$

- Condition met, ODE is exact

- Exact ODE implies solution will be of this form

$$\left[F(x, y) = C1, M(x, y) = F'(x, y), N(x, y) = \frac{\partial}{\partial y} F(x, y) \right]$$

- Solve for $F(x, y)$ by integrating $M(x, y)$ with respect to x

$$F(x, y) = \int \frac{2x}{y^3} dx + _F1(y)$$

- Evaluate integral

$$F(x, y) = \frac{x^2}{y^3} + _F1(y)$$

- Take derivative of $F(x, y)$ with respect to y

$$N(x, y) = \frac{\partial}{\partial y} F(x, y)$$

- Compute derivative

$$\frac{1}{y^2} - \frac{3x^2}{y^4} = -\frac{3x^2}{y^4} + \frac{d}{dy} _F1(y)$$

- Isolate for $\frac{d}{dy} _F1(y)$

$$\frac{d}{dy} _F1(y) = \frac{1}{y^2}$$

- Solve for $_F1(y)$

$$_F1(y) = -\frac{1}{y}$$

- Substitute $_F1(y)$ into equation for $F(x, y)$

$$F(x, y) = \frac{x^2}{y^3} - \frac{1}{y}$$

- Substitute $F(x, y)$ into the solution of the ODE

$$\frac{x^2}{y^3} - \frac{1}{y} = C1$$

- Solve for y

$$\left\{ y = \frac{\left(12\sqrt{3}x\sqrt{27x^2C1^2-4}C1+108x^2C1^2-8\right)^{1/3}}{6C1} + \frac{2}{3C1\left(12\sqrt{3}x\sqrt{27x^2C1^2-4}C1+108x^2C1^2-8\right)^{1/3}} - \frac{1}{3C1}, y = -\left(\frac{1}{3C1}\right)^{3/2} \right.$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli

```

```

trying separable
trying inverse linear
trying homogeneous types:
trying homogeneous D
<- homogeneous successful`

```

Maple dsolve solution

Solving time : 0.011 (sec)

Leaf size : 313

```

dsolve(2*x/y(x)^3+(1/y(x))^2-3*x^2/y(x)^4)*diff(y(x),x) = 0,
y(x),singsol=all)

```

$$\begin{aligned}
 y &= \frac{1 + \frac{\left(12\sqrt{3}x\sqrt{27x^2c_1^2 - 4c_1 - 108x^2c_1^2 + 8}\right)^{1/3}}{2} + \frac{2}{\left(12\sqrt{3}x\sqrt{27x^2c_1^2 - 4c_1 - 108x^2c_1^2 + 8}\right)^{1/3}}}{3c_1} \\
 y &= \frac{(1 + i\sqrt{3}) \left(12\sqrt{3}x\sqrt{27x^2c_1^2 - 4c_1 - 108x^2c_1^2 + 8}\right)^{2/3} - 4i\sqrt{3} - 4 \left(12\sqrt{3}x\sqrt{27x^2c_1^2 - 4c_1 - 108x^2c_1^2 + 8}\right)^{1/3}}{12 \left(12\sqrt{3}x\sqrt{27x^2c_1^2 - 4c_1 - 108x^2c_1^2 + 8}\right)^{1/3} c_1} \\
 y &= \frac{(i\sqrt{3} - 1) \left(12\sqrt{3}x\sqrt{27x^2c_1^2 - 4c_1 - 108x^2c_1^2 + 8}\right)^{2/3} - 4i\sqrt{3} + 4 \left(12\sqrt{3}x\sqrt{27x^2c_1^2 - 4c_1 - 108x^2c_1^2 + 8}\right)^{1/3}}{12 \left(12\sqrt{3}x\sqrt{27x^2c_1^2 - 4c_1 - 108x^2c_1^2 + 8}\right)^{1/3} c_1}
 \end{aligned}$$

Mathematica DSolve solution

Solving time : 60.158 (sec)

Leaf size : 458

```
DSolve[{2*x/y[x]^3+(1/y[x]^2-3*x^2/y[x]^4)*D[y[x],x]==0,{}},
y[x],x,IncludeSingularSolutions->True]
```

$$y(x) \rightarrow \frac{1}{3} \left(\frac{\sqrt[3]{27e^{c_1}x^2 + 3\sqrt{81e^{2c_1}x^4 - 12e^{4c_1}x^2} - 2e^{3c_1}}}{\sqrt[3]{2}} \right. \\ \left. + \frac{\sqrt[3]{2}e^{2c_1}}{\sqrt[3]{27e^{c_1}x^2 + 3\sqrt{81e^{2c_1}x^4 - 12e^{4c_1}x^2} - 2e^{3c_1}}} - e^{c_1} \right)$$

$$y(x) \rightarrow \frac{i(\sqrt{3} + i) \sqrt[3]{27e^{c_1}x^2 + 3\sqrt{81e^{2c_1}x^4 - 12e^{4c_1}x^2} - 2e^{3c_1}}}{6\sqrt[3]{2}} \\ - \frac{i(\sqrt{3} - i) e^{2c_1}}{3 \cdot 2^{2/3} \sqrt[3]{27e^{c_1}x^2 + 3\sqrt{81e^{2c_1}x^4 - 12e^{4c_1}x^2} - 2e^{3c_1}}} - \frac{e^{c_1}}{3}$$

$$y(x) \rightarrow - \frac{i(\sqrt{3} - i) \sqrt[3]{27e^{c_1}x^2 + 3\sqrt{81e^{2c_1}x^4 - 12e^{4c_1}x^2} - 2e^{3c_1}}}{6\sqrt[3]{2}} \\ + \frac{i(\sqrt{3} + i) e^{2c_1}}{3 \cdot 2^{2/3} \sqrt[3]{27e^{c_1}x^2 + 3\sqrt{81e^{2c_1}x^4 - 12e^{4c_1}x^2} - 2e^{3c_1}}} - \frac{e^{c_1}}{3}$$

2.3.7 problem 8

Solved as first order Exact ode	264
Maple step by step solution	268
Maple trace	269
Maple dsolve solution	269
Mathematica DSolve solution	270

Internal problem ID [18227]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter IV. Methods of solution: First order equations. section 29. Problems at page 81

Problem number : 8

Date solved : Friday, December 20, 2024 at 10:46:21 AM

CAS classification : [_exact]

Solve

$$\left(T + \frac{1}{\sqrt{t^2 - T^2}}\right) T' = \frac{T}{t\sqrt{t^2 - T^2}} - t$$

Solved as first order Exact ode

Time used: 0.468 (sec)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \tag{B}$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(t, T) dt + N(t, T) dT = 0 \quad (1A)$$

Therefore

$$\begin{aligned} \left(T + \frac{1}{\sqrt{-T^2 + t^2}} \right) dT &= \left(\frac{T}{t\sqrt{-T^2 + t^2}} - t \right) dt \\ \left(-\frac{T}{t\sqrt{-T^2 + t^2}} + t \right) dt &+ \left(T + \frac{1}{\sqrt{-T^2 + t^2}} \right) dT = 0 \end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(t, T) &= -\frac{T}{t\sqrt{-T^2 + t^2}} + t \\ N(t, T) &= T + \frac{1}{\sqrt{-T^2 + t^2}} \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial T} = \frac{\partial N}{\partial t}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial T} &= \frac{\partial}{\partial T} \left(-\frac{T}{t\sqrt{-T^2 + t^2}} + t \right) \\ &= -\frac{t}{(-T^2 + t^2)^{3/2}} \end{aligned}$$

And

$$\begin{aligned} \frac{\partial N}{\partial t} &= \frac{\partial}{\partial t} \left(T + \frac{1}{\sqrt{-T^2 + t^2}} \right) \\ &= -\frac{t}{(-T^2 + t^2)^{3/2}} \end{aligned}$$

Since $\frac{\partial M}{\partial T} = \frac{\partial N}{\partial t}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(t, T)$

$$\frac{\partial \phi}{\partial t} = M \quad (1)$$

$$\frac{\partial \phi}{\partial T} = N \quad (2)$$

Integrating (1) w.r.t. t gives

$$\begin{aligned} \int \frac{\partial \phi}{\partial t} dt &= \int M dt \\ \int \frac{\partial \phi}{\partial t} dt &= \int -\frac{T}{t\sqrt{-T^2+t^2}} + t dt \\ \phi &= \frac{t^2\sqrt{-T^2} + 2T \ln\left(\frac{\sqrt{-T^2}\sqrt{-T^2+t^2}-T^2}{t}\right) + 2T \ln(2)}{2\sqrt{-T^2}} + f(T) \end{aligned} \quad (3)$$

Where $f(T)$ is used for the constant of integration since ϕ is a function of both t and T . Taking derivative of equation (3) w.r.t T gives

$$\begin{aligned} \frac{\partial \phi}{\partial T} &= \frac{-\frac{t^2 T}{\sqrt{-T^2}} + 2 \ln\left(\frac{\sqrt{-T^2}\sqrt{-T^2+t^2}-T^2}{t}\right) + \frac{2T\left(-\frac{\sqrt{-T^2+t^2}T}{\sqrt{-T^2}} - \frac{\sqrt{-T^2}T}{\sqrt{-T^2+t^2}} - 2T\right)}{\sqrt{-T^2}\sqrt{-T^2+t^2}-T^2} + 2 \ln(2)}{2\sqrt{-T^2}} \quad (4) \\ &+ \frac{\left(t^2\sqrt{-T^2} + 2T \ln\left(\frac{\sqrt{-T^2}\sqrt{-T^2+t^2}-T^2}{t}\right) + 2T \ln(2)\right) T}{2(-T^2)^{3/2}} + f'(T) \\ &= \frac{2\sqrt{-T^2}\sqrt{-T^2+t^2} - 2T^2 + t^2}{\sqrt{-T^2+t^2}(\sqrt{-T^2}\sqrt{-T^2+t^2}-T^2)} + f'(T) \end{aligned}$$

But equation (2) says that $\frac{\partial \phi}{\partial T} = T + \frac{1}{\sqrt{-T^2+t^2}}$. Therefore equation (4) becomes

$$T + \frac{1}{\sqrt{-T^2+t^2}} = \frac{2\sqrt{-T^2}\sqrt{-T^2+t^2} - 2T^2 + t^2}{\sqrt{-T^2+t^2}(\sqrt{-T^2}\sqrt{-T^2+t^2}-T^2)} + f'(T) \quad (5)$$

Solving equation (5) for $f'(T)$ gives

$$\begin{aligned} f'(T) &= -\frac{\sqrt{-T^2}T^3 - \sqrt{-T^2}Tt^2 + \sqrt{-T^2+t^2}T^3 + \sqrt{-T^2}\sqrt{-T^2+t^2} - T^2 + t^2}{\sqrt{-T^2+t^2}(\sqrt{-T^2}\sqrt{-T^2+t^2}-T^2)} \\ &= \frac{(T^3 + \sqrt{-T^2})\sqrt{-T^2+t^2} + (T-t)(T+t)(\sqrt{-T^2}T-1)}{\sqrt{-T^2+t^2}(T^2 - \sqrt{-T^2}\sqrt{-T^2+t^2})} \end{aligned}$$

Integrating the above w.r.t T results in

$$\int f'(T) dT = \int \left(\frac{(T^3 + \sqrt{-T^2}) \sqrt{-T^2 + t^2} + (T - t)(T + t)(\sqrt{-T^2} T - 1)}{\sqrt{-T^2 + t^2} (T^2 - \sqrt{-T^2} \sqrt{-T^2 + t^2})} \right) dT$$

$$f(T) = \frac{\sqrt{-T^2} \ln(T)}{T} + \frac{T^2}{2} + c_1$$

Where c_1 is constant of integration. Substituting result found above for $f(T)$ into equation (3) gives ϕ

$$\phi = \frac{t^2 \sqrt{-T^2} + 2T \ln \left(\frac{\sqrt{-T^2} \sqrt{-T^2 + t^2} - T^2}{t} \right) + 2T \ln(2)}{2\sqrt{-T^2}} + \frac{\sqrt{-T^2} \ln(T)}{T} + \frac{T^2}{2} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into the constant c_1 gives the solution as

$$c_1 = \frac{t^2 \sqrt{-T^2} + 2T \ln \left(\frac{\sqrt{-T^2} \sqrt{-T^2 + t^2} - T^2}{t} \right) + 2T \ln(2)}{2\sqrt{-T^2}} + \frac{\sqrt{-T^2} \ln(T)}{T} + \frac{T^2}{2}$$

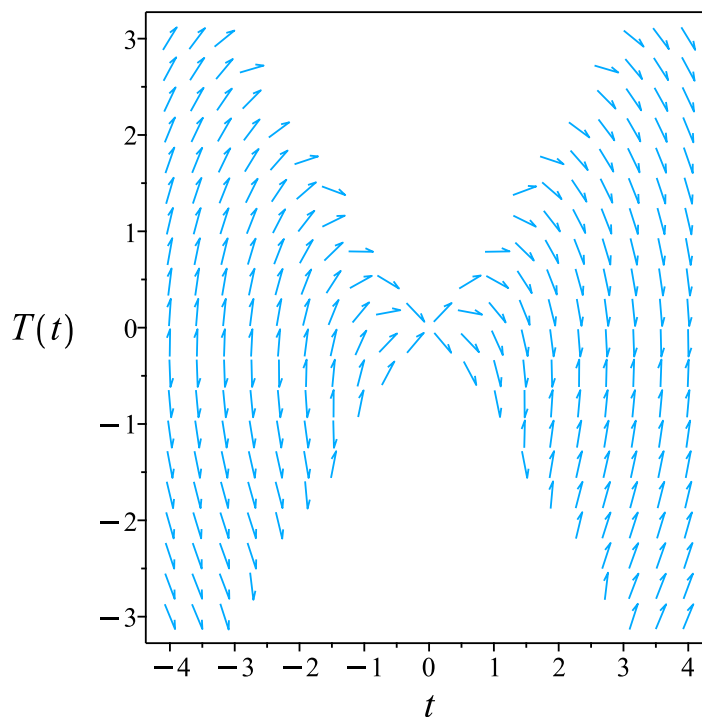


Figure 2.38: Slope field plot
 $\left(T + \frac{1}{\sqrt{t^2 - T^2}}\right) T' = \frac{T}{t\sqrt{t^2 - T^2}} - t$

Summary of solutions found

$$\frac{t^2\sqrt{-T^2} + 2T \ln\left(\frac{\sqrt{-T^2}\sqrt{t^2-T^2}-T^2}{t}\right) + 2T \ln(2)}{2\sqrt{-T^2}} + \frac{\sqrt{-T^2} \ln(T)}{T} + \frac{T^2}{2} = c_1$$

Maple step by step solution

Let's solve

$$\left(T + \frac{1}{\sqrt{t^2-T^2}}\right) T' = \frac{T}{t\sqrt{t^2-T^2}} - t$$

- Highest derivative means the order of the ODE is 1
- T'
- Check if ODE is exact
 - ODE is exact if the lhs is the total derivative of a C^2 function

$$F'(t, T) = 0$$

- Compute derivative of lhs

$$F'(t, T) + \left(\frac{\partial}{\partial T} F(t, T)\right) T' = 0$$

- Evaluate derivatives

$$-\frac{1}{t\sqrt{-T^2+t^2}} - \frac{T^2}{t(-T^2+t^2)^{3/2}} = -\frac{t}{(-T^2+t^2)^{3/2}}$$

- Simplify

$$-\frac{t}{(-T^2+t^2)^{3/2}} = -\frac{t}{(-T^2+t^2)^{3/2}}$$

- Condition met, ODE is exact

- Exact ODE implies solution will be of this form

$$[F(t, T) = C1, M(t, T) = F'(t, T), N(t, T) = \frac{\partial}{\partial T} F(t, T)]$$

- Solve for $F(t, T)$ by integrating $M(t, T)$ with respect to t

$$F(t, T) = \int \left(-\frac{T}{t\sqrt{-T^2+t^2}} + t\right) dt + _F1(T)$$

- Evaluate integral

$$F(t, T) = \frac{t^2}{2} + \frac{T \ln\left(\frac{-2T^2+2\sqrt{-T^2}\sqrt{-T^2+t^2}}{t}\right)}{\sqrt{-T^2}} + _F1(T)$$

- Take derivative of $F(t, T)$ with respect to T

$$N(t, T) = \frac{\partial}{\partial T} F(t, T)$$

- Compute derivative

$$T + \frac{1}{\sqrt{-T^2+t^2}} = \frac{\ln\left(\frac{-2T^2+2\sqrt{-T^2}\sqrt{-T^2+t^2}}{t}\right)}{\sqrt{-T^2}} + \frac{T^2 \ln\left(\frac{-2T^2+2\sqrt{-T^2}\sqrt{-T^2+t^2}}{t}\right)}{(-T^2)^{3/2}} + \frac{T\left(-4T - \frac{2\sqrt{-T^2+t^2}}{\sqrt{-T^2}} - \frac{2\sqrt{-T^2}T}{\sqrt{-T^2+t^2}}\right)}{\sqrt{-T^2}\left(-2T^2+2\sqrt{-T^2}\sqrt{-T^2+t^2}\right)}$$

- Isolate for $\frac{d}{dT}F1(T)$

$$\frac{d}{dT}F1(T) = T + \frac{1}{\sqrt{-T^2+t^2}} - \frac{\ln\left(\frac{-2T^2+2\sqrt{-T^2}\sqrt{-T^2+t^2}}{t}\right)}{\sqrt{-T^2}} - \frac{T^2 \ln\left(\frac{-2T^2+2\sqrt{-T^2}\sqrt{-T^2+t^2}}{t}\right)}{(-T^2)^{3/2}} - \frac{T\left(-4T - \frac{2\sqrt{-T^2+t^2}}{\sqrt{-T^2}}\right)}{\sqrt{-T^2}\left(-2T^2+2\sqrt{-T^2}\sqrt{-T^2+t^2}\right)}$$

- Solve for $F1(T)$

$$F1(T) = \frac{(-T^2)^{3/2}T^4 + T^6\sqrt{-T^2} - 2\sqrt{-T^2}T^4t^2 + 2\arctan\left(\frac{T}{\sqrt{-T^2+t^2}}\right)(-T^2)^{3/2}t^2 + 2\arctan\left(\frac{T}{\sqrt{-T^2+t^2}}\right)\sqrt{-T^2}T^2t^2 + 4\ln\left(\frac{-2T^2+2\sqrt{-T^2}\sqrt{-T^2+t^2}}{t}\right)}{4(-T^2)^{3/2}t^2}$$

- Substitute $F1(T)$ into equation for $F(t, T)$

$$F(t, T) = \frac{t^2}{2} + \frac{T \ln\left(\frac{-2T^2+2\sqrt{-T^2}\sqrt{-T^2+t^2}}{t}\right)}{\sqrt{-T^2}} + \frac{(-T^2)^{3/2}T^4 + T^6\sqrt{-T^2} - 2\sqrt{-T^2}T^4t^2 + 2\arctan\left(\frac{T}{\sqrt{-T^2+t^2}}\right)(-T^2)^{3/2}t^2 + 2\arctan\left(\frac{T}{\sqrt{-T^2+t^2}}\right)\sqrt{-T^2}T^2t^2 + 4\ln\left(\frac{-2T^2+2\sqrt{-T^2}\sqrt{-T^2+t^2}}{t}\right)}{4(-T^2)^{3/2}t^2}$$

- Substitute $F(t, T)$ into the solution of the ODE

$$\frac{t^2}{2} + \frac{T \ln\left(\frac{-2T^2+2\sqrt{-T^2}\sqrt{-T^2+t^2}}{t}\right)}{\sqrt{-T^2}} + \frac{(-T^2)^{3/2}T^4 + T^6\sqrt{-T^2} - 2\sqrt{-T^2}T^4t^2 + 2\arctan\left(\frac{T}{\sqrt{-T^2+t^2}}\right)(-T^2)^{3/2}t^2 + 2\arctan\left(\frac{T}{\sqrt{-T^2+t^2}}\right)\sqrt{-T^2}T^2t^2 + 4\ln\left(\frac{-2T^2+2\sqrt{-T^2}\sqrt{-T^2+t^2}}{t}\right)}{4(-T^2)^{3/2}t^2}$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying homogeneous types:
differential order: 1; looking for linear symmetries
trying exact
<- exact successful`

```

Maple dsolve solution

Solving time : 0.003 (sec)

Leaf size : 79

```

dsolve((T(t)+1/(t^2-T(t)^2)^(1/2))*diff(T(t),t) = T(t)/t/(t^2-T(t)^2)^(1/2)-t,
T(t),singsol=all)

```

$$\frac{\left(\frac{t^2}{2} + \frac{T^2}{2} + c_1\right) \sqrt{-T^2} + T \left(\ln \left(\frac{\sqrt{-T^2} \sqrt{t^2 - T^2} - T^2}{t} \right) + \ln(2) - \ln(T) \right)}{\sqrt{-T^2}} = 0$$

Mathematica DSolve solution

Solving time : 1.545 (sec)

Leaf size : 44

```
DSolve[{(T[t]+1/Sqrt[t^2-T[t]^2])*D[T[t],t]== T[t]/(t*Sqrt[t^2-T[t]^2])-t,{}},  
T[t],t,IncludeSingularSolutions->True]
```

$$\text{Solve} \left[-\arctan \left(\frac{\sqrt{t^2 - T(t)^2}}{T(t)} \right) + \frac{t^2}{2} + \frac{T(t)^2}{2} = c_1, T(t) \right]$$

2.4 Chapter IV. Methods of solution: First order equations. section 31. Problems at page 85

2.4.1	problem 1	272
2.4.2	problem 2	287
2.4.3	problem 3	295
2.4.4	problem 4	306
2.4.5	problem 5	320
2.4.6	problem 6	337
2.4.7	problem 7	346

2.4.1 problem 1

Solved as first order linear ode	272
Solved as first order separable ode	273
Solved as first order Exact ode	275
Solved using Lie symmetry for first order ode	280
Maple step by step solution	285
Maple trace	285
Maple dsolve solution	285
Mathematica DSolve solution	286

Internal problem ID [18228]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter IV. Methods of solution: First order equations. section 31. Problems at page 85

Problem number : 1

Date solved : Monday, December 23, 2024 at 09:16:29 PM

CAS classification : [_separable]

Solve

$$y' + xy = x$$

Solved as first order linear ode

Time used: 0.332 (sec)

In canonical form a linear first order is

$$y' + q(x)y = p(x)$$

Comparing the above to the given ode shows that

$$q(x) = x$$

$$p(x) = x$$

The integrating factor μ is

$$\mu = e^{\int x dx}$$

Therefore the solution is

$$y = \left(\int x e^{\int x dx} dx + c_1 \right) e^{-\int x dx}$$

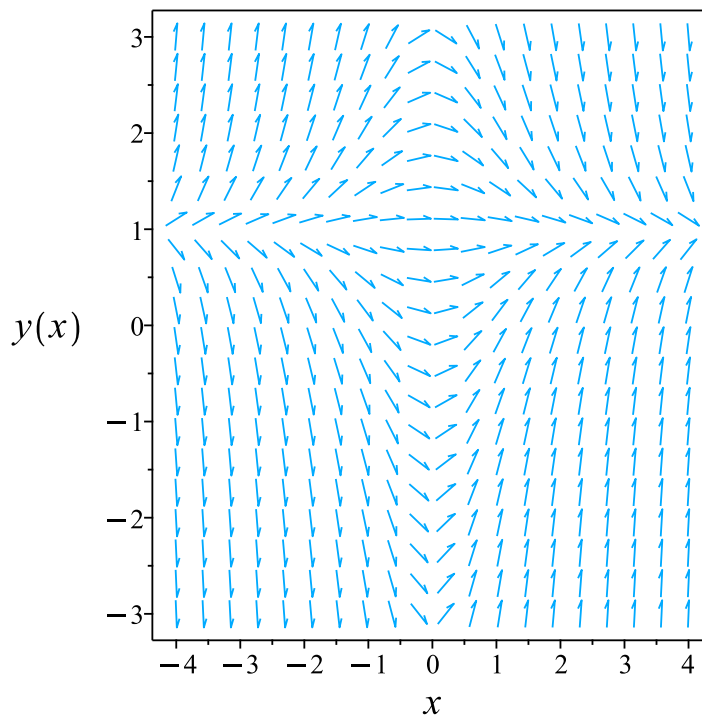


Figure 2.39: Slope field plot
 $y' + xy = x$

Summary of solutions found

$$y = \left(\int x e^{\int x dx} dx + c_1 \right) e^{-\int x dx}$$

Solved as first order separable ode

Time used: 0.176 (sec)

The ode $y' = -xy + x$ is separable as it can be written as

$$\begin{aligned} y' &= -xy + x \\ &= f(x)g(y) \end{aligned}$$

Where

$$\begin{aligned} f(x) &= x \\ g(y) &= -y + 1 \end{aligned}$$

Integrating gives

$$\begin{aligned}\int \frac{1}{g(y)} dy &= \int f(x) dx \\ \int \frac{1}{-y+1} dy &= \int x dx \\ -\ln(y-1) &= \frac{x^2}{2} + c_1\end{aligned}$$

We now need to find the singular solutions, these are found by finding for what values $g(y)$ is zero, since we had to divide by this above. Solving $g(y) = 0$ or $-y + 1 = 0$ for y gives

$$y = 1$$

Now we go over each such singular solution and check if it verifies the ode itself and any initial conditions given. If it does not then the singular solution will not be used.

Therefore the solutions found are

$$\begin{aligned}-\ln(y-1) &= \frac{x^2}{2} + c_1 \\ y &= 1\end{aligned}$$

Solving for y gives

$$\begin{aligned}y &= 1 \\ y &= e^{-\frac{x^2}{2} - c_1} + 1\end{aligned}$$

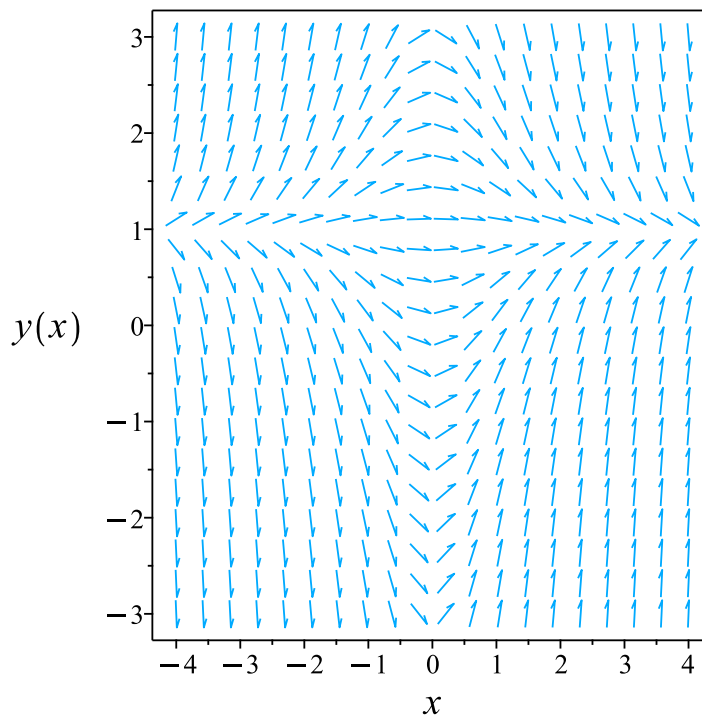


Figure 2.40: Slope field plot
 $y' + xy = x$

Summary of solutions found

$$y = 1$$

$$y = e^{-\frac{x^2}{2} - c_1} + 1$$

Solved as first order Exact ode

Time used: 0.226 (sec)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned}dy &= (-xy + x) dx \\ (xy - x) dx + dy &= 0\end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(x, y) &= xy - x \\ N(x, y) &= 1\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(xy - x) \\ &= x\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(1) \\ &= 0\end{aligned}$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned} A &= \frac{1}{N} \left(\frac{\partial M}{\partial y} - \frac{\partial N}{\partial x} \right) \\ &= 1((x) - (0)) \\ &= x \end{aligned}$$

Since A does not depend on y , then it can be used to find an integrating factor. The integrating factor μ is

$$\begin{aligned} \mu &= e^{\int A \, dx} \\ &= e^{\int x \, dx} \end{aligned}$$

The result of integrating gives

$$\begin{aligned} \mu &= e^{\frac{x^2}{2}} \\ &= e^{\frac{x^2}{2}} \end{aligned}$$

M and N are multiplied by this integrating factor, giving new M and new N which are called \bar{M} and \bar{N} for now so not to confuse them with the original M and N .

$$\begin{aligned} \bar{M} &= \mu M \\ &= e^{\frac{x^2}{2}} (xy - x) \\ &= x(y - 1) e^{\frac{x^2}{2}} \end{aligned}$$

And

$$\begin{aligned} \bar{N} &= \mu N \\ &= e^{\frac{x^2}{2}} (1) \\ &= e^{\frac{x^2}{2}} \end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned} \bar{M} + \bar{N} \frac{dy}{dx} &= 0 \\ \left(x(y - 1) e^{\frac{x^2}{2}} \right) + \left(e^{\frac{x^2}{2}} \right) \frac{dy}{dx} &= 0 \end{aligned}$$

The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = \overline{M} \quad (1)$$

$$\frac{\partial \phi}{\partial y} = \overline{N} \quad (2)$$

Integrating (2) w.r.t. y gives

$$\int \frac{\partial \phi}{\partial y} dy = \int \overline{N} dy$$

$$\int \frac{\partial \phi}{\partial y} dy = \int e^{\frac{x^2}{2}} dy$$

$$\phi = e^{\frac{x^2}{2}} y + f(x) \quad (3)$$

Where $f(x)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t x gives

$$\frac{\partial \phi}{\partial x} = x e^{\frac{x^2}{2}} y + f'(x) \quad (4)$$

But equation (1) says that $\frac{\partial \phi}{\partial x} = x(y - 1) e^{\frac{x^2}{2}}$. Therefore equation (4) becomes

$$x(y - 1) e^{\frac{x^2}{2}} = x e^{\frac{x^2}{2}} y + f'(x) \quad (5)$$

Solving equation (5) for $f'(x)$ gives

$$f'(x) = -x e^{\frac{x^2}{2}}$$

Integrating the above w.r.t x gives

$$\int f'(x) dx = \int \left(-x e^{\frac{x^2}{2}} \right) dx$$

$$f(x) = -e^{\frac{x^2}{2}} + c_1$$

Where c_1 is constant of integration. Substituting result found above for $f(x)$ into equation (3) gives ϕ

$$\phi = e^{\frac{x^2}{2}} y - e^{\frac{x^2}{2}} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into the constant c_1 gives the solution as

$$c_1 = e^{\frac{x^2}{2}} y - e^{\frac{x^2}{2}}$$

Solving for y gives

$$y = \left(e^{\frac{x^2}{2}} + c_1 \right) e^{-\frac{x^2}{2}}$$

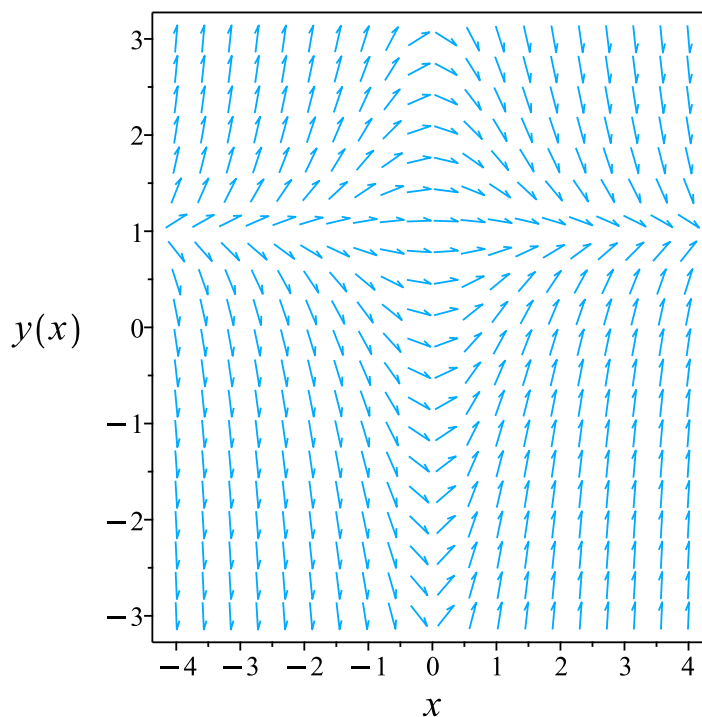


Figure 2.41: Slope field plot
 $y' + xy = x$

Summary of solutions found

$$y = \left(e^{\frac{x^2}{2}} + c_1 \right) e^{-\frac{x^2}{2}}$$

Solved using Lie symmetry for first order ode

Time used: 0.501 (sec)

Writing the ode as

$$\begin{aligned}y' &= -xy + x \\y' &= \omega(x, y)\end{aligned}$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

To determine ξ, η then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$\xi = xa_2 + ya_3 + a_1 \quad (1\text{E})$$

$$\eta = xb_2 + yb_3 + b_1 \quad (2\text{E})$$

Where the unknown coefficients are

$$\{a_1, a_2, a_3, b_1, b_2, b_3\}$$

Substituting equations (1E,2E) and ω into (A) gives

$$\begin{aligned}b_2 + (-xy + x)(b_3 - a_2) - (-xy + x)^2 a_3 \\ - (-y + 1)(xa_2 + ya_3 + a_1) + x(xb_2 + yb_3 + b_1) = 0\end{aligned} \quad (5\text{E})$$

Putting the above in normal form gives

$$\begin{aligned}-x^2y^2a_3 + 2x^2ya_3 - x^2a_3 + x^2b_2 + 2xya_2 + y^2a_3 \\ - 2xa_2 + xb_1 + xb_3 + ya_1 - ya_3 - a_1 + b_2 = 0\end{aligned}$$

Setting the numerator to zero gives

$$\begin{aligned}-x^2y^2a_3 + 2x^2ya_3 - x^2a_3 + x^2b_2 + 2xya_2 + y^2a_3 \\ - 2xa_2 + xb_1 + xb_3 + ya_1 - ya_3 - a_1 + b_2 = 0\end{aligned} \quad (6\text{E})$$

Looking at the above PDE shows the following are all the terms with $\{x, y\}$ in them.

$$\{x, y\}$$

The following substitution is now made to be able to collect on all terms with $\{x, y\}$ in them

$$\{x = v_1, y = v_2\}$$

The above PDE (6E) now becomes

$$\begin{aligned} -a_3v_1^2v_2^2 + 2a_3v_1^2v_2 + 2a_2v_1v_2 - a_3v_1^2 + a_3v_2^2 + b_2v_1^2 \\ + a_1v_2 - 2a_2v_1 - a_3v_2 + b_1v_1 + b_3v_1 - a_1 + b_2 = 0 \end{aligned} \quad (7E)$$

Collecting the above on the terms v_i introduced, and these are

$$\{v_1, v_2\}$$

Equation (7E) now becomes

$$\begin{aligned} -a_3v_1^2v_2^2 + 2a_3v_1^2v_2 + (-a_3 + b_2)v_1^2 + 2a_2v_1v_2 \\ + (-2a_2 + b_1 + b_3)v_1 + a_3v_2^2 + (a_1 - a_3)v_2 - a_1 + b_2 = 0 \end{aligned} \quad (8E)$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$\begin{aligned} a_3 &= 0 \\ 2a_2 &= 0 \\ -a_3 &= 0 \\ 2a_3 &= 0 \\ -a_1 + b_2 &= 0 \\ a_1 - a_3 &= 0 \\ -a_3 + b_2 &= 0 \\ -2a_2 + b_1 + b_3 &= 0 \end{aligned}$$

Solving the above equations for the unknowns gives

$$\begin{aligned} a_1 &= 0 \\ a_2 &= 0 \\ a_3 &= 0 \\ b_1 &= -b_3 \\ b_2 &= 0 \\ b_3 &= b_3 \end{aligned}$$

Substituting the above solution in the anstaz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$\begin{aligned}\xi &= 0 \\ \eta &= y - 1\end{aligned}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS \quad (1)$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = x$$

S is found from

$$\begin{aligned}S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{y-1} dy\end{aligned}$$

Which results in

$$S = \ln(y - 1)$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = -xy + x$$

Evaluating all the partial derivatives gives

$$\begin{aligned}R_x &= 1 \\R_y &= 0 \\S_x &= 0 \\S_y &= \frac{1}{y-1}\end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = -x \tag{2A}$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = -R$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S .

Since the ode has the form $\frac{d}{dR}S(R) = f(R)$, then we only need to integrate $f(R)$.

$$\begin{aligned}\int dS &= \int -R dR \\S(R) &= -\frac{R^2}{2} + c_2\end{aligned}$$

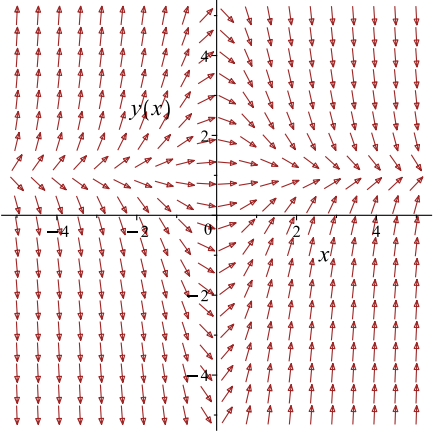
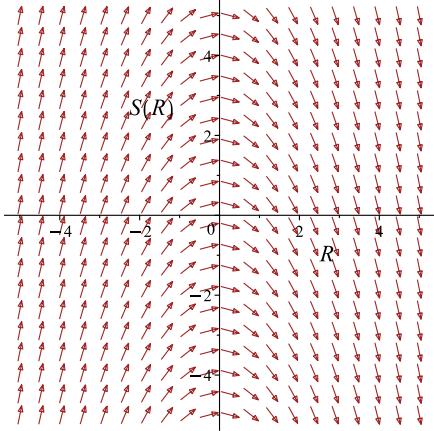
To complete the solution, we just need to transform the above back to x, y coordinates. This results in

$$\ln(y-1) = -\frac{x^2}{2} + c_2$$

Which gives

$$y = e^{-\frac{x^2}{2} + c_2} + 1$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = -xy + x$ 	$R = x$ $S = \ln(y - 1)$	$\frac{dS}{dR} = -R$ 

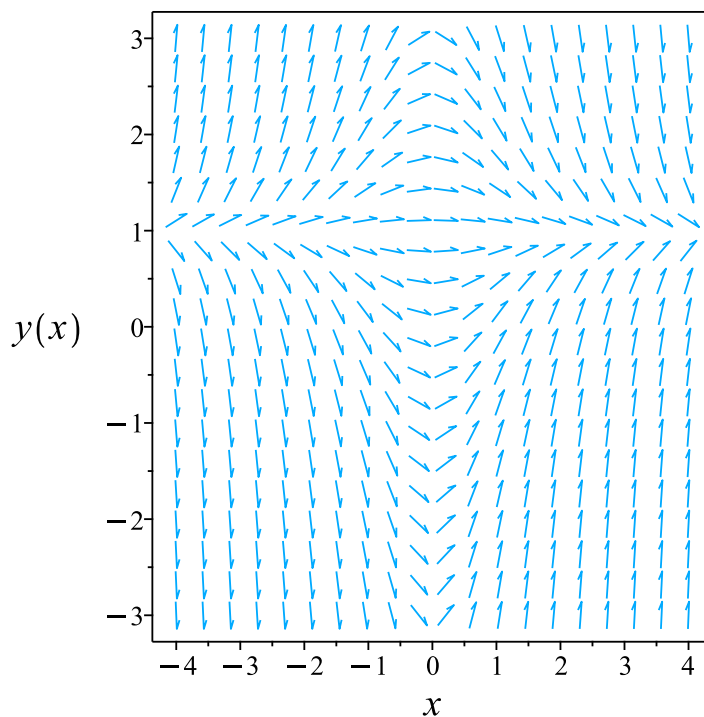


Figure 2.42: Slope field plot
 $y' + xy = x$

Summary of solutions found

$$y = e^{-\frac{x^2}{2} + c_2} + 1$$

Maple step by step solution

Let's solve

$$y' + xy = x$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Solve for the highest derivative

$$y' = -xy + x$$

- Separate variables

$$\frac{y'}{y-1} = -x$$

- Integrate both sides with respect to x

$$\int \frac{y'}{y-1} dx = \int -x dx + C1$$

- Evaluate integral

$$\ln(y-1) = -\frac{x^2}{2} + C1$$

- Solve for y

$$y = e^{-\frac{x^2}{2} + C1} + 1$$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`
```

Maple dsolve solution

Solving time : 0.002 (sec)

Leaf size : 14

```
dsolve(diff(y(x),x)+x*y(x) = x,
        y(x),singsol=all)
```

$$y = 1 + e^{-\frac{x^2}{2}} c_1$$

Mathematica DSolve solution

Solving time : 0.324 (sec)

Leaf size : 24

```
DSolve[{D[y[x],x]+x*y[x]==x,{}},  
y[x],x,IncludeSingularSolutions->True]
```

$$y(x) \rightarrow 1 + c_1 e^{-\frac{x^2}{2}}$$

$$y(x) \rightarrow 1$$

2.4.2 problem 2

Solved as first order linear ode	287
Solved as first order Exact ode	289
Maple step by step solution	292
Maple trace	293
Maple dsolve solution	293
Mathematica DSolve solution	294

Internal problem ID [18229]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter IV. Methods of solution: First order equations. section 31. Problems at page 85

Problem number : 2

Date solved : Monday, December 23, 2024 at 09:16:31 PM

CAS classification : [_linear]

Solve

$$y' + \frac{y}{x} = \sin(x)$$

Solved as first order linear ode

Time used: 0.119 (sec)

In canonical form a linear first order is

$$y' + q(x)y = p(x)$$

Comparing the above to the given ode shows that

$$q(x) = \frac{1}{x}$$

$$p(x) = \sin(x)$$

The integrating factor μ is

$$\begin{aligned} \mu &= e^{\int q dx} \\ &= e^{\int \frac{1}{x} dx} \\ &= x \end{aligned}$$

The ode becomes

$$\frac{d}{dx}(\mu y) = \mu p$$

$$\frac{d}{dx}(\mu y) = (\mu) (\sin(x))$$

$$\frac{d}{dx}(xy) = (x) (\sin(x))$$

$$d(xy) = (\sin(x) x) dx$$

Integrating gives

$$\begin{aligned} xy &= \int \sin(x) x dx \\ &= \sin(x) - \cos(x) x + c_1 \end{aligned}$$

Dividing throughout by the integrating factor x gives the final solution

$$y = \frac{\sin(x) - \cos(x) x + c_1}{x}$$

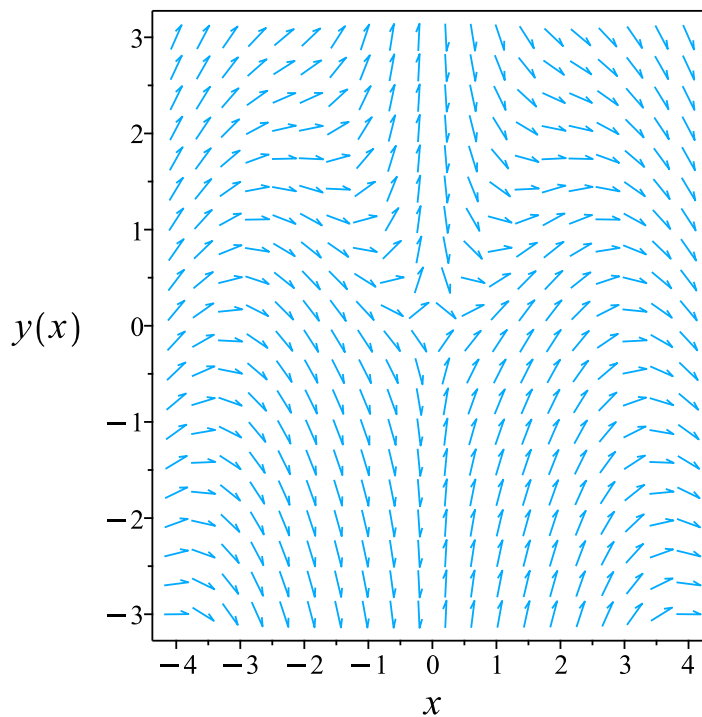


Figure 2.43: Slope field plot
 $y' + \frac{y}{x} = \sin(x)$

Summary of solutions found

$$y = \frac{\sin(x) - \cos(x)x + c_1}{x}$$

Solved as first order Exact ode

Time used: 0.190 (sec)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx}\phi(x, y) = 0$$

Hence

$$\frac{\partial\phi}{\partial x} + \frac{\partial\phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial\phi}{\partial x} &= M \\ \frac{\partial\phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (\text{1A})$$

Therefore

$$\begin{aligned} (x) dy &= (\sin(x)x - y) dx \\ (-\sin(x)x + y) dx &+ (x) dy = 0 \end{aligned} \quad (\text{2A})$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(x, y) &= -\sin(x)x + y \\N(x, y) &= x\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(-\sin(x)x + y) \\ &= 1\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(x) \\ &= 1\end{aligned}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \tag{1}$$

$$\frac{\partial \phi}{\partial y} = N \tag{2}$$

Integrating (2) w.r.t. y gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial y} dy &= \int N dy \\ \int \frac{\partial \phi}{\partial y} dy &= \int x dy \\ \phi &= xy + f(x)\end{aligned} \tag{3}$$

Where $f(x)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t x gives

$$\frac{\partial \phi}{\partial x} = y + f'(x) \tag{4}$$

But equation (1) says that $\frac{\partial \phi}{\partial x} = -\sin(x)x + y$. Therefore equation (4) becomes

$$-\sin(x)x + y = y + f'(x) \quad (5)$$

Solving equation (5) for $f'(x)$ gives

$$f'(x) = -\sin(x)x$$

Integrating the above w.r.t x gives

$$\begin{aligned} \int f'(x) dx &= \int (-\sin(x)x) dx \\ f(x) &= \cos(x)x - \sin(x) + c_1 \end{aligned}$$

Where c_1 is constant of integration. Substituting result found above for $f(x)$ into equation (3) gives ϕ

$$\phi = xy + \cos(x)x - \sin(x) + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into the constant c_1 gives the solution as

$$c_1 = xy + \cos(x)x - \sin(x)$$

Solving for y gives

$$y = -\frac{\cos(x)x - \sin(x) - c_1}{x}$$

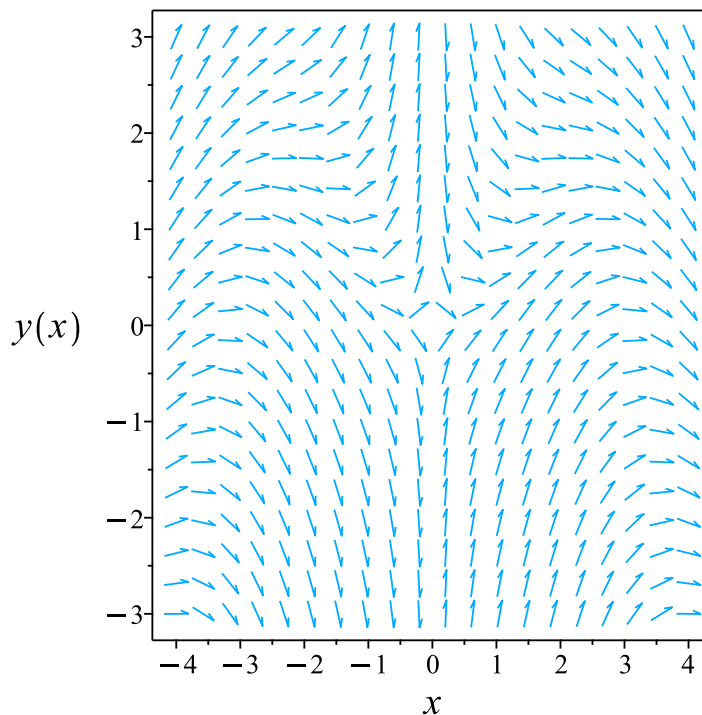


Figure 2.44: Slope field plot
 $y' + \frac{y}{x} = \sin(x)$

Summary of solutions found

$$y = -\frac{\cos(x)x - \sin(x) - c_1}{x}$$

Maple step by step solution

Let's solve

$$y' + \frac{y}{x} = \sin(x)$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Solve for the highest derivative

$$y' = -\frac{y}{x} + \sin(x)$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE

$$y' + \frac{y}{x} = \sin(x)$$

- The ODE is linear; multiply by an integrating factor $\mu(x)$

$$\mu(x) \left(y' + \frac{y}{x} \right) = \mu(x) \sin(x)$$

- Assume the lhs of the ODE is the total derivative $\frac{d}{dx}(y\mu(x))$

$$\mu(x) \left(y' + \frac{y}{x} \right) = y' \mu(x) + y \mu'(x)$$
- Isolate $\mu'(x)$

$$\mu'(x) = \frac{\mu(x)}{x}$$
- Solve to find the integrating factor

$$\mu(x) = x$$
- Integrate both sides with respect to x

$$\int \left(\frac{d}{dx}(y\mu(x)) \right) dx = \int \mu(x) \sin(x) dx + C1$$
- Evaluate the integral on the lhs

$$y\mu(x) = \int \mu(x) \sin(x) dx + C1$$
- Solve for y

$$y = \frac{\int \mu(x) \sin(x) dx + C1}{\mu(x)}$$
- Substitute $\mu(x) = x$

$$y = \frac{\int \sin(x)x dx + C1}{x}$$
- Evaluate the integrals on the rhs

$$y = \frac{\sin(x) - \cos(x)x + C1}{x}$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`

```

Maple dsolve solution

Solving time : 0.002 (sec)

Leaf size : 17

```

dsolve(diff(y(x),x)+y(x)/x = sin(x),
        y(x),singsol=all)

```

$$y = \frac{\sin(x) - \cos(x)x + c_1}{x}$$

Mathematica DSolve solution

Solving time : 0.068 (sec)

Leaf size : 19

```
DSolve[{D[y[x],x]+y[x]/x==Sin[x],{}},  
y[x],x,IncludeSingularSolutions->True]
```

$$y(x) \rightarrow \frac{\sin(x) - x \cos(x) + c_1}{x}$$

2.4.3 problem 3

Solved as first order Bernoulli ode	295
Solved as first order Exact ode	299
Maple step by step solution	304
Maple trace	304
Maple dsolve solution	304
Mathematica DSolve solution	305

Internal problem ID [18230]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter IV. Methods of solution: First order equations. section 31. Problems at page 85

Problem number : 3

Date solved : Monday, December 23, 2024 at 09:16:32 PM

CAS classification : [_Bernoulli]

Solve

$$y' + \frac{y}{x} = \frac{\sin(x)}{y^3}$$

Solved as first order Bernoulli ode

Time used: 0.516 (sec)

In canonical form, the ODE is

$$\begin{aligned} y' &= F(x, y) \\ &= \frac{-y^4 + \sin(x)x}{xy^3} \end{aligned}$$

This is a Bernoulli ODE.

$$y' = \left(-\frac{1}{x}\right)y + (\sin(x))\frac{1}{y^3} \quad (1)$$

The standard Bernoulli ODE has the form

$$y' = f_0(x)y + f_1(x)y^n \quad (2)$$

Comparing this to (1) shows that

$$\begin{aligned} f_0 &= -\frac{1}{x} \\ f_1 &= \sin(x) \end{aligned}$$

The first step is to divide the above equation by y^n which gives

$$\frac{y'}{y^n} = f_0(x)y^{1-n} + f_1(x) \quad (3)$$

The next step is use the substitution $v = y^{1-n}$ in equation (3) which generates a new ODE in $v(x)$ which will be linear and can be easily solved using an integrating factor. Backsubstitution then gives the solution $y(x)$ which is what we want.

This method is now applied to the ODE at hand. Comparing the ODE (1) With (2) Shows that

$$\begin{aligned} f_0(x) &= -\frac{1}{x} \\ f_1(x) &= \sin(x) \\ n &= -3 \end{aligned}$$

Dividing both sides of ODE (1) by $y^n = \frac{1}{y^3}$ gives

$$y'y^3 = -\frac{y^4}{x} + \sin(x) \quad (4)$$

Let

$$\begin{aligned} v &= y^{1-n} \\ &= y^4 \end{aligned} \quad (5)$$

Taking derivative of equation (5) w.r.t x gives

$$v' = 4y^3y' \quad (6)$$

Substituting equations (5) and (6) into equation (4) gives

$$\begin{aligned} \frac{v'(x)}{4} &= -\frac{v(x)}{x} + \sin(x) \\ v' &= -\frac{4v}{x} + 4\sin(x) \end{aligned} \quad (7)$$

The above now is a linear ODE in $v(x)$ which is now solved.

In canonical form a linear first order is

$$v'(x) + q(x)v(x) = p(x)$$

Comparing the above to the given ode shows that

$$\begin{aligned} q(x) &= \frac{4}{x} \\ p(x) &= 4\sin(x) \end{aligned}$$

The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int q dx} \\ &= e^{\int \frac{4}{x} dx} \\ &= x^4\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dx}(\mu v) &= \mu p \\ \frac{d}{dx}(\mu v) &= (\mu) (4 \sin(x)) \\ \frac{d}{dx}(v x^4) &= (x^4) (4 \sin(x)) \\ d(v x^4) &= (4 \sin(x) x^4) dx\end{aligned}$$

Integrating gives

$$\begin{aligned}v x^4 &= \int 4 \sin(x) x^4 dx \\ &= -4 \cos(x) x^4 + 16 \sin(x) x^3 + 48 \cos(x) x^2 - 96 \cos(x) - 96 \sin(x) x + c_1\end{aligned}$$

Dividing throughout by the integrating factor x^4 gives the final solution

$$v(x) = \frac{4(-x^4 + 12x^2 - 24) \cos(x) + 16(x^3 - 6x) \sin(x) + c_1}{x^4}$$

The substitution $v = y^{1-n}$ is now used to convert the above solution back to y which results in

$$y^4 = \frac{4(-x^4 + 12x^2 - 24) \cos(x) + 16(x^3 - 6x) \sin(x) + c_1}{x^4}$$

Solving for y gives

$$\begin{aligned}y &= \frac{(-4 \cos(x) x^4 + 16 \sin(x) x^3 + 48 \cos(x) x^2 - 96 \cos(x) - 96 \sin(x) x + c_1)^{1/4}}{x} \\ y &= -\frac{i(-4 \cos(x) x^4 + 16 \sin(x) x^3 + 48 \cos(x) x^2 - 96 \cos(x) - 96 \sin(x) x + c_1)^{1/4}}{x} \\ y &= \frac{i(-4 \cos(x) x^4 + 16 \sin(x) x^3 + 48 \cos(x) x^2 - 96 \cos(x) - 96 \sin(x) x + c_1)^{1/4}}{x} \\ y &= -\frac{(-4 \cos(x) x^4 + 16 \sin(x) x^3 + 48 \cos(x) x^2 - 96 \cos(x) - 96 \sin(x) x + c_1)^{1/4}}{x}\end{aligned}$$

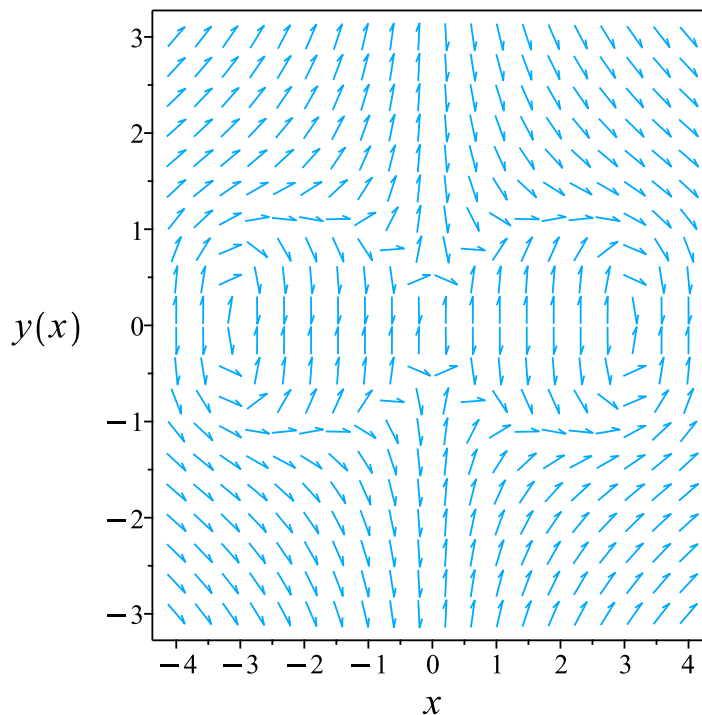


Figure 2.45: Slope field plot

$$y' + \frac{y}{x} = \frac{\sin(x)}{y^3}$$

Summary of solutions found

$$y = \frac{(-4 \cos(x) x^4 + 16 \sin(x) x^3 + 48 \cos(x) x^2 - 96 \cos(x) - 96 \sin(x) x + c_1)^{1/4}}{x}$$

$$y = -\frac{i(-4 \cos(x) x^4 + 16 \sin(x) x^3 + 48 \cos(x) x^2 - 96 \cos(x) - 96 \sin(x) x + c_1)^{1/4}}{x}$$

$$y = \frac{i(-4 \cos(x) x^4 + 16 \sin(x) x^3 + 48 \cos(x) x^2 - 96 \cos(x) - 96 \sin(x) x + c_1)^{1/4}}{x}$$

$$y = -\frac{(-4 \cos(x) x^4 + 16 \sin(x) x^3 + 48 \cos(x) x^2 - 96 \cos(x) - 96 \sin(x) x + c_1)^{1/4}}{x}$$

Solved as first order Exact ode

Time used: 0.444 (sec)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (\text{1A})$$

Therefore

$$\begin{aligned} (x y^3) dy &= (-y^4 + \sin(x) x) dx \\ (y^4 - \sin(x) x) dx + (x y^3) dy &= 0 \end{aligned} \quad (\text{2A})$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= y^4 - \sin(x) x \\ N(x, y) &= x y^3 \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(y^4 - \sin(x)x) \\ &= 4y^3\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(xy^3) \\ &= y^3\end{aligned}$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned}A &= \frac{1}{N} \left(\frac{\partial M}{\partial y} - \frac{\partial N}{\partial x} \right) \\ &= \frac{1}{xy^3} ((4y^3) - (y^3)) \\ &= \frac{3}{x}\end{aligned}$$

Since A does not depend on y , then it can be used to find an integrating factor. The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int A dx} \\ &= e^{\int \frac{3}{x} dx}\end{aligned}$$

The result of integrating gives

$$\begin{aligned}\mu &= e^{3\ln(x)} \\ &= x^3\end{aligned}$$

M and N are multiplied by this integrating factor, giving new M and new N which are called \bar{M} and \bar{N} for now so not to confuse them with the original M and N .

$$\begin{aligned}\bar{M} &= \mu M \\ &= x^3(y^4 - \sin(x)x) \\ &= (y^4 - \sin(x)x)x^3\end{aligned}$$

And

$$\begin{aligned}\bar{N} &= \mu N \\ &= x^3(x y^3) \\ &= x^4 y^3\end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned}\bar{M} + \bar{N} \frac{dy}{dx} &= 0 \\ ((y^4 - \sin(x)x)x^3) + (x^4 y^3) \frac{dy}{dx} &= 0\end{aligned}$$

The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = \bar{M} \quad (1)$$

$$\frac{\partial \phi}{\partial y} = \bar{N} \quad (2)$$

Integrating (1) w.r.t. x gives

$$\int \frac{\partial \phi}{\partial x} dx = \int \bar{M} dx$$

$$\int \frac{\partial \phi}{\partial x} dx = \int (y^4 - \sin(x)x)x^3 dx$$

$$\phi = (x^4 - 12x^2 + 24) \cos(x) + 4(-x^3 + 6x) \sin(x) + \frac{y^4 x^4}{4} + f(y) \quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = x^4 y^3 + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = x^4 y^3$. Therefore equation (4) becomes

$$x^4 y^3 = x^4 y^3 + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(y)$ into equation (3) gives ϕ

$$\phi = (x^4 - 12x^2 + 24) \cos(x) + 4(-x^3 + 6x) \sin(x) + \frac{y^4 x^4}{4} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into the constant c_1 gives the solution as

$$c_1 = (x^4 - 12x^2 + 24) \cos(x) + 4(-x^3 + 6x) \sin(x) + \frac{y^4 x^4}{4}$$

Solving for y gives

$$y = \frac{(-4 \cos(x) x^4 + 16 \sin(x) x^3 + 48 \cos(x) x^2 - 96 \sin(x) x - 96 \cos(x) + 4c_1)^{1/4}}{x}$$

$$y = -\frac{i(-4 \cos(x) x^4 + 16 \sin(x) x^3 + 48 \cos(x) x^2 - 96 \sin(x) x - 96 \cos(x) + 4c_1)^{1/4}}{x}$$

$$y = \frac{i(-4 \cos(x) x^4 + 16 \sin(x) x^3 + 48 \cos(x) x^2 - 96 \sin(x) x - 96 \cos(x) + 4c_1)^{1/4}}{x}$$

$$y = -\frac{(-4 \cos(x) x^4 + 16 \sin(x) x^3 + 48 \cos(x) x^2 - 96 \sin(x) x - 96 \cos(x) + 4c_1)^{1/4}}{x}$$

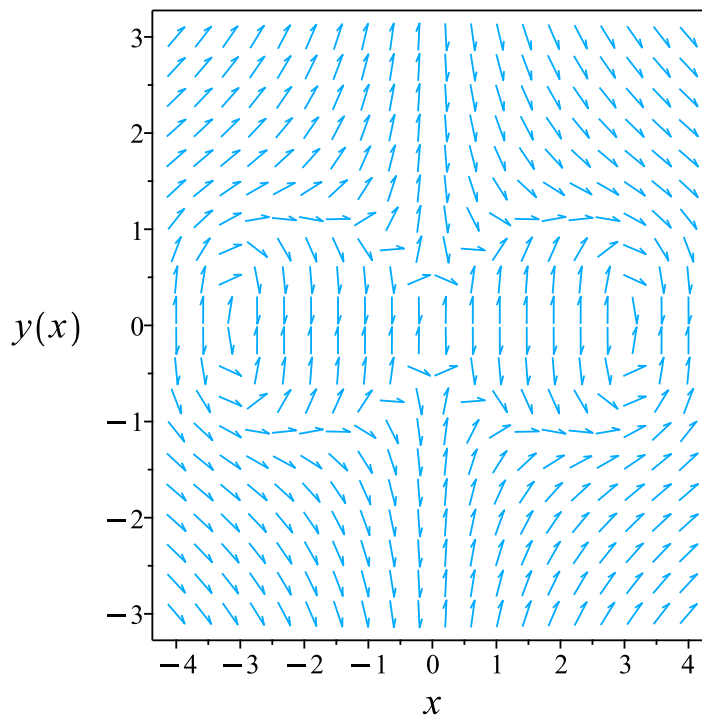


Figure 2.46: Slope field plot

$$y' + \frac{y}{x} = \frac{\sin(x)}{y^3}$$

Summary of solutions found

$$y = \frac{(-4 \cos(x) x^4 + 16 \sin(x) x^3 + 48 \cos(x) x^2 - 96 \sin(x) x - 96 \cos(x) + 4c_1)^{1/4}}{x}$$

$$y = -\frac{i(-4 \cos(x) x^4 + 16 \sin(x) x^3 + 48 \cos(x) x^2 - 96 \sin(x) x - 96 \cos(x) + 4c_1)^{1/4}}{x}$$

$$y = \frac{i(-4 \cos(x) x^4 + 16 \sin(x) x^3 + 48 \cos(x) x^2 - 96 \sin(x) x - 96 \cos(x) + 4c_1)^{1/4}}{x}$$

$$y = -\frac{(-4 \cos(x) x^4 + 16 \sin(x) x^3 + 48 \cos(x) x^2 - 96 \sin(x) x - 96 \cos(x) + 4c_1)^{1/4}}{x}$$

Maple step by step solution

Let's solve

$$y' + \frac{y}{x} = \frac{\sin(x)}{y^3}$$

- Highest derivative means the order of the ODE is 1

y'

- Solve for the highest derivative

$$y' = -\frac{y}{x} + \frac{\sin(x)}{y^3}$$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
<- Bernoulli successful`
```

Maple dsolve solution

Solving time : 0.026 (sec)

Leaf size : 156

```
dsolve(diff(y(x),x)+y(x)/x = sin(x)/y(x)^3,
        y(x),singsol=all)
```

$$y = \frac{(4(-x^4 + 12x^2 - 24) \cos(x) + 16(x^3 - 6x) \sin(x) + c_1)^{1/4}}{x}$$

$$y = -\frac{(4(-x^4 + 12x^2 - 24) \cos(x) + 16(x^3 - 6x) \sin(x) + c_1)^{1/4}}{x}$$

$$y = -\frac{i(4(-x^4 + 12x^2 - 24) \cos(x) + 16(x^3 - 6x) \sin(x) + c_1)^{1/4}}{x}$$

$$y = \frac{i(4(-x^4 + 12x^2 - 24) \cos(x) + 16(x^3 - 6x) \sin(x) + c_1)^{1/4}}{x}$$

Mathematica DSolve solution

Solving time : 0.531 (sec)

Leaf size : 114

```
DSolve[{D[y[x],x]+y[x]/x==Sin[x]/y[x]^2,{}},
y[x],x,IncludeSingularSolutions->True]
```

$$y(x) \rightarrow \frac{\sqrt[3]{9(x^2 - 2)\sin(x) - 3x(x^2 - 6)\cos(x) + c_1}}{x}$$

$$y(x) \rightarrow -\frac{\sqrt[3]{-1}\sqrt[3]{9(x^2 - 2)\sin(x) - 3x(x^2 - 6)\cos(x) + c_1}}{x}$$

$$y(x) \rightarrow \frac{(-1)^{2/3}\sqrt[3]{9(x^2 - 2)\sin(x) - 3x(x^2 - 6)\cos(x) + c_1}}{x}$$

2.4.4 problem 4

Solved as first order linear ode	306
Solved as first order homogeneous class D2 ode	308
Solved as first order Exact ode	309
Solved using Lie symmetry for first order ode	313
Maple step by step solution	317
Maple trace	318
Maple dsolve solution	319
Mathematica DSolve solution	319

Internal problem ID [18231]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter IV. Methods of solution: First order equations. section 31. Problems at page 85

Problem number : 4

Date solved : Monday, December 23, 2024 at 09:17:08 PM

CAS classification : [_linear]

Solve

$$p' = \frac{p + at^3 - 2pt^2}{t(-t^2 + 1)}$$

Solved as first order linear ode

Time used: 0.121 (sec)

In canonical form a linear first order is

$$p' + q(t)p = p(t)$$

Comparing the above to the given ode shows that

$$q(t) = -\frac{2t^2 - 1}{t^3 - t}$$

$$p(t) = -\frac{at^2}{t^2 - 1}$$

The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int q dt} \\ &= e^{\int -\frac{2t^2-1}{t^3-t} dt} \\ &= \frac{1}{t\sqrt{t+1}\sqrt{t-1}}\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dt}(\mu p) &= \mu p \\ \frac{d}{dt}(\mu p) &= (\mu) \left(-\frac{at^2}{t^2-1} \right) \\ \frac{d}{dt} \left(\frac{p}{t\sqrt{t+1}\sqrt{t-1}} \right) &= \left(\frac{1}{t\sqrt{t+1}\sqrt{t-1}} \right) \left(-\frac{at^2}{t^2-1} \right) \\ d \left(\frac{p}{t\sqrt{t+1}\sqrt{t-1}} \right) &= \left(-\frac{at}{(t^2-1)\sqrt{t+1}\sqrt{t-1}} \right) dt\end{aligned}$$

Integrating gives

$$\begin{aligned}\frac{p}{t\sqrt{t+1}\sqrt{t-1}} &= \int -\frac{at}{(t^2-1)\sqrt{t+1}\sqrt{t-1}} dt \\ &= \frac{\sqrt{t-1}\sqrt{t+1}a}{t^2-1} + c_1\end{aligned}$$

Dividing throughout by the integrating factor $\frac{1}{t\sqrt{t+1}\sqrt{t-1}}$ gives the final solution

$$p = \frac{t(\sqrt{t-1}\sqrt{t+1}a + c_1(t^2-1))\sqrt{t-1}\sqrt{t+1}}{t^2-1}$$

Summary of solutions found

$$p = \frac{t(\sqrt{t-1}\sqrt{t+1}a + c_1(t^2-1))\sqrt{t-1}\sqrt{t+1}}{t^2-1}$$

Solved as first order homogeneous class D2 ode

Time used: 0.281 (sec)

Applying change of variables $p = u(t)t$, then the ode becomes

$$u'(t)t + u(t) = \frac{u(t)t + at^3 - 2u(t)t^3}{t(-t^2 + 1)}$$

Which is now solved The ode $u'(t) = \frac{t(u(t)-a)}{t^2-1}$ is separable as it can be written as

$$\begin{aligned} u'(t) &= \frac{t(u(t) - a)}{t^2 - 1} \\ &= f(t)g(u) \end{aligned}$$

Where

$$\begin{aligned} f(t) &= \frac{t}{t^2 - 1} \\ g(u) &= u - a \end{aligned}$$

Integrating gives

$$\begin{aligned} \int \frac{1}{g(u)} du &= \int f(t) dt \\ \int \frac{1}{u - a} du &= \int \frac{t}{t^2 - 1} dt \\ \ln(-u(t) + a) &= \ln(\sqrt{t^2 - 1}) + c_1 \end{aligned}$$

We now need to find the singular solutions, these are found by finding for what values $g(u)$ is zero, since we had to divide by this above. Solving $g(u) = 0$ or $u - a = 0$ for $u(t)$ gives

$$u(t) = a$$

Now we go over each such singular solution and check if it verifies the ode itself and any initial conditions given. If it does not then the singular solution will not be used.

Therefore the solutions found are

$$\begin{aligned} \ln(-u(t) + a) &= \ln(\sqrt{t^2 - 1}) + c_1 \\ u(t) &= a \end{aligned}$$

Solving for $u(t)$ gives

$$\begin{aligned} u(t) &= a \\ u(t) &= -e^{c_1} \sqrt{t^2 - 1} + a \end{aligned}$$

Converting $u(t) = a$ back to p gives

$$p = at$$

Converting $u(t) = -e^{c_1\sqrt{t^2-1}} + a$ back to p gives

$$p = \left(-e^{c_1\sqrt{t^2-1}} + a\right) t$$

Summary of solutions found

$$p = at$$

$$p = \left(-e^{c_1\sqrt{t^2-1}} + a\right) t$$

Solved as first order Exact ode

Time used: 0.318 (sec)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx}\phi(x, y) = 0$$

Hence

$$\frac{\partial\phi}{\partial x} + \frac{\partial\phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial\phi}{\partial x} &= M \\ \frac{\partial\phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition

$\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(t, p) dt + N(t, p) dp = 0 \quad (1A)$$

Therefore

$$\begin{aligned} dp &= \left(\frac{a t^3 - 2p t^2 + p}{t(-t^2 + 1)} \right) dt \\ \left(-\frac{a t^3 - 2p t^2 + p}{t(-t^2 + 1)} \right) dt + dp &= 0 \end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(t, p) &= -\frac{a t^3 - 2p t^2 + p}{t(-t^2 + 1)} \\ N(t, p) &= 1 \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial p} = \frac{\partial N}{\partial t}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial p} &= \frac{\partial}{\partial p} \left(-\frac{a t^3 - 2p t^2 + p}{t(-t^2 + 1)} \right) \\ &= \frac{-2t^2 + 1}{t^3 - t} \end{aligned}$$

And

$$\begin{aligned} \frac{\partial N}{\partial t} &= \frac{\partial}{\partial t}(1) \\ &= 0 \end{aligned}$$

Since $\frac{\partial M}{\partial p} \neq \frac{\partial N}{\partial t}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned} A &= \frac{1}{N} \left(\frac{\partial M}{\partial p} - \frac{\partial N}{\partial t} \right) \\ &= 1 \left(\left(-\frac{-2t^2 + 1}{t(-t^2 + 1)} \right) - (0) \right) \\ &= \frac{-2t^2 + 1}{t^3 - t} \end{aligned}$$

Since A does not depend on p , then it can be used to find an integrating factor. The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int A dt} \\ &= e^{\int \frac{-2t^2+1}{t^3-t} dt}\end{aligned}$$

The result of integrating gives

$$\begin{aligned}\mu &= e^{-\ln(t) - \frac{\ln(t+1)}{2} - \frac{\ln(t-1)}{2}} \\ &= \frac{1}{t\sqrt{t+1}\sqrt{t-1}}\end{aligned}$$

M and N are multiplied by this integrating factor, giving new M and new N which are called \bar{M} and \bar{N} for now so not to confuse them with the original M and N .

$$\begin{aligned}\bar{M} &= \mu M \\ &= \frac{1}{t\sqrt{t+1}\sqrt{t-1}} \left(-\frac{at^3 - 2pt^2 + p}{t(-t^2 + 1)} \right) \\ &= \frac{at^3 - 2pt^2 + p}{t^2(t^2 - 1)\sqrt{t+1}\sqrt{t-1}}\end{aligned}$$

And

$$\begin{aligned}\bar{N} &= \mu N \\ &= \frac{1}{t\sqrt{t+1}\sqrt{t-1}}(1) \\ &= \frac{1}{t\sqrt{t+1}\sqrt{t-1}}\end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned}\bar{M} + \bar{N} \frac{dp}{dt} &= 0 \\ \left(\frac{at^3 - 2pt^2 + p}{t^2(t^2 - 1)\sqrt{t+1}\sqrt{t-1}} \right) + \left(\frac{1}{t\sqrt{t+1}\sqrt{t-1}} \right) \frac{dp}{dt} &= 0\end{aligned}$$

The following equations are now set up to solve for the function $\phi(t, p)$

$$\frac{\partial \phi}{\partial t} = \bar{M} \tag{1}$$

$$\frac{\partial \phi}{\partial p} = \bar{N} \tag{2}$$

Integrating (2) w.r.t. p gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial p} dp &= \int \bar{N} dp \\ \int \frac{\partial \phi}{\partial p} dp &= \int \frac{1}{t\sqrt{t+1}\sqrt{t-1}} dp \\ \phi &= \frac{p}{t\sqrt{t+1}\sqrt{t-1}} + f(t)\end{aligned}\quad (3)$$

Where $f(t)$ is used for the constant of integration since ϕ is a function of both t and p . Taking derivative of equation (3) w.r.t t gives

$$\begin{aligned}\frac{\partial \phi}{\partial t} &= -\frac{p}{t^2\sqrt{t+1}\sqrt{t-1}} - \frac{p}{2t(t+1)^{3/2}\sqrt{t-1}} - \frac{p}{2t\sqrt{t+1}(t-1)^{3/2}} + f'(t) \\ &= -\frac{2p(t^2 - \frac{1}{2})}{(t-1)^{3/2}(t+1)^{3/2}t^2} + f'(t)\end{aligned}\quad (4)$$

But equation (1) says that $\frac{\partial \phi}{\partial t} = \frac{at^3 - 2pt^2 + p}{t^2(t^2-1)\sqrt{t+1}\sqrt{t-1}}$. Therefore equation (4) becomes

$$\frac{at^3 - 2pt^2 + p}{t^2(t^2-1)\sqrt{t+1}\sqrt{t-1}} = -\frac{2p(t^2 - \frac{1}{2})}{(t-1)^{3/2}(t+1)^{3/2}t^2} + f'(t)\quad (5)$$

Solving equation (5) for $f'(t)$ gives

$$f'(t) = \frac{at}{(t+1)^{3/2}(t-1)^{3/2}}$$

Integrating the above w.r.t t gives

$$\begin{aligned}\int f'(t) dt &= \int \left(\frac{at}{(t+1)^{3/2}(t-1)^{3/2}} \right) dt \\ f(t) &= -\frac{a}{\sqrt{t-1}\sqrt{t+1}} + c_1\end{aligned}$$

Where c_1 is constant of integration. Substituting result found above for $f(t)$ into equation (3) gives ϕ

$$\phi = \frac{p}{t\sqrt{t+1}\sqrt{t-1}} - \frac{a}{\sqrt{t-1}\sqrt{t+1}} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into the constant c_1 gives the solution as

$$c_1 = \frac{p}{t\sqrt{t+1}\sqrt{t-1}} - \frac{a}{\sqrt{t-1}\sqrt{t+1}}$$

Solving for p gives

$$p = c_1 t \sqrt{t+1} \sqrt{t-1} + at$$

Summary of solutions found

$$p = c_1 t \sqrt{t+1} \sqrt{t-1} + at$$

Solved using Lie symmetry for first order ode

Time used: 0.551 (sec)

Writing the ode as

$$p' = \frac{-at^3 + 2pt^2 - p}{t(t^2 - 1)}$$

$$p' = \omega(t, p)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_t + \omega(\eta_p - \xi_t) - \omega^2 \xi_p - \omega_t \xi - \omega_p \eta = 0 \quad (\text{A})$$

To determine ξ, η then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$\xi = pa_3 + ta_2 + a_1 \quad (\text{1E})$$

$$\eta = pb_3 + tb_2 + b_1 \quad (\text{2E})$$

Where the unknown coefficients are

$$\{a_1, a_2, a_3, b_1, b_2, b_3\}$$

Substituting equations (1E,2E) and ω into (A) gives

$$b_2 + \frac{(-at^3 + 2pt^2 - p)(b_3 - a_2)}{t(t^2 - 1)} - \frac{(-at^3 + 2pt^2 - p)^2 a_3}{t^2(t^2 - 1)^2}$$

$$- \left(\frac{-3at^2 + 4pt}{t(t^2 - 1)} - \frac{-at^3 + 2pt^2 - p}{t^2(t^2 - 1)} - \frac{2(-at^3 + 2pt^2 - p)}{(t^2 - 1)^2} \right) (pa_3$$

$$+ ta_2 + a_1) - \frac{(2t^2 - 1)(pb_3 + tb_2 + b_1)}{t(t^2 - 1)} = 0 \quad (\text{5E})$$

Putting the above in normal form gives

$$\frac{a^2 t^6 a_3 - 4apt^5 a_3 - a t^6 a_2 + a t^6 b_3 + 2p^2 t^4 a_3 + t^6 b_2 + 4apt^3 a_3 + 3a t^4 a_2 - a t^4 b_3 - 2pt^4 a_1 + 2t^5 b_1 + 2a t^4 a_1 - 2t^5 b_1 - 2a t^3 a_1 + 3p^2 t^2 a_3 + 2p t^3 a_2 + t^4 b_2 - p t^2 a_1 + 3t^3 b_1 + pa_1 - tb_1}{t^2 (t^2 - 1)^2} = 0$$

Setting the numerator to zero gives

$$\begin{aligned} & -a^2 t^6 a_3 + 4apt^5 a_3 + a t^6 a_2 - a t^6 b_3 - 2p^2 t^4 a_3 - t^6 b_2 - 4apt^3 a_3 - 3a t^4 a_2 + a t^4 b_3 \\ & + 2p t^4 a_1 - 2t^5 b_1 - 2a t^3 a_1 + 3p^2 t^2 a_3 + 2p t^3 a_2 + t^4 b_2 - p t^2 a_1 + 3t^3 b_1 + pa_1 - tb_1 \\ & = 0 \end{aligned} \quad (6E)$$

Looking at the above PDE shows the following are all the terms with $\{p, t\}$ in them.

$$\{p, t\}$$

The following substitution is now made to be able to collect on all terms with $\{p, t\}$ in them

$$\{p = v_1, t = v_2\}$$

The above PDE (6E) now becomes

$$\begin{aligned} & -a^2 a_3 v_2^6 + aa_2 v_2^6 + 4aa_3 v_1 v_2^5 - ab_3 v_2^6 - 2a_3 v_1^2 v_2^4 - b_2 v_2^6 \\ & - 3aa_2 v_2^4 - 4aa_3 v_1 v_2^3 + ab_3 v_2^4 + 2a_1 v_1 v_2^4 - 2b_1 v_2^5 - 2aa_1 v_2^3 \\ & + 2a_2 v_1 v_2^3 + 3a_3 v_1^2 v_2^2 + b_2 v_2^4 - a_1 v_1 v_2^2 + 3b_1 v_2^3 + a_1 v_1 - b_1 v_2 = 0 \end{aligned} \quad (7E)$$

Collecting the above on the terms v_i introduced, and these are

$$\{v_1, v_2\}$$

Equation (7E) now becomes

$$\begin{aligned} & -2a_3 v_1^2 v_2^4 + 3a_3 v_1^2 v_2^2 + 4aa_3 v_1 v_2^5 + 2a_1 v_1 v_2^4 + (-4aa_3 + 2a_2) v_1 v_2^3 \\ & - a_1 v_1 v_2^2 + a_1 v_1 + (-a^2 a_3 + aa_2 - ab_3 - b_2) v_2^6 - 2b_1 v_2^5 \\ & + (-3aa_2 + ab_3 + b_2) v_2^4 + (-2aa_1 + 3b_1) v_2^3 - b_1 v_2 = 0 \end{aligned} \quad (8E)$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$\begin{aligned}
 a_1 &= 0 \\
 -a_1 &= 0 \\
 2a_1 &= 0 \\
 -2a_3 &= 0 \\
 3a_3 &= 0 \\
 -2b_1 &= 0 \\
 -b_1 &= 0 \\
 4aa_3 &= 0 \\
 -2aa_1 + 3b_1 &= 0 \\
 -4aa_3 + 2a_2 &= 0 \\
 -3aa_2 + ab_3 + b_2 &= 0 \\
 -a^2a_3 + aa_2 - ab_3 - b_2 &= 0
 \end{aligned}$$

Solving the above equations for the unknowns gives

$$\begin{aligned}
 a_1 &= 0 \\
 a_2 &= 0 \\
 a_3 &= 0 \\
 b_1 &= 0 \\
 b_2 &= -ab_3 \\
 b_3 &= b_3
 \end{aligned}$$

Substituting the above solution in the ansatz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$\begin{aligned}
 \xi &= 0 \\
 \eta &= -at + p
 \end{aligned}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(t, p) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dt}{\xi} = \frac{dp}{\eta} = dS \tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial t} + \eta \frac{\partial}{\partial p}\right) S(t, p) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = t$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{-at + p} dy \end{aligned}$$

Which results in

$$S = \ln(-at + p)$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_t + \omega(t, p)S_p}{R_t + \omega(t, p)R_p} \quad (2)$$

Where in the above R_t, R_p, S_t, S_p are all partial derivatives and $\omega(t, p)$ is the right hand side of the original ode given by

$$\omega(t, p) = \frac{-at^3 + 2pt^2 - p}{t(t^2 - 1)}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_t &= 1 \\ R_p &= 0 \\ S_t &= \frac{a}{at - p} \\ S_p &= \frac{1}{-at + p} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = \frac{2t^2 - 1}{t^3 - t} \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for t, p in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = \frac{2R^2 - 1}{R^3 - R}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S .

Since the ode has the form $\frac{d}{dR}S(R) = f(R)$, then we only need to integrate $f(R)$.

$$\int dS = \int \frac{2R^2 - 1}{R(R^2 - 1)} dR$$

$$S(R) = \ln(R) + \frac{\ln(R+1)}{2} + \frac{\ln(R-1)}{2} + c_2$$

To complete the solution, we just need to transform the above back to t, p coordinates. This results in

$$\ln(-at + p) = \ln(t) + \frac{\ln(t+1)}{2} + \frac{\ln(t-1)}{2} + c_2$$

Which gives

$$p = t \left(e^{-\frac{\ln(t+1)}{2} - \frac{\ln(t-1)}{2} - c_2} a + 1 \right) e^{\ln(\sqrt{t+1}) + \ln(\sqrt{t-1}) + c_2}$$

Summary of solutions found

$$p = t \left(e^{-\frac{\ln(t+1)}{2} - \frac{\ln(t-1)}{2} - c_2} a + 1 \right) e^{\ln(\sqrt{t+1}) + \ln(\sqrt{t-1}) + c_2}$$

Maple step by step solution

Let's solve

$$p' = \frac{p+at^3-2pt^2}{t(-t^2+1)}$$

- Highest derivative means the order of the ODE is 1

$$p'$$

- Solve for the highest derivative

$$p' = \frac{p+at^3-2pt^2}{t(-t^2+1)}$$

- Collect w.r.t. p and simplify

$$p' = \frac{(2t^2-1)p}{t(t^2-1)} - \frac{at^2}{t^2-1}$$

- Group terms with p on the lhs of the ODE and the rest on the rhs of the ODE

$$p' - \frac{(2t^2-1)p}{t(t^2-1)} = -\frac{at^2}{t^2-1}$$

- The ODE is linear; multiply by an integrating factor $\mu(t)$

$$\mu(t) \left(p' - \frac{(2t^2-1)p}{t(t^2-1)} \right) = -\frac{\mu(t)at^2}{t^2-1}$$

- Assume the lhs of the ODE is the total derivative $\frac{d}{dt}(p\mu(t))$

$$\mu(t) \left(p' - \frac{(2t^2-1)p}{t(t^2-1)} \right) = p'\mu(t) + p\mu'(t)$$

- Isolate $\mu'(t)$

$$\mu'(t) = -\frac{\mu(t)(2t^2-1)}{t(t^2-1)}$$

- Solve to find the integrating factor

$$\mu(t) = \frac{1}{t\sqrt{t+1}\sqrt{t-1}}$$

- Integrate both sides with respect to t

$$\int \left(\frac{d}{dt}(p\mu(t)) \right) dt = \int -\frac{\mu(t)at^2}{t^2-1} dt + C1$$

- Evaluate the integral on the lhs

$$p\mu(t) = \int -\frac{\mu(t)at^2}{t^2-1} dt + C1$$

- Solve for p

$$p = \frac{\int -\frac{\mu(t)at^2}{t^2-1} dt + C1}{\mu(t)}$$

- Substitute $\mu(t) = \frac{1}{t\sqrt{t+1}\sqrt{t-1}}$

$$p = t\sqrt{t+1}\sqrt{t-1} \left(\int -\frac{at}{(t^2-1)\sqrt{t+1}\sqrt{t-1}} dt + C1 \right)$$

- Evaluate the integrals on the rhs

$$p = t\sqrt{t+1}\sqrt{t-1} \left(\frac{\sqrt{t-1}\sqrt{t+1}a}{t^2-1} + C1 \right)$$

- Simplify

$$p = \frac{t(\sqrt{t-1}\sqrt{t+1}a + C1(t^2-1))\sqrt{t-1}\sqrt{t+1}}{t^2-1}$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`

```

Maple dsolve solution

Solving time : 0.003 (sec)

Leaf size : 20

```
dsolve(diff(p(t),t) = (p(t)+a*t^3-2*p(t)*t^2)/t/(-t^2+1),  
p(t),singsol=all)
```

$$p = t \left(c_1 \sqrt{t+1} \sqrt{t-1} + a \right)$$

Mathematica DSolve solution

Solving time : 0.046 (sec)

Leaf size : 23

```
DSolve[{D[p[t],t]==(p[t]+a*t^3-2*p[t]*t^2)/(t*(1-t^2)),{}},  
p[t],t,IncludeSingularSolutions->True]
```

$$p(t) \rightarrow t \left(a + c_1 \sqrt{1-t^2} \right)$$

2.4.5 problem 5

Solved as first order Bernoulli ode	320
Solved as first order Exact ode	323
Solved using Lie symmetry for first order ode	328
Maple step by step solution	335
Maple trace	336
Maple dsolve solution	336
Mathematica DSolve solution	336

Internal problem ID [18232]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter IV. Methods of solution: First order equations. section 31. Problems at page 85

Problem number : 5

Date solved : Monday, December 23, 2024 at 09:17:10 PM

CAS classification : [_Bernoulli]

Solve

$$(T \ln(t) - 1)T = tT'$$

Solved as first order Bernoulli ode

Time used: 0.125 (sec)

In canonical form, the ODE is

$$\begin{aligned} T' &= F(t, T) \\ &= \frac{(T \ln(t) - 1)T}{t} \end{aligned}$$

This is a Bernoulli ODE.

$$T' = \left(-\frac{1}{t}\right)T + \left(\frac{\ln(t)}{t}\right)T^2 \quad (1)$$

The standard Bernoulli ODE has the form

$$T' = f_0(t)T + f_1(t)T^n \quad (2)$$

Comparing this to (1) shows that

$$\begin{aligned} f_0 &= -\frac{1}{t} \\ f_1 &= \frac{\ln(t)}{t} \end{aligned}$$

The first step is to divide the above equation by T^n which gives

$$\frac{T'}{T^n} = f_0(t)T^{1-n} + f_1(t) \quad (3)$$

The next step is use the substitution $v = T^{1-n}$ in equation (3) which generates a new ODE in $v(t)$ which will be linear and can be easily solved using an integrating factor. Backsubstitution then gives the solution $T(t)$ which is what we want.

This method is now applied to the ODE at hand. Comparing the ODE (1) With (2) Shows that

$$\begin{aligned} f_0(t) &= -\frac{1}{t} \\ f_1(t) &= \frac{\ln(t)}{t} \\ n &= 2 \end{aligned}$$

Dividing both sides of ODE (1) by $T^n = T^2$ gives

$$T' \frac{1}{T^2} = -\frac{1}{tT} + \frac{\ln(t)}{t} \quad (4)$$

Let

$$\begin{aligned} v &= T^{1-n} \\ &= \frac{1}{T} \end{aligned} \quad (5)$$

Taking derivative of equation (5) w.r.t t gives

$$v' = -\frac{1}{T^2}T' \quad (6)$$

Substituting equations (5) and (6) into equation (4) gives

$$\begin{aligned} -v'(t) &= -\frac{v(t)}{t} + \frac{\ln(t)}{t} \\ v' &= \frac{v}{t} - \frac{\ln(t)}{t} \end{aligned} \quad (7)$$

The above now is a linear ODE in $v(t)$ which is now solved.

In canonical form a linear first order is

$$v'(t) + q(t)v(t) = p(t)$$

Comparing the above to the given ode shows that

$$q(t) = -\frac{1}{t}$$

$$p(t) = -\frac{\ln(t)}{t}$$

The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int q dt} \\ &= e^{\int -\frac{1}{t} dt} \\ &= \frac{1}{t}\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dt}(\mu v) &= \mu p \\ \frac{d}{dt}(\mu v) &= (\mu) \left(-\frac{\ln(t)}{t} \right) \\ \frac{d}{dt} \left(\frac{v}{t} \right) &= \left(\frac{1}{t} \right) \left(-\frac{\ln(t)}{t} \right) \\ d \left(\frac{v}{t} \right) &= \left(-\frac{\ln(t)}{t^2} \right) dt\end{aligned}$$

Integrating gives

$$\begin{aligned}\frac{v}{t} &= \int -\frac{\ln(t)}{t^2} dt \\ &= \frac{\ln(t)}{t} + \frac{1}{t} + c_1\end{aligned}$$

Dividing throughout by the integrating factor $\frac{1}{t}$ gives the final solution

$$v(t) = c_1 t + \ln(t) + 1$$

The substitution $v = T^{1-n}$ is now used to convert the above solution back to T which results in

$$\frac{1}{T} = c_1 t + \ln(t) + 1$$

Solving for T gives

$$T = \frac{1}{c_1 t + \ln(t) + 1}$$

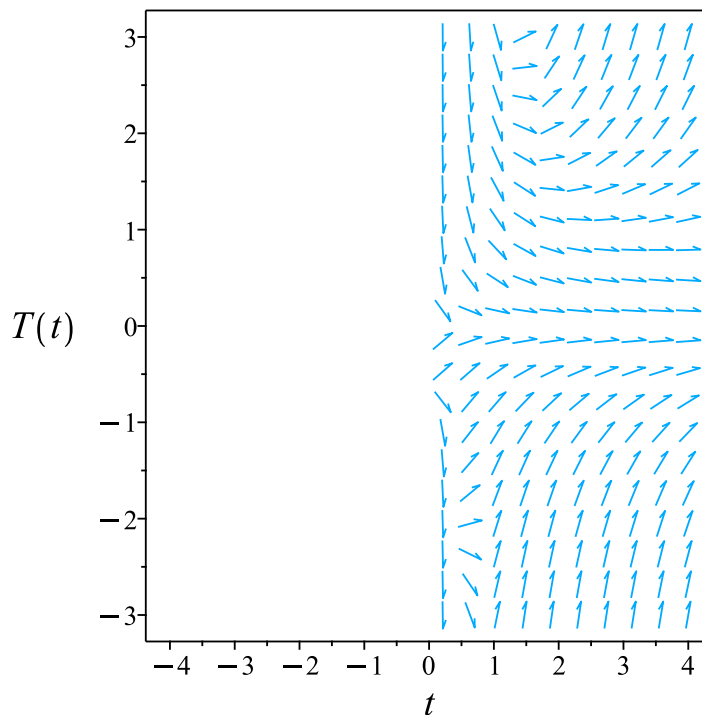


Figure 2.47: Slope field plot
 $(T \ln(t) - 1)T = tT'$

Summary of solutions found

$$T = \frac{1}{c_1 t + \ln(t) + 1}$$

Solved as first order Exact ode

Time used: 0.319 (sec)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(t, T) dt + N(t, T) dT = 0 \quad (1A)$$

Therefore

$$\begin{aligned}(-t) dT &= -(T \ln(t) - 1) T dt \\ ((T \ln(t) - 1) T) dt + (-t) dT &= 0\end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(t, T) &= (T \ln(t) - 1) T \\ N(t, T) &= -t\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial T} = \frac{\partial N}{\partial t}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial T} &= \frac{\partial}{\partial T}((T \ln(t) - 1) T) \\ &= 2T \ln(t) - 1\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial t} &= \frac{\partial}{\partial t}(-t) \\ &= -1\end{aligned}$$

Since $\frac{\partial M}{\partial T} \neq \frac{\partial N}{\partial t}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned} A &= \frac{1}{N} \left(\frac{\partial M}{\partial T} - \frac{\partial N}{\partial t} \right) \\ &= -\frac{1}{t} ((2T \ln(t) - 1) - (-1)) \\ &= -\frac{2T \ln(t)}{t} \end{aligned}$$

Since A depends on T , it can not be used to obtain an integrating factor. We will now try a second method to find an integrating factor. Let

$$\begin{aligned} B &= \frac{1}{M} \left(\frac{\partial N}{\partial t} - \frac{\partial M}{\partial T} \right) \\ &= \frac{1}{(T \ln(t) - 1) T} ((-1) - (2T \ln(t) - 1)) \\ &= -\frac{2 \ln(t)}{T \ln(t) - 1} \end{aligned}$$

Since B depends on t , it can not be used to obtain an integrating factor. We will now try a third method to find an integrating factor. Let

$$R = \frac{\frac{\partial N}{\partial t} - \frac{\partial M}{\partial T}}{xM - yN}$$

R is now checked to see if it is a function of only $t = tT$. Therefore

$$\begin{aligned} R &= \frac{\frac{\partial N}{\partial t} - \frac{\partial M}{\partial T}}{xM - yN} \\ &= \frac{(-1) - (2T \ln(t) - 1)}{t((T \ln(t) - 1) T) - T(-t)} \\ &= -\frac{2}{tT} \end{aligned}$$

Replacing all powers of terms tT by t gives

$$R = -\frac{2}{t}$$

Since R depends on t only, then it can be used to find an integrating factor. Let the integrating factor be μ then

$$\begin{aligned} \mu &= e^{\int R dt} \\ &= e^{\int (-\frac{2}{t}) dt} \end{aligned}$$

The result of integrating gives

$$\begin{aligned}\mu &= e^{-2\ln(t)} \\ &= \frac{1}{t^2}\end{aligned}$$

Now t is replaced back with tT giving

$$\mu = \frac{1}{t^2 T^2}$$

Multiplying M and N by this integrating factor gives new M and new N which are called \bar{M} and \bar{N} so not to confuse them with the original M and N

$$\begin{aligned}\bar{M} &= \mu M \\ &= \frac{1}{t^2 T^2} ((T \ln(t) - 1) T) \\ &= \frac{T \ln(t) - 1}{T t^2}\end{aligned}$$

And

$$\begin{aligned}\bar{N} &= \mu N \\ &= \frac{1}{t^2 T^2} (-t) \\ &= -\frac{1}{t T^2}\end{aligned}$$

A modified ODE is now obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned}\bar{M} + \bar{N} \frac{dT}{dt} &= 0 \\ \left(\frac{T \ln(t) - 1}{T t^2} \right) + \left(-\frac{1}{t T^2} \right) \frac{dT}{dt} &= 0\end{aligned}$$

The following equations are now set up to solve for the function $\phi(t, T)$

$$\frac{\partial \phi}{\partial t} = \bar{M} \tag{1}$$

$$\frac{\partial \phi}{\partial T} = \bar{N} \tag{2}$$

Integrating (1) w.r.t. t gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial t} dt &= \int \bar{M} dt \\ \int \frac{\partial \phi}{\partial t} dt &= \int \frac{T \ln(t) - 1}{T t^2} dt \\ \phi &= \frac{-T \ln(t) - T + 1}{tT} + f(T)\end{aligned} \tag{3}$$

Where $f(T)$ is used for the constant of integration since ϕ is a function of both t and T . Taking derivative of equation (3) w.r.t T gives

$$\begin{aligned}\frac{\partial \phi}{\partial T} &= \frac{-\ln(t) - 1}{tT} - \frac{-T \ln(t) - T + 1}{tT^2} + f'(T) \\ &= -\frac{1}{tT^2} + f'(T)\end{aligned}\quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial T} = -\frac{1}{tT^2}$. Therefore equation (4) becomes

$$-\frac{1}{tT^2} = -\frac{1}{tT^2} + f'(T)\quad (5)$$

Solving equation (5) for $f'(T)$ gives

$$f'(T) = 0$$

Therefore

$$f(T) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(T)$ into equation (3) gives ϕ

$$\phi = \frac{-T \ln(t) - T + 1}{tT} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into the constant c_1 gives the solution as

$$c_1 = \frac{-T \ln(t) - T + 1}{tT}$$

Solving for T gives

$$T = \frac{1}{c_1 t + \ln(t) + 1}$$

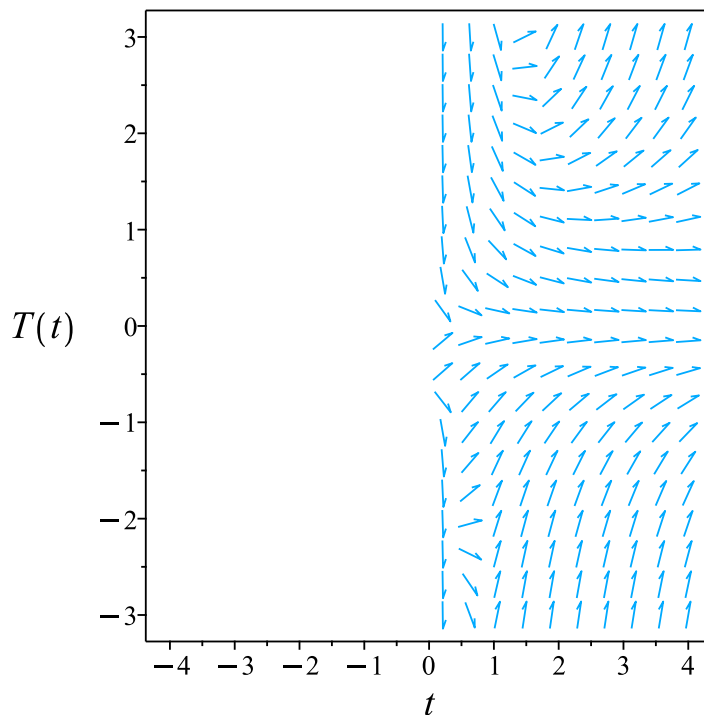


Figure 2.48: Slope field plot
 $(T \ln(t) - 1)T = tT'$

Summary of solutions found

$$T = \frac{1}{c_1 t + \ln(t) + 1}$$

Solved using Lie symmetry for first order ode

Time used: 1.423 (sec)

Writing the ode as

$$T' = \frac{(T \ln(t) - 1)T}{t}$$

$$T' = \omega(t, T)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_t + \omega(\eta_T - \xi_t) - \omega^2 \xi_T - \omega_t \xi - \omega_T \eta = 0 \quad (\text{A})$$

To determine ξ, η then (A) is solved using ansatz. Making bivariate polynomials of degree 2 to use as anstaz gives

$$\xi = T^2 a_6 + tT a_5 + t^2 a_4 + T a_3 + t a_2 + a_1 \quad (1\text{E})$$

$$\eta = T^2 b_6 + tT b_5 + t^2 b_4 + T b_3 + t b_2 + b_1 \quad (2\text{E})$$

Where the unknown coefficients are

$$\{a_1, a_2, a_3, a_4, a_5, a_6, b_1, b_2, b_3, b_4, b_5, b_6\}$$

Substituting equations (1E,2E) and ω into (A) gives

$$\begin{aligned} & Tb_5 + 2tb_4 + b_2 + \frac{(T \ln(t) - 1)T(-Ta_5 + 2Tb_6 - 2ta_4 + tb_5 - a_2 + b_3)}{t} \\ & - \frac{(T \ln(t) - 1)^2 T^2(2Ta_6 + ta_5 + a_3)}{t^2} \\ & - \left(\frac{T^2}{t^2} - \frac{(T \ln(t) - 1)T}{t^2} \right) (T^2a_6 + tTa_5 + t^2a_4 + Ta_3 + ta_2 + a_1) \\ & - \left(\frac{T \ln(t)}{t} + \frac{T \ln(t) - 1}{t} \right) (T^2b_6 + tTb_5 + t^2b_4 + Tb_3 + tb_2 + b_1) = 0 \end{aligned} \quad (5E)$$

Putting the above in normal form gives

$$\begin{aligned} & \frac{-Tb_5t^2 + T^3ta_5 + T^2t^2a_4 + T^2ta_5 + T^2tb_6 - Tt^2a_4 + 2 \ln(t)^2 T^5a_6 - 5 \ln(t) T^4a_6 + 2T^2a_3 + T^2ta_2 + \dots}{\dots} \\ & = 0 \end{aligned}$$

Setting the numerator to zero gives

$$\begin{aligned} & Tb_5t^2 - T^3ta_5 - T^2t^2a_4 - T^2ta_5 - T^2tb_6 + Tt^2a_4 - 2 \ln(t)^2 T^5a_6 \\ & + 5 \ln(t) T^4a_6 - 2T^2a_3 - T^2ta_2 - \ln(t)^2 T^4a_3 + 3 \ln(t) T^3a_3 \\ & + \ln(t) T^2a_1 + 2b_2t^2 - T^3a_3 - T^2a_1 - Ta_1 + tb_1 + 3t^3b_4 - T^4a_6 \\ & - 3T^3a_6 - \ln(t)^2 T^4ta_5 + 2 \ln(t) T^3ta_5 - \ln(t) T^2t^2a_4 - \ln(t) T^2t^2b_5 \\ & - 2 \ln(t) Tt^3b_4 - \ln(t) T^2tb_3 - 2 \ln(t) Tt^2b_2 - 2 \ln(t) Ttb_1 = 0 \end{aligned} \quad (6E)$$

Looking at the above PDE shows the following are all the terms with $\{T, t\}$ in them.

$$\{T, t, \ln(t)\}$$

The following substitution is now made to be able to collect on all terms with $\{T, t\}$ in them

$$\{T = v_1, t = v_2, \ln(t) = v_3\}$$

The above PDE (6E) now becomes

$$\begin{aligned}
& -v_3^2 v_1^4 v_2 a_5 - 2v_3^2 v_1^5 a_6 - v_3^2 v_1^4 a_3 - v_3 v_1^2 v_2^2 a_4 + 2v_3 v_1^3 v_2 a_5 + 5v_3 v_1^4 a_6 \\
& - 2v_3 v_1 v_2^3 b_4 - v_3 v_1^2 v_2^2 b_5 + 3v_3 v_1^3 a_3 - v_1^2 v_2^2 a_4 - v_1^3 v_2 a_5 - v_1^4 a_6 - 2v_3 v_1 v_2^2 b_2 \\
& - v_3 v_1^2 v_2 b_3 + v_3 v_1^2 a_1 - v_1^2 v_2 a_2 - v_1^3 a_3 + v_1 v_2^2 a_4 - v_1^2 v_2 a_5 - 3v_1^3 a_6 - 2v_3 v_1 v_2 b_1 \\
& + 3v_2^3 b_4 + v_1 b_5 v_2^2 - v_1^2 v_2 b_6 - v_1^2 a_1 - 2v_1^2 a_3 + 2b_2 v_2^2 - v_1 a_1 + v_2 b_1 = 0
\end{aligned} \tag{7E}$$

Collecting the above on the terms v_i introduced, and these are

$$\{v_1, v_2, v_3\}$$

Equation (7E) now becomes

$$\begin{aligned}
& -2v_3^2 v_1^5 a_6 - v_3^2 v_1^4 v_2 a_5 - v_3^2 v_1^4 a_3 + 5v_3 v_1^4 a_6 - v_1^4 a_6 + 2v_3 v_1^3 v_2 a_5 - v_1^3 v_2 a_5 \\
& + 3v_3 v_1^3 a_3 + (-a_3 - 3a_6) v_1^3 + (-a_4 - b_5) v_1^2 v_2^2 v_3 - v_1^2 v_2^2 a_4 - v_3 v_1^2 v_2 b_3 \\
& + (-a_2 - a_5 - b_6) v_1^2 v_2 + v_3 v_1^2 a_1 + (-a_1 - 2a_3) v_1^2 - 2v_3 v_1 v_2^3 b_4 - 2v_3 v_1 v_2^2 b_2 \\
& + (a_4 + b_5) v_1 v_2^2 - 2v_3 v_1 v_2 b_1 - v_1 a_1 + 3v_2^3 b_4 + 2b_2 v_2^2 + v_2 b_1 = 0
\end{aligned} \tag{8E}$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$a_1 = 0$$

$$b_1 = 0$$

$$-a_1 = 0$$

$$-a_3 = 0$$

$$3a_3 = 0$$

$$-a_4 = 0$$

$$-a_5 = 0$$

$$2a_5 = 0$$

$$-2a_6 = 0$$

$$-a_6 = 0$$

$$5a_6 = 0$$

$$-2b_1 = 0$$

$$-2b_2 = 0$$

$$2b_2 = 0$$

$$-b_3 = 0$$

$$-2b_4 = 0$$

$$3b_4 = 0$$

$$-a_1 - 2a_3 = 0$$

$$-a_3 - 3a_6 = 0$$

$$-a_4 - b_5 = 0$$

$$a_4 + b_5 = 0$$

$$-a_2 - a_5 - b_6 = 0$$

Solving the above equations for the unknowns gives

$$\begin{aligned}
 a_1 &= 0 \\
 a_2 &= -b_6 \\
 a_3 &= 0 \\
 a_4 &= 0 \\
 a_5 &= 0 \\
 a_6 &= 0 \\
 b_1 &= 0 \\
 b_2 &= 0 \\
 b_3 &= 0 \\
 b_4 &= 0 \\
 b_5 &= 0 \\
 b_6 &= b_6
 \end{aligned}$$

Substituting the above solution in the ansatz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$\begin{aligned}
 \xi &= -t \\
 \eta &= T^2
 \end{aligned}$$

Shifting is now applied to make $\xi = 0$ in order to simplify the rest of the computation

$$\begin{aligned}
 \eta &= \eta - \omega(t, T) \xi \\
 &= T^2 - \left(\frac{(T \ln(t) - 1) T}{t} \right) (-t) \\
 &= T^2 \ln(t) + T^2 - T \\
 \xi &= 0
 \end{aligned}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(t, T) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dt}{\xi} = \frac{dT}{\eta} = dS \tag{1}$$

The above comes from the requirements that $(\xi \frac{\partial}{\partial t} + \eta \frac{\partial}{\partial T}) S(t, T) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the

canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = t$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{T^2 \ln(t) + T^2 - T} dy \end{aligned}$$

Which results in

$$S = -\ln(T) + \ln(T \ln(t) + T - 1)$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_t + \omega(t, T)S_T}{R_t + \omega(t, T)R_T} \quad (2)$$

Where in the above R_t, R_T, S_t, S_T are all partial derivatives and $\omega(t, T)$ is the right hand side of the original ode given by

$$\omega(t, T) = \frac{(T \ln(t) - 1)T}{t}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_t &= 1 \\ R_T &= 0 \\ S_t &= \frac{T}{t(T \ln(t) + T - 1)} \\ S_T &= \frac{1}{T(T \ln(t) + T - 1)} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = \frac{1}{t} \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for t, T in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = \frac{1}{R}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S .

Since the ode has the form $\frac{d}{dR}S(R) = f(R)$, then we only need to integrate $f(R)$.

$$\int dS = \int \frac{1}{R} dR$$

$$S(R) = \ln(R) + c_2$$

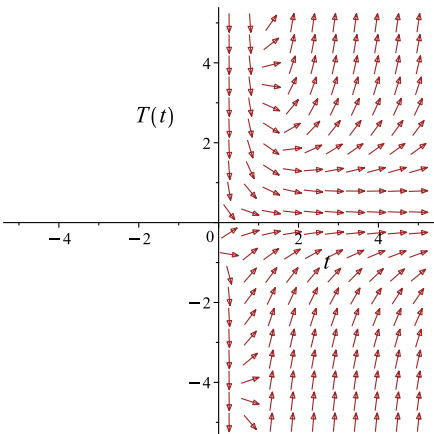
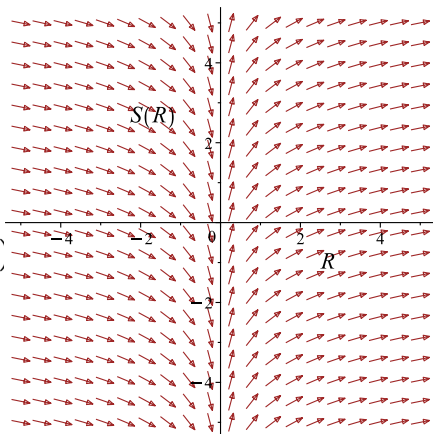
To complete the solution, we just need to transform the above back to t, T coordinates. This results in

$$-\ln(T) + \ln(T \ln(t) + T - 1) = \ln(t) + c_2$$

Which gives

$$T = -\frac{1}{e^{c_2 t} - \ln(t) - 1}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in t, T coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dT}{dt} = \frac{(T \ln(t) - 1)T}{t}$ 	$R = t$ $S = -\ln(T) + \ln(T \ln(t) - 1)$	$\frac{dS}{dR} = \frac{1}{R}$ 

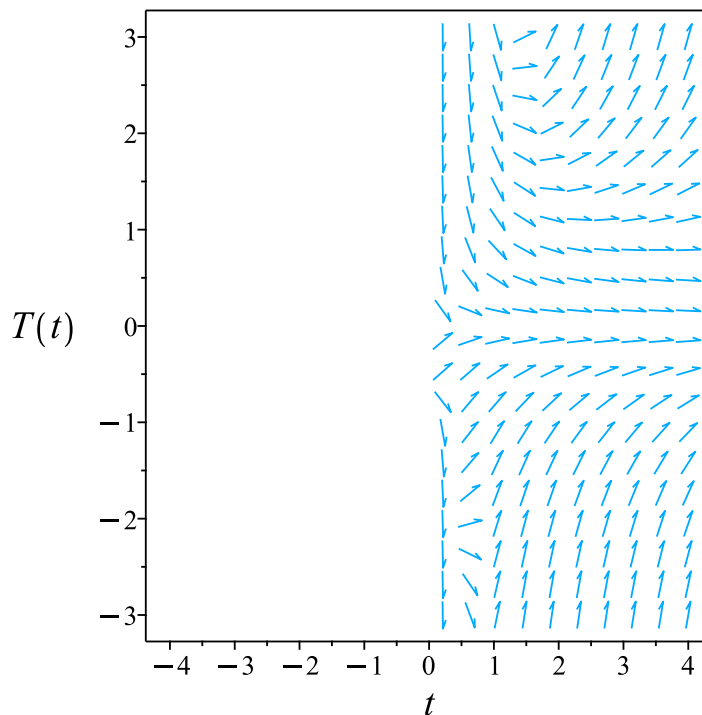


Figure 2.49: Slope field plot
 $(T \ln(t) - 1)T = tT'$

Summary of solutions found

$$T = -\frac{1}{e^{c_2 t} - \ln(t) - 1}$$

Maple step by step solution

Let's solve

$$(T \ln(t) - 1)T = tT'$$

- Highest derivative means the order of the ODE is 1
- Solve for the highest derivative

$$T' = \frac{(T \ln(t) - 1)T}{t}$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
<- Bernoulli successful`

```

Maple dsolve solution

Solving time : 0.003 (sec)

Leaf size : 13

```

dsolve((T(t)*ln(t)-1)*T(t) = t*diff(T(t),t),
        T(t),singsol=all)

```

$$T = \frac{1}{c_1 t + \ln(t) + 1}$$

Mathematica DSolve solution

Solving time : 0.149 (sec)

Leaf size : 20

```

DSolve[{(T[t]*Log[t]-1)*T[t]==t*D[T[t],t],{}}],
        T[t],t,IncludeSingularSolutions->True]

```

$$T(t) \rightarrow \frac{1}{\log(t) + c_1 t + 1}$$

$$T(t) \rightarrow 0$$

2.4.6 problem 6

Solved as first order linear ode	337
Solved as first order Exact ode	339
Maple step by step solution	343
Maple trace	344
Maple dsolve solution	345
Mathematica DSolve solution	345

Internal problem ID [18233]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter IV. Methods of solution: First order equations. section 31. Problems at page 85

Problem number : 6

Date solved : Monday, December 23, 2024 at 09:17:12 PM

CAS classification : [_linear]

Solve

$$y' + y \cos(x) = \frac{\sin(2x)}{2}$$

Solved as first order linear ode

Time used: 0.164 (sec)

In canonical form a linear first order is

$$y' + q(x)y = p(x)$$

Comparing the above to the given ode shows that

$$\begin{aligned} q(x) &= \cos(x) \\ p(x) &= \frac{\sin(2x)}{2} \end{aligned}$$

The integrating factor μ is

$$\begin{aligned} \mu &= e^{\int q dx} \\ &= e^{\int \cos(x) dx} \\ &= e^{\sin(x)} \end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dx}(\mu y) &= \mu p \\ \frac{d}{dx}(\mu y) &= (\mu) \left(\frac{\sin(2x)}{2} \right) \\ \frac{d}{dx}(y e^{\sin(x)}) &= (e^{\sin(x)}) \left(\frac{\sin(2x)}{2} \right) \\ d(y e^{\sin(x)}) &= \left(\frac{\sin(2x) e^{\sin(x)}}{2} \right) dx\end{aligned}$$

Integrating gives

$$\begin{aligned}y e^{\sin(x)} &= \int \frac{\sin(2x) e^{\sin(x)}}{2} dx \\ &= \sin(x) e^{\sin(x)} - e^{\sin(x)} + c_1\end{aligned}$$

Dividing throughout by the integrating factor $e^{\sin(x)}$ gives the final solution

$$y = \sin(x) + e^{-\sin(x)} c_1 - 1$$

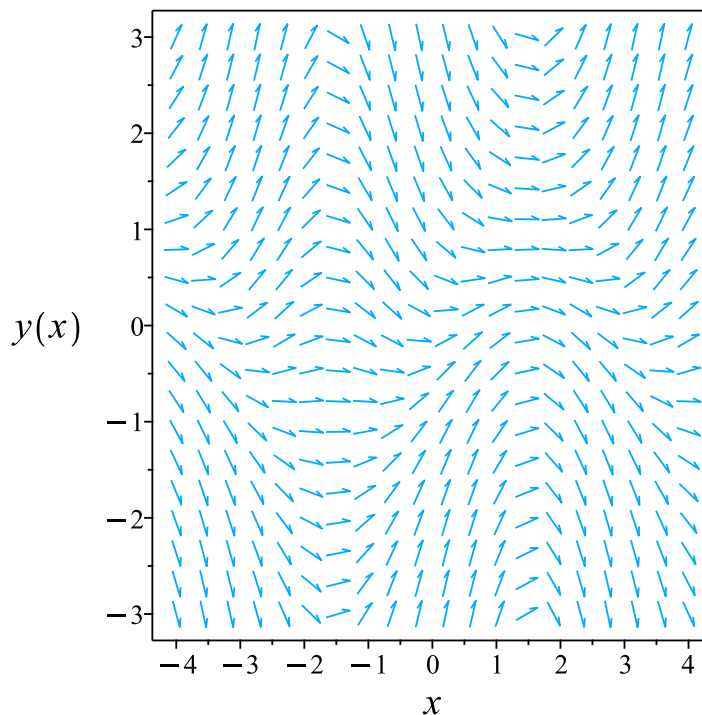


Figure 2.50: Slope field plot
 $y' + y \cos(x) = \frac{\sin(2x)}{2}$

Summary of solutions found

$$y = \sin(x) + e^{-\sin(x)}c_1 - 1$$

Solved as first order Exact ode

Time used: 0.154 (sec)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx}\phi(x, y) = 0$$

Hence

$$\frac{\partial\phi}{\partial x} + \frac{\partial\phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial\phi}{\partial x} &= M \\ \frac{\partial\phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (\text{1A})$$

Therefore

$$\begin{aligned} dy &= \left(-y \cos(x) + \frac{\sin(2x)}{2} \right) dx \\ \left(y \cos(x) - \frac{\sin(2x)}{2} \right) dx + dy &= 0 \end{aligned} \quad (\text{2A})$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(x, y) &= y \cos(x) - \frac{\sin(2x)}{2} \\N(x, y) &= 1\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} \left(y \cos(x) - \frac{\sin(2x)}{2} \right) \\&= \cos(x)\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(1) \\&= 0\end{aligned}$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned}A &= \frac{1}{N} \left(\frac{\partial M}{\partial y} - \frac{\partial N}{\partial x} \right) \\&= 1((\cos(x)) - (0)) \\&= \cos(x)\end{aligned}$$

Since A does not depend on y , then it can be used to find an integrating factor. The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int A \, dx} \\&= e^{\int \cos(x) \, dx}\end{aligned}$$

The result of integrating gives

$$\begin{aligned}\mu &= e^{\sin(x)} \\&= e^{\sin(x)}\end{aligned}$$

M and N are multiplied by this integrating factor, giving new M and new N which are called \overline{M} and \overline{N} for now so not to confuse them with the original M and N .

$$\begin{aligned}\overline{M} &= \mu M \\ &= e^{\sin(x)} \left(y \cos(x) - \frac{\sin(2x)}{2} \right) \\ &= \cos(x) (-\sin(x) + y) e^{\sin(x)}\end{aligned}$$

And

$$\begin{aligned}\overline{N} &= \mu N \\ &= e^{\sin(x)}(1) \\ &= e^{\sin(x)}\end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned}\overline{M} + \overline{N} \frac{dy}{dx} &= 0 \\ (\cos(x) (-\sin(x) + y) e^{\sin(x)}) + (e^{\sin(x)}) \frac{dy}{dx} &= 0\end{aligned}$$

The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = \overline{M} \tag{1}$$

$$\frac{\partial \phi}{\partial y} = \overline{N} \tag{2}$$

Integrating (2) w.r.t. y gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial y} dy &= \int \overline{N} dy \\ \int \frac{\partial \phi}{\partial y} dy &= \int e^{\sin(x)} dy \\ \phi &= y e^{\sin(x)} + f(x)\end{aligned} \tag{3}$$

Where $f(x)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t x gives

$$\frac{\partial \phi}{\partial x} = \cos(x) e^{\sin(x)} y + f'(x) \tag{4}$$

But equation (1) says that $\frac{\partial \phi}{\partial x} = \cos(x)(-\sin(x) + y)e^{\sin(x)}$. Therefore equation (4) becomes

$$\cos(x)(-\sin(x) + y)e^{\sin(x)} = \cos(x)e^{\sin(x)}y + f'(x) \quad (5)$$

Solving equation (5) for $f'(x)$ gives

$$f'(x) = -\cos(x)e^{\sin(x)}\sin(x)$$

Integrating the above w.r.t x gives

$$\begin{aligned} \int f'(x) dx &= \int (-\cos(x)e^{\sin(x)}\sin(x)) dx \\ f(x) &= -\sin(x)e^{\sin(x)} + e^{\sin(x)} + c_1 \end{aligned}$$

Where c_1 is constant of integration. Substituting result found above for $f(x)$ into equation (3) gives ϕ

$$\phi = ye^{\sin(x)} - \sin(x)e^{\sin(x)} + e^{\sin(x)} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into the constant c_1 gives the solution as

$$c_1 = ye^{\sin(x)} - \sin(x)e^{\sin(x)} + e^{\sin(x)}$$

Solving for y gives

$$y = e^{-\sin(x)}(\sin(x)e^{\sin(x)} - e^{\sin(x)} + c_1)$$

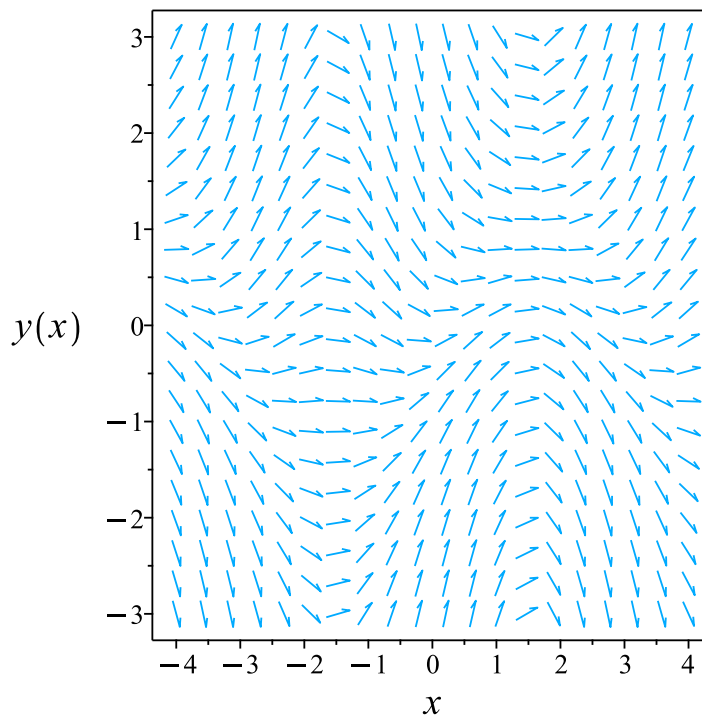


Figure 2.51: Slope field plot
 $y' + y \cos(x) = \frac{\sin(2x)}{2}$

Summary of solutions found

$$y = e^{-\sin(x)} (\sin(x) e^{\sin(x)} - e^{\sin(x)} + c_1)$$

Maple step by step solution

Let's solve

$$y' + y \cos(x) = \frac{\sin(2x)}{2}$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Solve for the highest derivative

$$y' = -y \cos(x) + \frac{\sin(2x)}{2}$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE

$$y' + y \cos(x) = \frac{\sin(2x)}{2}$$

- The ODE is linear; multiply by an integrating factor $\mu(x)$

$$\mu(x) (y' + y \cos(x)) = \frac{\mu(x) \sin(2x)}{2}$$

- Assume the lhs of the ODE is the total derivative $\frac{d}{dx}(y\mu(x))$

$$\mu(x)(y' + y \cos(x)) = y'\mu(x) + y\mu'(x)$$
- Isolate $\mu'(x)$

$$\mu'(x) = \mu(x) \cos(x)$$
- Solve to find the integrating factor

$$\mu(x) = e^{\sin(x)}$$
- Integrate both sides with respect to x

$$\int \left(\frac{d}{dx}(y\mu(x)) \right) dx = \int \frac{\mu(x)\sin(2x)}{2} dx + C1$$
- Evaluate the integral on the lhs

$$y\mu(x) = \int \frac{\mu(x)\sin(2x)}{2} dx + C1$$
- Solve for y

$$y = \frac{\int \frac{\mu(x)\sin(2x)}{2} dx + C1}{\mu(x)}$$
- Substitute $\mu(x) = e^{\sin(x)}$

$$y = \frac{\int \frac{\sin(2x)e^{\sin(x)}}{2} dx + C1}{e^{\sin(x)}}$$
- Evaluate the integrals on the rhs

$$y = \frac{\sin(x)e^{\sin(x)} - e^{\sin(x)} + C1}{e^{\sin(x)}}$$
- Simplify

$$y = \sin(x) + e^{-\sin(x)} C1 - 1$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`

```

Maple dsolve solution

Solving time : 0.002 (sec)

Leaf size : 15

```
dsolve(diff(y(x),x)+y(x)*cos(x) = 1/2*sin(2*x),  
        y(x),singsol=all)
```

$$y = \sin(x) + e^{-\sin(x)}c_1 - 1$$

Mathematica DSolve solution

Solving time : 0.064 (sec)

Leaf size : 18

```
DSolve[{D[y[x],x]+y[x]*Cos[x]==1/2*Sin[2*x],{}},  
        y[x],x,IncludeSingularSolutions->True]
```

$$y(x) \rightarrow \sin(x) + c_1 e^{-\sin(x)} - 1$$

2.4.7 problem 7

Solved as first order Bernoulli ode	346
Maple step by step solution	349
Maple trace	350
Maple dsolve solution	350
Mathematica DSolve solution	350

Internal problem ID [18234]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter IV. Methods of solution: First order equations. section 31. Problems at page 85

Problem number : 7

Date solved : Monday, December 23, 2024 at 09:17:15 PM

CAS classification : [_Bernoulli]

Solve

$$y - \cos(x) y' = y^2 \cos(x) (1 - \sin(x))$$

Solved as first order Bernoulli ode

Time used: 0.448 (sec)

In canonical form, the ODE is

$$\begin{aligned} y' &= F(x, y) \\ &= \frac{y(y \cos(x) \sin(x) - y \cos(x) + 1)}{\cos(x)} \end{aligned}$$

This is a Bernoulli ODE.

$$y' = \left(\frac{1}{\cos(x)} \right) y + \left(\frac{\cos(x) \sin(x) - \cos(x)}{\cos(x)} \right) y^2 \quad (1)$$

The standard Bernoulli ODE has the form

$$y' = f_0(x)y + f_1(x)y^n \quad (2)$$

Comparing this to (1) shows that

$$\begin{aligned} f_0 &= \frac{1}{\cos(x)} \\ f_1 &= \frac{\cos(x) \sin(x) - \cos(x)}{\cos(x)} \end{aligned}$$

The first step is to divide the above equation by y^n which gives

$$\frac{y'}{y^n} = f_0(x)y^{1-n} + f_1(x) \quad (3)$$

The next step is use the substitution $v = y^{1-n}$ in equation (3) which generates a new ODE in $v(x)$ which will be linear and can be easily solved using an integrating factor. Backsubstitution then gives the solution $y(x)$ which is what we want.

This method is now applied to the ODE at hand. Comparing the ODE (1) With (2) Shows that

$$\begin{aligned} f_0(x) &= \frac{1}{\cos(x)} \\ f_1(x) &= \frac{\cos(x)\sin(x) - \cos(x)}{\cos(x)} \\ n &= 2 \end{aligned}$$

Dividing both sides of ODE (1) by $y^n = y^2$ gives

$$y' \frac{1}{y^2} = \frac{1}{\cos(x)y} + \frac{\cos(x)\sin(x) - \cos(x)}{\cos(x)} \quad (4)$$

Let

$$\begin{aligned} v &= y^{1-n} \\ &= \frac{1}{y} \end{aligned} \quad (5)$$

Taking derivative of equation (5) w.r.t x gives

$$v' = -\frac{1}{y^2}y' \quad (6)$$

Substituting equations (5) and (6) into equation (4) gives

$$\begin{aligned} -v'(x) &= \frac{v(x)}{\cos(x)} + \frac{\cos(x)\sin(x) - \cos(x)}{\cos(x)} \\ v' &= -\frac{v}{\cos(x)} - \frac{\cos(x)\sin(x) - \cos(x)}{\cos(x)} \end{aligned} \quad (7)$$

The above now is a linear ODE in $v(x)$ which is now solved.

In canonical form a linear first order is

$$v'(x) + q(x)v(x) = p(x)$$

Comparing the above to the given ode shows that

$$\begin{aligned}q(x) &= \sec(x) \\p(x) &= 1 - \sin(x)\end{aligned}$$

The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int q dx} \\&= e^{\int \sec(x) dx} \\&= \sec(x) + \tan(x)\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dx}(\mu v) &= \mu p \\ \frac{d}{dx}(\mu v) &= (\mu)(1 - \sin(x)) \\ \frac{d}{dx}(v(\sec(x) + \tan(x))) &= (\sec(x) + \tan(x))(1 - \sin(x)) \\ d(v(\sec(x) + \tan(x))) &= ((1 - \sin(x))(\sec(x) + \tan(x))) dx\end{aligned}$$

Integrating gives

$$\begin{aligned}v(\sec(x) + \tan(x)) &= \int (1 - \sin(x))(\sec(x) + \tan(x)) dx \\ &= \sin(x) + c_1\end{aligned}$$

Dividing throughout by the integrating factor $\sec(x) + \tan(x)$ gives the final solution

$$v(x) = \frac{(\sin(x) + c_1)(\cos(x) - \sin(x) + 1)}{\cos(x) + \sin(x) + 1}$$

The substitution $v = y^{1-n}$ is now used to convert the above solution back to y which results in

$$\frac{1}{y} = \frac{(\sin(x) + c_1)(\cos(x) - \sin(x) + 1)}{\cos(x) + \sin(x) + 1}$$

Solving for y gives

$$y = \frac{\cos(x) + \sin(x) + 1}{\cos(x) \sin(x) + \cos(x) c_1 - \sin(x)^2 - c_1 \sin(x) + \sin(x) + c_1}$$

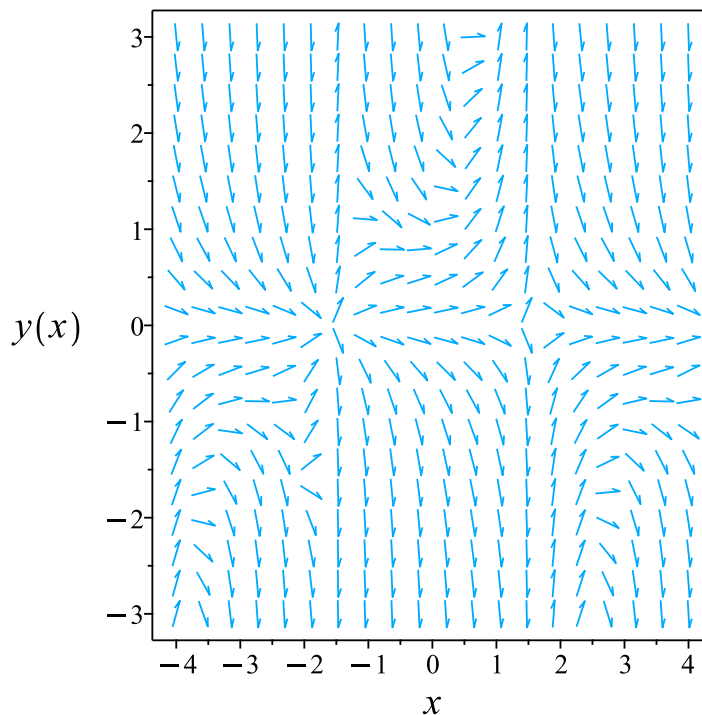


Figure 2.52: Slope field plot
 $y - \cos(x)y' = y^2 \cos(x)(1 - \sin(x))$

Summary of solutions found

$$y = \frac{\cos(x) + \sin(x) + 1}{\cos(x) \sin(x) + \cos(x) c_1 - \sin(x)^2 - c_1 \sin(x) + \sin(x) + c_1}$$

Maple step by step solution

Let's solve

$$y - \cos(x)y' = y^2 \cos(x)(1 - \sin(x))$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Solve for the highest derivative

$$y' = -\frac{-y + y^2 \cos(x)(1 - \sin(x))}{\cos(x)}$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
<- Bernoulli successful`

```

Maple dsolve solution

Solving time : 0.003 (sec)

Leaf size : 27

```

dsolve(y(x)-cos(x)*diff(y(x),x) = y(x)^2*cos(x)*(1-sin(x)),
        y(x),singsol=all)

```

$$y = \frac{\cos(x) + \sin(x) + 1}{(\sin(x) + c_1)(\cos(x) - \sin(x) + 1)}$$

Mathematica DSolve solution

Solving time : 0.372 (sec)

Leaf size : 41

```

DSolve[{y[x]-Cos[x]*D[y[x],x]==y[x]^2*Cos[x]*(1-Sin[x]),{}}],
        y[x],x,IncludeSingularSolutions->True]

```

$$y(x) \rightarrow \frac{e^{2\operatorname{arctanh}(\tan(\frac{x}{2}))}}{\cos(x)e^{2\operatorname{arctanh}(\tan(\frac{x}{2}))} + c_1}$$

$$y(x) \rightarrow 0$$

2.5 Chapter IV. Methods of solution: First order equations. section 32. Problems at page 89

2.5.1	problem 2	352
2.5.2	problem 3	356
2.5.3	problem 4	362
2.5.4	problem 5	365
2.5.5	problem 7	376
2.5.6	problem 8	379

2.5.1 problem 2

Solved as first order ode of type dAlembert	352
Maple step by step solution	354
Maple trace	355
Maple dsolve solution	355
Mathematica DSolve solution	355

Internal problem ID [18235]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter IV. Methods of solution: First order equations. section 32. Problems at page 89

Problem number : 2

Date solved : Monday, December 23, 2024 at 09:17:22 PM

CAS classification : [_rational, _dAlembert]

Solve

$$xy'^2 - y + 2y' = 0$$

Solved as first order ode of type dAlembert

Time used: 0.145 (sec)

Let $p = y'$ the ode becomes

$$x p^2 + 2p - y = 0$$

Solving for y from the above results in

$$y = x p^2 + 2p \tag{1}$$

This has the form

$$y = x f(p) + g(p) \tag{*}$$

Where f, g are functions of $p = y'(x)$. The above ode is dAlembert ode which is now solved.

Taking derivative of (*) w.r.t. x gives

$$\begin{aligned}
 p &= f + (x f' + g') \frac{dp}{dx} \\
 p - f &= (x f' + g') \frac{dp}{dx}
 \end{aligned}
 \tag{2}$$

Comparing the form $y = xf + g$ to (1A) shows that

$$\begin{aligned} f &= p^2 \\ g &= 2p \end{aligned}$$

Hence (2) becomes

$$-p^2 + p = (2xp + 2)p'(x) \quad (2A)$$

The singular solution is found by setting $\frac{dp}{dx} = 0$ in the above which gives

$$-p^2 + p = 0$$

Solving the above for p results in

$$\begin{aligned} p_1 &= 0 \\ p_2 &= 1 \end{aligned}$$

Substituting these in (1A) and keeping singular solution that verifies the ode gives

$$\begin{aligned} y &= 0 \\ y &= x + 2 \end{aligned}$$

The general solution is found when $\frac{dp}{dx} \neq 0$. From eq. (2A). This results in

$$p'(x) = \frac{-p(x)^2 + p(x)}{2p(x)x + 2} \quad (3)$$

Inverting the above ode gives

$$\frac{d}{dp}x(p) = \frac{2x(p)p + 2}{-p^2 + p} \quad (4)$$

This ODE is now solved for $x(p)$. The integrating factor is

$$\begin{aligned} \mu &= e^{\int \frac{2}{p-1} dp} \\ \mu &= (p-1)^2 \\ \mu &= (p-1)^2 \end{aligned} \quad (5)$$

Integrating gives

$$\begin{aligned} x(p) &= \frac{1}{\mu} \left(\int \mu \left(-\frac{2}{p(p-1)} \right) dp + c_1 \right) \\ &= \frac{1}{\mu} \left(\frac{-2p + 2 \ln(p) + c_1}{(p-1)^2} + c_1 \right) \\ &= \frac{-2p + 2 \ln(p) + c_1}{(p-1)^2} \end{aligned} \quad (5)$$

Now we need to eliminate p between the above solution and (1A). The first method is to solve for p from Eq. (1A) and substitute the result into Eq. (5). The Second method is to solve for p from Eq. (5) and substitute the result into (1A).

Eliminating p from the following two equations

$$x = \frac{-2p + 2 \ln(p) + c_1}{(p - 1)^2}$$

$$y = xp^2 + 2p$$

results in

$$p = e^{\text{RootOf}(-xe^{2-Z} + 2xe^{-Z} - 2e^{-Z} + c_1 + 2_Z - x)}$$

Substituting the above into Eq (1A) and simplifying gives

$$y = xe^{2\text{RootOf}(-xe^{2-Z} + 2xe^{-Z} - 2e^{-Z} + c_1 + 2_Z - x)} + 2e^{\text{RootOf}(-xe^{2-Z} + 2xe^{-Z} - 2e^{-Z} + c_1 + 2_Z - x)}$$

Summary of solutions found

$$y = 0$$

$$y = x + 2$$

$$y = xe^{2\text{RootOf}(-xe^{2-Z} + 2xe^{-Z} - 2e^{-Z} + c_1 + 2_Z - x)} + 2e^{\text{RootOf}(-xe^{2-Z} + 2xe^{-Z} - 2e^{-Z} + c_1 + 2_Z - x)}$$

Maple step by step solution

Let's solve

$$xy'^2 - y + 2y' = 0$$

- Highest derivative means the order of the ODE is 1
- y'
- Solve for the highest derivative
- $\left[y' = \frac{-1 + \sqrt{xy+1}}{x}, y' = -\frac{1 + \sqrt{xy+1}}{x} \right]$
- Solve the equation $y' = \frac{-1 + \sqrt{xy+1}}{x}$
- Solve the equation $y' = -\frac{1 + \sqrt{xy+1}}{x}$
- Set of solutions
- $\{workingODE, workingODE\}$

Maple trace

```

`Methods for first order ODEs:
-> Solving 1st order ODE of high degree, 1st attempt
trying 1st order WeierstrassP solution for high degree ODE
trying 1st order WeierstrassPPrime solution for high degree ODE
trying 1st order JacobiSN solution for high degree ODE
trying 1st order ODE linearizable_by_differentiation
trying differential order: 1; missing variables
trying dAlembert
<- dAlembert successful`

```

Maple dsolve solution

Solving time : 0.030 (sec)
 Leaf size : 65

```

dsolve(x*diff(y(x),x)^2-y(x)+2*diff(y(x),x) = 0,
        y(x),singsol=all)

```

$$y = 2xe^{\text{RootOf}(-xe^{2-Z}+2xe^{-Z}-2e^{-Z}+c_1+2_Z-x)} + 2\text{RootOf}(-xe^{2-Z}+2xe^{-Z}-2e^{-Z}+c_1+2_Z-x) + c_1 - x$$

Mathematica DSolve solution

Solving time : 10.342 (sec)
 Leaf size : 50

```

DSolve[{x*D[y[x],x]^2-y[x]+2*D[y[x],x]==0,{}},
        y[x],x,IncludeSingularSolutions->True]

```

$$\text{Solve} \left[\left\{ x = \frac{2 \log(K[1]) - 2K[1]}{(K[1] - 1)^2} + \frac{c_1}{(K[1] - 1)^2}, y(x) = xK[1]^2 + 2K[1] \right\}, \{y(x), K[1]\} \right]$$

2.5.2 problem 3

Maple step by step solution	358
Maple trace	360
Maple dsolve solution	361
Mathematica DSolve solution	361

Internal problem ID [18236]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter IV. Methods of solution: First order equations. section 32. Problems at page 89

Problem number : 3

Date solved : Monday, December 23, 2024 at 09:17:24 PM

CAS classification : [_quadrature]

Solve

$$2y'^3 + y'^2 - y = 0$$

Solving for the derivative gives these ODE's to solve

$$y' = \frac{(-1 + 54y + 6\sqrt{-3y + 81y^2})^{1/3}}{6} + \frac{1}{6(-1 + 54y + 6\sqrt{-3y + 81y^2})^{1/3}} - \frac{1}{6} \quad (1)$$

$$y' = -\frac{(-1 + 54y + 6\sqrt{-3y + 81y^2})^{1/3}}{12} - \frac{1}{12(-1 + 54y + 6\sqrt{-3y + 81y^2})^{1/3}} - \frac{1}{6} + \frac{i\sqrt{3} \left(\frac{(-1+54y+6\sqrt{-3y+81y^2})^{1/3}}{6} - \frac{1}{6(-1+54y+6\sqrt{-3y+81y^2})^{1/3}} \right)}{2} \quad (2)$$

$$y' = -\frac{(-1 + 54y + 6\sqrt{-3y + 81y^2})^{1/3}}{12} - \frac{1}{12(-1 + 54y + 6\sqrt{-3y + 81y^2})^{1/3}} - \frac{1}{6} - \frac{i\sqrt{3} \left(\frac{(-1+54y+6\sqrt{-3y+81y^2})^{1/3}}{6} - \frac{1}{6(-1+54y+6\sqrt{-3y+81y^2})^{1/3}} \right)}{2} \quad (3)$$

Now each of the above is solved separately.

Solving Eq. (1)

Unable to integrate (or intergal too complicated), and since no initial conditions are given, then the result can be written as

$$\int^y \frac{6(-1 + 54\tau + 6\sqrt{81\tau^2 - 3\tau})^{1/3}}{(-1 + 54\tau + 6\sqrt{81\tau^2 - 3\tau})^{2/3} - (-1 + 54\tau + 6\sqrt{81\tau^2 - 3\tau})^{1/3} + 1} d\tau = x + c_1$$

We now need to find the singular solutions, these are found by finding for what values $\left(\frac{(-1+54y+6\sqrt{81y^2-3y})^{1/3}}{6} + \frac{1}{6(-1+54y+6\sqrt{81y^2-3y})^{1/3}} - \frac{1}{6}\right)$ is zero. These give

$$y = \text{RootOf} \left(- \left(-1 + 54_Z + 6\sqrt{3} \sqrt{27_Z^2 - _Z} \right)^{2/3} + \left(-1 + 54_Z + 6\sqrt{3} \sqrt{27_Z^2 - _Z} \right)^{1/3} - 1 \right)$$

Now we go over each such singular solution and check if it verifies the ode itself and any initial conditions given. If it does not then the singular solution will not be used.

The solution $\text{RootOf} \left(- \left(-1 + 54_Z + 6\sqrt{3} \sqrt{27_Z^2 - _Z} \right)^{2/3} + \left(-1 + 54_Z + 6\sqrt{3} \sqrt{27_Z^2 - _Z} \right)^{1/3} - 1 \right)$ will not be used

Solving Eq. (2)

Unable to integrate (or intergal too complicated), and since no initial conditions are given, then the result can be written as

$$\int^y \frac{12(-1 + 54\tau + 6\sqrt{81\tau^2 - 3\tau})^{1/3}}{i(-1 + 54\tau + 6\sqrt{81\tau^2 - 3\tau})^{2/3} \sqrt{3} - (-1 + 54\tau + 6\sqrt{81\tau^2 - 3\tau})^{2/3} - i\sqrt{3} - 2(-1 + 54\tau + 6\sqrt{81\tau^2 - 3\tau})^{1/3}} d\tau = x + c_1$$

We now need to find the singular solutions, these are found by finding for what values

$$\left(- \frac{(-1+54y+6\sqrt{81y^2-3y})^{1/3}}{12} - \frac{1}{12(-1+54y+6\sqrt{81y^2-3y})^{1/3}} - \frac{1}{6} + \frac{i\sqrt{3} \left(\frac{(-1+54y+6\sqrt{81y^2-3y})^{1/3}}{6} - \frac{1}{6(-1+54y+6\sqrt{81y^2-3y})^{1/3}} \right)}{2} \right)$$

is zero. These give

$$y = \text{RootOf} \left(i \left(-1 + 54_Z + 6\sqrt{3} \sqrt{27_Z^2 - _Z} \right)^{2/3} \sqrt{3} - \left(-1 + 54_Z + 6\sqrt{3} \sqrt{27_Z^2 - _Z} \right)^{2/3} - i\sqrt{3} - 2 \left(-1 + 54_Z + 6\sqrt{3} \sqrt{27_Z^2 - _Z} \right)^{1/3} - 1 \right)$$

Now we go over each such singular solution and check if it verifies the ode itself and any initial conditions given. If it does not then the singular solution will not be used.

The solution $\text{RootOf} \left(i \left(-1 + 54_Z + 6\sqrt{3} \sqrt{27_Z^2 - _Z} \right)^{2/3} \sqrt{3} - \left(-1 + 54_Z + 6\sqrt{3} \sqrt{27_Z^2 - _Z} \right)^{2/3} \sqrt{3} + 2 \left(-1 + 54_Z + 6\sqrt{3} \sqrt{27_Z^2 - _Z} \right)^{1/3} + 1 \right)$ will not be used

Solving Eq. (3)

Unable to integrate (or intergal too complicated), and since no initial conditions are given, then the result can be written as

$$\int^y - \frac{12(-1 + 54\tau + 6\sqrt{81\tau^2 - 3\tau})^{1/3}}{i(-1 + 54\tau + 6\sqrt{81\tau^2 - 3\tau})^{2/3} \sqrt{3} + (-1 + 54\tau + 6\sqrt{81\tau^2 - 3\tau})^{2/3} - i\sqrt{3} + 2(-1 + 54\tau + 6\sqrt{81\tau^2 - 3\tau})^{1/3} + 1} d\tau$$

We now need to find the singular solutions, these are found by finding for what values

$$\left(- \frac{(-1+54y+6\sqrt{81y^2-3y})^{1/3}}{12} - \frac{1}{12(-1+54y+6\sqrt{81y^2-3y})^{1/3}} - \frac{1}{6} - \frac{i\sqrt{3} \left(\frac{(-1+54y+6\sqrt{81y^2-3y})^{1/3}}{6} - \frac{1}{6(-1+54y+6\sqrt{81y^2-3y})^{1/3}} \right)}{2} \right)$$

is zero. These give

$$y = \text{RootOf} \left(i \left(-1 + 54_Z + 6\sqrt{3} \sqrt{27_Z^2 - _Z} \right)^{2/3} \sqrt{3} + \left(-1 + 54_Z + 6\sqrt{3} \sqrt{27_Z^2 - _Z} \right)^{2/3} - i\sqrt{3} + 2 \left(-1 + 54_Z + 6\sqrt{3} \sqrt{27_Z^2 - _Z} \right)^{1/3} + 1 \right)$$

Now we go over each such singular solution and check if it verifies the ode itself and any initial conditions given. If it does not then the singular solution will not be used.

The solution $\text{RootOf} \left(i \left(-1 + 54_Z + 6\sqrt{3} \sqrt{27_Z^2 - _Z} \right)^{2/3} \sqrt{3} + \left(-1 + 54_Z + 6\sqrt{3} \sqrt{27_Z^2 - _Z} \right)^{2/3} - i\sqrt{3} + 2 \left(-1 + 54_Z + 6\sqrt{3} \sqrt{27_Z^2 - _Z} \right)^{1/3} + 1 \right)$ will not be used

Maple step by step solution

Let's solve

$$2y'^3 + y'^2 - y = 0$$

- Highest derivative means the order of the ODE is 1

y'

- Solve for the highest derivative

$$\left[y' = \frac{(-1+54y+6\sqrt{-3y+81y^2})^{1/3}}{6} + \frac{1}{6(-1+54y+6\sqrt{-3y+81y^2})^{1/3}} - \frac{1}{6}; y' = -\frac{(-1+54y+6\sqrt{-3y+81y^2})^{1/3}}{12} - \frac{1}{12} \right]$$

□ Solve the equation $y' = \frac{(-1+54y+6\sqrt{-3y+81y^2})^{1/3}}{6} + \frac{1}{6(-1+54y+6\sqrt{-3y+81y^2})^{1/3}} - \frac{1}{6}$

○ Separate variables

$$\frac{y'}{\frac{(-1+54y+6\sqrt{-3y+81y^2})^{1/3}}{6} + \frac{1}{6(-1+54y+6\sqrt{-3y+81y^2})^{1/3}} - \frac{1}{6}} = 1$$

○ Integrate both sides with respect to x

$$\int \frac{y'}{\frac{(-1+54y+6\sqrt{-3y+81y^2})^{1/3}}{6} + \frac{1}{6(-1+54y+6\sqrt{-3y+81y^2})^{1/3}} - \frac{1}{6}} dx = \int 1 dx + C_1$$

○ Cannot compute integral

$$\int \frac{y'}{\frac{(-1+54y+6\sqrt{-3y+81y^2})^{1/3}}{6} + \frac{1}{6(-1+54y+6\sqrt{-3y+81y^2})^{1/3}} - \frac{1}{6}} dx = x + C_1$$

□ Solve the equation $y' = -\frac{(-1+54y+6\sqrt{-3y+81y^2})^{1/3}}{12} - \frac{1}{12(-1+54y+6\sqrt{-3y+81y^2})^{1/3}} - \frac{1}{6} - \frac{I\sqrt{3} \left(\frac{-1+54y+6\sqrt{-3y+81y^2}}{6} \right)}{2}$

○ Separate variables

$$\frac{y'}{-\frac{(-1+54y+6\sqrt{-3y+81y^2})^{1/3}}{12} - \frac{1}{12(-1+54y+6\sqrt{-3y+81y^2})^{1/3}} - \frac{1}{6} - \frac{I\sqrt{3} \left(\frac{(-1+54y+6\sqrt{-3y+81y^2})^{1/3}}{6} - \frac{1}{6(-1+54y+6\sqrt{-3y+81y^2})^{1/3}} \right)}{2}} = 1$$

○ Integrate both sides with respect to x

$$\int \frac{y'}{-\frac{(-1+54y+6\sqrt{-3y+81y^2})^{1/3}}{12} - \frac{1}{12(-1+54y+6\sqrt{-3y+81y^2})^{1/3}} - \frac{1}{6} - \frac{I\sqrt{3} \left(\frac{(-1+54y+6\sqrt{-3y+81y^2})^{1/3}}{6} - \frac{1}{6(-1+54y+6\sqrt{-3y+81y^2})^{1/3}} \right)}{2}} dx = \int 1 dx + C_1$$

○ Cannot compute integral

$$\int \frac{y'}{-\frac{(-1+54y+6\sqrt{-3y+81y^2})^{1/3}}{12} - \frac{1}{12(-1+54y+6\sqrt{-3y+81y^2})^{1/3}} - \frac{1}{6} - \frac{I\sqrt{3} \left(\frac{(-1+54y+6\sqrt{-3y+81y^2})^{1/3}}{6} - \frac{1}{6(-1+54y+6\sqrt{-3y+81y^2})^{1/3}} \right)}{2}} dx = x + C_1$$

□ Solve the equation $y' = -\frac{(-1+54y+6\sqrt{-3y+81y^2})^{1/3}}{12} - \frac{1}{12(-1+54y+6\sqrt{-3y+81y^2})^{1/3}} - \frac{1}{6} + \frac{I\sqrt{3}}{6} \left(\frac{(-1+54y+6\sqrt{-3y+81y^2})^{1/3}}{6} - \frac{1}{6(-1+54y+6\sqrt{-3y+81y^2})^{1/3}} \right)$

- Separate variables

$$\frac{y'}{\left(\frac{(-1+54y+6\sqrt{-3y+81y^2})^{1/3}}{12} - \frac{1}{12(-1+54y+6\sqrt{-3y+81y^2})^{1/3}} - \frac{1}{6} + \frac{I\sqrt{3}}{6} \left(\frac{(-1+54y+6\sqrt{-3y+81y^2})^{1/3}}{6} - \frac{1}{6(-1+54y+6\sqrt{-3y+81y^2})^{1/3}} \right) \right)}$$

- Integrate both sides with respect to x

$$\int \frac{y'}{\left(\frac{(-1+54y+6\sqrt{-3y+81y^2})^{1/3}}{12} - \frac{1}{12(-1+54y+6\sqrt{-3y+81y^2})^{1/3}} - \frac{1}{6} + \frac{I\sqrt{3}}{6} \left(\frac{(-1+54y+6\sqrt{-3y+81y^2})^{1/3}}{6} - \frac{1}{6(-1+54y+6\sqrt{-3y+81y^2})^{1/3}} \right) \right)} dx$$

- Cannot compute integral

$$\int \frac{y'}{\left(\frac{(-1+54y+6\sqrt{-3y+81y^2})^{1/3}}{12} - \frac{1}{12(-1+54y+6\sqrt{-3y+81y^2})^{1/3}} - \frac{1}{6} + \frac{I\sqrt{3}}{6} \left(\frac{(-1+54y+6\sqrt{-3y+81y^2})^{1/3}}{6} - \frac{1}{6(-1+54y+6\sqrt{-3y+81y^2})^{1/3}} \right) \right)} dx$$

- Set of solutions

$$\left\{ \int \frac{y'}{\left(\frac{(-1+54y+6\sqrt{-3y+81y^2})^{1/3}}{6} + \frac{1}{6(-1+54y+6\sqrt{-3y+81y^2})^{1/3}} - \frac{1}{6} \right)} dx = x + C1, \int \frac{y'}{\left(\frac{(-1+54y+6\sqrt{-3y+81y^2})^{1/3}}{12} - \frac{1}{12(-1+54y+6\sqrt{-3y+81y^2})^{1/3}} - \frac{1}{6} + \frac{I\sqrt{3}}{6} \left(\frac{(-1+54y+6\sqrt{-3y+81y^2})^{1/3}}{6} - \frac{1}{6(-1+54y+6\sqrt{-3y+81y^2})^{1/3}} \right) \right)} dx \right.$$

Maple trace

```

Methods for first order ODEs:
*** Sublevel 2 ***
Methods for first order ODEs:
-> Solving 1st order ODE of high degree, 1st attempt
trying 1st order WeierstrassP solution for high degree ODE
trying 1st order WeierstrassPPrime solution for high degree ODE
trying 1st order JacobiSN solution for high degree ODE
trying 1st order ODE linearizable_by_differentiation
trying differential order: 1; missing variables
    
```

```
<- differential order: 1; missing x successful`
```

Maple dsolve solution

Solving time : 0.019 (sec)

Leaf size : 387

```
dsolve(2*diff(y(x),x)^3+diff(y(x),x)^2-y(x) = 0,
      y(x),singsol=all)
```

$y = 0$

$$-6\sqrt{3} \left(\int^y \frac{(18\sqrt{27a^2 - a} + (54a - 1)\sqrt{3})^{1/3}}{3^{2/3} - \sqrt{3} (18\sqrt{27a^2 - a} + (54a - 1)\sqrt{3})^{1/3} + 3^{1/3} (18\sqrt{27a^2 - a} + (54a - 1)\sqrt{3})^{1/3}} \right.$$

$$-72 \left(\int^y \frac{(18\sqrt{27a^2 - a} + (54a - 1)\sqrt{3})^{1/3}}{(i3^{5/6} + 3^{1/3} - 2 \cdot 3^{1/6} (18\sqrt{27a^2 - a} + (54a - 1)\sqrt{3})^{1/3}) (3^{1/3} + 3^{1/6} (18\sqrt{27a^2 - a} + (54a - 1)\sqrt{3})^{1/3})} d_a \right) + (-$$

$$-72 \left(\int^y \frac{(18\sqrt{27a^2 - a} + (54a - 1)\sqrt{3})^{1/3}}{(i3^{5/6} - 3^{1/3} + 2 \cdot 3^{1/6} (18\sqrt{27a^2 - a} + (54a - 1)\sqrt{3})^{1/3}) (3^{1/3} + 3^{1/6} (18\sqrt{27a^2 - a} + (54a - 1)\sqrt{3})^{1/3})} d_a \right) + ($$

$$\frac{\sqrt{3} + 3i}{-\sqrt{3} + 3i}$$

Mathematica DSolve solution

Solving time : 0.0 (sec)

Leaf size : 0

```
DSolve[{2*D[y[x],x]^3+D[y[x],x]^2-y[x]==0,{}},
      y[x],x,IncludeSingularSolutions->True]
```

Timed out

2.5.3 problem 4

Solved as first order quadrature ode	362
Maple step by step solution	363
Maple trace	364
Maple dsolve solution	364
Mathematica DSolve solution	364

Internal problem ID [18237]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter IV. Methods of solution: First order equations. section 32. Problems at page 89

Problem number : 4

Date solved : Monday, December 23, 2024 at 09:18:40 PM

CAS classification : [_quadrature]

Solve

$$y' = e^{z-y'}$$

Solved as first order quadrature ode

Time used: 0.092 (sec)

Since the ode has the form $y' = f(z)$, then we only need to integrate $f(z)$.

$$\int dy = \int \text{LambertW}(e^z) dz$$

$$y = \frac{\text{LambertW}(e^z)^2}{2} + \text{LambertW}(e^z) + c_1$$

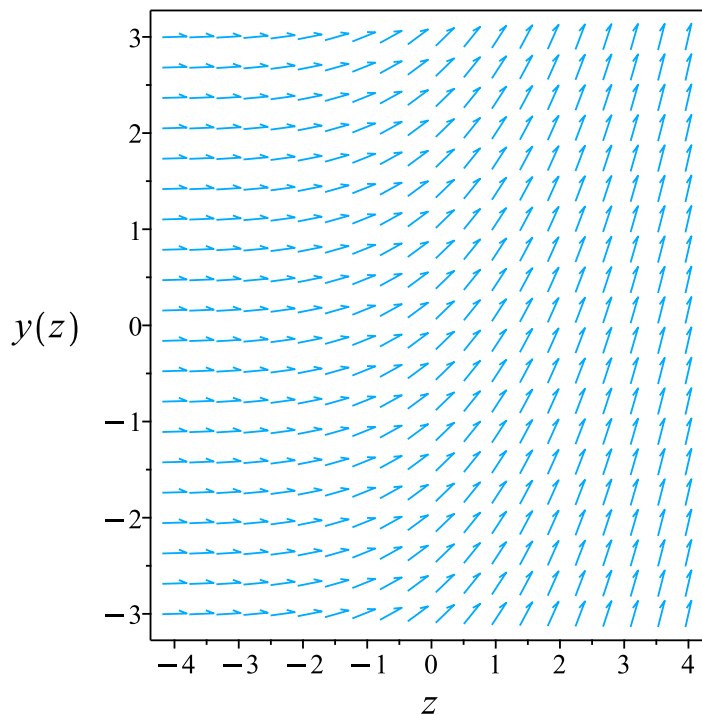


Figure 2.53: Slope field plot
 $y' = e^{z-y'}$

Summary of solutions found

$$y = \frac{\text{LambertW}(e^z)^2}{2} + \text{LambertW}(e^z) + c_1$$

Maple step by step solution

Let's solve

$$y' = e^{z-y'}$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Solve for the highest derivative

$$y' = \text{LambertW}(e^z)$$

- Integrate both sides with respect to z

$$\int y' dz = \int \text{LambertW}(e^z) dz + C1$$

- Evaluate integral

$$y = \frac{\text{LambertW}(e^z)^2}{2} + \text{LambertW}(e^z) + C1$$

- Solve for y

$$y = \frac{\text{LambertW}(e^z)^2}{2} + \text{LambertW}(e^z) + C1$$

Maple trace

```

`Methods for first order ODEs:
-> Solving 1st order ODE of high degree, 1st attempt
trying 1st order WeierstrassP solution for high degree ODE
trying 1st order WeierstrassPPrime solution for high degree ODE
trying 1st order JacobiSN solution for high degree ODE
trying 1st order ODE linearizable_by_differentiation
trying differential order: 1; missing variables
<- differential order: 1; missing y(x) successful`

```

Maple dsolve solution

Solving time : 0.207 (sec)

Leaf size : 16

```

dsolve(diff(y(z),z) = exp(z-diff(y(z),z)),
        y(z),singsol=all)

```

$$y = \frac{\text{LambertW}(e^z)^2}{2} + \text{LambertW}(e^z) + c_1$$

Mathematica DSolve solution

Solving time : 0.053 (sec)

Leaf size : 22

```

DSolve[{D[y[z],z]==Exp[z-D[y[z],z]],{}}
        y[z],z,IncludeSingularSolutions->True]

```

$$y(z) \rightarrow \frac{1}{2}W(e^z)^2 + W(e^z) + c_1$$

2.5.4 problem 5

Solved as first order isobaric ode	365
Solved using Lie symmetry for first order ode	368
Maple step by step solution	373
Maple trace	374
Maple dsolve solution	374
Mathematica DSolve solution	375

Internal problem ID [18238]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter IV. Methods of solution: First order equations. section 32. Problems at page 89

Problem number : 5

Date solved : Monday, December 23, 2024 at 09:18:41 PM

CAS classification : [[_homogeneous, 'class G']]

Solve

$$\sqrt{t^2 + T} = T'$$

Solved as first order isobaric ode

Time used: 2.435 (sec)

Solving for T' gives

$$T' = \sqrt{t^2 + T} \tag{1}$$

Each of the above ode's is now solved An ode $T' = f(t, T)$ is isobaric if

$$f(tt, t^m T) = t^{m-1} f(t, T) \tag{1}$$

Where here

$$f(t, T) = \sqrt{t^2 + T} \tag{2}$$

m is the order of isobaric. Substituting (2) into (1) and solving for m gives

$$m = 2$$

Since the ode is isobaric of order $m = 2$, then the substitution

$$\begin{aligned} T &= ut^m \\ &= ut^2 \end{aligned}$$

Converts the ODE to a separable in $u(t)$. Performing this substitution gives

$$2tu(t) + t^2u'(t) = \sqrt{t^2 + t^2u(t)}$$

The ode $u'(t) = \frac{\sqrt{1+u(t)}-2u(t)}{t}$ is separable as it can be written as

$$\begin{aligned} u'(t) &= \frac{\sqrt{1+u(t)}-2u(t)}{t} \\ &= f(t)g(u) \end{aligned}$$

Where

$$\begin{aligned} f(t) &= \frac{1}{t} \\ g(u) &= \sqrt{u+1} - 2u \end{aligned}$$

Integrating gives

$$\begin{aligned} \int \frac{1}{g(u)} du &= \int f(t) dt \\ \int \frac{1}{\sqrt{u+1} - 2u} du &= \int \frac{1}{t} dt \\ -\frac{\ln(2u(t) - \sqrt{1+u(t)})}{2} + \frac{\sqrt{17} \operatorname{arctanh}\left(\frac{(4\sqrt{1+u(t)}-1)\sqrt{17}}{17}\right)}{17} &= \ln(t) + c_1 \end{aligned}$$

We now need to find the singular solutions, these are found by finding for what values $g(u)$ is zero, since we had to divide by this above. Solving $g(u) = 0$ or $\sqrt{u+1} - 2u = 0$ for $u(t)$ gives

$$u(t) = \frac{1}{8} + \frac{\sqrt{17}}{8}$$

Now we go over each such singular solution and check if it verifies the ode itself and any initial conditions given. If it does not then the singular solution will not be used.

Therefore the solutions found are

$$\begin{aligned} -\frac{\ln(2u(t) - \sqrt{1+u(t)})}{2} + \frac{\sqrt{17} \operatorname{arctanh}\left(\frac{(4\sqrt{1+u(t)}-1)\sqrt{17}}{17}\right)}{17} &= \ln(t) + c_1 \\ u(t) &= \frac{1}{8} + \frac{\sqrt{17}}{8} \end{aligned}$$

Converting $-\frac{\ln(2u(t)-\sqrt{1+u(t)})}{2} + \frac{\sqrt{17} \operatorname{arctanh}\left(\frac{(4\sqrt{1+u(t)}-1)\sqrt{17}}{17}\right)}{17} = \ln(t) + c_1$ back to T gives

$$-\frac{\ln\left(\frac{2T}{t^2} - \sqrt{1 + \frac{T}{t^2}}\right)}{2} + \frac{\sqrt{17} \operatorname{arctanh}\left(\frac{(4\sqrt{1 + \frac{T}{t^2}} - 1)\sqrt{17}}{17}\right)}{17} = \ln(t) + c_1$$

Converting $u(t) = \frac{1}{8} + \frac{\sqrt{17}}{8}$ back to T gives

$$\frac{T}{t^2} = \frac{1}{8} + \frac{\sqrt{17}}{8}$$

Solving for T gives

$$-\frac{\ln\left(\frac{2T}{t^2} - \sqrt{1 + \frac{T}{t^2}}\right)}{2} + \frac{\sqrt{17} \operatorname{arctanh}\left(\frac{(4\sqrt{1 + \frac{T}{t^2}} - 1)\sqrt{17}}{17}\right)}{17} = \ln(t) + c_1$$

$$T = \frac{(1 + \sqrt{17})t^2}{8}$$

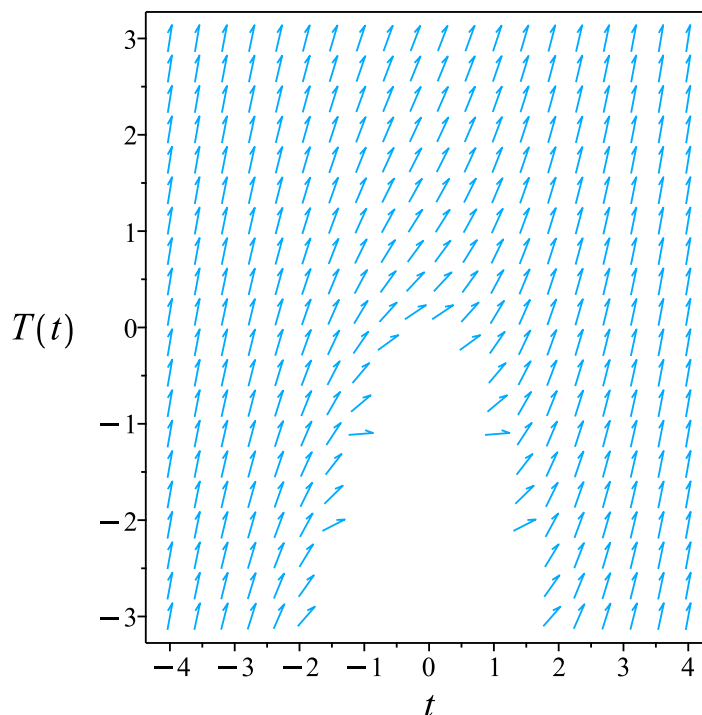


Figure 2.54: Slope field plot
 $\sqrt{t^2 + T} = T'$

Summary of solutions found

$$-\frac{\ln\left(\frac{2T}{t^2} - \sqrt{1 + \frac{T}{t^2}}\right)}{2} + \frac{\sqrt{17} \operatorname{arctanh}\left(\frac{(4\sqrt{1 + \frac{T}{t^2}} - 1)\sqrt{17}}{17}\right)}{17} = \ln(t) + c_1$$

$$T = \frac{(1 + \sqrt{17})t^2}{8}$$

Solved using Lie symmetry for first order ode

Time used: 1.726 (sec)

Writing the ode as

$$T' = \sqrt{t^2 + T}$$

$$T' = \omega(t, T)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_t + \omega(\eta_T - \xi_t) - \omega^2 \xi_T - \omega_t \xi - \omega_T \eta = 0 \quad (\text{A})$$

To determine ξ, η then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$\xi = Ta_3 + ta_2 + a_1 \quad (\text{1E})$$

$$\eta = Tb_3 + tb_2 + b_1 \quad (\text{2E})$$

Where the unknown coefficients are

$$\{a_1, a_2, a_3, b_1, b_2, b_3\}$$

Substituting equations (1E,2E) and ω into (A) gives

$$b_2 + \sqrt{t^2 + T}(b_3 - a_2) - (t^2 + T)a_3 - \frac{t(Ta_3 + ta_2 + a_1)}{\sqrt{t^2 + T}} - \frac{Tb_3 + tb_2 + b_1}{2\sqrt{t^2 + T}} = 0 \quad (\text{5E})$$

Putting the above in normal form gives

$$\frac{2\sqrt{t^2 + T}t^2a_3 + 2\sqrt{t^2 + T}Ta_3 + 2Tta_3 + 4t^2a_2 - 2t^2b_3 - 2b_2\sqrt{t^2 + T} + 2Ta_2 - Tb_3 + 2ta_1 + tb_2 + b_1}{2\sqrt{t^2 + T}} = 0$$

Setting the numerator to zero gives

$$\begin{aligned} & -2\sqrt{t^2 + T}t^2a_3 - 2\sqrt{t^2 + T}Ta_3 - 2Tta_3 - 4t^2a_2 + 2t^2b_3 \\ & + 2b_2\sqrt{t^2 + T} - 2Ta_2 + Tb_3 - 2ta_1 - tb_2 - b_1 = 0 \end{aligned} \quad (6E)$$

Simplifying the above gives

$$\begin{aligned} & -2\sqrt{t^2 + T}t^2a_3 - 2(t^2 + T)a_2 + 2(t^2 + T)b_3 - 2\sqrt{t^2 + T}Ta_3 \\ & - 2Tta_3 - 2t^2a_2 + 2b_2\sqrt{t^2 + T} - Tb_3 - 2ta_1 - tb_2 - b_1 = 0 \end{aligned} \quad (6E)$$

Since the PDE has radicals, simplifying gives

$$\begin{aligned} & -2\sqrt{t^2 + T}t^2a_3 - 2\sqrt{t^2 + T}Ta_3 - 2Tta_3 - 4t^2a_2 + 2t^2b_3 \\ & + 2b_2\sqrt{t^2 + T} - 2Ta_2 + Tb_3 - 2ta_1 - tb_2 - b_1 = 0 \end{aligned}$$

Looking at the above PDE shows the following are all the terms with $\{T, t\}$ in them.

$$\{T, t, \sqrt{t^2 + T}\}$$

The following substitution is now made to be able to collect on all terms with $\{T, t\}$ in them

$$\{T = v_1, t = v_2, \sqrt{t^2 + T} = v_3\}$$

The above PDE (6E) now becomes

$$\begin{aligned} & -2v_3v_2^2a_3 - 4v_2^2a_2 - 2v_1v_2a_3 - 2v_3v_1a_3 + 2v_2^2b_3 \\ & - 2v_2a_1 - 2v_1a_2 - v_2b_2 + 2b_2v_3 + v_1b_3 - b_1 = 0 \end{aligned} \quad (7E)$$

Collecting the above on the terms v_i introduced, and these are

$$\{v_1, v_2, v_3\}$$

Equation (7E) now becomes

$$\begin{aligned} & -2v_1v_2a_3 - 2v_3v_1a_3 + (-2a_2 + b_3)v_1 - 2v_3v_2^2a_3 \\ & + (-4a_2 + 2b_3)v_2^2 + (-2a_1 - b_2)v_2 + 2b_2v_3 - b_1 = 0 \end{aligned} \quad (8E)$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$\begin{aligned} -2a_3 &= 0 \\ -b_1 &= 0 \\ 2b_2 &= 0 \\ -2a_1 - b_2 &= 0 \\ -4a_2 + 2b_3 &= 0 \\ -2a_2 + b_3 &= 0 \end{aligned}$$

Solving the above equations for the unknowns gives

$$\begin{aligned} a_1 &= 0 \\ a_2 &= a_2 \\ a_3 &= 0 \\ b_1 &= 0 \\ b_2 &= 0 \\ b_3 &= 2a_2 \end{aligned}$$

Substituting the above solution in the ansatz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$\begin{aligned} \xi &= t \\ \eta &= 2T \end{aligned}$$

Shifting is now applied to make $\xi = 0$ in order to simplify the rest of the computation

$$\begin{aligned} \eta &= \eta - \omega(t, T) \xi \\ &= 2T - \left(\sqrt{t^2 + T} \right) (t) \\ &= -\sqrt{t^2 + T} t + 2T \\ \xi &= 0 \end{aligned}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(t, T) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dt}{\xi} = \frac{dT}{\eta} = dS \tag{1}$$

The above comes from the requirements that $(\xi \frac{\partial}{\partial t} + \eta \frac{\partial}{\partial T}) S(t, T) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = t$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{-\sqrt{t^2 + T}t + 2T} dy \end{aligned}$$

Which results in

$$S = -\frac{\ln(\sqrt{t^2 + T}t + 2T)}{4} - \frac{\sqrt{17} \operatorname{arctanh}\left(\frac{(4\sqrt{t^2 + T} + t)\sqrt{17}}{17t}\right)}{34} + \frac{\ln(-\sqrt{t^2 + T}t + 2T)}{4} - \frac{\sqrt{17} \operatorname{arctanh}\left(\frac{(4\sqrt{t^2 + T} - t)\sqrt{17}}{17t}\right)}{34}$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_t + \omega(t, T)S_T}{R_t + \omega(t, T)R_T} \quad (2)$$

Where in the above R_t, R_T, S_t, S_T are all partial derivatives and $\omega(t, T)$ is the right hand side of the original ode given by

$$\omega(t, T) = \sqrt{t^2 + T}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_t &= 1 \\ R_T &= 0 \\ S_t &= \frac{t^6 + 2Tt^4 - 3T^2t^2 - 4T^3}{(\sqrt{t^2 + T}t - 2T)^2 (\sqrt{t^2 + T}t + 2T) \sqrt{t^2 + T}} \\ S_T &= \frac{(t + 2\sqrt{t^2 + T})T + t^3}{(-t^4 - Tt^2 + 4T^2) \sqrt{t^2 + T}} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = 0 \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for t, T in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = 0$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S .

Since the ode has the form $\frac{d}{dR}S(R) = f(R)$, then we only need to integrate $f(R)$.

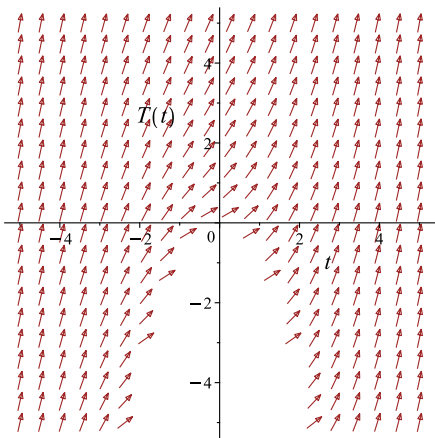
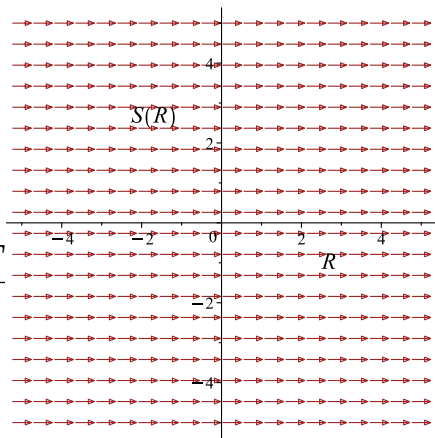
$$\int dS = \int 0 dR + c_2$$

$$S(R) = c_2$$

To complete the solution, we just need to transform the above back to t, T coordinates. This results in

$$-\frac{\ln(\sqrt{t^2 + T}t + 2T)}{4} - \frac{\sqrt{17} \operatorname{arctanh}\left(\frac{(4\sqrt{t^2 + T} + t)\sqrt{17}}{17t}\right)}{34} + \frac{\ln(-\sqrt{t^2 + T}t + 2T)}{4} + \frac{\sqrt{17} \operatorname{arctanh}\left(\frac{(t - 4\sqrt{t^2 + T})\sqrt{17}}{17t}\right)}{34}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in t, T coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dT}{dt} = \sqrt{t^2 + T}$ 	$R = t$ $S = -\frac{\ln(\sqrt{t^2 + T}t + 2T)}{4}$	$\frac{dS}{dR} = 0$ 

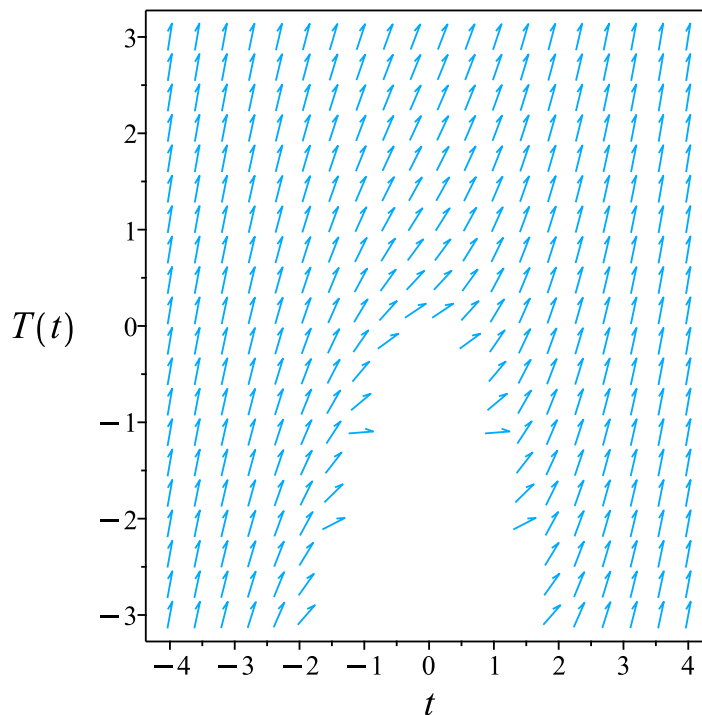


Figure 2.55: Slope field plot
 $\sqrt{t^2 + T} = T'$

Summary of solutions found

$$\begin{aligned}
 & -\frac{\ln(\sqrt{t^2 + T}t + 2T)}{4} - \frac{\sqrt{17} \operatorname{arctanh}\left(\frac{(4\sqrt{t^2 + T} + t)\sqrt{17}}{17t}\right)}{34} \\
 & + \frac{\ln(-\sqrt{t^2 + T}t + 2T)}{4} + \frac{\sqrt{17} \operatorname{arctanh}\left(\frac{(t - 4\sqrt{t^2 + T})\sqrt{17}}{17t}\right)}{34} \\
 & + \frac{\ln(-t^4 - Tt^2 + 4T^2)}{4} - \frac{\sqrt{17} \operatorname{arctanh}\left(\frac{(-t^2 + 8T)\sqrt{17}}{17t^2}\right)}{34} = c_2
 \end{aligned}$$

Maple step by step solution

Let's solve

$$\sqrt{t^2 + T} = T'$$

- Highest derivative means the order of the ODE is 1
 T'
- Solve for the highest derivative

$$T' = \sqrt{t^2 + T}$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying homogeneous types:
trying homogeneous G
1st order, trying the canonical coordinates of the invariance group
<- 1st order, canonical coordinates successful
<- homogeneous successful`

```

Maple dsolve solution

Solving time : 0.005 (sec)

Leaf size : 136

```

dsolve((t^2+T(t))^(1/2) = diff(T(t),t),
       T(t),singsol=all)

```

$$\begin{aligned}
& -17 \ln(-t^4 - Tt^2 + 4T^2) - 17 \ln(-\sqrt{t^2 + T}t + 2T) + 17 \ln(\sqrt{t^2 + T}t + 2T) \\
& + \left(2 \operatorname{arctanh} \left(\frac{(4\sqrt{t^2 + T} + t)\sqrt{17}}{17t} \right) - 2 \operatorname{arctanh} \left(\frac{(t - 4\sqrt{t^2 + T})\sqrt{17}}{17t} \right) \right. \\
& \left. - 2 \operatorname{arctanh} \left(\frac{(t^2 - 8T)\sqrt{17}}{17t^2} \right) \right) \sqrt{17} - c_1 = 0
\end{aligned}$$

Mathematica DSolve solution

Solving time : 0.256 (sec)

Leaf size : 135

```
DSolve[{Sqrt[t^2+T[t]]==D[T[t],t],{}}],
      T[t],t,IncludeSingularSolutions->True]
```

$$\text{Solve} \left[\begin{aligned} & \frac{1}{34} \left(-34 \log \left(\sqrt{t^2 + T(t)} - t \right) \right. \\ & - \left(\sqrt{17} - 17 \right) \log \left(2 \left(\sqrt{17} - 4 \right) t \sqrt{t^2 + T(t)} - 2 \left(\sqrt{17} - 4 \right) t^2 - \left(\sqrt{17} - 3 \right) T(t) \right) \\ & \left. + \left(17 + \sqrt{17} \right) \log \left(2 \left(4 + \sqrt{17} \right) t \sqrt{t^2 + T(t)} - 2 \left(4 + \sqrt{17} \right) t^2 - \left(3 + \sqrt{17} \right) T(t) \right) \right] = c_1, T(t) \end{aligned} \right]$$

2.5.5 problem 7

Maple step by step solution	377
Maple trace	377
Maple dsolve solution	378
Mathematica DSolve solution	378

Internal problem ID [18239]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter IV. Methods of solution: First order equations. section 32. Problems at page 89

Problem number : 7

Date solved : Monday, December 23, 2024 at 09:18:46 PM

CAS classification : [_quadrature]

Solve

$$(x^2 - 1) y'^2 = 1$$

Solving for the derivative gives these ODE's to solve

$$y' = \frac{1}{\sqrt{x^2 - 1}} \tag{1}$$

$$y' = -\frac{1}{\sqrt{x^2 - 1}} \tag{2}$$

Now each of the above is solved separately.

Solving Eq. (1)

Since the ode has the form $y' = f(x)$, then we only need to integrate $f(x)$.

$$\begin{aligned} \int dy &= \int \frac{1}{\sqrt{x^2 - 1}} dx \\ y &= \ln \left(x + \sqrt{x^2 - 1} \right) + c_1 \end{aligned}$$

Solving Eq. (2)

Since the ode has the form $y' = f(x)$, then we only need to integrate $f(x)$.

$$\begin{aligned} \int dy &= \int -\frac{1}{\sqrt{x^2 - 1}} dx \\ y &= -\ln \left(x + \sqrt{x^2 - 1} \right) + c_2 \end{aligned}$$

Maple step by step solution

Let's solve

$$(x^2 - 1)y'^2 = 1$$

- Highest derivative means the order of the ODE is 1

y'

- Solve for the highest derivative

$$\left[y' = \frac{1}{\sqrt{x^2-1}}, y' = -\frac{1}{\sqrt{x^2-1}} \right]$$

- Solve the equation $y' = \frac{1}{\sqrt{x^2-1}}$

- Integrate both sides with respect to x

$$\int y' dx = \int \frac{1}{\sqrt{x^2-1}} dx + _C1$$

- Evaluate integral

$$y = \ln(x + \sqrt{x^2 - 1}) + _C1$$

- Solve for y

$$y = \ln(x + \sqrt{x^2 - 1}) + _C1$$

- Solve the equation $y' = -\frac{1}{\sqrt{x^2-1}}$

- Integrate both sides with respect to x

$$\int y' dx = \int -\frac{1}{\sqrt{x^2-1}} dx + _C1$$

- Evaluate integral

$$y = -\ln(x + \sqrt{x^2 - 1}) + _C1$$

- Solve for y

$$y = -\ln(x + \sqrt{x^2 - 1}) + _C1$$

- Set of solutions

$$\{y = -\ln(x + \sqrt{x^2 - 1}) + C1, y = \ln(x + \sqrt{x^2 - 1}) + C1\}$$

Maple trace

```
`Methods for first order ODEs:
```

```
-> Solving 1st order ODE of high degree, 1st attempt
```

```
trying 1st order WeierstrassP solution for high degree ODE
```

```
trying 1st order WeierstrassPPrime solution for high degree ODE
```

```
trying 1st order JacobiSN solution for high degree ODE
```

```
trying 1st order ODE linearizable_by_differentiation
```

```
trying differential order: 1; missing variables
```

```
<- differential order: 1; missing y(x) successful`
```

Maple dsolve solution

Solving time : 0.016 (sec)

Leaf size : 33

```
dsolve((x^2-1)*diff(y(x),x)^2 = 1,  
        y(x),singsol=all)
```

$$y = \ln(x + \sqrt{x^2 - 1}) + c_1$$

$$y = -\ln(x + \sqrt{x^2 - 1}) + c_1$$

Mathematica DSolve solution

Solving time : 0.02 (sec)

Leaf size : 41

```
DSolve[{(x^2-1)*D[y[x],x]^2==1,{}},  
        y[x],x,IncludeSingularSolutions->True]
```

$$y(x) \rightarrow -\operatorname{arctanh}\left(\frac{x}{\sqrt{x^2 - 1}}\right) + c_1$$

$$y(x) \rightarrow \operatorname{arctanh}\left(\frac{x}{\sqrt{x^2 - 1}}\right) + c_1$$

2.5.6 problem 8

Solved as first order homogeneous class C ode	379
Solved using Lie symmetry for first order ode	381
Solved as first order ode of type Riccati	386
Maple step by step solution	390
Maple trace	391
Maple dsolve solution	391
Mathematica DSolve solution	391

Internal problem ID [18240]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter IV. Methods of solution: First order equations. section 32. Problems at page 89

Problem number : 8

Date solved : Monday, December 23, 2024 at 09:18:47 PM

CAS classification : [[_homogeneous, 'class C'], _Riccati]

Solve

$$y' = (x + y)^2$$

Solved as first order homogeneous class C ode

Time used: 0.064 (sec)

Let

$$z = x + y \tag{1}$$

Then

$$z'(x) = 1 + y'$$

Therefore

$$y' = z'(x) - 1$$

Hence the given ode can now be written as

$$z'(x) - 1 = z^2$$

This is separable first order ode. Integrating

$$\int dx = \int \frac{1}{z^2 + 1} dz$$
$$x + c_1 = \arctan(z)$$

Replacing z back by its value from (1) then the above gives the solution as Solving for y gives

$$y = -x + \tan(x + c_1)$$

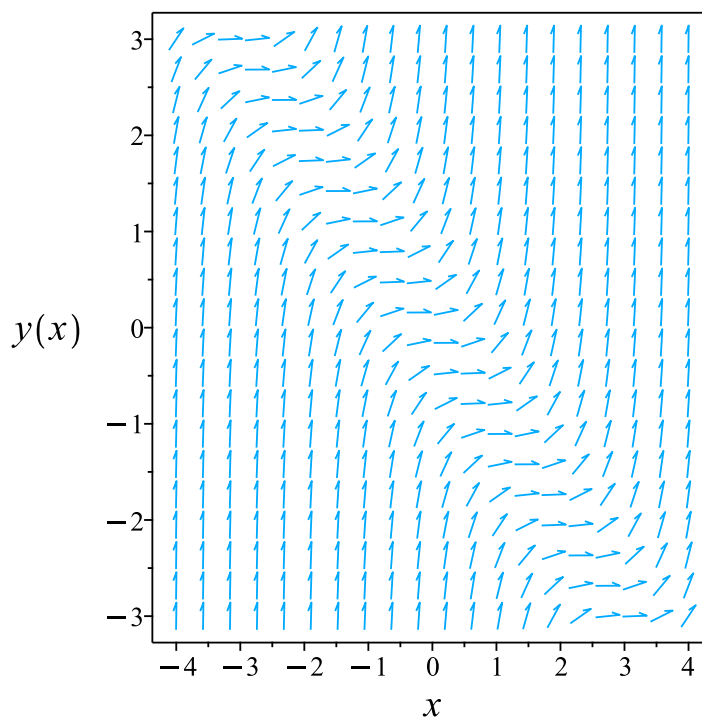


Figure 2.56: Slope field plot
 $y' = (x + y)^2$

Summary of solutions found

$$y = -x + \tan(x + c_1)$$

Solved using Lie symmetry for first order ode

Time used: 0.607 (sec)

Writing the ode as

$$y' = (x + y)^2$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

To determine ξ, η then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$\xi = xa_2 + ya_3 + a_1 \quad (\text{1E})$$

$$\eta = xb_2 + yb_3 + b_1 \quad (\text{2E})$$

Where the unknown coefficients are

$$\{a_1, a_2, a_3, b_1, b_2, b_3\}$$

Substituting equations (1E,2E) and ω into (A) gives

$$b_2 + (x + y)^2 (b_3 - a_2) - (x + y)^4 a_3 - (2x + 2y)(xa_2 + ya_3 + a_1) - (2x + 2y)(xb_2 + yb_3 + b_1) = 0 \quad (\text{5E})$$

Putting the above in normal form gives

$$-x^4 a_3 - 4x^3 y a_3 - 6x^2 y^2 a_3 - 4x y^3 a_3 - y^4 a_3 - 3x^2 a_2 - 2x^2 b_2 + x^2 b_3 - 4x y a_2 - 2x y a_3 - 2x y b_2 - y^2 a_2 - 2y^2 a_3 - y^2 b_3 - 2x a_1 - 2x b_1 - 2y a_1 - 2y b_1 + b_2 = 0$$

Setting the numerator to zero gives

$$-x^4 a_3 - 4x^3 y a_3 - 6x^2 y^2 a_3 - 4x y^3 a_3 - y^4 a_3 - 3x^2 a_2 - 2x^2 b_2 + x^2 b_3 - 4x y a_2 - 2x y a_3 - 2x y b_2 - y^2 a_2 - 2y^2 a_3 - y^2 b_3 - 2x a_1 - 2x b_1 - 2y a_1 - 2y b_1 + b_2 = 0 \quad (\text{6E})$$

Looking at the above PDE shows the following are all the terms with $\{x, y\}$ in them.

$$\{x, y\}$$

The following substitution is now made to be able to collect on all terms with $\{x, y\}$ in them

$$\{x = v_1, y = v_2\}$$

The above PDE (6E) now becomes

$$\begin{aligned} -a_3v_1^4 - 4a_3v_1^3v_2 - 6a_3v_1^2v_2^2 - 4a_3v_1v_2^3 - a_3v_2^4 - 3a_2v_1^2 - 4a_2v_1v_2 - a_2v_2^2 - 2a_3v_1v_2 \\ - 2a_3v_2^2 - 2b_2v_1^2 - 2b_2v_1v_2 + b_3v_1^2 - b_3v_2^2 - 2a_1v_1 - 2a_1v_2 - 2b_1v_1 - 2b_1v_2 + b_2 = 0 \end{aligned} \quad (7E)$$

Collecting the above on the terms v_i introduced, and these are

$$\{v_1, v_2\}$$

Equation (7E) now becomes

$$\begin{aligned} -a_3v_1^4 - 4a_3v_1^3v_2 - 6a_3v_1^2v_2^2 + (-3a_2 - 2b_2 + b_3)v_1^2 \\ - 4a_3v_1v_2^3 + (-4a_2 - 2a_3 - 2b_2)v_1v_2 + (-2a_1 - 2b_1)v_1 \\ - a_3v_2^4 + (-a_2 - 2a_3 - b_3)v_2^2 + (-2a_1 - 2b_1)v_2 + b_2 = 0 \end{aligned} \quad (8E)$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$\begin{aligned} b_2 &= 0 \\ -6a_3 &= 0 \\ -4a_3 &= 0 \\ -a_3 &= 0 \\ -2a_1 - 2b_1 &= 0 \\ -4a_2 - 2a_3 - 2b_2 &= 0 \\ -3a_2 - 2b_2 + b_3 &= 0 \\ -a_2 - 2a_3 - b_3 &= 0 \end{aligned}$$

Solving the above equations for the unknowns gives

$$\begin{aligned} a_1 &= -b_1 \\ a_2 &= 0 \\ a_3 &= 0 \\ b_1 &= b_1 \\ b_2 &= 0 \\ b_3 &= 0 \end{aligned}$$

Substituting the above solution in the anstaz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$\begin{aligned}\xi &= -1 \\ \eta &= 1\end{aligned}$$

Shifting is now applied to make $\xi = 0$ in order to simplify the rest of the computation

$$\begin{aligned}\eta &= \eta - \omega(x, y) \xi \\ &= 1 - ((x + y)^2) (-1) \\ &= x^2 + 2xy + y^2 + 1 \\ \xi &= 0\end{aligned}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS \quad (1)$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = x$$

S is found from

$$\begin{aligned}S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{x^2 + 2xy + y^2 + 1} dy\end{aligned}$$

Which results in

$$S = \arctan(x + y)$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = (x + y)^2$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 1 \\ R_y &= 0 \\ S_x &= \frac{1}{1 + (x + y)^2} \\ S_y &= \frac{1}{1 + (x + y)^2} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = 1 \tag{2A}$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = 1$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S .

Since the ode has the form $\frac{d}{dR}S(R) = f(R)$, then we only need to integrate $f(R)$.

$$\begin{aligned} \int dS &= \int 1 dR \\ S(R) &= R + c_2 \end{aligned}$$

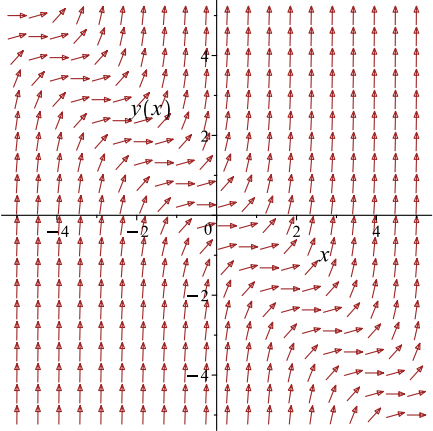
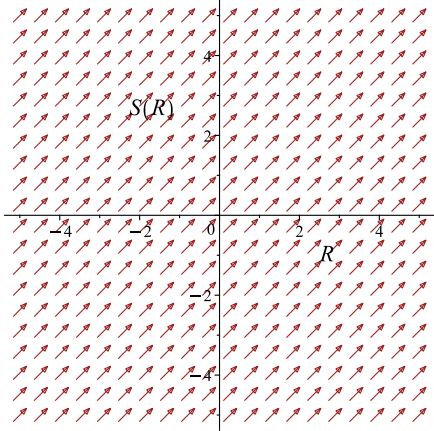
To complete the solution, we just need to transform the above back to x, y coordinates. This results in

$$\arctan(x + y) = x + c_2$$

Which gives

$$y = -x + \tan(x + c_2)$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = (x + y)^2$ 	$R = x$ $S = \arctan(x + y)$	$\frac{dS}{dR} = 1$ 

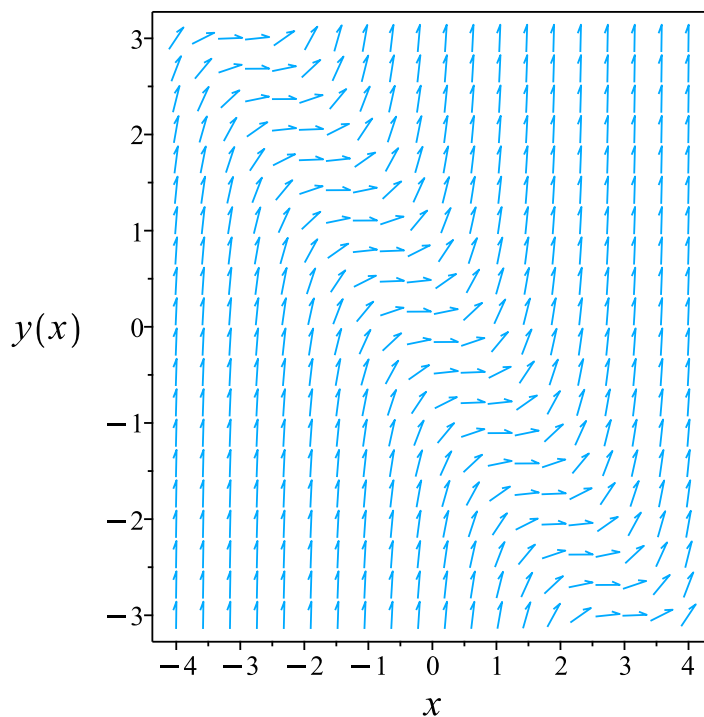


Figure 2.57: Slope field plot $y' = (x + y)^2$

Summary of solutions found

$$y = -x + \tan(x + c_2)$$

Solved as first order ode of type Riccati

Time used: 0.252 (sec)

In canonical form the ODE is

$$\begin{aligned} y' &= F(x, y) \\ &= (x + y)^2 \end{aligned}$$

This is a Riccati ODE. Comparing the ODE to solve

$$y' = x^2 + 2xy + y^2$$

With Riccati ODE standard form

$$y' = f_0(x) + f_1(x)y + f_2(x)y^2$$

Shows that $f_0(x) = x^2$, $f_1(x) = 2x$ and $f_2(x) = 1$. Let

$$\begin{aligned} y &= \frac{-u'}{f_2 u} \\ &= \frac{-u'}{u} \end{aligned} \tag{1}$$

Using the above substitution in the given ODE results (after some simplification) in a second order ODE to solve for $u(x)$ which is

$$f_2 u''(x) - (f_2' + f_1 f_2) u'(x) + f_2^2 f_0 u(x) = 0 \tag{2}$$

But

$$\begin{aligned} f_2' &= 0 \\ f_1 f_2 &= 2x \\ f_2^2 f_0 &= x^2 \end{aligned}$$

Substituting the above terms back in equation (2) gives

$$u''(x) - 2xu'(x) + x^2u(x) = 0$$

In normal form the given ode is written as

$$\frac{d^2u}{dx^2} + p(x) \left(\frac{du}{dx} \right) + q(x) u = 0 \quad (2)$$

Where

$$\begin{aligned} p(x) &= -2x \\ q(x) &= x^2 \end{aligned}$$

Calculating the Liouville ode invariant Q given by

$$\begin{aligned} Q &= q - \frac{p'}{2} - \frac{p^2}{4} \\ &= x^2 - \frac{(-2x)'}{2} - \frac{(-2x)^2}{4} \\ &= x^2 - \frac{(-2)}{2} - \frac{(4x^2)}{4} \\ &= x^2 - (-1) - x^2 \\ &= 1 \end{aligned}$$

Since the Liouville ode invariant does not depend on the independent variable x then the transformation

$$u = v(x) z(x) \quad (3)$$

is used to change the original ode to a constant coefficients ode in v . In (3) the term $z(x)$ is given by

$$\begin{aligned} z(x) &= e^{-\int \frac{p(x)}{2} dx} \\ &= e^{-\int \frac{-2x}{2} dx} \\ &= e^{\frac{x^2}{2}} \end{aligned} \quad (5)$$

Hence (3) becomes

$$u = v(x) e^{\frac{x^2}{2}} \quad (4)$$

Applying this change of variable to the original ode results in

$$e^{\frac{x^2}{2}} \left(v(x) + \frac{d^2}{dx^2} v(x) \right) = 0$$

Which is now solved for $v(x)$.

The above ode can be simplified to

$$v(x) + \frac{d^2}{dx^2}v(x) = 0$$

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Av''(x) + Bv'(x) + Cv(x) = 0$$

Where in the above $A = 1, B = 0, C = 1$. Let the solution be $v(x) = e^{\lambda x}$. Substituting this into the ODE gives

$$\lambda^2 e^{x\lambda} + e^{x\lambda} = 0 \quad (1)$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$\lambda^2 + 1 = 0 \quad (2)$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = 0, C = 1$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{0}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{0^2 - (4)(1)(1)} \\ &= \pm i \end{aligned}$$

Hence

$$\lambda_1 = +i$$

$$\lambda_2 = -i$$

Which simplifies to

$$\lambda_1 = i$$

$$\lambda_2 = -i$$

Since roots are complex conjugate of each others, then let the roots be

$$\lambda_{1,2} = \alpha \pm i\beta$$

Where $\alpha = 0$ and $\beta = 1$. Therefore the final solution, when using Euler relation, can be written as

$$v(x) = e^{\alpha x} (c_1 \cos(\beta x) + c_2 \sin(\beta x))$$

Which becomes

$$v(x) = e^0(\cos(x) c_1 + c_2 \sin(x))$$

Or

$$v(x) = \cos(x) c_1 + c_2 \sin(x)$$

Will add steps showing solving for IC soon.

Now that $v(x)$ is known, then

$$\begin{aligned} u &= v(x) z(x) \\ &= (\cos(x) c_1 + c_2 \sin(x)) (z(x)) \end{aligned} \tag{7}$$

But from (5)

$$z(x) = e^{\frac{x^2}{2}}$$

Hence (7) becomes

$$u = (\cos(x) c_1 + c_2 \sin(x)) e^{\frac{x^2}{2}}$$

Will add steps showing solving for IC soon.

Taking derivative gives

$$u'(x) = (-c_1 \sin(x) + c_2 \cos(x)) e^{\frac{x^2}{2}} + (\cos(x) c_1 + c_2 \sin(x)) x e^{\frac{x^2}{2}}$$

Doing change of constants, the solution becomes

$$y = -\frac{\left((-c_3 \sin(x) + \cos(x)) e^{\frac{x^2}{2}} + (\cos(x) c_3 + \sin(x)) x e^{\frac{x^2}{2}}\right) e^{-\frac{x^2}{2}}}{\cos(x) c_3 + \sin(x)}$$

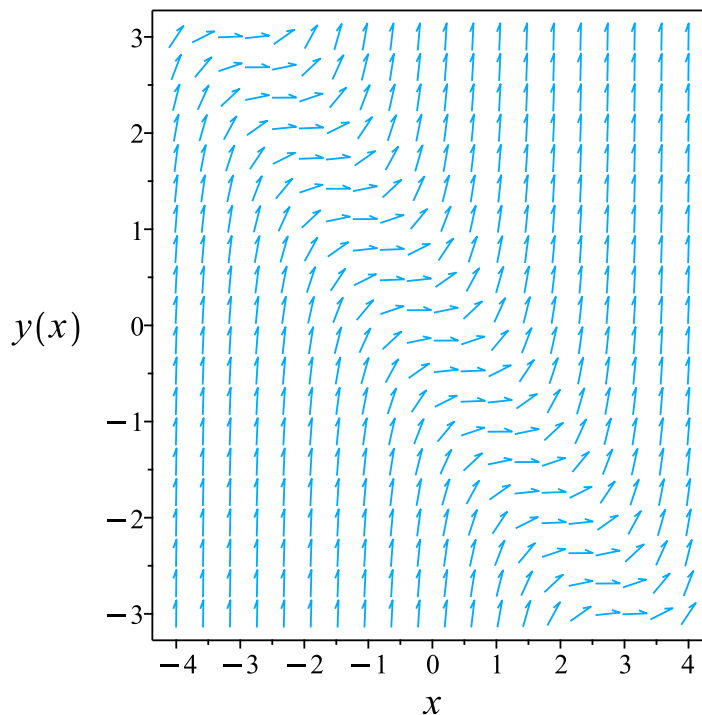


Figure 2.58: Slope field plot
 $y' = (x + y)^2$

Summary of solutions found

$$y = -\frac{\left((-c_3 \sin(x) + \cos(x)) e^{\frac{x^2}{2}} + (\cos(x) c_3 + \sin(x)) x e^{\frac{x^2}{2}}\right) e^{-\frac{x^2}{2}}}{\cos(x) c_3 + \sin(x)}$$

Maple step by step solution

Let's solve

$$y' = (x + y)^2$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Solve for the highest derivative

$$y' = (x + y)^2$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
trying inverse linear
trying homogeneous types:
trying homogeneous C
1st order, trying the canonical coordinates of the invariance group
  -> Calling odsolve with the ODE`, diff(y(x), x) = -1, y(x)`      *** Sublevel 2 ***
      Methods for first order ODEs:
      --- Trying classification methods ---
      trying a quadrature
      trying 1st order linear
      <- 1st order linear successful
<- 1st order, canonical coordinates successful
<- homogeneous successful`

```

Maple dsolve solution

Solving time : 0.006 (sec)

Leaf size : 16

```

dsolve(diff(y(x),x) = (x+y(x))^2,
       y(x),singsol=all)

```

$$y = -x - \tan(c_1 - x)$$

Mathematica DSolve solution

Solving time : 0.559 (sec)

Leaf size : 14

```

DSolve[{D[y[x],x]==(x+y[x])^2,{}},
       y[x],x,IncludeSingularSolutions->True]

```

$$y(x) \rightarrow -x + \tan(x + c_1)$$

2.6 Chapter IV. Methods of solution: First order equations. section 33. Problems at page 91

2.6.1	problem 1	393
2.6.2	problem 2 (eq 39)	405
2.6.3	problem 3 (eq 41)	410
2.6.4	problem 4 (eq 50)	415
2.6.5	problem 8 (eq 68)	424
2.6.6	problem 8 (eq 69)	441
2.6.7	problem 9 (a)	461
2.6.8	problem 9 (b)	478
2.6.9	problem 9 (c)	482
2.6.10	problem 9 (d)	485
2.6.11	problem 9 (e)	502
2.6.12	problem 10 (a)	528
2.6.13	problem 10 (b)	542
2.6.14	problem 10 (c)	562

2.6.1 problem 1

Solved as second order linear constant coeff ode	393
Solved as second order can be made integrable	394
Solved as second order ode using Kovacic algorithm	397
Solved as second order ode adjoint method	400
Maple step by step solution	403
Maple trace	404
Maple dsolve solution	404
Mathematica DSolve solution	404

Internal problem ID [18241]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter IV. Methods of solution: First order equations. section 33. Problems at page 91

Problem number : 1

Date solved : Monday, December 23, 2024 at 09:18:48 PM

CAS classification : [[_2nd_order, _missing_x]]

Solve

$$\theta'' = -p^2\theta$$

Solved as second order linear constant coeff ode

Time used: 0.285 (sec)

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$A\theta''(t) + B\theta'(t) + C\theta(t) = 0$$

Where in the above $A = 1, B = 0, C = p^2$. Let the solution be $\theta = e^{\lambda t}$. Substituting this into the ODE gives

$$\lambda^2 e^{t\lambda} + p^2 e^{t\lambda} = 0 \tag{1}$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda t}$ gives

$$\lambda^2 + p^2 = 0 \tag{2}$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = 0, C = p^2$ into the above gives

$$\begin{aligned}\lambda_{1,2} &= \frac{0}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{0^2 - (4)(1)(p^2)} \\ &= \pm \sqrt{-p^2}\end{aligned}$$

Hence

$$\begin{aligned}\lambda_1 &= +\sqrt{-p^2} \\ \lambda_2 &= -\sqrt{-p^2}\end{aligned}$$

Which simplifies to

$$\begin{aligned}\lambda_1 &= \sqrt{-p^2} \\ \lambda_2 &= -\sqrt{-p^2}\end{aligned}$$

Since roots are real and distinct, then the solution is

$$\begin{aligned}\theta &= c_1 e^{\lambda_1 t} + c_2 e^{\lambda_2 t} \\ \theta &= c_1 e^{(\sqrt{-p^2})t} + c_2 e^{(-\sqrt{-p^2})t}\end{aligned}$$

Or

$$\theta = c_1 e^{t\sqrt{-p^2}} + c_2 e^{-t\sqrt{-p^2}}$$

Will add steps showing solving for IC soon.

Summary of solutions found

$$\theta = c_1 e^{t\sqrt{-p^2}} + c_2 e^{-t\sqrt{-p^2}}$$

Solved as second order can be made integrable

Time used: 2.397 (sec)

Multiplying the ode by θ' gives

$$\theta' \theta'' + p^2 \theta' \theta = 0$$

Integrating the above w.r.t t gives

$$\begin{aligned}\int (\theta' \theta'' + p^2 \theta' \theta) dt &= 0 \\ \frac{\theta'^2}{2} + \frac{p^2 \theta^2}{2} &= c_1\end{aligned}$$

Which is now solved for θ . Solving for the derivative gives these ODE's to solve

$$\theta' = \sqrt{-p^2\theta^2 + 2c_1} \quad (1)$$

$$\theta' = -\sqrt{-p^2\theta^2 + 2c_1} \quad (2)$$

Now each of the above is solved separately.

Solving Eq. (1)

Integrating gives

$$\int \frac{1}{\sqrt{-p^2\theta^2 + 2c_1}} d\theta = dt$$

$$\frac{\arctan\left(\frac{p\theta}{\sqrt{-p^2\theta^2 + 2c_1}}\right)}{p} = t + c_2$$

Singular solutions are found by solving

$$\sqrt{-p^2\theta^2 + 2c_1} = 0$$

for θ . This is because we had to divide by this in the above step. This gives the following singular solution(s), which also have to satisfy the given ODE.

$$\theta = \frac{\sqrt{2}\sqrt{c_1}}{p}$$

$$\theta = -\frac{\sqrt{2}\sqrt{c_1}}{p}$$

Solving for θ gives

$$\theta = \frac{\sqrt{2}\sqrt{c_1}}{p}$$

$$\theta = \frac{\tan(c_2p + tp)\sqrt{2}\sqrt{\frac{c_1}{\tan^2(c_2p + tp) + 1}}}{p}$$

$$\theta = -\frac{\sqrt{2}\sqrt{c_1}}{p}$$

Solving Eq. (2)

Integrating gives

$$\int -\frac{1}{\sqrt{-p^2\theta^2 + 2c_1}} d\theta = dt$$

$$-\frac{\arctan\left(\frac{p\theta}{\sqrt{-p^2\theta^2 + 2c_1}}\right)}{p} = t + c_3$$

Singular solutions are found by solving

$$-\sqrt{-p^2\theta^2 + 2c_1} = 0$$

for θ . This is because we had to divide by this in the above step. This gives the following singular solution(s), which also have to satisfy the given ODE.

$$\theta = \frac{\sqrt{2}\sqrt{c_1}}{p}$$

$$\theta = -\frac{\sqrt{2}\sqrt{c_1}}{p}$$

Solving for θ gives

$$\theta = \frac{\sqrt{2}\sqrt{c_1}}{p}$$

$$\theta = -\frac{\sqrt{2}\sqrt{c_1}}{p}$$

$$\theta = -\frac{\tan(c_3p + tp)\sqrt{2}\sqrt{\frac{c_1}{\tan^2(c_3p + tp) + 1}}}{p}$$

Will add steps showing solving for IC soon.

The solution

$$\theta = \frac{\sqrt{2}\sqrt{c_1}}{p}$$

was found not to satisfy the ode or the IC. Hence it is removed. The solution

$$\theta = -\frac{\sqrt{2}\sqrt{c_1}}{p}$$

was found not to satisfy the ode or the IC. Hence it is removed.

Summary of solutions found

$$\theta = \frac{\tan(c_2 p + tp) \sqrt{2} \sqrt{\frac{c_1}{\tan(c_2 p + tp)^2 + 1}}}{p}$$

$$\theta = - \frac{\tan(c_3 p + tp) \sqrt{2} \sqrt{\frac{c_1}{\tan(c_3 p + tp)^2 + 1}}}{p}$$

Solved as second order ode using Kovacic algorithm

Time used: 0.091 (sec)

Writing the ode as

$$\theta'' + p^2 \theta = 0 \quad (1)$$

$$A\theta'' + B\theta' + C\theta = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$A = 1$$

$$B = 0 \quad (3)$$

$$C = p^2$$

Applying the Liouville transformation on the dependent variable gives

$$z(t) = \theta e^{\int \frac{B}{2A} dt}$$

Then (2) becomes

$$z''(t) = r z(t) \quad (4)$$

Where r is given by

$$r = \frac{s}{t} \quad (5)$$

$$= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2}$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{-p^2}{1} \quad (6)$$

Comparing the above to (5) shows that

$$\begin{aligned}s &= -p^2 \\ t &= 1\end{aligned}$$

Therefore eq. (4) becomes

$$z''(t) = (-p^2) z(t) \quad (7)$$

Equation (7) is now solved. After finding $z(t)$ then θ is found using the inverse transformation

$$\theta = z(t) e^{-\int \frac{B}{2A} dt}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 2.31: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned}O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - 0 \\ &= 0\end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = -p^2$ is not a function of t , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(t) = e^{t\sqrt{-p^2}}$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in θ is found from

$$\theta_1 = z_1 e^{\int -\frac{1}{2} \frac{B}{A} dt}$$

Since $B = 0$ then the above reduces to

$$\begin{aligned} \theta_1 &= z_1 \\ &= e^{t\sqrt{-p^2}} \end{aligned}$$

Which simplifies to

$$\theta_1 = e^{t\sqrt{-p^2}}$$

The second solution θ_2 to the original ode is found using reduction of order

$$\theta_2 = \theta_1 \int \frac{e^{\int -\frac{B}{A} dt}}{\theta_1^2} dt$$

Since $B = 0$ then the above becomes

$$\begin{aligned} \theta_2 &= \theta_1 \int \frac{1}{\theta_1^2} dt \\ &= e^{t\sqrt{-p^2}} \int \frac{1}{e^{2t\sqrt{-p^2}}} dt \\ &= e^{t\sqrt{-p^2}} \left(\frac{\sqrt{-p^2} e^{-2t\sqrt{-p^2}}}{2p^2} \right) \end{aligned}$$

Therefore the solution is

$$\begin{aligned} \theta &= c_1 \theta_1 + c_2 \theta_2 \\ &= c_1 \left(e^{t\sqrt{-p^2}} \right) + c_2 \left(e^{t\sqrt{-p^2}} \left(\frac{\sqrt{-p^2} e^{-2t\sqrt{-p^2}}}{2p^2} \right) \right) \end{aligned}$$

Will add steps showing solving for IC soon.

Summary of solutions found

$$\theta = c_1 e^{t\sqrt{-p^2}} + \frac{c_2 e^{-t\sqrt{-p^2}} \sqrt{-p^2}}{2p^2}$$

Solved as second order ode adjoint method

Time used: 0.242 (sec)

In normal form the ode

$$\theta'' = -p^2\theta \quad (1)$$

Becomes

$$\theta'' + p(t)\theta' + q(t)\theta = r(t) \quad (2)$$

Where

$$\begin{aligned} p(t) &= 0 \\ q(t) &= p^2 \\ r(t) &= 0 \end{aligned}$$

The Lagrange adjoint ode is given by

$$\begin{aligned} \xi'' - (\xi p)' + \xi q &= 0 \\ \xi'' - (0)' + (p^2\xi(t)) &= 0 \\ \xi''(t) + p^2\xi(t) &= 0 \end{aligned}$$

Which is solved for $\xi(t)$. This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$A\xi''(t) + B\xi'(t) + C\xi(t) = 0$$

Where in the above $A = 1, B = 0, C = p^2$. Let the solution be $\xi = e^{\lambda t}$. Substituting this into the ODE gives

$$\lambda^2 e^{t\lambda} + p^2 e^{t\lambda} = 0 \quad (1)$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda t}$ gives

$$\lambda^2 + p^2 = 0 \quad (2)$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = 0, C = p^2$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{0}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{0^2 - (4)(1)(p^2)} \\ &= \pm \sqrt{-p^2} \end{aligned}$$

Hence

$$\begin{aligned} \lambda_1 &= +\sqrt{-p^2} \\ \lambda_2 &= -\sqrt{-p^2} \end{aligned}$$

Which simplifies to

$$\begin{aligned} \lambda_1 &= \sqrt{-p^2} \\ \lambda_2 &= -\sqrt{-p^2} \end{aligned}$$

Since roots are real and distinct, then the solution is

$$\begin{aligned} \xi &= c_1 e^{\lambda_1 t} + c_2 e^{\lambda_2 t} \\ \xi &= c_1 e^{(\sqrt{-p^2})t} + c_2 e^{(-\sqrt{-p^2})t} \end{aligned}$$

Or

$$\xi = c_1 e^{t\sqrt{-p^2}} + c_2 e^{-t\sqrt{-p^2}}$$

Will add steps showing solving for IC soon.

The original ode now reduces to first order ode

$$\begin{aligned} \xi(t) \theta' - \theta \xi'(t) + \xi(t) p(t) \theta &= \int \xi(t) r(t) dt \\ \theta' + \theta \left(p(t) - \frac{\xi'(t)}{\xi(t)} \right) &= \frac{\int \xi(t) r(t) dt}{\xi(t)} \end{aligned}$$

Or

$$\theta' - \frac{\theta \left(c_1 \sqrt{-p^2} e^{t\sqrt{-p^2}} - c_2 e^{-t\sqrt{-p^2}} \sqrt{-p^2} \right)}{c_1 e^{t\sqrt{-p^2}} + c_2 e^{-t\sqrt{-p^2}}} = 0$$

Which is now a first order ode. This is now solved for θ . In canonical form a linear first order is

$$\theta' + q(t)\theta = p(t)$$

Comparing the above to the given ode shows that

$$q(t) = -\frac{\sqrt{-p^2} \left(c_1 e^{2t\sqrt{-p^2}} - c_2 \right)}{c_1 e^{2t\sqrt{-p^2}} + c_2}$$

$$p(t) = 0$$

The integrating factor μ is

$$\begin{aligned} \mu &= e^{\int q dt} \\ &= e^{\int -\frac{\sqrt{-p^2} \left(c_1 e^{2t\sqrt{-p^2}} - c_2 \right)}{c_1 e^{2t\sqrt{-p^2}} + c_2} dt} \\ &= \frac{\sqrt{e^{2t\sqrt{-p^2}}}}{c_1 e^{2t\sqrt{-p^2}} + c_2} \end{aligned}$$

The ode becomes

$$\frac{d}{dt} \mu \theta = 0$$

$$\frac{d}{dt} \left(\frac{\theta \sqrt{e^{2t\sqrt{-p^2}}}}{c_1 e^{2t\sqrt{-p^2}} + c_2} \right) = 0$$

Integrating gives

$$\begin{aligned} \frac{\theta \sqrt{e^{2t\sqrt{-p^2}}}}{c_1 e^{2t\sqrt{-p^2}} + c_2} &= \int 0 dt + c_3 \\ &= c_3 \end{aligned}$$

Dividing throughout by the integrating factor $\frac{\sqrt{e^{2t\sqrt{-p^2}}}}{c_1 e^{2t\sqrt{-p^2}} + c_2}$ gives the final solution

$$\theta = \frac{\left(c_1 e^{2t\sqrt{-p^2}} + c_2 \right) c_3}{\sqrt{e^{2t\sqrt{-p^2}}}}$$

Hence, the solution found using Lagrange adjoint equation method is

$$\theta = \frac{\left(c_1 e^{2t\sqrt{-p^2}} + c_2 \right) c_3}{\sqrt{e^{2t\sqrt{-p^2}}}}$$

The constants can be merged to give

$$\theta = \frac{c_1 e^{2t\sqrt{-p^2}} + c_2}{\sqrt{e^{2t\sqrt{-p^2}}}}$$

Will add steps showing solving for IC soon.

Summary of solutions found

$$\theta = \frac{c_1 e^{2t\sqrt{-p^2}} + c_2}{\sqrt{e^{2t\sqrt{-p^2}}}}$$

Maple step by step solution

Let's solve

$$\theta'' = -p^2\theta$$

- Highest derivative means the order of the ODE is 2

$$\theta''$$

- Group terms with θ on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$\theta'' + p^2\theta = 0$$

- Characteristic polynomial of ODE

$$p^2 + r^2 = 0$$

- Use quadratic formula to solve for r

$$r = \frac{0 \pm (\sqrt{-4p^2})}{2}$$

- Roots of the characteristic polynomial

$$r = (\sqrt{-p^2}, -\sqrt{-p^2})$$

- 1st solution of the ODE

$$\theta_1(t) = e^{t\sqrt{-p^2}}$$

- 2nd solution of the ODE

$$\theta_2(t) = e^{-t\sqrt{-p^2}}$$

- General solution of the ODE

$$\theta = C1\theta_1(t) + C2\theta_2(t)$$

- Substitute in solutions

$$\theta = C1 e^{t\sqrt{-p^2}} + C2 e^{-t\sqrt{-p^2}}$$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
<- constant coefficients successful`

```

Maple dsolve solution

Solving time : 0.006 (sec)

Leaf size : 17

```

dsolve(diff(diff(theta(t),t),t) = -p^2*theta(t),
        theta(t),singsol=all)

```

$$\theta = c_1 \sin(tp) + c_2 \cos(tp)$$

Mathematica DSolve solution

Solving time : 0.029 (sec)

Leaf size : 20

```

DSolve[{D[theta[t],{t,2}]== -p^2*theta[t],{}}],
        theta[t],t,IncludeSingularSolutions->True]

```

$$\theta(t) \rightarrow c_1 \cos(pt) + c_2 \sin(pt)$$

2.6.2 problem 2 (eq 39)

Solved as first order quadrature ode	405
Solved as first order Exact ode	406
Maple step by step solution	408
Maple trace	409
Maple dsolve solution	409
Mathematica DSolve solution	409

Internal problem ID [18242]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter IV. Methods of solution: First order equations. section 33. Problems at page 91

Problem number : 2 (eq 39)

Date solved : Monday, December 23, 2024 at 09:18:51 PM

CAS classification : [_quadrature]

Solve

$$\sec(\theta)^2 = \frac{ms'}{k}$$

Solved as first order quadrature ode

Time used: 0.131 (sec)

Since the ode has the form $s' = f(\theta)$, then we only need to integrate $f(\theta)$.

$$\int ds = \int \frac{\sec(\theta)^2 k}{m} d\theta$$

$$s = \frac{k \tan(\theta)}{m} + c_1$$

Summary of solutions found

$$s = \frac{k \tan(\theta)}{m} + c_1$$

Solved as first order Exact ode

Time used: 0.088 (sec)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx}\phi(x, y) = 0$$

Hence

$$\frac{\partial\phi}{\partial x} + \frac{\partial\phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial\phi}{\partial x} &= M \\ \frac{\partial\phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(\theta, s) d\theta + N(\theta, s) ds = 0 \quad (\text{1A})$$

Therefore

$$\begin{aligned} \left(-\frac{m}{k}\right) ds &= (-\sec(\theta)^2) d\theta \\ (\sec(\theta)^2) d\theta + \left(-\frac{m}{k}\right) ds &= 0 \end{aligned} \quad (\text{2A})$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(\theta, s) &= \sec(\theta)^2 \\ N(\theta, s) &= -\frac{m}{k} \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial s} = \frac{\partial N}{\partial \theta}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial s} &= \frac{\partial}{\partial s} (\sec(\theta)^2) \\ &= 0\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial \theta} &= \frac{\partial}{\partial \theta} \left(-\frac{m}{k}\right) \\ &= 0\end{aligned}$$

Since $\frac{\partial M}{\partial s} = \frac{\partial N}{\partial \theta}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(\theta, s)$

$$\frac{\partial \phi}{\partial \theta} = M \tag{1}$$

$$\frac{\partial \phi}{\partial s} = N \tag{2}$$

Integrating (1) w.r.t. θ gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial \theta} d\theta &= \int M d\theta \\ \int \frac{\partial \phi}{\partial \theta} d\theta &= \int \sec(\theta)^2 d\theta \\ \phi &= \tan(\theta) + f(s)\end{aligned} \tag{3}$$

Where $f(s)$ is used for the constant of integration since ϕ is a function of both θ and s . Taking derivative of equation (3) w.r.t s gives

$$\frac{\partial \phi}{\partial s} = 0 + f'(s) \tag{4}$$

But equation (2) says that $\frac{\partial \phi}{\partial s} = -\frac{m}{k}$. Therefore equation (4) becomes

$$-\frac{m}{k} = 0 + f'(s) \tag{5}$$

Solving equation (5) for $f'(s)$ gives

$$f'(s) = -\frac{m}{k}$$

Integrating the above w.r.t s gives

$$\int f'(s) ds = \int \left(-\frac{m}{k}\right) ds$$

$$f(s) = -\frac{ms}{k} + c_1$$

Where c_1 is constant of integration. Substituting result found above for $f(s)$ into equation (3) gives ϕ

$$\phi = \tan(\theta) - \frac{ms}{k} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into the constant c_1 gives the solution as

$$c_1 = \tan(\theta) - \frac{ms}{k}$$

Solving for s gives

$$s = \frac{(-c_1 + \tan(\theta)) k}{m}$$

Summary of solutions found

$$s = \frac{(-c_1 + \tan(\theta)) k}{m}$$

Maple step by step solution

Let's solve

$$\sec(\theta)^2 = \frac{ms'}{k}$$

- Highest derivative means the order of the ODE is 1
- s'
- Separate variables

$$s' = \frac{\sec(\theta)^2 k}{m}$$

- Integrate both sides with respect to θ

$$\int s' d\theta = \int \frac{\sec(\theta)^2 k}{m} d\theta + C1$$

- Evaluate integral

$$s = \frac{k \tan(\theta)}{m} + C1$$

- Solve for s

$$s = \frac{k \tan(\theta) + C1m}{m}$$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
<- quadrature successful`
```

Maple dsolve solution

Solving time : 0.002 (sec)

Leaf size : 13

```
dsolve(sec(theta)^2 = m/k*diff(s(theta),theta),
        s(theta),singsol=all)
```

$$s = \frac{k \tan(\theta)}{m} + c_1$$

Mathematica DSolve solution

Solving time : 0.009 (sec)

Leaf size : 15

```
DSolve[{Sec[theta]^2==m/k*D[s[theta],theta],{}}
        s[theta],theta,IncludeSingularSolutions->True]
```

$$s(\theta) \rightarrow \frac{k \tan(\theta)}{m} + c_1$$

2.6.3 problem 3 (eq 41)

Solved as second order missing y ode	410
Maple step by step solution	412
Maple trace	413
Maple dsolve solution	414
Mathematica DSolve solution	414

Internal problem ID [18243]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter IV. Methods of solution: First order equations. section 33. Problems at page 91

Problem number : 3 (eq 41)

Date solved : Monday, December 23, 2024 at 09:18:52 PM

CAS classification : [[_2nd_order, _missing_x]]

Solve

$$y'' = \frac{m\sqrt{1+y'^2}}{k}$$

Solved as second order missing y ode

Time used: 0.320 (sec)

This is second order ode with missing dependent variable y . Let

$$p(x) = y'$$

Then

$$p'(x) = y''$$

Hence the ode becomes

$$p'(x) - \frac{m\sqrt{1+p(x)^2}}{k} = 0$$

Which is now solve for $p(x)$ as first order ode. Integrating gives

$$\int \frac{k}{m\sqrt{p^2+1}} dp = dx$$

$$\frac{k \operatorname{arcsinh}(p)}{m} = x + c_1$$

Singular solutions are found by solving

$$\frac{m\sqrt{p^2 + 1}}{k} = 0$$

for $p(x)$. This is because we had to divide by this in the above step. This gives the following singular solution(s), which also have to satisfy the given ODE.

$$p(x) = -i$$

$$p(x) = i$$

Solving for $p(x)$ gives

$$p(x) = -i$$

$$p(x) = i$$

$$p(x) = \sinh\left(\frac{m(x + c_1)}{k}\right)$$

For solution (1) found earlier, since $p = y'$ then we now have a new first order ode to solve which is

$$y' = -i$$

Since the ode has the form $y' = f(x)$, then we only need to integrate $f(x)$.

$$\int dy = \int -i dx$$

$$y = -ix + c_2$$

For solution (2) found earlier, since $p = y'$ then we now have a new first order ode to solve which is

$$y' = i$$

Since the ode has the form $y' = f(x)$, then we only need to integrate $f(x)$.

$$\int dy = \int i dx$$

$$y = ix + c_3$$

For solution (3) found earlier, since $p = y'$ then we now have a new first order ode to solve which is

$$y' = \sinh\left(\frac{m(x + c_1)}{k}\right)$$

Since the ode has the form $y' = f(x)$, then we only need to integrate $f(x)$.

$$\int dy = \int \sinh\left(\frac{m(x+c_1)}{k}\right) dx$$

$$y = \frac{k \cosh\left(\frac{mx}{k} + \frac{mc_1}{k}\right)}{m} + c_4$$

$$y = \frac{k \cosh\left(\frac{m(x+c_1)}{k}\right)}{m} + c_4$$

Will add steps showing solving for IC soon.

Summary of solutions found

$$y = -ix + c_2$$

$$y = ix + c_3$$

$$y = \frac{k \cosh\left(\frac{m(x+c_1)}{k}\right)}{m} + c_4$$

Maple step by step solution

Let's solve

$$y'' = \frac{m\sqrt{1+y'^2}}{k}$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Make substitution $u = y'$ to reduce order of ODE

$$u'(x) = \frac{m\sqrt{1+u(x)^2}}{k}$$

- Solve for the highest derivative

$$u'(x) = \frac{m\sqrt{1+u(x)^2}}{k}$$

- Separate variables

$$\frac{u'(x)}{\sqrt{1+u(x)^2}} = \frac{m}{k}$$

- Integrate both sides with respect to x

$$\int \frac{u'(x)}{\sqrt{1+u(x)^2}} dx = \int \frac{m}{k} dx + C1$$

- Evaluate integral

- $$\operatorname{arcsinh}(u(x)) = \frac{mx}{k} + C1$$
- Solve for $u(x)$

$$u(x) = \sinh\left(\frac{C1k+xm}{k}\right)$$
- Solve 1st ODE for $u(x)$

$$u(x) = \sinh\left(\frac{C1k+xm}{k}\right)$$
- Make substitution $u = y'$

$$y' = \sinh\left(\frac{C1k+xm}{k}\right)$$
- Integrate both sides to solve for y

$$\int y' dx = \int \sinh\left(\frac{C1k+xm}{k}\right) dx + C2$$
- Compute integrals
$$y = \frac{k \cosh\left(\frac{mx}{k} + C1\right)}{m} + C2$$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying 2nd order Liouville
trying 2nd order WeierstrassP
trying 2nd order JacobiSN
differential order: 2; trying a linearization to 3rd order
trying 2nd order ODE linearizable_by_differentiation
-> Calling odsolve with the ODE`, diff(diff(diff(y(x), x), x), x)-m^2*(diff(y(x), x))/
  Methods for third order ODEs:
  --- Trying classification methods ---
  trying a quadrature
  checking if the LODE has constant coefficients
  <- constant coefficients successful
trying 2nd order, 2 integrating factors of the form mu(x,y)
trying differential order: 2; missing variables
`, `-> Computing symmetries using: way = 3
-> Calling odsolve with the ODE`, diff(_b(_a), _a) = m*(1+_b(_a)^2)^(1/2)/k, _b(_a), F
  symmetry methods on request
`, `1st order, trying reduction of order with given symmetries: `[1, 0]

```

Maple dsolve solution

Solving time : 0.307 (sec)

Leaf size : 36

```
dsolve(diff(diff(y(x),x),x) = m/k*(1+diff(y(x),x)^2)^(1/2),
        y(x),singsol=all)
```

$$y = -ix + c_1$$

$$y = ix + c_1$$

$$y = c_2 + \frac{k \cosh\left(\frac{m(x+c_1)}{k}\right)}{m}$$

Mathematica DSolve solution

Solving time : 45.711 (sec)

Leaf size : 58

```
DSolve[{D[y[x],{x,2}]==m/k*Sqrt[1+D[y[x],x]^2],{}}],
        y[x],x,IncludeSingularSolutions->True]
```

$$y(x) \rightarrow c_2 - \frac{k}{m \sqrt{\operatorname{sech}^2\left(\frac{mx}{k} + c_1\right)}}$$

$$y(x) \rightarrow \frac{k}{m \sqrt{\operatorname{sech}^2\left(\frac{mx}{k} + c_1\right)}} + c_2$$

2.6.4 problem 4 (eq 50)

Solved as second order missing x ode	415
Solved as second order can be made integrable	418
Maple step by step solution	419
Maple trace	422
Maple dsolve solution	422
Mathematica DSolve solution	423

Internal problem ID [18244]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter IV. Methods of solution: First order equations. section 33. Problems at page 91

Problem number : 4 (eq 50)

Date solved : Monday, December 23, 2024 at 09:18:54 PM

CAS classification :

[[_2nd_order, _missing_x], [_2nd_order, _reducible, _mu_x_y1]]

Solve

$$\phi'' = \frac{4\pi n c}{\sqrt{v_0^2 + \frac{2e(\phi - V_0)}{m}}}$$

Solved as second order missing x ode

Time used: 4.773 (sec)

This is missing independent variable second order ode. Solved by reduction of order by using substitution which makes the dependent variable ϕ an independent variable. Using

$$\phi' = p$$

Then

$$\begin{aligned}\phi'' &= \frac{dp}{dx} \\ &= \frac{dp}{d\phi} \frac{d\phi}{dx} \\ &= p \frac{dp}{d\phi}\end{aligned}$$

Hence the ode becomes

$$p(\phi) \left(\frac{d}{d\phi} p(\phi) \right) = \frac{4\pi n c}{\sqrt{v_0^2 + \frac{2e(\phi - V_0)}{m}}}$$

Which is now solved as first order ode for $p(\phi)$.

The ode $p' = \frac{4\pi n c}{p \sqrt{-\frac{v_0^2 m + 2eV_0 - 2e\phi}{m}}}$ is separable as it can be written as

$$\begin{aligned} p' &= \frac{4\pi n c}{p \sqrt{-\frac{v_0^2 m + 2eV_0 - 2e\phi}{m}}} \\ &= f(\phi)g(p) \end{aligned}$$

Where

$$\begin{aligned} f(\phi) &= \frac{4\pi n c}{\sqrt{-\frac{v_0^2 m + 2eV_0 - 2e\phi}{m}}} \\ g(p) &= \frac{1}{p} \end{aligned}$$

Integrating gives

$$\begin{aligned} \int \frac{1}{g(p)} dp &= \int f(\phi) d\phi \\ \int p dp &= \int \frac{4\pi n c}{\sqrt{-\frac{v_0^2 m + 2eV_0 - 2e\phi}{m}}} d\phi \\ \frac{p^2}{2} &= \frac{4\sqrt{\frac{(-2V_0 + 2\phi)e + v_0^2 m}{m}} \pi n c m}{e} + c_1 \end{aligned}$$

Solving for p gives

$$\begin{aligned} p &= \frac{\sqrt{2} \sqrt{e \left(4\sqrt{-\frac{v_0^2 m + 2eV_0 - 2e\phi}{m}} \pi n c m + c_1 e \right)}}{e} \\ p &= \frac{\sqrt{2} \sqrt{e \left(4\sqrt{-\frac{v_0^2 m + 2eV_0 - 2e\phi}{m}} \pi n c m + c_1 e \right)}}{e} \end{aligned}$$

For solution (1) found earlier, since $p = \phi'$ then we now have a new first order ode to solve which is

$$\phi' = \frac{\sqrt{2} \sqrt{e \left(4 \sqrt{-\frac{v_0^2 m - 2e\phi + 2eV_0}{m}} \pi n c m + c_1 e \right)}}{e}$$

Integrating gives

$$\int \frac{e\sqrt{2}}{2 \sqrt{e \left(4 \sqrt{-\frac{v_0^2 m + 2eV_0 - 2e\phi}{m}} \pi n c m + c_1 e \right)}} d\phi = dx$$

$$\frac{\sqrt{2} \sqrt{4ecm\pi \sqrt{\frac{(-2V_0 + 2\phi)e + v_0^2 m}{m}} + e^2 c_1} \left(2 \sqrt{\frac{(-2V_0 + 2\phi)e + v_0^2 m}{m}} \pi n c m - c_1 e \right)}{24e n^2 m c^2 \pi^2} = x + c_2$$

For solution (2) found earlier, since $p = \phi'$ then we now have a new first order ode to solve which is

$$\phi' = -\frac{\sqrt{2} \sqrt{e \left(4 \sqrt{-\frac{v_0^2 m - 2e\phi + 2eV_0}{m}} \pi n c m + c_1 e \right)}}{e}$$

Integrating gives

$$\int -\frac{e\sqrt{2}}{2 \sqrt{e \left(4 \sqrt{-\frac{v_0^2 m + 2eV_0 - 2e\phi}{m}} \pi n c m + c_1 e \right)}} d\phi = dx$$

$$\frac{\sqrt{4ecm\pi \sqrt{\frac{(-2V_0 + 2\phi)e + v_0^2 m}{m}} + e^2 c_1} \left(\sqrt{\frac{(-2V_0 + 2\phi)e + v_0^2 m}{m}} \pi n c m - \frac{c_1 e}{2} \right) \sqrt{2}}{12e n^2 m c^2 \pi^2} = x + c_3$$

Will add steps showing solving for IC soon.

Summary of solutions found

$$\frac{\sqrt{2} \sqrt{4ecm\pi \sqrt{\frac{(-2V_0 + 2\phi)e + v_0^2 m}{m}} + e^2 c_1} \left(2 \sqrt{\frac{(-2V_0 + 2\phi)e + v_0^2 m}{m}} \pi n c m - c_1 e \right)}{24e n^2 m c^2 \pi^2} = x + c_2$$

$$\frac{\sqrt{4ecm\pi \sqrt{\frac{(-2V_0 + 2\phi)e + v_0^2 m}{m}} + e^2 c_1} \left(\sqrt{\frac{(-2V_0 + 2\phi)e + v_0^2 m}{m}} \pi n c m - \frac{c_1 e}{2} \right) \sqrt{2}}{12e n^2 m c^2 \pi^2} = x + c_3$$

Solved as second order can be made integrable

Time used: 5.401 (sec)

Multiplying the ode by ϕ' gives

$$\phi' \phi'' - \frac{4\phi' \pi n c}{\sqrt{\frac{v_0^2 m + 2e\phi - 2eV_0}{m}}} = 0$$

Integrating the above w.r.t x gives

$$\int \left(\phi' \phi'' - \frac{4\phi' \pi n c}{\sqrt{\frac{v_0^2 m + 2e\phi - 2eV_0}{m}}} \right) dx = 0$$

$$\frac{\phi'^2}{2} - \frac{4\pi n c \sqrt{\frac{2e\phi}{m} + \frac{v_0^2 m - 2eV_0}{m}} m}{e} = c_1$$

Which is now solved for ϕ . Solving for the derivative gives these ODE's to solve

$$\phi' = \frac{\sqrt{2} \sqrt{e \left(4 \sqrt{\frac{v_0^2 m + 2e\phi - 2eV_0}{m}} \pi n c m + c_1 e \right)}}{e} \quad (1)$$

$$\phi' = - \frac{\sqrt{2} \sqrt{e \left(4 \sqrt{\frac{v_0^2 m + 2e\phi - 2eV_0}{m}} \pi n c m + c_1 e \right)}}{e} \quad (2)$$

Now each of the above is solved separately.

Solving Eq. (1)

Integrating gives

$$\int \frac{e\sqrt{2}}{2\sqrt{e \left(4 \sqrt{\frac{v_0^2 m - 2eV_0 + 2e\phi}{m}} \pi n c m + c_1 e \right)}} d\phi = dx$$

$$\frac{\sqrt{2} \sqrt{4ecm\pi \sqrt{\frac{(-2V_0 + 2\phi)e + v_0^2 m}{m}} + e^2 c_1} \left(2\sqrt{\frac{(-2V_0 + 2\phi)e + v_0^2 m}{m}} \pi n c m - c_1 e \right)}{24e n^2 m c^2 \pi^2} = x + c_2$$

Solving Eq. (2)

Integrating gives

$$\int -\frac{e\sqrt{2}}{2\sqrt{e\left(4\sqrt{\frac{v_0^2 m - 2eV_0 + 2e\phi}{m}} \pi ncm + c_1 e\right)}} d\phi = dx$$

$$-\frac{\sqrt{4ecmn\pi\sqrt{\frac{(-2V_0 + 2\phi)e + v_0^2 m}{m}} + e^2 c_1} \left(\sqrt{\frac{(-2V_0 + 2\phi)e + v_0^2 m}{m}} \pi ncm - \frac{c_1 e}{2}\right) \sqrt{2}}{12e n^2 m c^2 \pi^2} = x + c_3$$

Will add steps showing solving for IC soon.

Summary of solutions found

$$\frac{\sqrt{2} \sqrt{4ecmn\pi\sqrt{\frac{(-2V_0 + 2\phi)e + v_0^2 m}{m}} + e^2 c_1} \left(2\sqrt{\frac{(-2V_0 + 2\phi)e + v_0^2 m}{m}} \pi ncm - c_1 e\right)}{24e n^2 m c^2 \pi^2} = x + c_2$$

$$-\frac{\sqrt{4ecmn\pi\sqrt{\frac{(-2V_0 + 2\phi)e + v_0^2 m}{m}} + e^2 c_1} \left(\sqrt{\frac{(-2V_0 + 2\phi)e + v_0^2 m}{m}} \pi ncm - \frac{c_1 e}{2}\right) \sqrt{2}}{12e n^2 m c^2 \pi^2} = x + c_3$$

Maple step by step solution

Let's solve

$$\phi'' = \frac{4\pi n c}{\sqrt{v_0^2 + \frac{2e(\phi - V_0)}{m}}}$$

- Highest derivative means the order of the ODE is 2

$$\phi''$$

- Define new dependent variable u

$$u(x) = \phi'$$

- Compute ϕ''

$$u'(x) = \phi''$$

- Use chain rule on the lhs

$$\phi' \left(\frac{d}{d\phi} u(\phi) \right) = \phi''$$

- Substitute in the definition of u

$$u(\phi) \left(\frac{d}{d\phi} u(\phi) \right) = \phi''$$

- Make substitutions $\phi' = u(\phi)$, $\phi'' = u(\phi) \left(\frac{d}{d\phi} u(\phi) \right)$ to reduce order of ODE

$$u(\phi) \left(\frac{d}{d\phi} u(\phi) \right) = \frac{4\pi nc}{\sqrt{v_0^2 + \frac{2e(\phi - V_0)}{m}}}$$

- Integrate both sides with respect to ϕ

$$\int u(\phi) \left(\frac{d}{d\phi} u(\phi) \right) d\phi = \int \frac{4\pi nc}{\sqrt{v_0^2 + \frac{2e(\phi - V_0)}{m}}} d\phi + C1$$

- Evaluate integral

$$\frac{u(\phi)^2}{2} = -\frac{4(-v_0^2 m + 2eV_0 - 2e\phi)\pi nc}{e\sqrt{-\frac{v_0^2 m + 2eV_0 - 2e\phi}{m}}} + C1$$

- Solve for $u(\phi)$

$$\left\{ u(\phi) = \frac{\sqrt{-2e\sqrt{-\frac{v_0^2 m + 2eV_0 - 2e\phi}{m}} \left(-4\pi nc m v_0^2 + 8V_0 \pi nc e - 8\phi \pi nc e - C1 e \sqrt{-\frac{v_0^2 m + 2eV_0 - 2e\phi}{m}} \right)}}{e\sqrt{-\frac{v_0^2 m + 2eV_0 - 2e\phi}{m}}}, u(\phi) = -\sqrt{-2e\sqrt{-\frac{v_0^2 m + 2eV_0 - 2e\phi}{m}} \left(-4\pi nc m v_0^2 + 8V_0 \pi nc e - 8\phi \pi nc e - C1 e \sqrt{-\frac{v_0^2 m + 2eV_0 - 2e\phi}{m}} \right)} \right.$$

- Solve 1st ODE for $u(\phi)$

$$u(\phi) = \frac{\sqrt{-2e\sqrt{-\frac{v_0^2 m + 2eV_0 - 2e\phi}{m}} \left(-4\pi nc m v_0^2 + 8V_0 \pi nc e - 8\phi \pi nc e - C1 e \sqrt{-\frac{v_0^2 m + 2eV_0 - 2e\phi}{m}} \right)}}{e\sqrt{-\frac{v_0^2 m + 2eV_0 - 2e\phi}{m}}}$$

- Revert to original variables with substitution $u(\phi) = \phi'$, $\phi = \phi$

$$\phi' = \frac{\sqrt{-2e\sqrt{-\frac{v_0^2 m - 2e\phi + 2eV_0}{m}} \left(-4\pi nc m v_0^2 + 8V_0 \pi nc e - 8\phi \pi nc e - C1 e \sqrt{-\frac{v_0^2 m - 2e\phi + 2eV_0}{m}} \right)}}{e\sqrt{-\frac{v_0^2 m - 2e\phi + 2eV_0}{m}}}$$

- Solve for the highest derivative

$$\phi' = \frac{\sqrt{-2e\sqrt{-\frac{v_0^2 m - 2e\phi + 2eV_0}{m}} \left(-4\pi nc m v_0^2 + 8V_0 \pi nc e - 8\phi \pi nc e - C1 e \sqrt{-\frac{v_0^2 m - 2e\phi + 2eV_0}{m}} \right)}}{e\sqrt{-\frac{v_0^2 m - 2e\phi + 2eV_0}{m}}}$$

- Separate variables

$$\frac{\phi' \sqrt{-\frac{v_0^2 m - 2e\phi + 2eV_0}{m}}}{\sqrt{-2e\sqrt{-\frac{v_0^2 m - 2e\phi + 2eV_0}{m}} \left(-4\pi nc m v_0^2 + 8V_0 \pi nc e - 8\phi \pi nc e - C1 e \sqrt{-\frac{v_0^2 m - 2e\phi + 2eV_0}{m}} \right)}} = \frac{1}{e}$$

- Integrate both sides with respect to x

$$\int \frac{\phi' \sqrt{-\frac{v_0^2 m - 2e\phi + 2eV_0}{m}}}{\sqrt{-2e\sqrt{-\frac{v_0^2 m - 2e\phi + 2eV_0}{m}} \left(-4\pi nc m v_0^2 + 8V_0 \pi nc e - 8\phi \pi nc e - C1 e \sqrt{-\frac{v_0^2 m - 2e\phi + 2eV_0}{m}} \right)}} dx = \int \frac{1}{e} dx + C2$$

- Evaluate integral

$$\frac{\left(-4\pi ncm v_0^2+8V_0\pi nce-8\phi\pi nce-C1e\sqrt{-\frac{-v_0^2m-2e\phi+2eV_0}{m}}\right)\sqrt{\frac{(-v_0^2m-2e\phi+2eV_0)\left(4\sqrt{-\frac{-v_0^2m-2e\phi+2eV_0}{m}}\pi ncm+C1\right)}{m}}}{12\sqrt{-2e\sqrt{-\frac{-v_0^2m-2e\phi+2eV_0}{m}}\left(-4\pi ncm v_0^2+8V_0\pi nce-8\phi\pi nce-C1e\sqrt{-\frac{-v_0^2m-2e\phi+2eV_0}{m}}\right)}e\sqrt{\sqrt{-\frac{-v_0^2m-2e\phi+2eV_0}{m}}\left(-4\pi ncm v_0^2\right)}}$$

- Solve for ϕ

$$\phi = \frac{-v_0^2m^2n^2c^2+2eV_0n^2mc^2+\left(\frac{e\left(576\pi^4C_2^2c^4e^2m^2n^4+1152\pi^4C_2c^4em^2n^4x+576\pi^4c^4m^2n^4x^2+24\sqrt{2}c^2mn^2\sqrt{288C_2^4\pi^4c^4e^4m^2n^4+1152\pi^4C_2^2c^4e^2m^2n^4x+576\pi^4c^4m^2n^4x^2}\right)}{24\sqrt{2}c^2mn^2}\right)}{24\sqrt{2}c^2mn^2}$$

- Solve 2nd ODE for $u(\phi)$

$$u(\phi) = -\frac{\sqrt{-2e\sqrt{-\frac{-v_0^2m+2eV_0-2e\phi}{m}}\left(-4\pi ncm v_0^2+8V_0\pi nce-8\phi\pi nce-C1e\sqrt{-\frac{-v_0^2m+2eV_0-2e\phi}{m}}\right)}}{e\sqrt{-\frac{-v_0^2m+2eV_0-2e\phi}{m}}}$$

- Revert to original variables with substitution $u(\phi) = \phi', \phi = \phi$

$$\phi' = -\frac{\sqrt{-2e\sqrt{-\frac{-v_0^2m-2e\phi+2eV_0}{m}}\left(-4\pi ncm v_0^2+8V_0\pi nce-8\phi\pi nce-C1e\sqrt{-\frac{-v_0^2m-2e\phi+2eV_0}{m}}\right)}}{e\sqrt{-\frac{-v_0^2m-2e\phi+2eV_0}{m}}}$$

- Solve for the highest derivative

$$\phi' = -\frac{\sqrt{-2e\sqrt{-\frac{-v_0^2m-2e\phi+2eV_0}{m}}\left(-4\pi ncm v_0^2+8V_0\pi nce-8\phi\pi nce-C1e\sqrt{-\frac{-v_0^2m-2e\phi+2eV_0}{m}}\right)}}{e\sqrt{-\frac{-v_0^2m-2e\phi+2eV_0}{m}}}$$

- Separate variables

$$\frac{\phi'\sqrt{-\frac{-v_0^2m-2e\phi+2eV_0}{m}}}{\sqrt{-2e\sqrt{-\frac{-v_0^2m-2e\phi+2eV_0}{m}}\left(-4\pi ncm v_0^2+8V_0\pi nce-8\phi\pi nce-C1e\sqrt{-\frac{-v_0^2m-2e\phi+2eV_0}{m}}\right)}} = -\frac{1}{e}$$

- Integrate both sides with respect to x

$$\int \frac{\phi'\sqrt{-\frac{-v_0^2m-2e\phi+2eV_0}{m}}}{\sqrt{-2e\sqrt{-\frac{-v_0^2m-2e\phi+2eV_0}{m}}\left(-4\pi ncm v_0^2+8V_0\pi nce-8\phi\pi nce-C1e\sqrt{-\frac{-v_0^2m-2e\phi+2eV_0}{m}}\right)}} dx = \int -\frac{1}{e} dx + C_2$$

- Evaluate integral

$$\frac{\left(-4\pi ncm v_0^2+8V_0\pi nce-8\phi\pi nce-C1e\sqrt{-\frac{-v_0^2m-2e\phi+2eV_0}{m}}\right)\sqrt{\frac{(-v_0^2m-2e\phi+2eV_0)\left(4\sqrt{-\frac{-v_0^2m-2e\phi+2eV_0}{m}}\pi ncm+C1\right)}{m}}}{12\sqrt{-2e\sqrt{-\frac{-v_0^2m-2e\phi+2eV_0}{m}}\left(-4\pi ncm v_0^2+8V_0\pi nce-8\phi\pi nce-C1e\sqrt{-\frac{-v_0^2m-2e\phi+2eV_0}{m}}\right)}e\sqrt{\sqrt{-\frac{-v_0^2m-2e\phi+2eV_0}{m}}\left(-4\pi ncm v_0^2\right)}}$$

- Solve for ϕ

Mathematica DSolve solution

Solving time : 73.891 (sec)

Leaf size : 2754

```
DSolve[{D[phi[x],{x,2}]==4*Pi*n*c/Sqrt[v0^2+2*e/m*(phi[x]-V0)],{}},  
phi[x],x,IncludeSingularSolutions->True]
```

Too large to display

2.6.5 problem 8 (eq 68)

Solved as first order separable ode	424
Solved as first order Exact ode	426
Solved using Lie symmetry for first order ode	430
Solved as first order ode of type Riccati	435
Maple step by step solution	439
Maple trace	439
Maple dsolve solution	440
Mathematica DSolve solution	440

Internal problem ID [18245]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter IV. Methods of solution: First order equations. section 33. Problems at page 91

Problem number : 8 (eq 68)

Date solved : Monday, December 23, 2024 at 09:19:07 PM

CAS classification : [_separable]

Solve

$$y' = x(ay^2 + b)$$

Solved as first order separable ode

Time used: 0.303 (sec)

The ode $y' = x(ay^2 + b)$ is separable as it can be written as

$$\begin{aligned} y' &= x(ay^2 + b) \\ &= f(x)g(y) \end{aligned}$$

Where

$$\begin{aligned} f(x) &= x \\ g(y) &= ay^2 + b \end{aligned}$$

Integrating gives

$$\int \frac{1}{g(y)} dy = \int f(x) dx$$

$$\int \frac{1}{ay^2 + b} dy = \int x dx$$

$$\frac{\arctan\left(\frac{ay}{\sqrt{ba}}\right)}{\sqrt{ba}} = \frac{x^2}{2} + c_1$$

We now need to find the singular solutions, these are found by finding for what values $g(y)$ is zero, since we had to divide by this above. Solving $g(y) = 0$ or $ay^2 + b = 0$ for y gives

$$y = \frac{\sqrt{-ba}}{a}$$

$$y = -\frac{\sqrt{-ba}}{a}$$

Now we go over each such singular solution and check if it verifies the ode itself and any initial conditions given. If it does not then the singular solution will not be used.

Therefore the solutions found are

$$\frac{\arctan\left(\frac{ay}{\sqrt{ba}}\right)}{\sqrt{ba}} = \frac{x^2}{2} + c_1$$

$$y = \frac{\sqrt{-ba}}{a}$$

$$y = -\frac{\sqrt{-ba}}{a}$$

Solving for y gives

$$y = \frac{\sqrt{-ba}}{a}$$

$$y = \frac{\tan\left(\frac{x^2\sqrt{ba}}{2} + c_1\sqrt{ba}\right)\sqrt{ba}}{a}$$

$$y = -\frac{\sqrt{-ba}}{a}$$

Summary of solutions found

$$y = \frac{\sqrt{-ba}}{a}$$

$$y = \frac{\tan\left(\frac{x^2\sqrt{ba}}{2} + c_1\sqrt{ba}\right)\sqrt{ba}}{a}$$

$$y = -\frac{\sqrt{-ba}}{a}$$

Solved as first order Exact ode

Time used: 0.373 (sec)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx}\phi(x, y) = 0$$

Hence

$$\frac{\partial\phi}{\partial x} + \frac{\partial\phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial\phi}{\partial x} &= M \\ \frac{\partial\phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might

or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned} dy &= (x(a y^2 + b)) dx \\ (-x(a y^2 + b)) dx + dy &= 0 \end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= -x(a y^2 + b) \\ N(x, y) &= 1 \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(-x(a y^2 + b)) \\ &= -2xay \end{aligned}$$

And

$$\begin{aligned} \frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(1) \\ &= 0 \end{aligned}$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned} A &= \frac{1}{N} \left(\frac{\partial M}{\partial y} - \frac{\partial N}{\partial x} \right) \\ &= 1((-2xay) - (0)) \\ &= -2xay \end{aligned}$$

Since A depends on y , it can not be used to obtain an integrating factor. We will now try a second method to find an integrating factor. Let

$$\begin{aligned} B &= \frac{1}{M} \left(\frac{\partial N}{\partial x} - \frac{\partial M}{\partial y} \right) \\ &= -\frac{1}{x(a y^2 + b)}((0) - (-2xay)) \\ &= -\frac{2ay}{a y^2 + b} \end{aligned}$$

Since B does not depend on x , it can be used to obtain an integrating factor. Let the integrating factor be μ . Then

$$\begin{aligned}\mu &= e^{\int B \, dy} \\ &= e^{\int -\frac{2ay}{ay^2+b} \, dy}\end{aligned}$$

The result of integrating gives

$$\begin{aligned}\mu &= e^{-\ln(ay^2+b)} \\ &= \frac{1}{ay^2+b}\end{aligned}$$

M and N are now multiplied by this integrating factor, giving new M and new N which are called \bar{M} and \bar{N} so not to confuse them with the original M and N .

$$\begin{aligned}\bar{M} &= \mu M \\ &= \frac{1}{ay^2+b}(-x(ay^2+b)) \\ &= -x\end{aligned}$$

And

$$\begin{aligned}\bar{N} &= \mu N \\ &= \frac{1}{ay^2+b}(1) \\ &= \frac{1}{ay^2+b}\end{aligned}$$

So now a modified ODE is obtained from the original ODE which will be exact and can be solved using the standard method. The modified ODE is

$$\begin{aligned}\bar{M} + \bar{N} \frac{dy}{dx} &= 0 \\ (-x) + \left(\frac{1}{ay^2+b} \right) \frac{dy}{dx} &= 0\end{aligned}$$

The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = \bar{M} \tag{1}$$

$$\frac{\partial \phi}{\partial y} = \bar{N} \tag{2}$$

Integrating (1) w.r.t. x gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int \bar{M} dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int -x dx \\ \phi &= -\frac{x^2}{2} + f(y)\end{aligned}\quad (3)$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\frac{\partial \phi}{\partial y} = 0 + f'(y) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial y} = \frac{1}{ay^2+b}$. Therefore equation (4) becomes

$$\frac{1}{ay^2+b} = 0 + f'(y) \quad (5)$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = \frac{1}{ay^2+b}$$

Integrating the above w.r.t y gives

$$\begin{aligned}\int f'(y) dy &= \int \left(\frac{1}{ay^2+b} \right) dy \\ f(y) &= \frac{\arctan\left(\frac{ay}{\sqrt{ba}}\right)}{\sqrt{ba}} + c_1\end{aligned}$$

Where c_1 is constant of integration. Substituting result found above for $f(y)$ into equation (3) gives ϕ

$$\phi = -\frac{x^2}{2} + \frac{\arctan\left(\frac{ay}{\sqrt{ba}}\right)}{\sqrt{ba}} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into the constant c_1 gives the solution as

$$c_1 = -\frac{x^2}{2} + \frac{\arctan\left(\frac{ay}{\sqrt{ba}}\right)}{\sqrt{ba}}$$

Solving for y gives

$$y = \frac{\tan\left(\frac{x^2\sqrt{ba}}{2} + c_1\sqrt{ba}\right)\sqrt{ba}}{a}$$

Summary of solutions found

$$y = \frac{\tan\left(\frac{x^2\sqrt{ba}}{2} + c_1\sqrt{ba}\right)\sqrt{ba}}{a}$$

Solved using Lie symmetry for first order ode

Time used: 1.043 (sec)

Writing the ode as

$$y' = x(a y^2 + b)$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

To determine ξ, η then (A) is solved using ansatz. Making bivariate polynomials of degree 2 to use as anstaz gives

$$\xi = x^2 a_4 + x y a_5 + y^2 a_6 + x a_2 + y a_3 + a_1 \quad (\text{1E})$$

$$\eta = x^2 b_4 + x y b_5 + y^2 b_6 + x b_2 + y b_3 + b_1 \quad (\text{2E})$$

Where the unknown coefficients are

$$\{a_1, a_2, a_3, a_4, a_5, a_6, b_1, b_2, b_3, b_4, b_5, b_6\}$$

Substituting equations (1E,2E) and ω into (A) gives

$$\begin{aligned}
& 2xb_4 + yb_5 + b_2 + x(ay^2 + b)(-2xa_4 + xb_5 - ya_5 + 2yb_6 - a_2 + b_3) \\
& - x^2(ay^2 + b)^2(xa_5 + 2ya_6 + a_3) \\
& - (ay^2 + b)(x^2a_4 + xya_5 + y^2a_6 + xa_2 + ya_3 + a_1) \\
& - 2xay(x^2b_4 + xyb_5 + y^2b_6 + xb_2 + yb_3 + b_1) = 0
\end{aligned} \tag{5E}$$

Putting the above in normal form gives

$$\begin{aligned}
& -a^2x^3y^4a_5 - 2a^2x^2y^5a_6 - a^2x^2y^4a_3 - 2abx^3y^2a_5 - 4abx^2y^3a_6 - 2abx^2y^2a_3 \\
& - 2ax^3yb_4 - 3ax^2y^2a_4 - ax^2y^2b_5 - 2ax^3y^3a_5 - ay^4a_6 - b^2x^3a_5 - 2b^2x^2ya_6 \\
& - 2ax^2yb_2 - 2ax^2y^2a_2 - ax^2y^2b_3 - ay^3a_3 - b^2x^2a_3 - 2axyb_1 - ay^2a_1 - 3bx^2a_4 \\
& + bx^2b_5 - 2bxya_5 + 2bxyb_6 - by^2a_6 - 2bxa_2 + bxb_3 - bya_3 - ba_1 + 2xb_4 + yb_5 + b_2 = 0
\end{aligned}$$

Setting the numerator to zero gives

$$\begin{aligned}
& -a^2x^3y^4a_5 - 2a^2x^2y^5a_6 - a^2x^2y^4a_3 - 2abx^3y^2a_5 - 4abx^2y^3a_6 \\
& - 2abx^2y^2a_3 - 2ax^3yb_4 - 3ax^2y^2a_4 - ax^2y^2b_5 - 2ax^3y^3a_5 - ay^4a_6 \\
& - b^2x^3a_5 - 2b^2x^2ya_6 - 2ax^2yb_2 - 2ax^2y^2a_2 - ax^2y^2b_3 - ay^3a_3 \\
& - b^2x^2a_3 - 2axyb_1 - ay^2a_1 - 3bx^2a_4 + bx^2b_5 - 2bxya_5 + 2bxyb_6 \\
& - by^2a_6 - 2bxa_2 + bxb_3 - bya_3 - ba_1 + 2xb_4 + yb_5 + b_2 = 0
\end{aligned} \tag{6E}$$

Looking at the above PDE shows the following are all the terms with $\{x, y\}$ in them.

$$\{x, y\}$$

The following substitution is now made to be able to collect on all terms with $\{x, y\}$ in them

$$\{x = v_1, y = v_2\}$$

The above PDE (6E) now becomes

$$\begin{aligned}
& -a^2a_5v_1^3v_2^4 - 2a^2a_6v_1^2v_2^5 - a^2a_3v_1^2v_2^4 - 2aba_5v_1^3v_2^2 - 4aba_6v_1^2v_2^3 \\
& - 2aba_3v_1^2v_2^2 - 3aa_4v_1^2v_2^2 - 2aa_5v_1v_2^3 - aa_6v_2^4 - 2ab_4v_1^3v_2 - ab_5v_1^2v_2^2 \\
& - b^2a_5v_1^3 - 2b^2a_6v_1^2v_2 - 2aa_2v_1v_2^2 - aa_3v_2^3 - 2ab_2v_1^2v_2 - ab_3v_1v_2^2 \\
& - b^2a_3v_1^2 - aa_1v_2^2 - 2ab_1v_1v_2 - 3ba_4v_1^2 - 2ba_5v_1v_2 - ba_6v_2^2 + bb_5v_1^2 \\
& + 2bb_6v_1v_2 - 2ba_2v_1 - ba_3v_2 + bb_3v_1 - ba_1 + 2b_4v_1 + b_5v_2 + b_2 = 0
\end{aligned} \tag{7E}$$

Collecting the above on the terms v_i introduced, and these are

$$\{v_1, v_2\}$$

Equation (7E) now becomes

$$\begin{aligned} & -a^2 a_5 v_1^3 v_2^4 - 2aba_5 v_1^3 v_2^2 - 2ab_4 v_1^3 v_2 - b^2 a_5 v_1^3 - 2a^2 a_6 v_1^2 v_2^5 - a^2 a_3 v_1^2 v_2^4 \\ & - 4aba_6 v_1^2 v_2^3 + (-2aba_3 - 3aa_4 - ab_5) v_1^2 v_2^2 + (-2b^2 a_6 - 2ab_2) v_1^2 v_2 \\ & + (-b^2 a_3 - 3ba_4 + bb_5) v_1^2 - 2aa_5 v_1 v_2^3 + (-2aa_2 - ab_3) v_1 v_2^2 \\ & + (-2ab_1 - 2ba_5 + 2bb_6) v_1 v_2 + (-2ba_2 + bb_3 + 2b_4) v_1 - aa_6 v_2^4 \\ & - aa_3 v_2^3 + (-aa_1 - ba_6) v_2^2 + (-ba_3 + b_5) v_2 - ba_1 + b_2 = 0 \end{aligned} \quad (8E)$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$\begin{aligned} & -aa_3 = 0 \\ & -2aa_5 = 0 \\ & -aa_6 = 0 \\ & -2ab_4 = 0 \\ & -a^2 a_3 = 0 \\ & -a^2 a_5 = 0 \\ & -2a^2 a_6 = 0 \\ & -b^2 a_5 = 0 \\ & -2aba_5 = 0 \\ & -4aba_6 = 0 \\ & -2aa_2 - ab_3 = 0 \\ & -ba_1 + b_2 = 0 \\ & -ba_3 + b_5 = 0 \\ & -aa_1 - ba_6 = 0 \\ & -2b^2 a_6 - 2ab_2 = 0 \\ & -2ba_2 + bb_3 + 2b_4 = 0 \\ & -b^2 a_3 - 3ba_4 + bb_5 = 0 \\ & -2aba_3 - 3aa_4 - ab_5 = 0 \\ & -2ab_1 - 2ba_5 + 2bb_6 = 0 \end{aligned}$$

Solving the above equations for the unknowns gives

$$\begin{aligned}
 a_1 &= 0 \\
 a_2 &= 0 \\
 a_3 &= 0 \\
 a_4 &= 0 \\
 a_5 &= 0 \\
 a_6 &= 0 \\
 b_1 &= b_1 \\
 b_2 &= 0 \\
 b_3 &= 0 \\
 b_4 &= 0 \\
 b_5 &= 0 \\
 b_6 &= \frac{ab_1}{b}
 \end{aligned}$$

Substituting the above solution in the anstaz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$\begin{aligned}
 \xi &= 0 \\
 \eta &= \frac{ay^2 + b}{b}
 \end{aligned}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS \tag{1}$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = x$$

S is found from

$$\begin{aligned}
 S &= \int \frac{1}{\eta} dy \\
 &= \int \frac{1}{\frac{ay^2+b}{b}} dy
 \end{aligned}$$

Which results in

$$S = \frac{b \arctan\left(\frac{ay}{\sqrt{ba}}\right)}{\sqrt{ba}}$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = x(ay^2 + b)$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 1 \\ R_y &= 0 \\ S_x &= 0 \\ S_y &= \frac{b}{ay^2 + b} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = bR \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = bR$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S .

Since the ode has the form $\frac{d}{dR}S(R) = f(R)$, then we only need to integrate $f(R)$.

$$\begin{aligned} \int dS &= \int bR dR \\ S(R) &= \frac{bR^2}{2} + c_2 \end{aligned}$$

To complete the solution, we just need to transform the above back to x, y coordinates. This results in

$$\frac{\sqrt{b} \arctan\left(\frac{\sqrt{a}y}{\sqrt{b}}\right)}{\sqrt{a}} = \frac{bx^2}{2} + c_2$$

Which gives

$$y = \frac{\sqrt{b} \tan\left(\frac{\sqrt{a}(bx^2+2c_2)}{2\sqrt{b}}\right)}{\sqrt{a}}$$

Summary of solutions found

$$y = \frac{\sqrt{b} \tan\left(\frac{\sqrt{a}(bx^2+2c_2)}{2\sqrt{b}}\right)}{\sqrt{a}}$$

Solved as first order ode of type Riccati

Time used: 0.482 (sec)

In canonical form the ODE is

$$\begin{aligned} y' &= F(x, y) \\ &= x(ay^2 + b) \end{aligned}$$

This is a Riccati ODE. Comparing the ODE to solve

$$y' = axy^2 + bx$$

With Riccati ODE standard form

$$y' = f_0(x) + f_1(x)y + f_2(x)y^2$$

Shows that $f_0(x) = bx$, $f_1(x) = 0$ and $f_2(x) = ax$. Let

$$\begin{aligned} y &= \frac{-u'}{f_2u} \\ &= \frac{-u'}{uax} \end{aligned} \tag{1}$$

Using the above substitution in the given ODE results (after some simplification) in a second order ODE to solve for $u(x)$ which is

$$f_2u''(x) - (f_2' + f_1f_2)u'(x) + f_2^2f_0u(x) = 0 \tag{2}$$

But

$$\begin{aligned}f_2' &= a \\f_1 f_2 &= 0 \\f_2^2 f_0 &= a^2 x^3 b\end{aligned}$$

Substituting the above terms back in equation (2) gives

$$axu''(x) - au'(x) + a^2x^3bu(x) = 0$$

In normal form the ode

$$ax \left(\frac{d^2u}{dx^2} \right) - a \left(\frac{du}{dx} \right) + a^2x^3bu = 0 \quad (1)$$

Becomes

$$\frac{d^2u}{dx^2} + p(x) \left(\frac{du}{dx} \right) + q(x)u = 0 \quad (2)$$

Where

$$\begin{aligned}p(x) &= -\frac{1}{x} \\q(x) &= abx^2\end{aligned}$$

Applying change of variables $\tau = g(x)$ to (2) gives

$$\frac{d^2}{d\tau^2}u(\tau) + p_1 \left(\frac{d}{d\tau}u(\tau) \right) + q_1u(\tau) = 0 \quad (3)$$

Where τ is the new independent variable, and

$$p_1(\tau) = \frac{\frac{d^2}{dx^2}\tau(x) + p(x) \left(\frac{d}{dx}\tau(x) \right)}{\left(\frac{d}{dx}\tau(x) \right)^2} \quad (4)$$

$$q_1(\tau) = \frac{q(x)}{\left(\frac{d}{dx}\tau(x) \right)^2} \quad (5)$$

Let $p_1 = 0$. Eq (4) simplifies to

$$\frac{d^2}{dx^2}\tau(x) + p(x) \left(\frac{d}{dx}\tau(x) \right) = 0$$

This ode is solved resulting in

$$\begin{aligned}
 \tau &= \int e^{-\int p(x)dx} dx \\
 &= \int e^{-\int -\frac{1}{x} dx} dx \\
 &= \int e^{\ln(x)} dx \\
 &= \int x dx \\
 &= \frac{x^2}{2}
 \end{aligned} \tag{6}$$

Using (6) to evaluate q_1 from (5) gives

$$\begin{aligned}
 q_1(\tau) &= \frac{q(x)}{\left(\frac{d}{dx}\tau(x)\right)^2} \\
 &= \frac{abx^2}{x^2} \\
 &= ba
 \end{aligned} \tag{7}$$

Substituting the above in (3) and noting that now $p_1 = 0$ results in

$$\begin{aligned}
 \frac{d^2}{d\tau^2}u(\tau) + q_1u(\tau) &= 0 \\
 \frac{d^2}{d\tau^2}u(\tau) + bau(\tau) &= 0
 \end{aligned}$$

The above ode is now solved for $u(\tau)$. This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Au''(\tau) + Bu'(\tau) + Cu(\tau) = 0$$

Where in the above $A = 1, B = 0, C = ba$. Let the solution be $u(\tau) = e^{\lambda\tau}$. Substituting this into the ODE gives

$$\lambda^2 e^{\lambda\tau} + ba e^{\lambda\tau} = 0 \tag{1}$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda\tau}$ gives

$$ba + \lambda^2 = 0 \tag{2}$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = 0, C = ba$ into the above gives

$$\begin{aligned}\lambda_{1,2} &= \frac{0}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{0^2 - (4)(1)(ba)} \\ &= \pm \sqrt{-ba}\end{aligned}$$

Hence

$$\begin{aligned}\lambda_1 &= +\sqrt{-ba} \\ \lambda_2 &= -\sqrt{-ba}\end{aligned}$$

Which simplifies to

$$\begin{aligned}\lambda_1 &= \sqrt{-ba} \\ \lambda_2 &= -\sqrt{-ba}\end{aligned}$$

Since roots are real and distinct, then the solution is

$$\begin{aligned}u(\tau) &= c_1 e^{\lambda_1 \tau} + c_2 e^{\lambda_2 \tau} \\ u(\tau) &= c_1 e^{(\sqrt{-ba})\tau} + c_2 e^{(-\sqrt{-ba})\tau}\end{aligned}$$

Or

$$u(\tau) = c_1 e^{\tau\sqrt{-ba}} + c_2 e^{-\tau\sqrt{-ba}}$$

Will add steps showing solving for IC soon.

The above solution is now transformed back to u using (6) which results in

$$u = c_1 e^{\frac{x^2\sqrt{-ba}}{2}} + c_2 e^{-\frac{x^2\sqrt{-ba}}{2}}$$

Will add steps showing solving for IC soon.

Taking derivative gives

$$u'(x) = c_1 x \sqrt{-ba} e^{\frac{x^2\sqrt{-ba}}{2}} - c_2 x \sqrt{-ba} e^{-\frac{x^2\sqrt{-ba}}{2}}$$

Doing change of constants, the solution becomes

$$y = -\frac{c_3 x \sqrt{-ba} e^{\frac{x^2\sqrt{-ba}}{2}} - x \sqrt{-ba} e^{-\frac{x^2\sqrt{-ba}}{2}}}{ax \left(c_3 e^{\frac{x^2\sqrt{-ba}}{2}} + e^{-\frac{x^2\sqrt{-ba}}{2}} \right)}$$

Summary of solutions found

$$y = -\frac{c_3 x \sqrt{-ba} e^{\frac{x^2 \sqrt{-ba}}{2}} - x \sqrt{-ba} e^{-\frac{x^2 \sqrt{-ba}}{2}}}{ax \left(c_3 e^{\frac{x^2 \sqrt{-ba}}{2}} + e^{-\frac{x^2 \sqrt{-ba}}{2}} \right)}$$

Maple step by step solution

Let's solve

$$y' = x(ay^2 + b)$$

- Highest derivative means the order of the ODE is 1

$$y'$$

- Solve for the highest derivative

$$y' = x(ay^2 + b)$$

- Separate variables

$$\frac{y'}{ay^2+b} = x$$

- Integrate both sides with respect to x

$$\int \frac{y'}{ay^2+b} dx = \int x dx + C1$$

- Evaluate integral

$$\frac{\arctan\left(\frac{ay}{\sqrt{ba}}\right)}{\sqrt{ba}} = \frac{x^2}{2} + C1$$

- Solve for y

$$y = \frac{\tan\left(\frac{x^2 \sqrt{ba}}{2} + C1 \sqrt{ba}\right) \sqrt{ba}}{a}$$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
<- separable successful`
```

Maple dsolve solution

Solving time : 0.006 (sec)

Leaf size : 28

```
dsolve(diff(y(x),x) = x*(a*y(x)^2+b),
        y(x),singsol=all)
```

$$y = \frac{\tan\left(\frac{\sqrt{ba}(x^2+2c_1)}{2}\right)\sqrt{ba}}{a}$$

Mathematica DSolve solution

Solving time : 6.63 (sec)

Leaf size : 75

```
DSolve[{D[y[x],x]==x*(a*y[x]^2+b)},{}],
        y[x],x,IncludeSingularSolutions->True]
```

$$y(x) \rightarrow \frac{\sqrt{b} \tan\left(\frac{1}{2}\sqrt{a}\sqrt{b}(x^2 + 2c_1)\right)}{\sqrt{a}}$$

$$y(x) \rightarrow -\frac{i\sqrt{b}}{\sqrt{a}}$$

$$y(x) \rightarrow \frac{i\sqrt{b}}{\sqrt{a}}$$

2.6.6 problem 8 (eq 69)

Solved as first order separable ode	441
Solved as first order Exact ode	443
Solved using Lie symmetry for first order ode	448
Solved as first order ode of type Riccati	454
Maple step by step solution	459
Maple trace	459
Maple dsolve solution	459
Mathematica DSolve solution	460

Internal problem ID [18246]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter IV. Methods of solution: First order equations. section 33. Problems at page 91

Problem number : 8 (eq 69)

Date solved : Monday, December 23, 2024 at 09:19:10 PM

CAS classification : [_separable]

Solve

$$n' = (n^2 + 1) x$$

Solved as first order separable ode

Time used: 0.145 (sec)

The ode $n' = (n^2 + 1) x$ is separable as it can be written as

$$\begin{aligned} n' &= (n^2 + 1) x \\ &= f(x)g(n) \end{aligned}$$

Where

$$\begin{aligned} f(x) &= x \\ g(n) &= n^2 + 1 \end{aligned}$$

Integrating gives

$$\int \frac{1}{g(n)} dn = \int f(x) dx$$
$$\int \frac{1}{n^2 + 1} dn = \int x dx$$
$$\arctan(n) = \frac{x^2}{2} + c_1$$

We now need to find the singular solutions, these are found by finding for what values $g(n)$ is zero, since we had to divide by this above. Solving $g(n) = 0$ or $n^2 + 1 = 0$ for n gives

$$n = -i$$
$$n = i$$

Now we go over each such singular solution and check if it verifies the ode itself and any initial conditions given. If it does not then the singular solution will not be used.

Therefore the solutions found are

$$\arctan(n) = \frac{x^2}{2} + c_1$$
$$n = -i$$
$$n = i$$

Solving for n gives

$$n = -i$$
$$n = i$$
$$n = \tan\left(\frac{x^2}{2} + c_1\right)$$

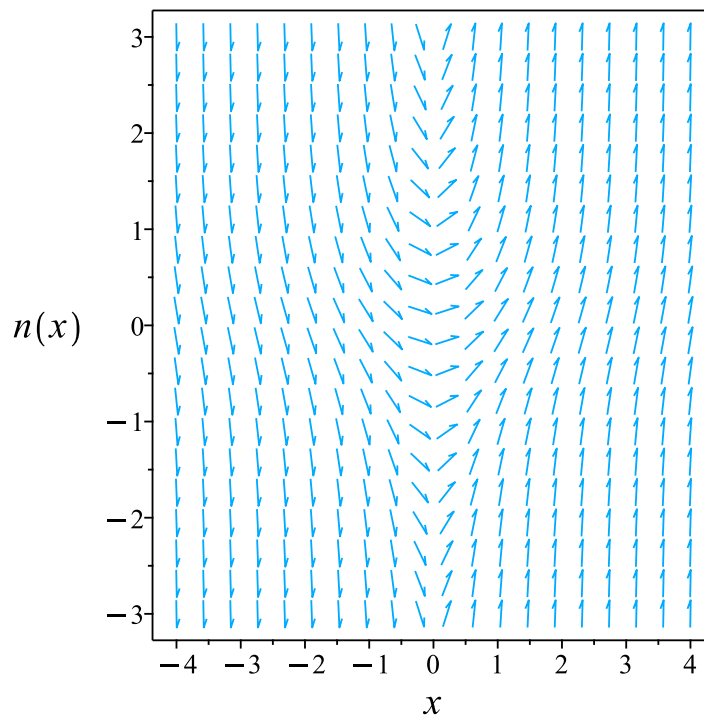


Figure 2.59: Slope field plot
 $n' = (n^2 + 1)x$

Summary of solutions found

$$n = -i$$

$$n = i$$

$$n = \tan\left(\frac{x^2}{2} + c_1\right)$$

Solved as first order Exact ode

Time used: 0.342 (sec)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \tag{A}$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx}\phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, n) dx + N(x, n) dn = 0 \quad (\text{1A})$$

Therefore

$$\begin{aligned} dn &= ((n^2 + 1) x) dx \\ (- (n^2 + 1) x) dx + dn &= 0 \end{aligned} \quad (\text{2A})$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, n) &= - (n^2 + 1) x \\ N(x, n) &= 1 \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial n} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial n} &= \frac{\partial}{\partial n} (-(n^2 + 1) x) \\ &= -2nx \end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(1) \\ &= 0\end{aligned}$$

Since $\frac{\partial M}{\partial n} \neq \frac{\partial N}{\partial x}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned}A &= \frac{1}{N} \left(\frac{\partial M}{\partial n} - \frac{\partial N}{\partial x} \right) \\ &= 1((-2nx) - (0)) \\ &= -2nx\end{aligned}$$

Since A depends on n , it can not be used to obtain an integrating factor. We will now try a second method to find an integrating factor. Let

$$\begin{aligned}B &= \frac{1}{M} \left(\frac{\partial N}{\partial x} - \frac{\partial M}{\partial n} \right) \\ &= -\frac{1}{(n^2 + 1)x} ((0) - (-2nx)) \\ &= -\frac{2n}{n^2 + 1}\end{aligned}$$

Since B does not depend on x , it can be used to obtain an integrating factor. Let the integrating factor be μ . Then

$$\begin{aligned}\mu &= e^{\int B \, dn} \\ &= e^{\int -\frac{2n}{n^2+1} \, dn}\end{aligned}$$

The result of integrating gives

$$\begin{aligned}\mu &= e^{-\ln(n^2+1)} \\ &= \frac{1}{n^2 + 1}\end{aligned}$$

M and N are now multiplied by this integrating factor, giving new M and new N which are called \bar{M} and \bar{N} so not to confuse them with the original M and N .

$$\begin{aligned}\bar{M} &= \mu M \\ &= \frac{1}{n^2 + 1} (-(n^2 + 1)x) \\ &= -x\end{aligned}$$

And

$$\begin{aligned}\bar{N} &= \mu N \\ &= \frac{1}{n^2 + 1} \quad (1) \\ &= \frac{1}{n^2 + 1}\end{aligned}$$

So now a modified ODE is obtained from the original ODE which will be exact and can be solved using the standard method. The modified ODE is

$$\begin{aligned}\bar{M} + \bar{N} \frac{dn}{dx} &= 0 \\ (-x) + \left(\frac{1}{n^2 + 1} \right) \frac{dn}{dx} &= 0\end{aligned}$$

The following equations are now set up to solve for the function $\phi(x, n)$

$$\frac{\partial \phi}{\partial x} = \bar{M} \quad (1)$$

$$\frac{\partial \phi}{\partial n} = \bar{N} \quad (2)$$

Integrating (1) w.r.t. x gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial x} dx &= \int \bar{M} dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int -x dx \\ \phi &= -\frac{x^2}{2} + f(n)\end{aligned} \quad (3)$$

Where $f(n)$ is used for the constant of integration since ϕ is a function of both x and n . Taking derivative of equation (3) w.r.t n gives

$$\frac{\partial \phi}{\partial n} = 0 + f'(n) \quad (4)$$

But equation (2) says that $\frac{\partial \phi}{\partial n} = \frac{1}{n^2 + 1}$. Therefore equation (4) becomes

$$\frac{1}{n^2 + 1} = 0 + f'(n) \quad (5)$$

Solving equation (5) for $f'(n)$ gives

$$f'(n) = \frac{1}{n^2 + 1}$$

Integrating the above w.r.t n gives

$$\int f'(n) \, dn = \int \left(\frac{1}{n^2 + 1} \right) \, dn$$
$$f(n) = \arctan(n) + c_1$$

Where c_1 is constant of integration. Substituting result found above for $f(n)$ into equation (3) gives ϕ

$$\phi = -\frac{x^2}{2} + \arctan(n) + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into the constant c_1 gives the solution as

$$c_1 = -\frac{x^2}{2} + \arctan(n)$$

Solving for n gives

$$n = \tan \left(\frac{x^2}{2} + c_1 \right)$$

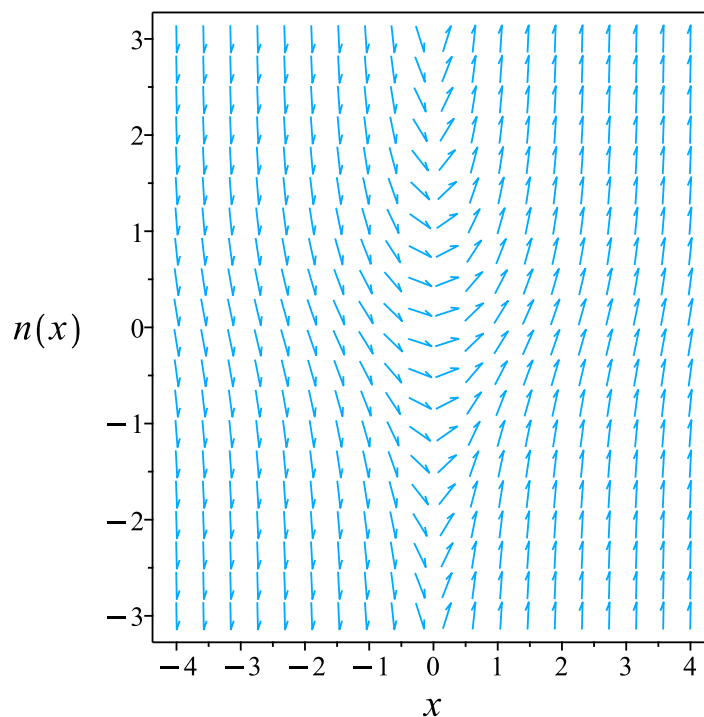


Figure 2.60: Slope field plot
 $n' = (n^2 + 1)x$

Summary of solutions found

$$n = \tan\left(\frac{x^2}{2} + c_1\right)$$

Solved using Lie symmetry for first order ode

Time used: 0.823 (sec)

Writing the ode as

$$\begin{aligned} n' &= (n^2 + 1)x \\ n' &= \omega(x, n) \end{aligned}$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_n - \xi_x) - \omega^2 \xi_n - \omega_x \xi - \omega_n \eta = 0 \quad (\text{A})$$

To determine ξ, η then (A) is solved using ansatz. Making bivariate polynomials of degree 2 to use as anstaz gives

$$\xi = n^2 a_6 + n x a_5 + x^2 a_4 + n a_3 + x a_2 + a_1 \quad (1\text{E})$$

$$\eta = n^2 b_6 + n x b_5 + x^2 b_4 + n b_3 + x b_2 + b_1 \quad (2\text{E})$$

Where the unknown coefficients are

$$\{a_1, a_2, a_3, a_4, a_5, a_6, b_1, b_2, b_3, b_4, b_5, b_6\}$$

Substituting equations (1E,2E) and ω into (A) gives

$$\begin{aligned} & nb_5 + 2xb_4 + b_2 + (n^2 + 1)x(-na_5 + 2nb_6 - 2xa_4 + xb_5 - a_2 + b_3) \\ & - (n^2 + 1)^2 x^2(2na_6 + xa_5 + a_3) \\ & - (n^2 + 1)(n^2a_6 + nxa_5 + x^2a_4 + na_3 + xa_2 + a_1) \\ & - 2nx(n^2b_6 + nxb_5 + x^2b_4 + nb_3 + xb_2 + b_1) = 0 \end{aligned} \quad (5E)$$

Putting the above in normal form gives

$$\begin{aligned} & -2n^5x^2a_6 - n^4x^3a_5 - n^4x^2a_3 - 4n^3x^2a_6 - 2n^2x^3a_5 - n^4a_6 - 2n^3xa_5 \\ & - 2n^2x^2a_3 - 3n^2x^2a_4 - n^2x^2b_5 - 2nx^3b_4 - n^3a_3 - 2n^2xa_2 - n^2xb_3 \\ & - 2nx^2a_6 - 2nx^2b_2 - x^3a_5 - n^2a_1 - n^2a_6 - 2nxa_5 - 2nxb_1 + 2nxb_6 \\ & - x^2a_3 - 3x^2a_4 + x^2b_5 - na_3 + nb_5 - 2xa_2 + xb_3 + 2xb_4 - a_1 + b_2 = 0 \end{aligned}$$

Setting the numerator to zero gives

$$\begin{aligned} & -2n^5x^2a_6 - n^4x^3a_5 - n^4x^2a_3 - 4n^3x^2a_6 - 2n^2x^3a_5 - n^4a_6 - 2n^3xa_5 \\ & - 2n^2x^2a_3 - 3n^2x^2a_4 - n^2x^2b_5 - 2nx^3b_4 - n^3a_3 - 2n^2xa_2 - n^2xb_3 \\ & - 2nx^2a_6 - 2nx^2b_2 - x^3a_5 - n^2a_1 - n^2a_6 - 2nxa_5 - 2nxb_1 + 2nxb_6 \\ & - x^2a_3 - 3x^2a_4 + x^2b_5 - na_3 + nb_5 - 2xa_2 + xb_3 + 2xb_4 - a_1 + b_2 = 0 \end{aligned} \quad (6E)$$

Looking at the above PDE shows the following are all the terms with $\{n, x\}$ in them.

$$\{n, x\}$$

The following substitution is now made to be able to collect on all terms with $\{n, x\}$ in them

$$\{n = v_1, x = v_2\}$$

The above PDE (6E) now becomes

$$\begin{aligned} & -a_5v_1^4v_2^3 - 2a_6v_1^5v_2^2 - a_3v_1^4v_2^2 - 2a_5v_1^2v_2^3 - 4a_6v_1^3v_2^2 - 2a_3v_1^2v_2^2 - 3a_4v_1^2v_2^2 \\ & - 2a_5v_1^3v_2 - a_6v_1^4 - 2b_4v_1v_2^3 - b_5v_1^2v_2^2 - 2a_2v_1^2v_2 - a_3v_1^3 - a_5v_2^3 - 2a_6v_1v_2^2 \\ & - 2b_2v_1v_2^2 - b_3v_1^2v_2 - a_1v_1^2 - a_3v_2^2 - 3a_4v_2^2 - 2a_5v_1v_2 - a_6v_1^2 - 2b_1v_1v_2 \\ & + b_5v_2^2 + 2b_6v_1v_2 - 2a_2v_2 - a_3v_1 + b_3v_2 + 2b_4v_2 + b_5v_1 - a_1 + b_2 = 0 \end{aligned} \quad (7E)$$

Collecting the above on the terms v_i introduced, and these are

$$\{v_1, v_2\}$$

Equation (7E) now becomes

$$\begin{aligned} & -2a_6v_1^5v_2^2 - a_5v_1^4v_2^3 - a_3v_1^4v_2^2 - a_6v_1^4 - 4a_6v_1^3v_2^2 - 2a_5v_1^3v_2 - a_3v_1^3 \\ & - 2a_5v_1^2v_2^3 + (-2a_3 - 3a_4 - b_5)v_1^2v_2^2 + (-2a_2 - b_3)v_1^2v_2 + (-a_1 - a_6)v_1^2 \quad (8E) \\ & - 2b_4v_1v_2^3 + (-2a_6 - 2b_2)v_1v_2^2 + (-2a_5 - 2b_1 + 2b_6)v_1v_2 + (-a_3 + b_5)v_1 \\ & - a_5v_2^3 + (-a_3 - 3a_4 + b_5)v_2^2 + (-2a_2 + b_3 + 2b_4)v_2 - a_1 + b_2 = 0 \end{aligned}$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$\begin{aligned} -a_3 &= 0 \\ -2a_5 &= 0 \\ -a_5 &= 0 \\ -4a_6 &= 0 \\ -2a_6 &= 0 \\ -a_6 &= 0 \\ -2b_4 &= 0 \\ -a_1 - a_6 &= 0 \\ -a_1 + b_2 &= 0 \\ -2a_2 - b_3 &= 0 \\ -a_3 + b_5 &= 0 \\ -2a_6 - 2b_2 &= 0 \\ -2a_2 + b_3 + 2b_4 &= 0 \\ -2a_3 - 3a_4 - b_5 &= 0 \\ -a_3 - 3a_4 + b_5 &= 0 \\ -2a_5 - 2b_1 + 2b_6 &= 0 \end{aligned}$$

Solving the above equations for the unknowns gives

$$a_1 = 0$$

$$a_2 = 0$$

$$a_3 = 0$$

$$a_4 = 0$$

$$a_5 = 0$$

$$a_6 = 0$$

$$b_1 = b_6$$

$$b_2 = 0$$

$$b_3 = 0$$

$$b_4 = 0$$

$$b_5 = 0$$

$$b_6 = b_6$$

Substituting the above solution in the anstaz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$\xi = 0$$

$$\eta = n^2 + 1$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, n) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dn}{\eta} = dS \quad (1)$$

The above comes from the requirements that $(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial n}) S(x, n) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = x$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{n^2 + 1} dy \end{aligned}$$

Which results in

$$S = \arctan(n)$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, n)S_n}{R_x + \omega(x, n)R_n} \quad (2)$$

Where in the above R_x, R_n, S_x, S_n are all partial derivatives and $\omega(x, n)$ is the right hand side of the original ode given by

$$\omega(x, n) = (n^2 + 1)x$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 1 \\ R_n &= 0 \\ S_x &= 0 \\ S_n &= \frac{1}{n^2 + 1} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = x \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, n in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = R$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S .

Since the ode has the form $\frac{d}{dR}S(R) = f(R)$, then we only need to integrate $f(R)$.

$$\begin{aligned} \int dS &= \int R dR \\ S(R) &= \frac{R^2}{2} + c_2 \end{aligned}$$

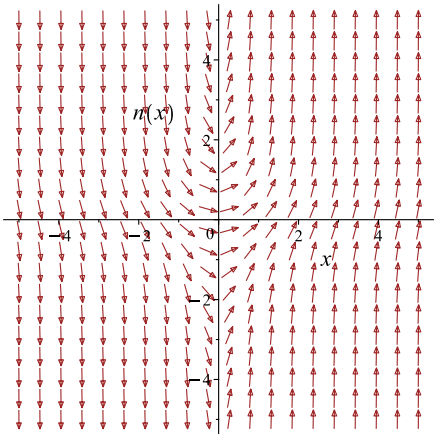
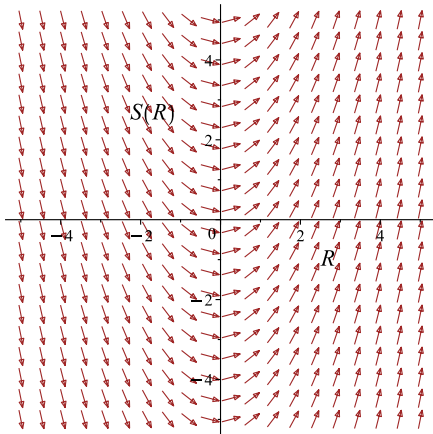
To complete the solution, we just need to transform the above back to x, n coordinates. This results in

$$\arctan(n) = \frac{x^2}{2} + c_2$$

Which gives

$$n = \tan\left(\frac{x^2}{2} + c_2\right)$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, n coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dn}{dx} = (n^2 + 1)x$ 	$R = x$ $S = \arctan(n)$	$\frac{dS}{dR} = R$ 

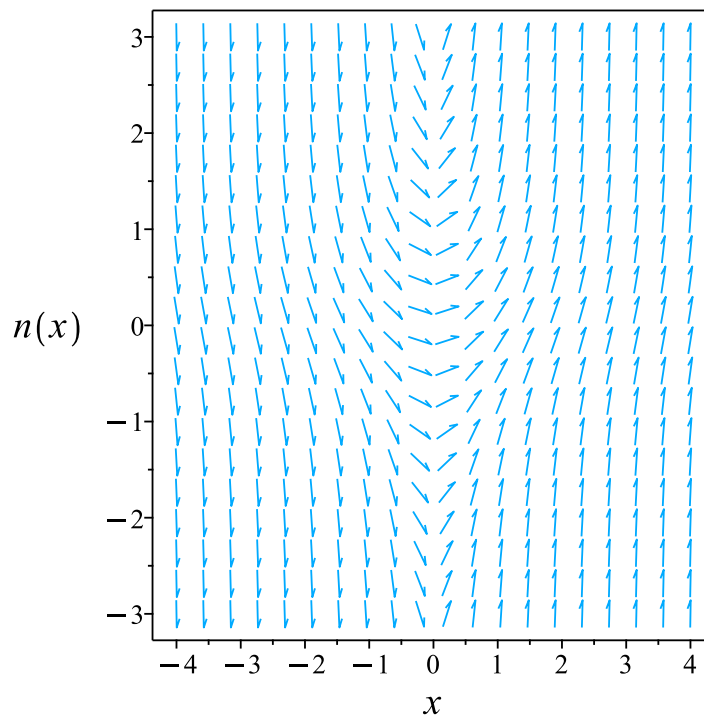


Figure 2.61: Slope field plot
 $n' = (n^2 + 1)x$

Summary of solutions found

$$n = \tan\left(\frac{x^2}{2} + c_2\right)$$

Solved as first order ode of type Riccati

Time used: 0.198 (sec)

In canonical form the ODE is

$$\begin{aligned} n' &= F(x, n) \\ &= (n^2 + 1)x \end{aligned}$$

This is a Riccati ODE. Comparing the ODE to solve

$$n' = n^2x + x$$

With Riccati ODE standard form

$$n' = f_0(x) + f_1(x)n + f_2(x)n^2$$

Shows that $f_0(x) = x$, $f_1(x) = 0$ and $f_2(x) = x$. Let

$$\begin{aligned} n &= \frac{-u'}{f_2 u} \\ &= \frac{-u'}{ux} \end{aligned} \quad (1)$$

Using the above substitution in the given ODE results (after some simplification) in a second order ODE to solve for $u(x)$ which is

$$f_2 u''(x) - (f_2' + f_1 f_2) u'(x) + f_2^2 f_0 u(x) = 0 \quad (2)$$

But

$$\begin{aligned} f_2' &= 1 \\ f_1 f_2 &= 0 \\ f_2^2 f_0 &= x^3 \end{aligned}$$

Substituting the above terms back in equation (2) gives

$$x u''(x) - u'(x) + x^3 u(x) = 0$$

In normal form the ode

$$x \left(\frac{d^2 u}{dx^2} \right) - \frac{du}{dx} + x^3 u = 0 \quad (1)$$

Becomes

$$\frac{d^2 u}{dx^2} + p(x) \left(\frac{du}{dx} \right) + q(x) u = 0 \quad (2)$$

Where

$$\begin{aligned} p(x) &= -\frac{1}{x} \\ q(x) &= x^2 \end{aligned}$$

Applying change of variables $\tau = g(x)$ to (2) gives

$$\frac{d^2}{d\tau^2} u(\tau) + p_1 \left(\frac{d}{d\tau} u(\tau) \right) + q_1 u(\tau) = 0 \quad (3)$$

Where τ is the new independent variable, and

$$p_1(\tau) = \frac{\frac{d^2}{dx^2} \tau(x) + p(x) \left(\frac{d}{dx} \tau(x) \right)}{\left(\frac{d}{dx} \tau(x) \right)^2} \quad (4)$$

$$q_1(\tau) = \frac{q(x)}{\left(\frac{d}{dx} \tau(x) \right)^2} \quad (5)$$

Let $p_1 = 0$. Eq (4) simplifies to

$$\frac{d^2}{dx^2}\tau(x) + p(x) \left(\frac{d}{dx}\tau(x) \right) = 0$$

This ode is solved resulting in

$$\begin{aligned} \tau &= \int e^{-\int p(x)dx} dx \\ &= \int e^{-\int -\frac{1}{x}dx} dx \\ &= \int e^{\ln(x)} dx \\ &= \int x dx \\ &= \frac{x^2}{2} \end{aligned} \tag{6}$$

Using (6) to evaluate q_1 from (5) gives

$$\begin{aligned} q_1(\tau) &= \frac{q(x)}{\left(\frac{d}{dx}\tau(x) \right)^2} \\ &= \frac{x^2}{x^2} \\ &= 1 \end{aligned} \tag{7}$$

Substituting the above in (3) and noting that now $p_1 = 0$ results in

$$\begin{aligned} \frac{d^2}{d\tau^2}u(\tau) + q_1u(\tau) &= 0 \\ \frac{d^2}{d\tau^2}u(\tau) + u(\tau) &= 0 \end{aligned}$$

The above ode is now solved for $u(\tau)$. This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Au''(\tau) + Bu'(\tau) + Cu(\tau) = 0$$

Where in the above $A = 1, B = 0, C = 1$. Let the solution be $u(\tau) = e^{\lambda\tau}$. Substituting this into the ODE gives

$$\lambda^2 e^{\tau\lambda} + e^{\tau\lambda} = 0 \tag{1}$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda\tau}$ gives

$$\lambda^2 + 1 = 0 \tag{2}$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = 0, C = 1$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{0}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{0^2 - (4)(1)(1)} \\ &= \pm i \end{aligned}$$

Hence

$$\begin{aligned} \lambda_1 &= +i \\ \lambda_2 &= -i \end{aligned}$$

Which simplifies to

$$\begin{aligned} \lambda_1 &= i \\ \lambda_2 &= -i \end{aligned}$$

Since roots are complex conjugate of each others, then let the roots be

$$\lambda_{1,2} = \alpha \pm i\beta$$

Where $\alpha = 0$ and $\beta = 1$. Therefore the final solution, when using Euler relation, can be written as

$$u(\tau) = e^{\alpha\tau} (c_1 \cos(\beta\tau) + c_2 \sin(\beta\tau))$$

Which becomes

$$u(\tau) = e^0 (c_1 \cos(\tau) + c_2 \sin(\tau))$$

Or

$$u(\tau) = c_1 \cos(\tau) + c_2 \sin(\tau)$$

Will add steps showing solving for IC soon.

The above solution is now transformed back to u using (6) which results in

$$u = c_1 \cos\left(\frac{x^2}{2}\right) + c_2 \sin\left(\frac{x^2}{2}\right)$$

Will add steps showing solving for IC soon.

Taking derivative gives

$$u'(x) = -c_1 x \sin\left(\frac{x^2}{2}\right) + c_2 x \cos\left(\frac{x^2}{2}\right)$$

Doing change of constants, the solution becomes

$$n = -\frac{-c_3 x \sin\left(\frac{x^2}{2}\right) + x \cos\left(\frac{x^2}{2}\right)}{x \left(c_3 \cos\left(\frac{x^2}{2}\right) + \sin\left(\frac{x^2}{2}\right)\right)}$$

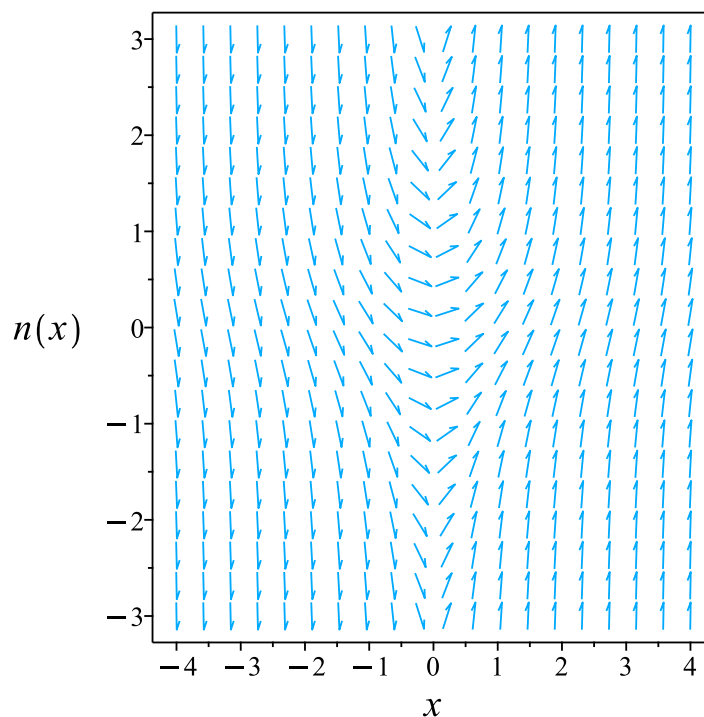


Figure 2.62: Slope field plot
 $n' = (n^2 + 1)x$

Summary of solutions found

$$n = -\frac{-c_3 x \sin\left(\frac{x^2}{2}\right) + x \cos\left(\frac{x^2}{2}\right)}{x \left(c_3 \cos\left(\frac{x^2}{2}\right) + \sin\left(\frac{x^2}{2}\right)\right)}$$

Maple step by step solution

Let's solve

$$n' = (n^2 + 1)x$$

- Highest derivative means the order of the ODE is 1

$$n'$$

- Solve for the highest derivative

$$n' = (n^2 + 1)x$$

- Separate variables

$$\frac{n'}{n^2+1} = x$$

- Integrate both sides with respect to x

$$\int \frac{n'}{n^2+1} dx = \int x dx + C1$$

- Evaluate integral

$$\arctan(n) = \frac{x^2}{2} + C1$$

- Solve for n

$$n = \tan\left(\frac{x^2}{2} + C1\right)$$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
<- separable successful`
```

Maple dsolve solution

Solving time : 0.005 (sec)

Leaf size : 12

```
dsolve(diff(n(x),x) = (n(x)^2+1)*x,
        n(x),singsol=all)
```

$$n = \tan\left(\frac{x^2}{2} + c_1\right)$$

Mathematica DSolve solution

Solving time : 0.197 (sec)

Leaf size : 30

```
DSolve[{D[n[x],x]==(n[x]^2+1)*x,{}},  
n[x],x,IncludeSingularSolutions->True]
```

$$n(x) \rightarrow \tan\left(\frac{x^2}{2} + c_1\right)$$

$$n(x) \rightarrow -i$$

$$n(x) \rightarrow i$$

2.6.7 problem 9 (a)

Solved as first order linear ode	461
Solved as first order separable ode	463
Solved as first order homogeneous class D2 ode	464
Solved as first order Exact ode	466
Solved using Lie symmetry for first order ode	471
Maple step by step solution	476
Maple trace	476
Maple dsolve solution	477
Mathematica DSolve solution	477

Internal problem ID [18247]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter IV. Methods of solution: First order equations. section 33. Problems at page 91

Problem number : 9 (a)

Date solved : Monday, December 23, 2024 at 09:19:12 PM

CAS classification : [_separable]

Solve

$$v' + \frac{2v}{u} = 3v$$

Solved as first order linear ode

Time used: 0.073 (sec)

In canonical form a linear first order is

$$v' + q(u)v = p(u)$$

Comparing the above to the given ode shows that

$$q(u) = -\frac{3u-2}{u}$$

$$p(u) = 0$$

The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int q \, du} \\ &= e^{\int -\frac{3u-2}{u} \, du} \\ &= u^2 e^{-3u}\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{du}\mu v &= 0 \\ \frac{d}{du}(v u^2 e^{-3u}) &= 0\end{aligned}$$

Integrating gives

$$\begin{aligned}v u^2 e^{-3u} &= \int 0 du + c_1 \\ &= c_1\end{aligned}$$

Dividing throughout by the integrating factor $u^2 e^{-3u}$ gives the final solution

$$v = \frac{e^{3u} c_1}{u^2}$$

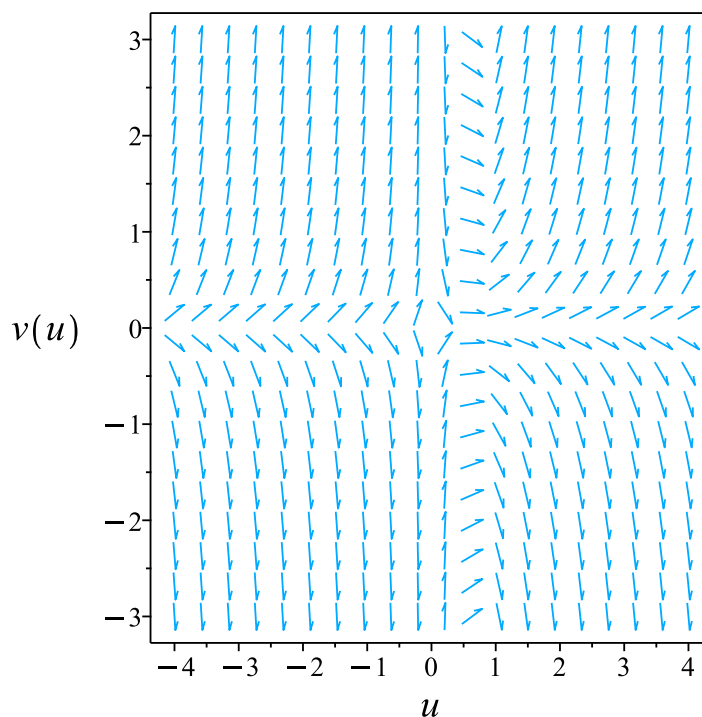


Figure 2.63: Slope field plot

$$v' + \frac{2v}{u} = 3v$$

Summary of solutions found

$$v = \frac{e^{3u} c_1}{u^2}$$

Solved as first order separable ode

Time used: 0.103 (sec)

The ode $v' = \frac{v(3u-2)}{u}$ is separable as it can be written as

$$\begin{aligned} v' &= \frac{v(3u-2)}{u} \\ &= f(u)g(v) \end{aligned}$$

Where

$$\begin{aligned} f(u) &= \frac{3u-2}{u} \\ g(v) &= v \end{aligned}$$

Integrating gives

$$\begin{aligned} \int \frac{1}{g(v)} dv &= \int f(u) du \\ \int \frac{1}{v} dv &= \int \frac{3u-2}{u} du \\ \ln(v) &= 3u + \ln\left(\frac{1}{u^2}\right) + c_1 \end{aligned}$$

We now need to find the singular solutions, these are found by finding for what values $g(v)$ is zero, since we had to divide by this above. Solving $g(v) = 0$ or $v = 0$ for v gives

$$v = 0$$

Now we go over each such singular solution and check if it verifies the ode itself and any initial conditions given. If it does not then the singular solution will not be used.

Therefore the solutions found are

$$\begin{aligned} \ln(v) &= 3u + \ln\left(\frac{1}{u^2}\right) + c_1 \\ v &= 0 \end{aligned}$$

Solving for v gives

$$\begin{aligned} v &= 0 \\ v &= \frac{e^{3u+c_1}}{u^2} \end{aligned}$$

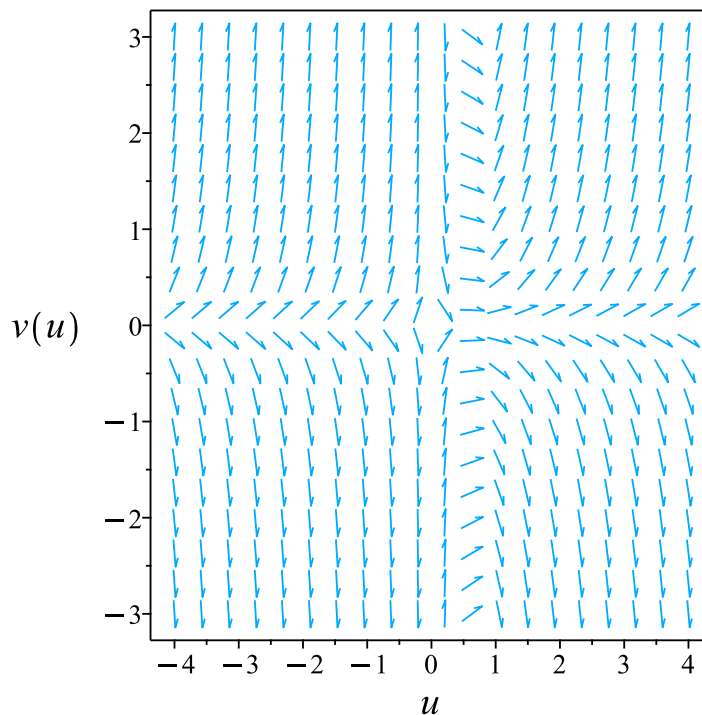


Figure 2.64: Slope field plot
 $v' + \frac{2v}{u} = 3v$

Summary of solutions found

$$v = 0$$

$$v = \frac{e^{3u+c_1}}{u^2}$$

Solved as first order homogeneous class D2 ode

Time used: 0.127 (sec)

Applying change of variables $v = u(u)u$, then the ode becomes

$$u'(u)u + 3u(u) = 3u(u)u$$

Which is now solved The ode $u'(u) = \frac{3u(u)(u-1)}{u}$ is separable as it can be written as

$$u'(u) = \frac{3u(u)(u-1)}{u}$$

$$= f(u)g(u)$$

Where

$$f(u) = \frac{3u - 3}{u}$$

$$g(u) = u$$

Integrating gives

$$\int \frac{1}{g(u)} du = \int f(u) du$$

$$\int \frac{1}{u} du = \int \frac{3u - 3}{u} du$$

$$\ln(u(u)) = 3u + \ln\left(\frac{1}{u^3}\right) + c_1$$

We now need to find the singular solutions, these are found by finding for what values $g(u)$ is zero, since we had to divide by this above. Solving $g(u) = 0$ or $u = 0$ for $u(u)$ gives

$$u(u) = 0$$

Now we go over each such singular solution and check if it verifies the ode itself and any initial conditions given. If it does not then the singular solution will not be used.

Therefore the solutions found are

$$\ln(u(u)) = 3u + \ln\left(\frac{1}{u^3}\right) + c_1$$

$$u(u) = 0$$

Solving for $u(u)$ gives

$$u(u) = 0$$

$$u(u) = \frac{e^{3u+c_1}}{u^3}$$

Converting $u(u) = 0$ back to v gives

$$v = 0$$

Converting $u(u) = \frac{e^{3u+c_1}}{u^3}$ back to v gives

$$v = \frac{e^{3u+c_1}}{u^2}$$

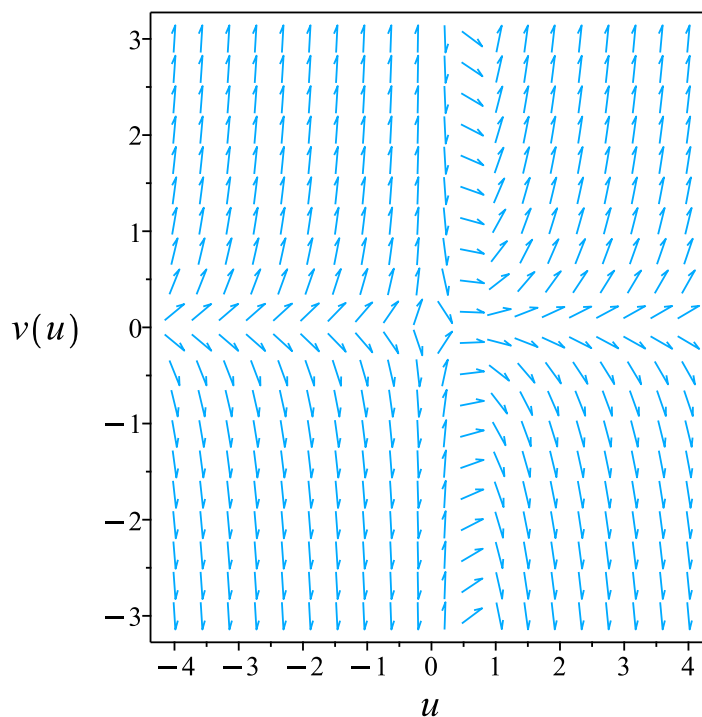


Figure 2.65: Slope field plot
 $v' + \frac{2v}{u} = 3v$

Summary of solutions found

$$v = 0$$

$$v = \frac{e^{3u+c_1}}{u^2}$$

Solved as first order Exact ode

Time used: 0.130 (sec)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(u, v) du + N(u, v) dv = 0 \quad (1A)$$

Therefore

$$\begin{aligned}dv &= \left(3v - \frac{2v}{u}\right) du \\ \left(\frac{2v}{u} - 3v\right) du + dv &= 0\end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(u, v) &= \frac{2v}{u} - 3v \\ N(u, v) &= 1\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial v} = \frac{\partial N}{\partial u}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial v} &= \frac{\partial}{\partial v} \left(\frac{2v}{u} - 3v\right) \\ &= \frac{2}{u} - 3\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial u} &= \frac{\partial}{\partial u}(1) \\ &= 0\end{aligned}$$

Since $\frac{\partial M}{\partial v} \neq \frac{\partial N}{\partial u}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned}A &= \frac{1}{N} \left(\frac{\partial M}{\partial v} - \frac{\partial N}{\partial u} \right) \\ &= 1 \left(\left(\frac{2}{u} - 3 \right) - (0) \right) \\ &= \frac{2}{u} - 3\end{aligned}$$

Since A does not depend on v , then it can be used to find an integrating factor. The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int A \, du} \\ &= e^{\int \frac{2}{u} - 3 \, du}\end{aligned}$$

The result of integrating gives

$$\begin{aligned}\mu &= e^{-3u + 2 \ln(u)} \\ &= u^2 e^{-3u}\end{aligned}$$

M and N are multiplied by this integrating factor, giving new M and new N which are called \overline{M} and \overline{N} for now so not to confuse them with the original M and N .

$$\begin{aligned}\overline{M} &= \mu M \\ &= u^2 e^{-3u} \left(\frac{2v}{u} - 3v \right) \\ &= -v(3u - 2) u e^{-3u}\end{aligned}$$

And

$$\begin{aligned}\overline{N} &= \mu N \\ &= u^2 e^{-3u} (1) \\ &= u^2 e^{-3u}\end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned}\bar{M} + \bar{N} \frac{dv}{du} &= 0 \\ (-v(3u - 2) u e^{-3u}) + (u^2 e^{-3u}) \frac{dv}{du} &= 0\end{aligned}$$

The following equations are now set up to solve for the function $\phi(u, v)$

$$\frac{\partial \phi}{\partial u} = \bar{M} \quad (1)$$

$$\frac{\partial \phi}{\partial v} = \bar{N} \quad (2)$$

Integrating (2) w.r.t. v gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial v} dv &= \int \bar{N} dv \\ \int \frac{\partial \phi}{\partial v} dv &= \int u^2 e^{-3u} dv \\ \phi &= v u^2 e^{-3u} + f(u)\end{aligned} \quad (3)$$

Where $f(u)$ is used for the constant of integration since ϕ is a function of both u and v . Taking derivative of equation (3) w.r.t u gives

$$\begin{aligned}\frac{\partial \phi}{\partial u} &= 2vu e^{-3u} - 3v u^2 e^{-3u} + f'(u) \\ &= -v(3u - 2) u e^{-3u} + f'(u)\end{aligned} \quad (4)$$

But equation (1) says that $\frac{\partial \phi}{\partial u} = -v(3u - 2) u e^{-3u}$. Therefore equation (4) becomes

$$-v(3u - 2) u e^{-3u} = -v(3u - 2) u e^{-3u} + f'(u) \quad (5)$$

Solving equation (5) for $f'(u)$ gives

$$f'(u) = 0$$

Therefore

$$f(u) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(u)$ into equation (3) gives ϕ

$$\phi = v u^2 e^{-3u} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into the constant c_1 gives the solution as

$$c_1 = v u^2 e^{-3u}$$

Solving for v gives

$$v = \frac{e^{3u} c_1}{u^2}$$

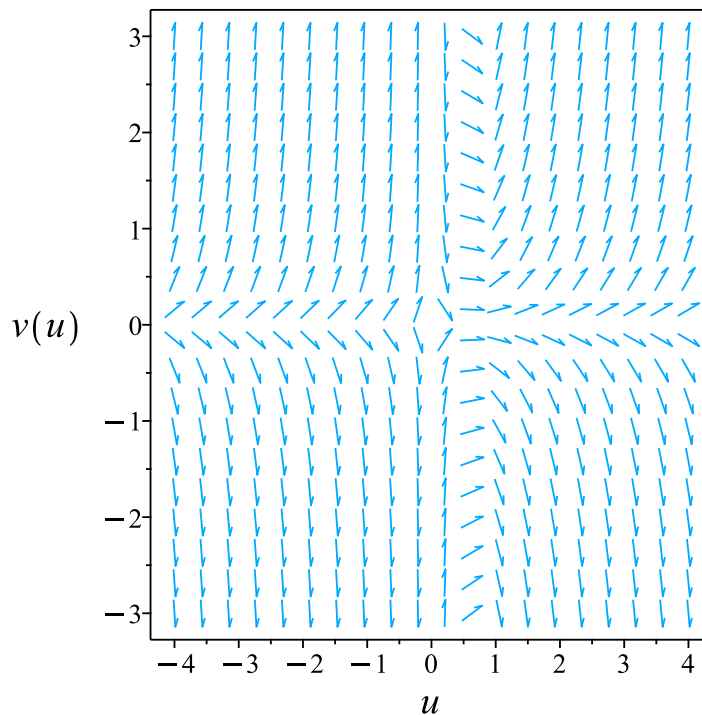


Figure 2.66: Slope field plot

$$v' + \frac{2v}{u} = 3v$$

Summary of solutions found

$$v = \frac{e^{3u} c_1}{u^2}$$

Solved using Lie symmetry for first order ode

Time used: 0.348 (sec)

Writing the ode as

$$v' = \frac{v(3u-2)}{u}$$

$$v' = \omega(u, v)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_u + \omega(\eta_v - \xi_u) - \omega^2 \xi_v - \omega_u \xi - \omega_v \eta = 0 \quad (\text{A})$$

To determine ξ, η then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$\xi = ua_2 + va_3 + a_1 \quad (\text{1E})$$

$$\eta = ub_2 + vb_3 + b_1 \quad (\text{2E})$$

Where the unknown coefficients are

$$\{a_1, a_2, a_3, b_1, b_2, b_3\}$$

Substituting equations (1E,2E) and ω into (A) gives

$$b_2 + \frac{v(3u-2)(b_3 - a_2)}{u} - \frac{v^2(3u-2)^2 a_3}{u^2} \quad (\text{5E})$$

$$- \left(\frac{3v}{u} - \frac{v(3u-2)}{u^2} \right) (ua_2 + va_3 + a_1) - \frac{(3u-2)(ub_2 + vb_3 + b_1)}{u} = 0$$

Putting the above in normal form gives

$$-\frac{9u^2v^2a_3 + 3u^3b_2 + 3u^2va_2 - 12uv^2a_3 + 3u^2b_1 - 3b_2u^2 + 6v^2a_3 - 2ub_1 + 2va_1}{u^2} = 0$$

Setting the numerator to zero gives

$$-9u^2v^2a_3 - 3u^3b_2 - 3u^2va_2 + 12uv^2a_3 - 3u^2b_1 + 3b_2u^2 - 6v^2a_3 + 2ub_1 - 2va_1 = 0 \quad (\text{6E})$$

Looking at the above PDE shows the following are all the terms with $\{u, v\}$ in them.

$$\{u, v\}$$

The following substitution is now made to be able to collect on all terms with $\{u, v\}$ in them

$$\{u = v_1, v = v_2\}$$

The above PDE (6E) now becomes

$$-9a_3v_1^2v_2^2 - 3a_2v_1^2v_2 + 12a_3v_1v_2^2 - 3b_2v_1^3 - 6a_3v_2^2 - 3b_1v_1^2 + 3b_2v_1^2 - 2a_1v_2 + 2b_1v_1 = 0 \quad (7E)$$

Collecting the above on the terms v_i introduced, and these are

$$\{v_1, v_2\}$$

Equation (7E) now becomes

$$-3b_2v_1^3 - 9a_3v_1^2v_2^2 - 3a_2v_1^2v_2 + (-3b_1 + 3b_2)v_1^2 + 12a_3v_1v_2^2 + 2b_1v_1 - 6a_3v_2^2 - 2a_1v_2 = 0 \quad (8E)$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$\begin{aligned} -2a_1 &= 0 \\ -3a_2 &= 0 \\ -9a_3 &= 0 \\ -6a_3 &= 0 \\ 12a_3 &= 0 \\ 2b_1 &= 0 \\ -3b_2 &= 0 \\ -3b_1 + 3b_2 &= 0 \end{aligned}$$

Solving the above equations for the unknowns gives

$$\begin{aligned} a_1 &= 0 \\ a_2 &= 0 \\ a_3 &= 0 \\ b_1 &= 0 \\ b_2 &= 0 \\ b_3 &= b_3 \end{aligned}$$

Substituting the above solution in the anstaz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$\begin{aligned}\xi &= 0 \\ \eta &= v\end{aligned}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(u, v) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{du}{\xi} = \frac{dv}{\eta} = dS \quad (1)$$

The above comes from the requirements that $(\xi \frac{\partial}{\partial u} + \eta \frac{\partial}{\partial v}) S(u, v) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = u$$

S is found from

$$\begin{aligned}S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{v} dy\end{aligned}$$

Which results in

$$S = \ln(v)$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_u + \omega(u, v)S_v}{R_u + \omega(u, v)R_v} \quad (2)$$

Where in the above R_u, R_v, S_u, S_v are all partial derivatives and $\omega(u, v)$ is the right hand side of the original ode given by

$$\omega(u, v) = \frac{v(3u - 2)}{u}$$

Evaluating all the partial derivatives gives

$$\begin{aligned}R_u &= 1 \\R_v &= 0 \\S_u &= 0 \\S_v &= \frac{1}{v}\end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = \frac{3u - 2}{u} \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for u, v in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = \frac{3R - 2}{R}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S .

Since the ode has the form $\frac{d}{dR}S(R) = f(R)$, then we only need to integrate $f(R)$.

$$\begin{aligned}\int dS &= \int \frac{3R - 2}{R} dR \\S(R) &= 3R - 2 \ln(R) + c_2\end{aligned}$$

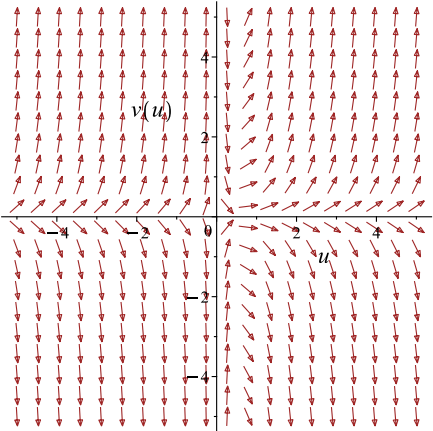
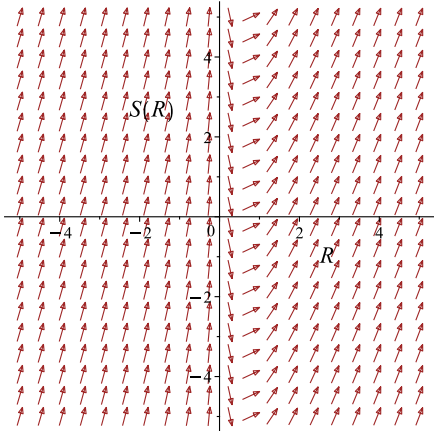
To complete the solution, we just need to transform the above back to u, v coordinates. This results in

$$\ln(v) = 3u - 2 \ln(u) + c_2$$

Which gives

$$v = \frac{e^{3u+c_2}}{u^2}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in u, v coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dv}{du} = \frac{v(3u-2)}{u}$ 	$R = u$ $S = \ln(v)$	$\frac{dS}{dR} = \frac{3R-2}{R}$ 

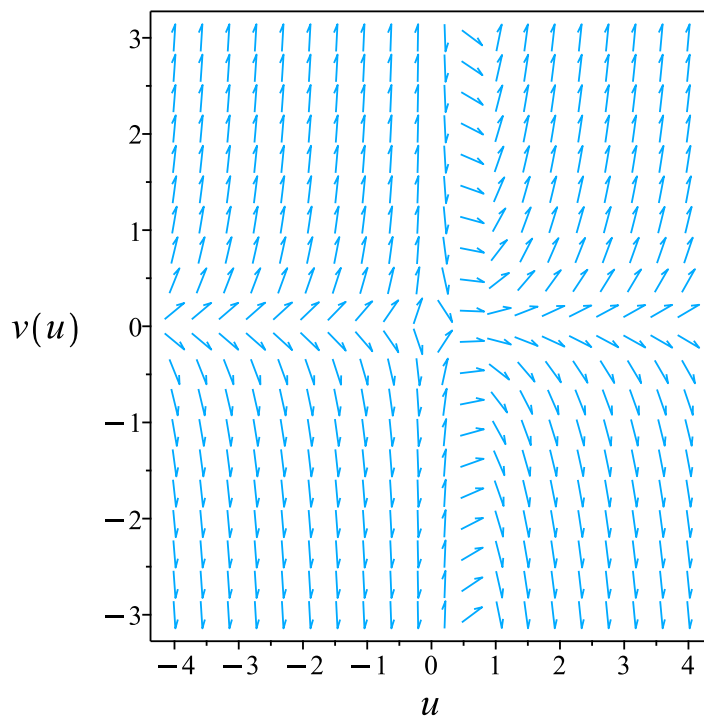


Figure 2.67: Slope field plot
 $v' + \frac{2v}{u} = 3v$

Summary of solutions found

$$v = \frac{e^{3u+c_2}}{u^2}$$

Maple step by step solution

Let's solve

$$v' + \frac{2v}{u} = 3v$$

- Highest derivative means the order of the ODE is 1

v'

- Solve for the highest derivative

$$v' = 3v - \frac{2v}{u}$$

- Separate variables

$$\frac{v'}{v} = \frac{3u-2}{u}$$

- Integrate both sides with respect to u

$$\int \frac{v'}{v} du = \int \frac{3u-2}{u} du + C1$$

- Evaluate integral

$$\ln(v) = 3u - 2 \ln(u) + C1$$

- Solve for v

$$v = \frac{e^{3u+C1}}{u^2}$$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`
```


Maple dsolve solution

Solving time : 0.002 (sec)

Leaf size : 13

```
dsolve(diff(v(u),u)+2*v(u)/u = 3*v(u),  
        v(u),singsol=all)
```

$$v = \frac{e^{3u}c_1}{u^2}$$

Mathematica DSolve solution

Solving time : 0.026 (sec)

Leaf size : 21

```
DSolve[{D[v[u],u]+2*v[u]/u==3*v[u],{}},  
        v[u],u,IncludeSingularSolutions->True]
```

$$v(u) \rightarrow \frac{c_1 e^{3u}}{u^2}$$

$$v(u) \rightarrow 0$$

2.6.8 problem 9 (b)

Solved as first order separable ode	478
Maple step by step solution	480
Maple trace	480
Maple dsolve solution	481
Mathematica DSolve solution	481

Internal problem ID [18248]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter IV. Methods of solution: First order equations. section 33. Problems at page 91

Problem number : 9 (b)

Date solved : Monday, December 23, 2024 at 09:19:14 PM

CAS classification : [_separable]

Solve

$$\sqrt{-u^2 + 1} v' = 2u\sqrt{1 - v^2}$$

Solved as first order separable ode

Time used: 0.352 (sec)

The ode $v' = \frac{2u\sqrt{1-v^2}}{\sqrt{-u^2+1}}$ is separable as it can be written as

$$\begin{aligned} v' &= \frac{2u\sqrt{1-v^2}}{\sqrt{-u^2+1}} \\ &= f(u)g(v) \end{aligned}$$

Where

$$\begin{aligned} f(u) &= \frac{2u}{\sqrt{-u^2+1}} \\ g(v) &= \sqrt{-v^2+1} \end{aligned}$$

Integrating gives

$$\begin{aligned} \int \frac{1}{g(v)} dv &= \int f(u) du \\ \int \frac{1}{\sqrt{-v^2+1}} dv &= \int \frac{2u}{\sqrt{-u^2+1}} du \\ \arcsin(v) &= -2\sqrt{-u^2+1} + c_1 \end{aligned}$$

We now need to find the singular solutions, these are found by finding for what values $g(v)$ is zero, since we had to divide by this above. Solving $g(v) = 0$ or $\sqrt{-v^2 + 1} = 0$ for v gives

$$v = -1$$

$$v = 1$$

Now we go over each such singular solution and check if it verifies the ode itself and any initial conditions given. If it does not then the singular solution will not be used.

Therefore the solutions found are

$$\arcsin(v) = -2\sqrt{-u^2 + 1} + c_1$$

$$v = -1$$

$$v = 1$$

Solving for v gives

$$v = -1$$

$$v = 1$$

$$v = \sin\left(-2\sqrt{-u^2 + 1} + c_1\right)$$

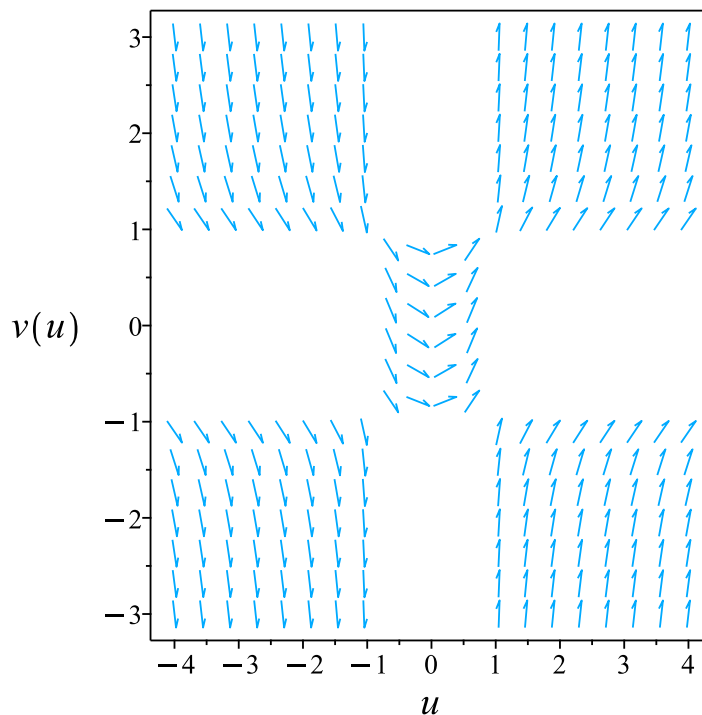


Figure 2.68: Slope field plot
 $\sqrt{-u^2 + 1} v' = 2u\sqrt{1 - v^2}$

Summary of solutions found

$$v = -1$$

$$v = 1$$

$$v = \sin\left(-2\sqrt{-u^2 + 1} + c_1\right)$$

Maple step by step solution

Let's solve

$$\sqrt{-u^2 + 1} v' = 2u\sqrt{1 - v^2}$$

- Highest derivative means the order of the ODE is 1

$$v'$$

- Solve for the highest derivative

$$v' = \frac{2u\sqrt{1-v^2}}{\sqrt{-u^2+1}}$$

- Separate variables

$$\frac{v'}{\sqrt{1-v^2}} = \frac{2u}{\sqrt{-u^2+1}}$$

- Integrate both sides with respect to u

$$\int \frac{v'}{\sqrt{1-v^2}} du = \int \frac{2u}{\sqrt{-u^2+1}} du + C1$$

- Evaluate integral

$$\arcsin(v) = \frac{2(u-1)(u+1)}{\sqrt{-u^2+1}} + C1$$

- Solve for v

$$v = \sin\left(\frac{C1\sqrt{-u^2+1}+2u^2-2}{\sqrt{-u^2+1}}\right)$$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
<- separable successful`
```

Maple dsolve solution

Solving time : 0.006 (sec)

Leaf size : 32

```
dsolve((-u^2+1)^(1/2)*diff(v(u),u) = 2*u*(1-v(u)^2)^(1/2),
      v(u),singsol=all)
```

$$v = \sin \left(\frac{2c_1 \sqrt{-u^2 + 1} + 2u^2 - 2}{\sqrt{-u^2 + 1}} \right)$$

Mathematica DSolve solution

Solving time : 6.503 (sec)

Leaf size : 110

```
DSolve[{Sqrt[1-u^2]*D[v[u],u]==2*u*Sqrt[1-v[u]^2],{}},
      v[u],u,IncludeSingularSolutions->True]
```

$$v(u) \rightarrow -\frac{\tan(2\sqrt{1-u^2}-c_1)}{\sqrt{\sec^2(2\sqrt{1-u^2}-c_1)}}$$

$$v(u) \rightarrow \frac{\tan(2\sqrt{1-u^2}-c_1)}{\sqrt{\sec^2(2\sqrt{1-u^2}-c_1)}}$$

$$v(u) \rightarrow -1$$

$$v(u) \rightarrow 1$$

2.6.9 problem 9 (c)

Solved as first order quadrature ode	482
Maple step by step solution	483
Maple trace	484
Maple dsolve solution	484
Mathematica DSolve solution	484

Internal problem ID [18249]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter IV. Methods of solution: First order equations. section 33. Problems at page 91

Problem number : 9 (c)

Date solved : Monday, December 23, 2024 at 09:19:17 PM

CAS classification : [_quadrature]

Solve

$$\sqrt{1 + v'} = \frac{e^u}{2}$$

Solved as first order quadrature ode

Time used: 0.067 (sec)

Since the ode has the form $v' = f(u)$, then we only need to integrate $f(u)$.

$$\int dv = \int \frac{e^{2u}}{4} - 1 du$$

$$v = -u + \frac{e^{2u}}{8} + c_1$$

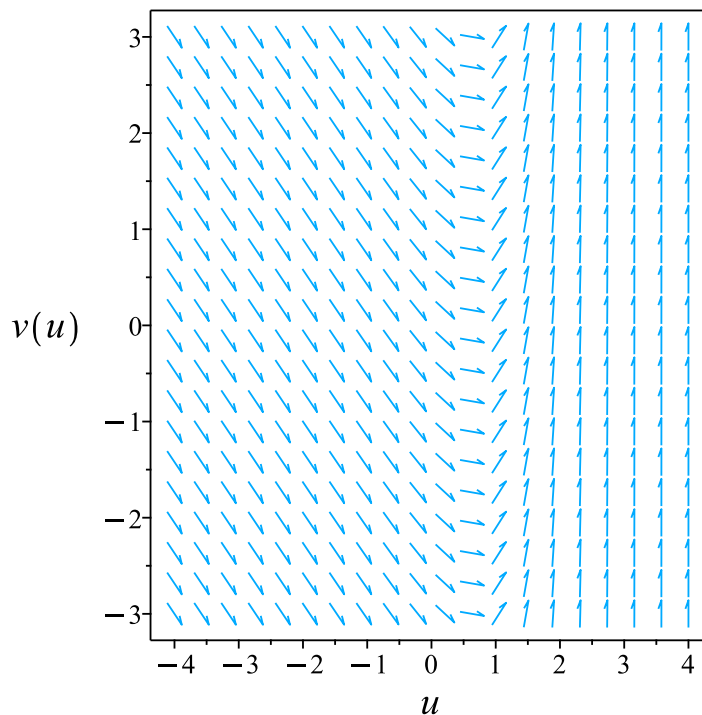


Figure 2.69: Slope field plot
 $\sqrt{1 + v'} = \frac{e^u}{2}$

Summary of solutions found

$$v = -u + \frac{e^{2u}}{8} + c_1$$

Maple step by step solution

Let's solve

$$\sqrt{1 + v'} = \frac{e^u}{2}$$

- Highest derivative means the order of the ODE is 1

v'

- Solve for the highest derivative

$$v' = \frac{(e^u)^2}{4} - 1$$

- Integrate both sides with respect to u

$$\int v' du = \int \left(\frac{(e^u)^2}{4} - 1 \right) du + C1$$

- Evaluate integral

$$v = -u + \frac{(e^u)^2}{8} + C1$$

- Solve for v

$$v = -u + \frac{(e^u)^2}{8} + C1$$

Maple trace

```

`Methods for first order ODEs:
-> Solving 1st order ODE of high degree, 1st attempt
trying 1st order WeierstrassP solution for high degree ODE
trying 1st order WeierstrassPPrime solution for high degree ODE
trying 1st order JacobiSN solution for high degree ODE
trying 1st order ODE linearizable_by_differentiation
trying differential order: 1; missing variables
<- differential order: 1; missing y(x) successful`

```

Maple dsolve solution

Solving time : 0.008 (sec)

Leaf size : 17

```

dsolve((1+diff(v(u),u))^(1/2) = 1/2*exp(u),
        v(u),singsol=all)

```

$$v = \frac{e^{2u}}{8} - \ln(e^u) + c_1$$

Mathematica DSolve solution

Solving time : 0.011 (sec)

Leaf size : 20

```

DSolve[{Sqrt[1+D[v[u],u]]==Exp[u]/2,{}},
        v[u],u,IncludeSingularSolutions->True]

```

$$v(u) \rightarrow -u + \frac{e^{2u}}{8} + c_1$$

2.6.10 problem 9 (d)

Solved as first order linear ode 485
 Solved as first order separable ode 487
 Solved as first order homogeneous class D2 ode 488
 Solved as first order Exact ode 490
 Solved using Lie symmetry for first order ode 495
 Maple step by step solution 500
 Maple trace 501
 Maple dsolve solution 501
 Mathematica DSolve solution 501

Internal problem ID [18250]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter IV. Methods of solution: First order equations. section 33. Problems at page 91

Problem number : 9 (d)

Date solved : Monday, December 23, 2024 at 09:19:17 PM

CAS classification : [_separable]

Solve

$$\frac{y'}{x} = y \sin(x^2 - 1) - \frac{2y}{\sqrt{x}}$$

Solved as first order linear ode

Time used: 0.503 (sec)

In canonical form a linear first order is

$$y' + q(x)y = p(x)$$

Comparing the above to the given ode shows that

$$q(x) = -(\sin(x^2 - 1)\sqrt{x} - 2)\sqrt{x}$$

$$p(x) = 0$$

The integrating factor μ is

$$\begin{aligned} \mu &= e^{\int q dx} \\ &= e^{\int -(\sin(x^2 - 1)\sqrt{x} - 2)\sqrt{x} dx} \\ &= e^{\frac{\cos(x^2 - 1)}{2} + \frac{4x^{3/2}}{3}} \end{aligned}$$

The ode becomes

$$\frac{d}{dx} \mu y = 0$$

$$\frac{d}{dx} \left(y e^{\frac{\cos(x^2-1)}{2} + \frac{4x^{3/2}}{3}} \right) = 0$$

Integrating gives

$$y e^{\frac{\cos(x^2-1)}{2} + \frac{4x^{3/2}}{3}} = \int 0 dx + c_1$$

$$= c_1$$

Dividing throughout by the integrating factor $e^{\frac{\cos(x^2-1)}{2} + \frac{4x^{3/2}}{3}}$ gives the final solution

$$y = e^{-\frac{\cos(x^2-1)}{2} - \frac{4x^{3/2}}{3}} c_1$$

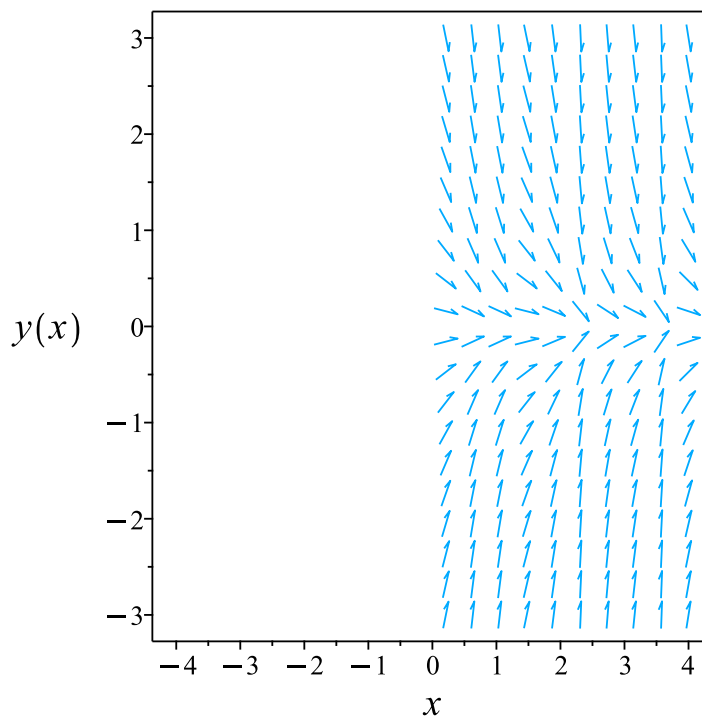


Figure 2.70: Slope field plot

$$\frac{y'}{y} = y \sin(x^2 - 1) - \frac{2y}{\sqrt{x}}$$

Summary of solutions found

$$y = e^{-\frac{\cos(x^2-1)}{2} - \frac{4x^{3/2}}{3}} c_1$$

Solved as first order separable ode

Time used: 0.149 (sec)

The ode $y' = y(\sin(x^2 - 1)\sqrt{x} - 2)\sqrt{x}$ is separable as it can be written as

$$\begin{aligned} y' &= y(\sin(x^2 - 1)\sqrt{x} - 2)\sqrt{x} \\ &= f(x)g(y) \end{aligned}$$

Where

$$\begin{aligned} f(x) &= (\sin(x^2 - 1)\sqrt{x} - 2)\sqrt{x} \\ g(y) &= y \end{aligned}$$

Integrating gives

$$\begin{aligned} \int \frac{1}{g(y)} dy &= \int f(x) dx \\ \int \frac{1}{y} dy &= \int (\sin(x^2 - 1)\sqrt{x} - 2)\sqrt{x} dx \\ \ln(y) &= -\frac{\cos(x^2 - 1)}{2} - \frac{4x^{3/2}}{3} + c_1 \end{aligned}$$

We now need to find the singular solutions, these are found by finding for what values $g(y)$ is zero, since we had to divide by this above. Solving $g(y) = 0$ or $y = 0$ for y gives

$$y = 0$$

Now we go over each such singular solution and check if it verifies the ode itself and any initial conditions given. If it does not then the singular solution will not be used.

Therefore the solutions found are

$$\begin{aligned} \ln(y) &= -\frac{\cos(x^2 - 1)}{2} - \frac{4x^{3/2}}{3} + c_1 \\ y &= 0 \end{aligned}$$

Solving for y gives

$$\begin{aligned} y &= 0 \\ y &= e^{-\frac{\cos(x^2 - 1)}{2} - \frac{4x^{3/2}}{3} + c_1} \end{aligned}$$

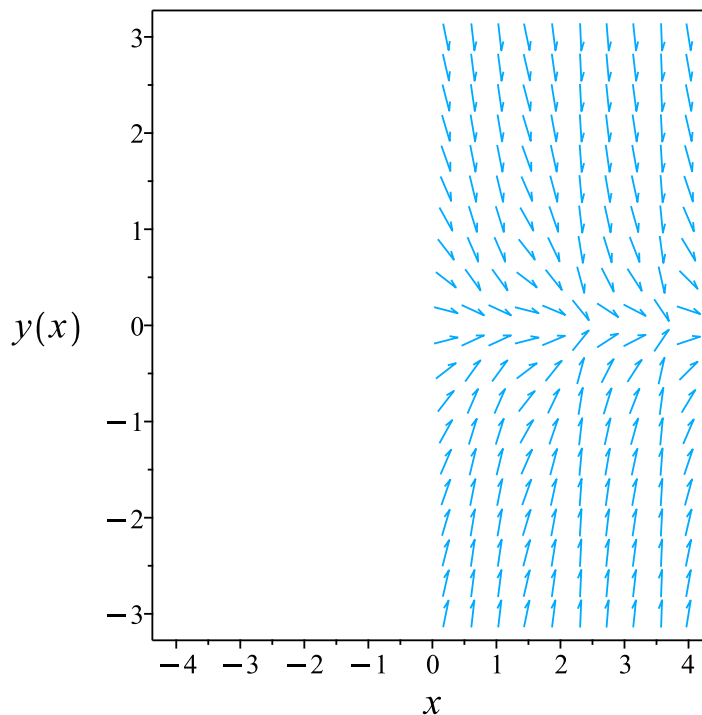


Figure 2.71: Slope field plot
 $\frac{y'}{x} = y \sin(x^2 - 1) - \frac{2y}{\sqrt{x}}$

Summary of solutions found

$$y = 0$$

$$y = e^{-\frac{\cos(x^2-1)}{2} - \frac{4x^{3/2}}{3} + c_1}$$

Solved as first order homogeneous class D2 ode

Time used: 0.294 (sec)

Applying change of variables $y = u(x)x$, then the ode becomes

$$\frac{u'(x)x + u(x)}{x} = u(x)x \sin(x^2 - 1) - 2u(x)\sqrt{x}$$

Which is now solved The ode $u'(x) = -\frac{u(x)(-\sin(x^2-1)x^2+2x^{3/2}+1)}{x}$ is separable as it can be written as

$$\begin{aligned} u'(x) &= -\frac{u(x)(-\sin(x^2-1)x^2+2x^{3/2}+1)}{x} \\ &= f(x)g(u) \end{aligned}$$

Where

$$f(x) = -\frac{-\sin(x^2 - 1)x^2 + 2x^{3/2} + 1}{x}$$

$$g(u) = u$$

Integrating gives

$$\int \frac{1}{g(u)} du = \int f(x) dx$$

$$\int \frac{1}{u} du = \int -\frac{-\sin(x^2 - 1)x^2 + 2x^{3/2} + 1}{x} dx$$

$$\ln(u(x)) = -\frac{\cos(x^2 - 1)}{2} + \ln\left(\frac{1}{x}\right) - \frac{4x^{3/2}}{3} + c_1$$

We now need to find the singular solutions, these are found by finding for what values $g(u)$ is zero, since we had to divide by this above. Solving $g(u) = 0$ or $u = 0$ for $u(x)$ gives

$$u(x) = 0$$

Now we go over each such singular solution and check if it verifies the ode itself and any initial conditions given. If it does not then the singular solution will not be used.

Therefore the solutions found are

$$\ln(u(x)) = -\frac{\cos(x^2 - 1)}{2} + \ln\left(\frac{1}{x}\right) - \frac{4x^{3/2}}{3} + c_1$$

$$u(x) = 0$$

Solving for $u(x)$ gives

$$u(x) = 0$$

$$u(x) = \frac{e^{-\frac{\cos(x^2 - 1)}{2} - \frac{4x^{3/2}}{3} + c_1}}{x}$$

Converting $u(x) = 0$ back to y gives

$$y = 0$$

Converting $u(x) = \frac{e^{-\frac{\cos(x^2 - 1)}{2} - \frac{4x^{3/2}}{3} + c_1}}{x}$ back to y gives

$$y = e^{-\frac{\cos(x^2 - 1)}{2} - \frac{4x^{3/2}}{3} + c_1}$$

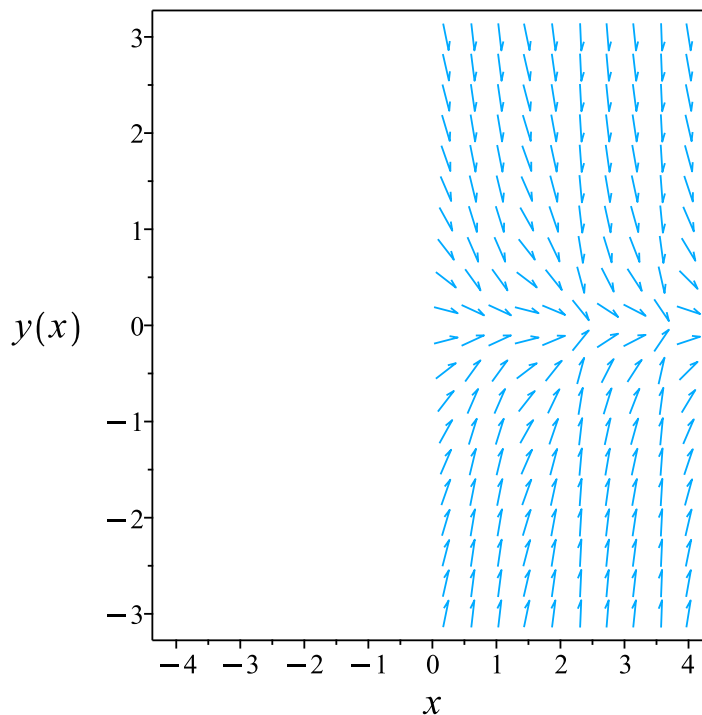


Figure 2.72: Slope field plot
 $\frac{y'}{x} = y \sin(x^2 - 1) - \frac{2y}{\sqrt{x}}$

Summary of solutions found

$$y = 0$$

$$y = e^{-\frac{\cos(x^2-1)}{2} - \frac{4x^{3/2}}{3} + c_1}$$

Solved as first order Exact ode

Time used: 0.523 (sec)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned}\left(\frac{1}{x}\right) dy &= \left(y \sin(x^2 - 1) - \frac{2y}{\sqrt{x}}\right) dx \\ \left(-y \sin(x^2 - 1) + \frac{2y}{\sqrt{x}}\right) dx &+ \left(\frac{1}{x}\right) dy = 0\end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(x, y) &= -y \sin(x^2 - 1) + \frac{2y}{\sqrt{x}} \\ N(x, y) &= \frac{1}{x}\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial y} &= \frac{\partial}{\partial y} \left(-y \sin(x^2 - 1) + \frac{2y}{\sqrt{x}}\right) \\ &= -\sin(x^2 - 1) + \frac{2}{\sqrt{x}}\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial x} &= \frac{\partial}{\partial x} \left(\frac{1}{x} \right) \\ &= -\frac{1}{x^2}\end{aligned}$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned}A &= \frac{1}{N} \left(\frac{\partial M}{\partial y} - \frac{\partial N}{\partial x} \right) \\ &= x \left(\left(-\sin(x^2 - 1) + \frac{2}{\sqrt{x}} \right) - \left(-\frac{1}{x^2} \right) \right) \\ &= x \left(-\sin(x^2 - 1) + \frac{2}{\sqrt{x}} + \frac{1}{x^2} \right)\end{aligned}$$

Since A does not depend on y , then it can be used to find an integrating factor. The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int A dx} \\ &= e^{\int x \left(-\sin(x^2 - 1) + \frac{2}{\sqrt{x}} + \frac{1}{x^2} \right) dx}\end{aligned}$$

The result of integrating gives

$$\begin{aligned}\mu &= e^{\frac{\cos(x^2 - 1)}{2} + \ln(x) + \frac{4x^{3/2}}{3}} \\ &= x e^{\frac{\cos(x^2 - 1)}{2} + \frac{4x^{3/2}}{3}}\end{aligned}$$

M and N are multiplied by this integrating factor, giving new M and new N which are called \bar{M} and \bar{N} for now so not to confuse them with the original M and N .

$$\begin{aligned}\bar{M} &= \mu M \\ &= x e^{\frac{\cos(x^2 - 1)}{2} + \frac{4x^{3/2}}{3}} \left(-y \sin(x^2 - 1) + \frac{2y}{\sqrt{x}} \right) \\ &= \left(-\sin(x^2 - 1) \sqrt{x} + 2 \right) \sqrt{x} y e^{\frac{\cos(x^2 - 1)}{2} + \frac{4x^{3/2}}{3}}\end{aligned}$$

And

$$\begin{aligned}\bar{N} &= \mu N \\ &= x e^{\frac{\cos(x^2 - 1)}{2} + \frac{4x^{3/2}}{3}} \left(\frac{1}{x} \right) \\ &= e^{\frac{\cos(x^2 - 1)}{2} + \frac{4x^{3/2}}{3}}\end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\overline{M} + \overline{N} \frac{dy}{dx} = 0$$

$$\left((-\sin(x^2 - 1) \sqrt{x} + 2) \sqrt{x} y e^{\frac{\cos(x^2 - 1)}{2} + \frac{4x^{3/2}}{3}} \right) + \left(e^{\frac{\cos(x^2 - 1)}{2} + \frac{4x^{3/2}}{3}} \right) \frac{dy}{dx} = 0$$

The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = \overline{M} \quad (1)$$

$$\frac{\partial \phi}{\partial y} = \overline{N} \quad (2)$$

Integrating (2) w.r.t. y gives

$$\int \frac{\partial \phi}{\partial y} dy = \int \overline{N} dy$$

$$\int \frac{\partial \phi}{\partial y} dy = \int e^{\frac{\cos(x^2 - 1)}{2} + \frac{4x^{3/2}}{3}} dy$$

$$\phi = y e^{\frac{\cos(x^2 - 1)}{2} + \frac{4x^{3/2}}{3}} + f(x) \quad (3)$$

Where $f(x)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t x gives

$$\frac{\partial \phi}{\partial x} = y(-\sin(x^2 - 1) x + 2\sqrt{x}) e^{\frac{\cos(x^2 - 1)}{2} + \frac{4x^{3/2}}{3}} + f'(x) \quad (4)$$

But equation (1) says that $\frac{\partial \phi}{\partial x} = (-\sin(x^2 - 1) \sqrt{x} + 2) \sqrt{x} y e^{\frac{\cos(x^2 - 1)}{2} + \frac{4x^{3/2}}{3}}$. Therefore equation (4) becomes

$$(-\sin(x^2 - 1) \sqrt{x} + 2) \sqrt{x} y e^{\frac{\cos(x^2 - 1)}{2} + \frac{4x^{3/2}}{3}} = y(-\sin(x^2 - 1) x + 2\sqrt{x}) e^{\frac{\cos(x^2 - 1)}{2} + \frac{4x^{3/2}}{3}} + f'(x) \quad (5)$$

Solving equation (5) for $f'(x)$ gives

$$f'(x) = 0$$

Therefore

$$f(x) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(x)$ into equation (3) gives ϕ

$$\phi = y e^{\frac{\cos(x^2-1)}{2} + \frac{4x^{3/2}}{3}} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into the constant c_1 gives the solution as

$$c_1 = y e^{\frac{\cos(x^2-1)}{2} + \frac{4x^{3/2}}{3}}$$

Solving for y gives

$$y = e^{-\frac{\cos(x^2-1)}{2} - \frac{4x^{3/2}}{3}} c_1$$

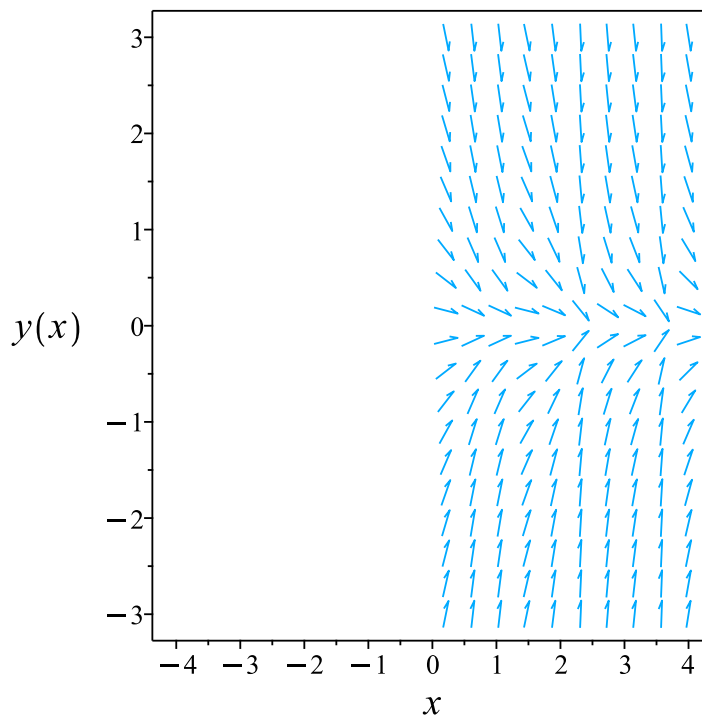


Figure 2.73: Slope field plot

$$\frac{y'}{x} = y \sin(x^2 - 1) - \frac{2y}{\sqrt{x}}$$

Summary of solutions found

$$y = e^{-\frac{\cos(x^2-1)}{2} - \frac{4x^{3/2}}{3}} c_1$$

Solved using Lie symmetry for first order ode

Time used: 0.849 (sec)

Writing the ode as

$$\begin{aligned}y' &= y(\sin(x^2 - 1)\sqrt{x} - 2)\sqrt{x} \\y' &= \omega(x, y)\end{aligned}$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

To determine ξ, η then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$\xi = xa_2 + ya_3 + a_1 \quad (\text{1E})$$

$$\eta = xb_2 + yb_3 + b_1 \quad (\text{2E})$$

Where the unknown coefficients are

$$\{a_1, a_2, a_3, b_1, b_2, b_3\}$$

Substituting equations (1E,2E) and ω into (A) gives

$$\begin{aligned}&b_2 + y(\sin(x^2 - 1)\sqrt{x} - 2)\sqrt{x}(b_3 - a_2) - y^2(\sin(x^2 - 1)\sqrt{x} - 2)^2 xa_3 \\&- \left(y \left(2x^{3/2} \cos(x^2 - 1) + \frac{\sin(x^2 - 1)}{2\sqrt{x}} \right) \sqrt{x} \right. \\&+ \left. \frac{y(\sin(x^2 - 1)\sqrt{x} - 2)}{2\sqrt{x}} \right) (xa_2 + ya_3 + a_1) \\&- (\sin(x^2 - 1)\sqrt{x} - 2)\sqrt{x}(xb_2 + yb_3 + b_1) = 0\end{aligned} \quad (\text{5E})$$

Putting the above in normal form gives

$$\sin(x^2 - 1)^2 x^{5/2} y^2 a_3 + 2 \cos(x^2 - 1) x^{7/2} y a_2 + 2 \cos(x^2 - 1) x^{5/2} y^2 a_3 + 2 \cos(x^2 - 1) x^{5/2} y a_1 - 4 \sin(x^2 - 1) x^{3/2} y^2 a_3 - 2 \cos(x^2 - 1) x^{7/2} y a_2 - 2 \cos(x^2 - 1) x^{5/2} y^2 a_3 - 2 \cos(x^2 - 1) x^{5/2} y a_1 - \sin(x^2 - 1) x^{5/2} b_2 - 2 \sin(x^2 - 1) x^{3/2} y a_2 + 4 \sin(x^2 - 1) x^2 y^2 a_3 - 4 y^2 x^{3/2} a_3 - \sin(x^2 - 1) x^{3/2} b_1 = 0$$

Setting the numerator to zero gives

$$-\sin(x^2 - 1)^2 x^{5/2} y^2 a_3 - 2 \cos(x^2 - 1) x^{7/2} y a_2 - 2 \cos(x^2 - 1) x^{5/2} y^2 a_3 \quad (\text{6E})$$

$$-2 \cos(x^2 - 1) x^{5/2} y a_1 - \sin(x^2 - 1) x^{5/2} b_2 - 2 \sin(x^2 - 1) x^{3/2} y a_2 + 4 \sin(x^2 - 1) x^2 y^2 a_3 - 4 y^2 x^{3/2} a_3 - \sin(x^2 - 1) x^{3/2} b_1 = 0$$

Simplifying the above gives

$$\begin{aligned}
 & -4y^2x^{3/2}a_3 + 2x^2b_2 + 3xya_2 + y^2a_3 + b_2\sqrt{x} + 2xb_1 + ya_1 - \frac{x^{5/2}y^2a_3}{2} + \frac{x^{5/2}y^2a_3 \cos(2x^2 - 2)}{2} \quad (6E) \\
 & -2 \cos(x^2 - 1)x^{7/2}ya_2 - 2 \cos(x^2 - 1)x^{5/2}y^2a_3 - 2 \cos(x^2 - 1)x^{5/2}ya_1 - \sin(x^2 - 1)x^{5/2}b_2 - 2 \sin(x^2 - 1)x
 \end{aligned}$$

Looking at the above PDE shows the following are all the terms with $\{x, y\}$ in them.

$$\{x, y, \sqrt{x}, x^{3/2}, x^{5/2}, x^{7/2}, \cos(x^2 - 1), \cos(2x^2 - 2), \sin(x^2 - 1)\}$$

The following substitution is now made to be able to collect on all terms with $\{x, y\}$ in them

$$\{x = v_1, y = v_2, \sqrt{x} = v_3, x^{3/2} = v_4, x^{5/2} = v_5, x^{7/2} = v_6, \cos(x^2 - 1) = v_7, \cos(2x^2 - 2) = v_8, \sin(x^2 - 1) = v_9\}$$

The above PDE (6E) now becomes

$$\begin{aligned}
 & -4v_2^2v_4a_3 + 2v_1^2b_2 + 3v_1v_2a_2 + v_2^2a_3 + b_2v_3 + 2v_1b_1 + v_2a_1 - \frac{1}{2}v_5v_2^2a_3 \quad (7E) \\
 & + \frac{1}{2}v_5v_2^2a_3v_8 - 2v_7v_6v_2a_2 - 2v_7v_5v_2^2a_3 - 2v_7v_5v_2a_1 - v_9v_5b_2 \\
 & - 2v_9v_4v_2a_2 + 4v_9v_1^2v_2^2a_3 - v_9v_4b_1 - v_9v_3v_2^2a_3 - v_9v_3v_2a_1 = 0
 \end{aligned}$$

Collecting the above on the terms v_i introduced, and these are

$$\{v_1, v_2, v_3, v_4, v_5, v_6, v_7, v_8, v_9\}$$

Equation (7E) now becomes

$$\begin{aligned}
 & -4v_2^2v_4a_3 + 2v_1^2b_2 + 3v_1v_2a_2 + v_2^2a_3 + b_2v_3 + 2v_1b_1 + v_2a_1 - \frac{1}{2}v_5v_2^2a_3 \quad (8E) \\
 & + \frac{1}{2}v_5v_2^2a_3v_8 - 2v_7v_6v_2a_2 - 2v_7v_5v_2^2a_3 - 2v_7v_5v_2a_1 - v_9v_5b_2 \\
 & - 2v_9v_4v_2a_2 + 4v_9v_1^2v_2^2a_3 - v_9v_4b_1 - v_9v_3v_2^2a_3 - v_9v_3v_2a_1 = 0
 \end{aligned}$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$\begin{aligned}a_1 &= 0 \\a_3 &= 0 \\b_2 &= 0 \\-2a_1 &= 0 \\-a_1 &= 0 \\-2a_2 &= 0 \\3a_2 &= 0 \\-4a_3 &= 0 \\-2a_3 &= 0 \\-a_3 &= 0 \\4a_3 &= 0 \\-\frac{a_3}{2} &= 0 \\ \frac{a_3}{2} &= 0 \\-b_1 &= 0 \\2b_1 &= 0 \\-b_2 &= 0 \\2b_2 &= 0\end{aligned}$$

Solving the above equations for the unknowns gives

$$\begin{aligned}a_1 &= 0 \\a_2 &= 0 \\a_3 &= 0 \\b_1 &= 0 \\b_2 &= 0 \\b_3 &= b_3\end{aligned}$$

Substituting the above solution in the anstaz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$\begin{aligned}\xi &= 0 \\ \eta &= y\end{aligned}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS \quad (1)$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = x$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{y} dy \end{aligned}$$

Which results in

$$S = \ln(y)$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = y(\sin(x^2 - 1) \sqrt{x} - 2) \sqrt{x}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= 1 \\ R_y &= 0 \\ S_x &= 0 \\ S_y &= \frac{1}{y} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = (\sin(x^2 - 1) \sqrt{x} - 2) \sqrt{x} \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = \left(\sin(R^2 - 1) \sqrt{R} - 2 \right) \sqrt{R}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S .

Since the ode has the form $\frac{d}{dR}S(R) = f(R)$, then we only need to integrate $f(R)$.

$$\int dS = \int \left(\sin(R^2 - 1) \sqrt{R} - 2 \right) \sqrt{R} dR$$

$$S(R) = -\frac{\cos(R^2 - 1)}{2} - \frac{4R^{3/2}}{3} + c_2$$

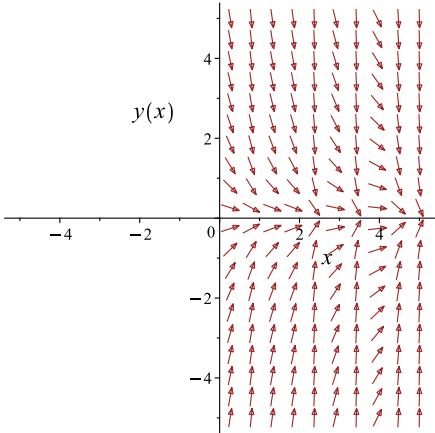
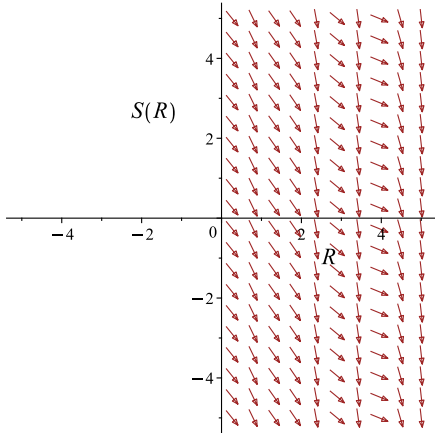
To complete the solution, we just need to transform the above back to x, y coordinates. This results in

$$\ln(y) = -\frac{\cos(x^2 - 1)}{2} - \frac{4x^{3/2}}{3} + c_2$$

Which gives

$$y = e^{-\frac{\cos(x^2 - 1)}{2} - \frac{4x^{3/2}}{3} + c_2}$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = y(\sin(x^2 - 1) \sqrt{x} - 2) \sqrt{x}$ 	$R = x$ $S = \ln(y)$	$\frac{dS}{dR} = \left(\sin(R^2 - 1) \sqrt{R} - 2 \right) \sqrt{R}$ 

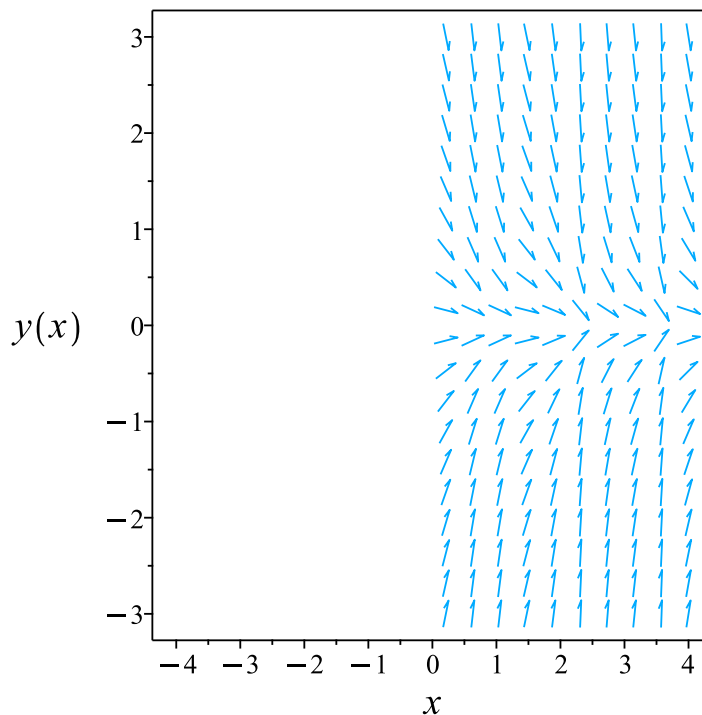


Figure 2.74: Slope field plot
 $\frac{y'}{x} = y \sin(x^2 - 1) - \frac{2y}{\sqrt{x}}$

Summary of solutions found

$$y = e^{-\frac{\cos(x^2-1)}{2} - \frac{4x^{3/2}}{3} + c_2}$$

Maple step by step solution

Let's solve

$$\frac{y'}{x} = y \sin(x^2 - 1) - \frac{2y}{\sqrt{x}}$$

- Highest derivative means the order of the ODE is 1

y'

- Solve for the highest derivative

$$y' = \left(y \sin(x^2 - 1) - \frac{2y}{\sqrt{x}} \right) x$$

- Separate variables

$$\frac{y'}{y} = (-2 + \sqrt{x} \sin((x - 1)(1 + x))) \sqrt{x}$$

- Integrate both sides with respect to x

$$\int \frac{y'}{y} dx = \int (-2 + \sqrt{x} \sin((x-1)(1+x))) \sqrt{x} dx + C1$$

- Evaluate integral

$$\ln(y) = -\frac{\cos((x-1)(1+x))}{2} - \frac{4x^{3/2}}{3} + C1$$

- Solve for y

$$y = e^{-\frac{\cos((x-1)(1+x))}{2} - \frac{4x^{3/2}}{3} + C1}$$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`
```

Maple dsolve solution

Solving time : 0.000 (sec)

Leaf size : 21

```
dsolve(1/x*diff(y(x),x) = y(x)*sin(x^2-1)-2*y(x)/x^(1/2),
        y(x),singsol=all)
```

$$y = e^{-\frac{\cos(x^2-1)}{2} - \frac{4x^{3/2}}{3}} c_1$$

Mathematica DSolve solution

Solving time : 0.1 (sec)

Leaf size : 37

```
DSolve[{1/x*D[y[x],x]==y[x]*Sin[x^2-1]-2*y[x]/Sqrt[x],{}},
        y[x],x,IncludeSingularSolutions->True]
```

$$y(x) \rightarrow c_1 e^{\frac{1}{6}(-8x^{3/2} - 3\cos(1-x^2))}$$

$$y(x) \rightarrow 0$$

2.6.11 problem 9 (e)

Solved as first order homogeneous class A ode	502
Solved as first order homogeneous class D2 ode	505
Solved as first order homogeneous class Maple C ode	508
Solved as first order Exact ode	512
Solved as first order isobaric ode	517
Solved using Lie symmetry for first order ode	520
Maple step by step solution	526
Maple trace	526
Maple dsolve solution	526
Mathematica DSolve solution	527

Internal problem ID [18251]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter IV. Methods of solution: First order equations. section 33. Problems at page 91

Problem number : 9 (e)

Date solved : Monday, December 23, 2024 at 09:19:20 PM

CAS classification :

[[_homogeneous, 'class A'], _rational, [_Abel, '2nd type', 'class A']]

Solve

$$y' = 1 + \frac{2y}{x - y}$$

Solved as first order homogeneous class A ode

Time used: 0.606 (sec)

In canonical form, the ODE is

$$\begin{aligned} y' &= F(x, y) \\ &= -\frac{x + y}{-x + y} \end{aligned} \tag{1}$$

An ode of the form $y' = \frac{M(x, y)}{N(x, y)}$ is called homogeneous if the functions $M(x, y)$ and $N(x, y)$ are both homogeneous functions and of the same order. Recall that a function $f(x, y)$ is homogeneous of order n if

$$f(t^n x, t^n y) = t^n f(x, y)$$

In this case, it can be seen that both $M = x + y$ and $N = x - y$ are both homogeneous and of the same order $n = 1$. Therefore this is a homogeneous ode. Since this ode is homogeneous, it is converted to separable ODE using the substitution $u = \frac{y}{x}$, or $y = ux$. Hence

$$\frac{dy}{dx} = \frac{du}{dx}x + u$$

Applying the transformation $y = ux$ to the above ODE in (1) gives

$$\begin{aligned} \frac{du}{dx}x + u &= \frac{-u - 1}{u - 1} \\ \frac{du}{dx} &= \frac{\frac{-u(x)-1}{u(x)-1} - u(x)}{x} \end{aligned}$$

Or

$$u'(x) - \frac{\frac{-u(x)-1}{u(x)-1} - u(x)}{x} = 0$$

Or

$$u'(x) xu(x) - u'(x) x + u(x)^2 + 1 = 0$$

Or

$$x(u(x) - 1) u'(x) + u(x)^2 + 1 = 0$$

Which is now solved as separable in $u(x)$.

The ode $u'(x) = -\frac{u(x)^2+1}{x(u(x)-1)}$ is separable as it can be written as

$$\begin{aligned} u'(x) &= -\frac{u(x)^2 + 1}{x(u(x) - 1)} \\ &= f(x)g(u) \end{aligned}$$

Where

$$\begin{aligned} f(x) &= -\frac{1}{x} \\ g(u) &= \frac{u^2 + 1}{u - 1} \end{aligned}$$

Integrating gives

$$\begin{aligned} \int \frac{1}{g(u)} du &= \int f(x) dx \\ \int \frac{u - 1}{u^2 + 1} du &= \int -\frac{1}{x} dx \\ \frac{\ln(u(x)^2 + 1)}{2} - \arctan(u(x)) &= \ln\left(\frac{1}{x}\right) + c_1 \end{aligned}$$

We now need to find the singular solutions, these are found by finding for what values $g(u)$ is zero, since we had to divide by this above. Solving $g(u) = 0$ or $\frac{u^2+1}{u-1} = 0$ for $u(x)$ gives

$$u(x) = -i$$

$$u(x) = i$$

Now we go over each such singular solution and check if it verifies the ode itself and any initial conditions given. If it does not then the singular solution will not be used.

Therefore the solutions found are

$$\frac{\ln(u(x)^2 + 1)}{2} - \arctan(u(x)) = \ln\left(\frac{1}{x}\right) + c_1$$

$$u(x) = -i$$

$$u(x) = i$$

Converting $\frac{\ln(u(x)^2+1)}{2} - \arctan(u(x)) = \ln\left(\frac{1}{x}\right) + c_1$ back to y gives

$$\frac{\ln\left(\frac{y^2+x^2}{x^2}\right)}{2} - \arctan\left(\frac{y}{x}\right) = \ln\left(\frac{1}{x}\right) + c_1$$

Converting $u(x) = -i$ back to y gives

$$y = -ix$$

Converting $u(x) = i$ back to y gives

$$y = ix$$

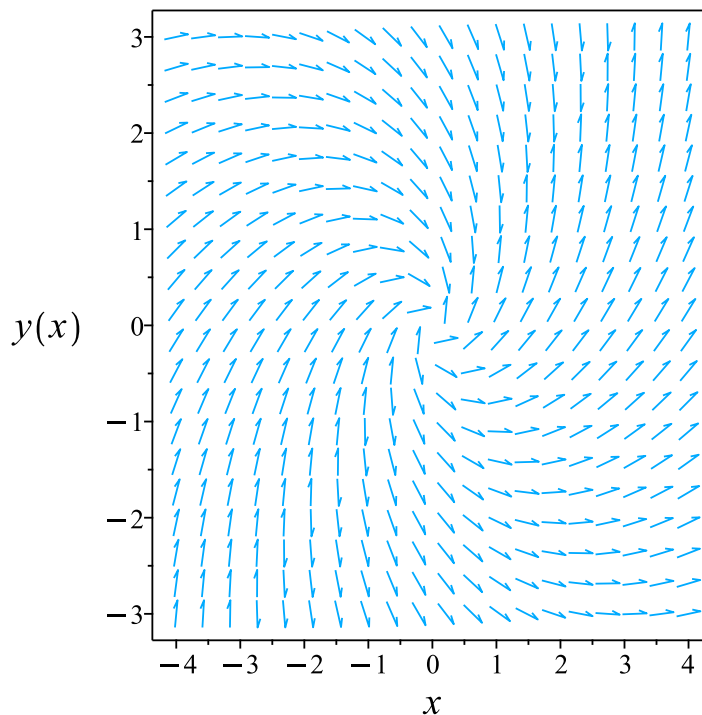


Figure 2.75: Slope field plot

$$y' = 1 + \frac{2y}{x-y}$$

Summary of solutions found

$$\frac{\ln\left(\frac{y^2+x^2}{x^2}\right)}{2} - \arctan\left(\frac{y}{x}\right) = \ln\left(\frac{1}{x}\right) + c_1$$

$$y = -ix$$

$$y = ix$$

Solved as first order homogeneous class D2 ode

Time used: 0.165 (sec)

Applying change of variables $y = u(x)x$, then the ode becomes

$$u'(x)x + u(x) = 1 + \frac{2u(x)x}{x - u(x)x}$$

Which is now solved The ode $u'(x) = -\frac{u(x)^2+1}{(u(x)-1)x}$ is separable as it can be written as

$$\begin{aligned} u'(x) &= -\frac{u(x)^2+1}{(u(x)-1)x} \\ &= f(x)g(u) \end{aligned}$$

Where

$$f(x) = -\frac{1}{x}$$

$$g(u) = \frac{u^2 + 1}{u - 1}$$

Integrating gives

$$\int \frac{1}{g(u)} du = \int f(x) dx$$

$$\int \frac{u - 1}{u^2 + 1} du = \int -\frac{1}{x} dx$$

$$\frac{\ln(u(x)^2 + 1)}{2} - \arctan(u(x)) = \ln\left(\frac{1}{x}\right) + c_1$$

We now need to find the singular solutions, these are found by finding for what values $g(u)$ is zero, since we had to divide by this above. Solving $g(u) = 0$ or $\frac{u^2+1}{u-1} = 0$ for $u(x)$ gives

$$u(x) = -i$$

$$u(x) = i$$

Now we go over each such singular solution and check if it verifies the ode itself and any initial conditions given. If it does not then the singular solution will not be used.

Therefore the solutions found are

$$\frac{\ln(u(x)^2 + 1)}{2} - \arctan(u(x)) = \ln\left(\frac{1}{x}\right) + c_1$$

$$u(x) = -i$$

$$u(x) = i$$

Converting $\frac{\ln(u(x)^2+1)}{2} - \arctan(u(x)) = \ln\left(\frac{1}{x}\right) + c_1$ back to y gives

$$\frac{\ln\left(\frac{y^2}{x^2} + 1\right)}{2} - \arctan\left(\frac{y}{x}\right) = \ln\left(\frac{1}{x}\right) + c_1$$

Converting $u(x) = -i$ back to y gives

$$y = -ix$$

Converting $u(x) = i$ back to y gives

$$y = ix$$

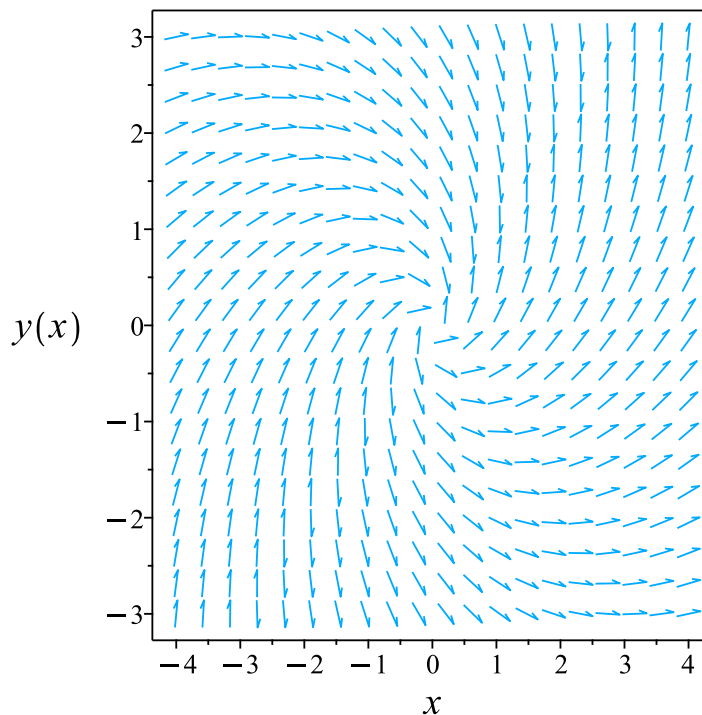


Figure 2.76: Slope field plot
 $y' = 1 + \frac{2y}{x-y}$

Summary of solutions found

$$\frac{\ln\left(\frac{y^2}{x^2} + 1\right)}{2} - \arctan\left(\frac{y}{x}\right) = \ln\left(\frac{1}{x}\right) + c_1$$

$$y = -ix$$

$$y = ix$$

Summary of solutions found

$$\frac{\ln\left(\frac{y^2+x^2}{x^2}\right)}{2} - \arctan\left(\frac{y}{x}\right) = \ln\left(\frac{1}{x}\right) + c_1$$

$$y = -ix$$

$$y = ix$$

Solved as first order homogeneous class Maple C ode

Time used: 0.750 (sec)

Let $Y = y - y_0$ and $X = x - x_0$ then the above is transformed to new ode in $Y(X)$

$$\frac{d}{dX}Y(X) = -\frac{x_0 + X + Y(X) + y_0}{-x_0 - X + Y(X) + y_0}$$

Solving for possible values of x_0 and y_0 which makes the above ode a homogeneous ode results in

$$x_0 = 0$$

$$y_0 = 0$$

Using these values now it is possible to easily solve for $Y(X)$. The above ode now becomes

$$\frac{d}{dX}Y(X) = -\frac{X + Y(X)}{-X + Y(X)}$$

In canonical form, the ODE is

$$\begin{aligned} Y' &= F(X, Y) \\ &= -\frac{X + Y}{-X + Y} \end{aligned} \quad (1)$$

An ode of the form $Y' = \frac{M(X, Y)}{N(X, Y)}$ is called homogeneous if the functions $M(X, Y)$ and $N(X, Y)$ are both homogeneous functions and of the same order. Recall that a function $f(X, Y)$ is homogeneous of order n if

$$f(t^n X, t^n Y) = t^n f(X, Y)$$

In this case, it can be seen that both $M = X + Y$ and $N = X - Y$ are both homogeneous and of the same order $n = 1$. Therefore this is a homogeneous ode. Since this ode is homogeneous, it is converted to separable ODE using the substitution $u = \frac{Y}{X}$, or $Y = uX$. Hence

$$\frac{dY}{dX} = \frac{du}{dX}X + u$$

Applying the transformation $Y = uX$ to the above ODE in (1) gives

$$\begin{aligned} \frac{du}{dX}X + u &= \frac{-u - 1}{u - 1} \\ \frac{du}{dX} &= \frac{\frac{-u(X)-1}{u(X)-1} - u(X)}{X} \end{aligned}$$

Or

$$\frac{d}{dX}u(X) - \frac{\frac{-u(X)-1}{u(X)-1} - u(X)}{X} = 0$$

Or

$$\left(\frac{d}{dX}u(X)\right)Xu(X) - \left(\frac{d}{dX}u(X)\right)X + u(X)^2 + 1 = 0$$

Or

$$X(u(X) - 1) \left(\frac{d}{dX}u(X)\right) + u(X)^2 + 1 = 0$$

Which is now solved as separable in $u(X)$.

The ode $\frac{d}{dX}u(X) = -\frac{u(X)^2+1}{X(u(X)-1)}$ is separable as it can be written as

$$\begin{aligned} \frac{d}{dX}u(X) &= -\frac{u(X)^2 + 1}{X(u(X) - 1)} \\ &= f(X)g(u) \end{aligned}$$

Where

$$\begin{aligned} f(X) &= -\frac{1}{X} \\ g(u) &= \frac{u^2 + 1}{u - 1} \end{aligned}$$

Integrating gives

$$\begin{aligned} \int \frac{1}{g(u)} du &= \int f(X) dX \\ \int \frac{u-1}{u^2+1} du &= \int -\frac{1}{X} dX \\ \frac{\ln(u(X)^2 + 1)}{2} - \arctan(u(X)) &= \ln\left(\frac{1}{X}\right) + c_1 \end{aligned}$$

We now need to find the singular solutions, these are found by finding for what values $g(u)$ is zero, since we had to divide by this above. Solving $g(u) = 0$ or $\frac{u^2+1}{u-1} = 0$ for $u(X)$ gives

$$\begin{aligned} u(X) &= -i \\ u(X) &= i \end{aligned}$$

Now we go over each such singular solution and check if it verifies the ode itself and any initial conditions given. If it does not then the singular solution will not be used.

Therefore the solutions found are

$$\frac{\ln(u(X)^2 + 1)}{2} - \arctan(u(X)) = \ln\left(\frac{1}{X}\right) + c_1$$

$$u(X) = -i$$

$$u(X) = i$$

Converting $\frac{\ln(u(X)^2+1)}{2} - \arctan(u(X)) = \ln\left(\frac{1}{X}\right) + c_1$ back to $Y(X)$ gives

$$\frac{\ln\left(\frac{Y(X)^2+X^2}{X^2}\right)}{2} - \arctan\left(\frac{Y(X)}{X}\right) = \ln\left(\frac{1}{X}\right) + c_1$$

Converting $u(X) = -i$ back to $Y(X)$ gives

$$Y(X) = -iX$$

Converting $u(X) = i$ back to $Y(X)$ gives

$$Y(X) = iX$$

Using the solution for $Y(X)$

$$\frac{\ln\left(\frac{Y(X)^2+X^2}{X^2}\right)}{2} - \arctan\left(\frac{Y(X)}{X}\right) = \ln\left(\frac{1}{X}\right) + c_1 \quad (\text{A})$$

And replacing back terms in the above solution using

$$Y = y + y_0$$

$$X = x + x_0$$

Or

$$Y = y$$

$$X = x$$

Then the solution in y becomes using EQ (A)

$$\frac{\ln\left(\frac{y^2+x^2}{x^2}\right)}{2} - \arctan\left(\frac{y}{x}\right) = \ln\left(\frac{1}{x}\right) + c_1$$

Using the solution for $Y(X)$

$$Y(X) = -iX \quad (\text{A})$$

And replacing back terms in the above solution using

$$\begin{aligned}Y &= y + y_0 \\X &= x + x_0\end{aligned}$$

Or

$$\begin{aligned}Y &= y \\X &= x\end{aligned}$$

Then the solution in y becomes using EQ (A)

$$y = -ix$$

Using the solution for $Y(X)$

$$Y(X) = iX \tag{A}$$

And replacing back terms in the above solution using

$$\begin{aligned}Y &= y + y_0 \\X &= x + x_0\end{aligned}$$

Or

$$\begin{aligned}Y &= y \\X &= x\end{aligned}$$

Then the solution in y becomes using EQ (A)

$$y = ix$$

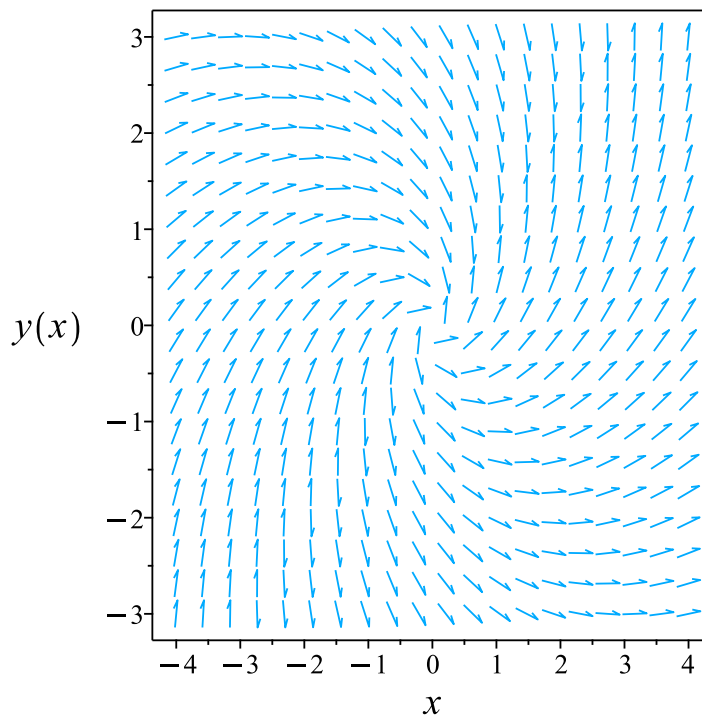


Figure 2.77: Slope field plot
 $y' = 1 + \frac{2y}{x-y}$

Solved as first order Exact ode

Time used: 0.260 (sec)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (1A)$$

Therefore

$$\begin{aligned} (-x + y) dy &= (-x - y) dx \\ (x + y) dx + (-x + y) dy &= 0 \end{aligned} \quad (2A)$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(x, y) &= x + y \\ N(x, y) &= -x + y \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\begin{aligned} \frac{\partial M}{\partial y} &= \frac{\partial}{\partial y}(x + y) \\ &= 1 \end{aligned}$$

And

$$\begin{aligned} \frac{\partial N}{\partial x} &= \frac{\partial}{\partial x}(-x + y) \\ &= -1 \end{aligned}$$

Since $\frac{\partial M}{\partial y} \neq \frac{\partial N}{\partial x}$, then the ODE is not exact. By inspection $\frac{1}{y^2+x^2}$ is an integrating factor. Therefore by multiplying $M = x + y$ and $N = -x + y$ by this integrating factor the

ode becomes exact. The new M, N are

$$M = \frac{x + y}{y^2 + x^2}$$

$$N = \frac{-x + y}{y^2 + x^2}$$

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\frac{\partial \phi}{\partial x} = M$$

$$\frac{\partial \phi}{\partial y} = N$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(x, y) dx + N(x, y) dy = 0 \quad (\text{1A})$$

Therefore

$$\left(\frac{-x + y}{x^2 + y^2} \right) dy = \left(-\frac{x + y}{x^2 + y^2} \right) dx$$

$$\left(\frac{x + y}{x^2 + y^2} \right) dx + \left(\frac{-x + y}{x^2 + y^2} \right) dy = 0 \quad (\text{2A})$$

Comparing (1A) and (2A) shows that

$$M(x, y) = \frac{x + y}{x^2 + y^2}$$

$$N(x, y) = \frac{-x + y}{x^2 + y^2}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Using result found above gives

$$\frac{\partial M}{\partial y} = \frac{\partial}{\partial y} \left(\frac{x + y}{x^2 + y^2} \right)$$

$$= \frac{x^2 - 2xy - y^2}{(x^2 + y^2)^2}$$

And

$$\frac{\partial N}{\partial x} = \frac{\partial}{\partial x} \left(\frac{-x + y}{x^2 + y^2} \right)$$

$$= \frac{x^2 - 2xy - y^2}{(x^2 + y^2)^2}$$

Since $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$, then the ODE is exact. The following equations are now set up to solve for the function $\phi(x, y)$

$$\frac{\partial \phi}{\partial x} = M \tag{1}$$

$$\frac{\partial \phi}{\partial y} = N \tag{2}$$

Integrating (1) w.r.t. x gives

$$\int \frac{\partial \phi}{\partial x} dx = \int M dx$$

$$\int \frac{\partial \phi}{\partial x} dx = \int \frac{x + y}{x^2 + y^2} dx$$

$$\phi = \frac{\ln(x^2 + y^2)}{2} + \arctan\left(\frac{x}{y}\right) + f(y) \tag{3}$$

Where $f(y)$ is used for the constant of integration since ϕ is a function of both x and y . Taking derivative of equation (3) w.r.t y gives

$$\begin{aligned}\frac{\partial\phi}{\partial y} &= \frac{y}{x^2 + y^2} - \frac{x}{y^2 \left(\frac{x^2}{y^2} + 1\right)} + f'(y) \\ &= \frac{-x + y}{x^2 + y^2} + f'(y)\end{aligned}\tag{4}$$

But equation (2) says that $\frac{\partial\phi}{\partial y} = \frac{-x+y}{x^2+y^2}$. Therefore equation (4) becomes

$$\frac{-x + y}{x^2 + y^2} = \frac{-x + y}{x^2 + y^2} + f'(y)\tag{5}$$

Solving equation (5) for $f'(y)$ gives

$$f'(y) = 0$$

Therefore

$$f(y) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(y)$ into equation (3) gives ϕ

$$\phi = \frac{\ln(x^2 + y^2)}{2} + \arctan\left(\frac{x}{y}\right) + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into the constant c_1 gives the solution as

$$c_1 = \frac{\ln(x^2 + y^2)}{2} + \arctan\left(\frac{x}{y}\right)$$

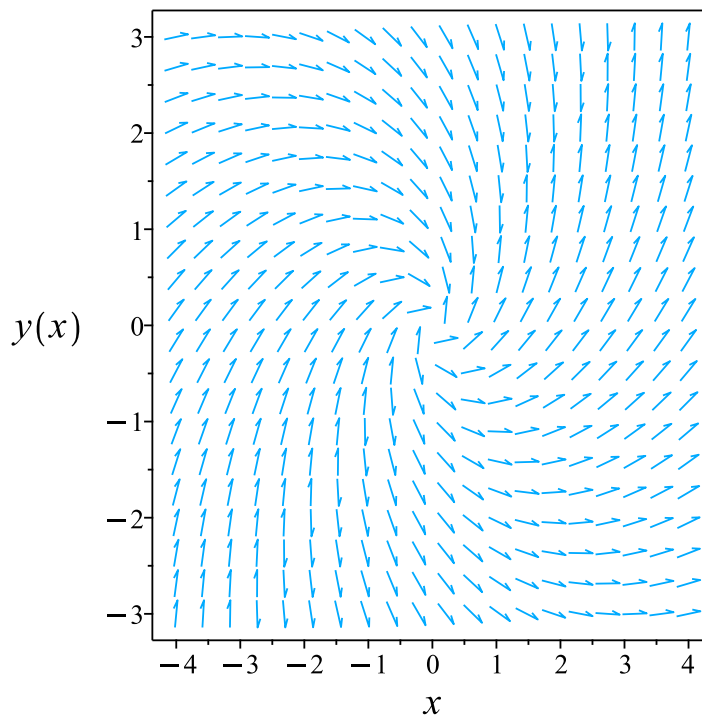


Figure 2.78: Slope field plot
 $y' = 1 + \frac{2y}{x-y}$

Summary of solutions found

$$\frac{\ln(y^2 + x^2)}{2} + \arctan\left(\frac{x}{y}\right) = c_1$$

Solved as first order isobaric ode

Time used: 0.467 (sec)

Solving for y' gives

$$y' = -\frac{x+y}{-x+y} \quad (1)$$

Each of the above ode's is now solved An ode $y' = f(x, y)$ is isobaric if

$$f(tx, t^m y) = t^{m-1} f(x, y) \quad (1)$$

Where here

$$f(x, y) = -\frac{x+y}{-x+y} \quad (2)$$

m is the order of isobaric. Substituting (2) into (1) and solving for m gives

$$m = 1$$

Since the ode is isobaric of order $m = 1$, then the substitution

$$\begin{aligned} y &= ux^m \\ &= ux \end{aligned}$$

Converts the ODE to a separable in $u(x)$. Performing this substitution gives

$$u(x) + xu'(x) = -\frac{x + xu(x)}{-x + xu(x)}$$

The ode $u'(x) = -\frac{u(x)^2+1}{(u(x)-1)x}$ is separable as it can be written as

$$\begin{aligned} u'(x) &= -\frac{u(x)^2 + 1}{(u(x) - 1)x} \\ &= f(x)g(u) \end{aligned}$$

Where

$$\begin{aligned} f(x) &= -\frac{1}{x} \\ g(u) &= \frac{u^2 + 1}{u - 1} \end{aligned}$$

Integrating gives

$$\begin{aligned} \int \frac{1}{g(u)} du &= \int f(x) dx \\ \int \frac{u-1}{u^2+1} du &= \int -\frac{1}{x} dx \\ \frac{\ln(u(x)^2+1)}{2} - \arctan(u(x)) &= \ln\left(\frac{1}{x}\right) + c_1 \end{aligned}$$

We now need to find the singular solutions, these are found by finding for what values $g(u)$ is zero, since we had to divide by this above. Solving $g(u) = 0$ or $\frac{u^2+1}{u-1} = 0$ for $u(x)$ gives

$$\begin{aligned} u(x) &= -i \\ u(x) &= i \end{aligned}$$

Now we go over each such singular solution and check if it verifies the ode itself and any initial conditions given. If it does not then the singular solution will not be used.

Therefore the solutions found are

$$\frac{\ln(u(x)^2 + 1)}{2} - \arctan(u(x)) = \ln\left(\frac{1}{x}\right) + c_1$$

$$u(x) = -i$$

$$u(x) = i$$

Converting $\frac{\ln(u(x)^2 + 1)}{2} - \arctan(u(x)) = \ln\left(\frac{1}{x}\right) + c_1$ back to y gives

$$\frac{\ln\left(\frac{y^2}{x^2} + 1\right)}{2} - \arctan\left(\frac{y}{x}\right) = \ln\left(\frac{1}{x}\right) + c_1$$

Converting $u(x) = -i$ back to y gives

$$\frac{y}{x} = -i$$

Converting $u(x) = i$ back to y gives

$$\frac{y}{x} = i$$

Solving for y gives

$$\frac{\ln\left(\frac{y^2}{x^2} + 1\right)}{2} - \arctan\left(\frac{y}{x}\right) = \ln\left(\frac{1}{x}\right) + c_1$$

$$y = -ix$$

$$y = ix$$

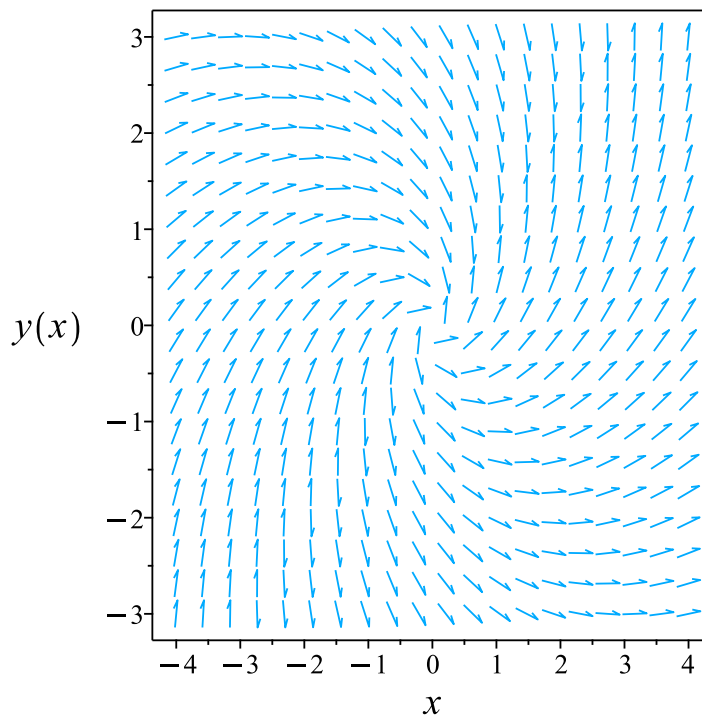


Figure 2.79: Slope field plot
 $y' = 1 + \frac{2y}{x-y}$

Summary of solutions found

$$\frac{\ln\left(\frac{y^2}{x^2} + 1\right)}{2} - \arctan\left(\frac{y}{x}\right) = \ln\left(\frac{1}{x}\right) + c_1$$

$$y = -ix$$

$$y = ix$$

Solved using Lie symmetry for first order ode

Time used: 0.420 (sec)

Writing the ode as

$$y' = -\frac{x+y}{-x+y}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2 \xi_y - \omega_x \xi - \omega_y \eta = 0 \quad (\text{A})$$

To determine ξ, η then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$\xi = xa_2 + ya_3 + a_1 \quad (1E)$$

$$\eta = xb_2 + yb_3 + b_1 \quad (2E)$$

Where the unknown coefficients are

$$\{a_1, a_2, a_3, b_1, b_2, b_3\}$$

Substituting equations (1E,2E) and ω into (A) gives

$$\begin{aligned} b_2 - \frac{(x+y)(b_3 - a_2)}{-x+y} - \frac{(x+y)^2 a_3}{(-x+y)^2} \\ - \left(-\frac{1}{-x+y} - \frac{x+y}{(-x+y)^2} \right) (xa_2 + ya_3 + a_1) \\ - \left(-\frac{1}{-x+y} + \frac{x+y}{(-x+y)^2} \right) (xb_2 + yb_3 + b_1) = 0 \end{aligned} \quad (5E)$$

Putting the above in normal form gives

$$\frac{x^2a_2 + x^2a_3 + x^2b_2 - x^2b_3 - 2xya_2 + 2xya_3 + 2xyb_2 + 2xyb_3 - y^2a_2 - y^2a_3 - y^2b_2 + y^2b_3 + 2xb_1 - 2ya_1}{(x-y)^2} = 0$$

Setting the numerator to zero gives

$$\begin{aligned} -x^2a_2 - x^2a_3 - x^2b_2 + x^2b_3 + 2xya_2 - 2xya_3 - 2xyb_2 \\ - 2xyb_3 + y^2a_2 + y^2a_3 + y^2b_2 - y^2b_3 - 2xb_1 + 2ya_1 = 0 \end{aligned} \quad (6E)$$

Looking at the above PDE shows the following are all the terms with $\{x, y\}$ in them.

$$\{x, y\}$$

The following substitution is now made to be able to collect on all terms with $\{x, y\}$ in them

$$\{x = v_1, y = v_2\}$$

The above PDE (6E) now becomes

$$\begin{aligned} & -a_2v_1^2 + 2a_2v_1v_2 + a_2v_2^2 - a_3v_1^2 - 2a_3v_1v_2 + a_3v_2^2 - b_2v_1^2 \\ & - 2b_2v_1v_2 + b_2v_2^2 + b_3v_1^2 - 2b_3v_1v_2 - b_3v_2^2 + 2a_1v_2 - 2b_1v_1 = 0 \end{aligned} \quad (7E)$$

Collecting the above on the terms v_i introduced, and these are

$$\{v_1, v_2\}$$

Equation (7E) now becomes

$$\begin{aligned} & (-a_2 - a_3 - b_2 + b_3)v_1^2 + (2a_2 - 2a_3 - 2b_2 - 2b_3)v_1v_2 \\ & - 2b_1v_1 + (a_2 + a_3 + b_2 - b_3)v_2^2 + 2a_1v_2 = 0 \end{aligned} \quad (8E)$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$\begin{aligned} 2a_1 &= 0 \\ -2b_1 &= 0 \\ -a_2 - a_3 - b_2 + b_3 &= 0 \\ a_2 + a_3 + b_2 - b_3 &= 0 \\ 2a_2 - 2a_3 - 2b_2 - 2b_3 &= 0 \end{aligned}$$

Solving the above equations for the unknowns gives

$$\begin{aligned} a_1 &= 0 \\ a_2 &= b_3 \\ a_3 &= -b_2 \\ b_1 &= 0 \\ b_2 &= b_2 \\ b_3 &= b_3 \end{aligned}$$

Substituting the above solution in the anstaz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$\begin{aligned} \xi &= x \\ \eta &= y \end{aligned}$$

Shifting is now applied to make $\xi = 0$ in order to simplify the rest of the computation

$$\begin{aligned}\eta &= \eta - \omega(x, y) \xi \\ &= y - \left(-\frac{x+y}{-x+y} \right) (x) \\ &= \frac{-x^2 - y^2}{x-y} \\ \xi &= 0\end{aligned}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS \quad (1)$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y} \right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = x$$

S is found from

$$\begin{aligned}S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{\frac{-x^2 - y^2}{x-y}} dy\end{aligned}$$

Which results in

$$S = \frac{\ln(x^2 + y^2)}{2} - \arctan\left(\frac{y}{x}\right)$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = -\frac{x+y}{-x+y}$$

Evaluating all the partial derivatives gives

$$\begin{aligned}R_x &= 1 \\R_y &= 0 \\S_x &= \frac{x+y}{x^2+y^2} \\S_y &= \frac{-x+y}{x^2+y^2}\end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = 0 \tag{2A}$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = 0$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S .

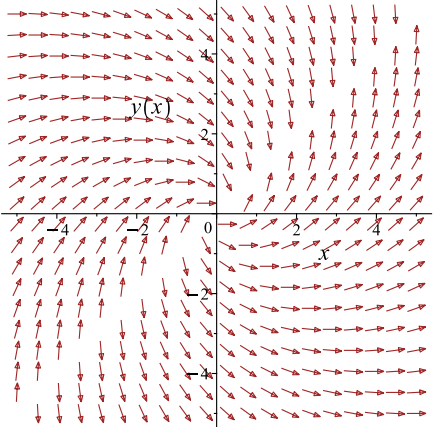
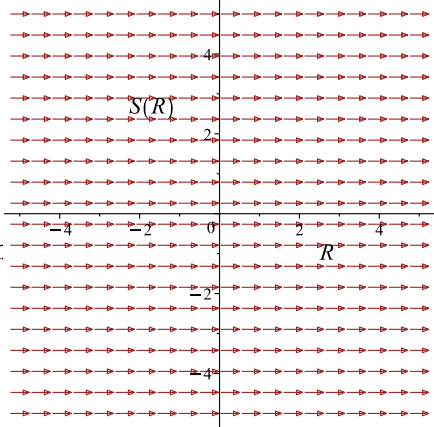
Since the ode has the form $\frac{d}{dR}S(R) = f(R)$, then we only need to integrate $f(R)$.

$$\begin{aligned}\int dS &= \int 0 dR + c_2 \\S(R) &= c_2\end{aligned}$$

To complete the solution, we just need to transform the above back to x, y coordinates. This results in

$$\frac{\ln(y^2 + x^2)}{2} - \arctan\left(\frac{y}{x}\right) = c_2$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in x, y coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dy}{dx} = -\frac{x+y}{-x+y}$ 	$R = x$ $S = \frac{\ln(x^2 + y^2)}{2} - \arctan\left(\frac{y}{x}\right)$	$\frac{dS}{dR} = 0$ 

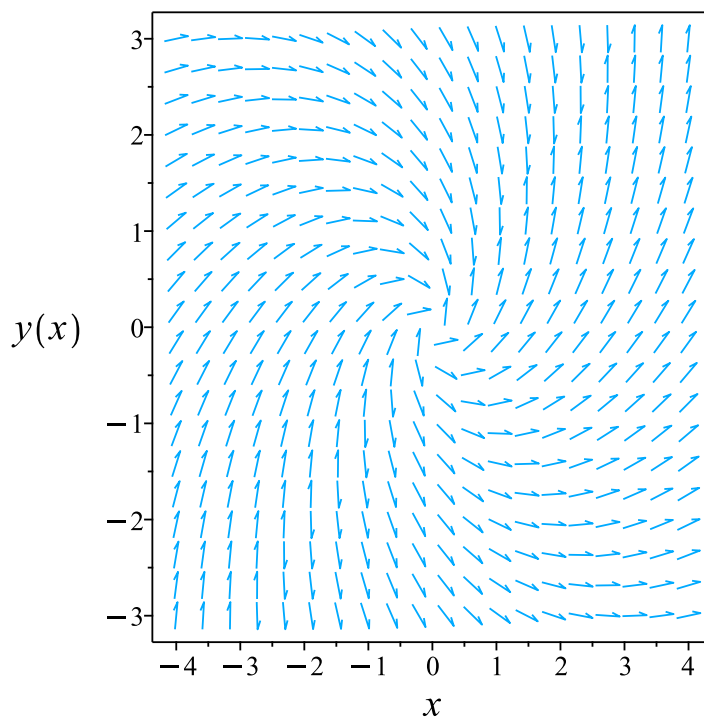


Figure 2.80: Slope field plot
 $y' = 1 + \frac{2y}{x-y}$

Summary of solutions found

$$\frac{\ln(y^2 + x^2)}{2} - \arctan\left(\frac{y}{x}\right) = c_2$$

Maple step by step solution

Let's solve

$$y' = 1 + \frac{2y}{x-y}$$

- Highest derivative means the order of the ODE is 1

y'

- Solve for the highest derivative

$$y' = 1 + \frac{2y}{x-y}$$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
trying inverse linear
trying homogeneous types:
trying homogeneous D
<- homogeneous successful`
```

Maple dsolve solution

Solving time : 0.013 (sec)

Leaf size : 24

```
dsolve(diff(y(x),x) = 1+2*y(x)/(x-y(x)),
        y(x),singsol=all)
```

$$y = \tan\left(\text{RootOf}\left(-2_Z + \ln\left(\sec(_Z)^2\right) + 2\ln(x) + 2c_1\right)\right) x$$

Mathematica DSolve solution

Solving time : 0.031 (sec)

Leaf size : 36

```
DSolve[{D[y[x],x]==1+2*y[x]/(x-y[x]),{}},  
y[x],x,IncludeSingularSolutions->True]
```

$$\text{Solve} \left[\frac{1}{2} \log \left(\frac{y(x)^2}{x^2} + 1 \right) - \arctan \left(\frac{y(x)}{x} \right) = -\log(x) + c_1, y(x) \right]$$

2.6.12 problem 10 (a)

Solved as first order linear ode	528
Solved as first order separable ode	530
Solved as first order Exact ode	531
Solved using Lie symmetry for first order ode	535
Maple step by step solution	540
Maple trace	541
Maple dsolve solution	541
Mathematica DSolve solution	541

Internal problem ID [18252]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter IV. Methods of solution: First order equations. section 33. Problems at page 91

Problem number : 10 (a)

Date solved : Monday, December 23, 2024 at 09:19:24 PM

CAS classification : [_separable]

Solve

$$v' + 2vu = 2u$$

Solved as first order linear ode

Time used: 0.100 (sec)

In canonical form a linear first order is

$$v' + q(u)v = p(u)$$

Comparing the above to the given ode shows that

$$q(u) = 2u$$

$$p(u) = 2u$$

The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int q \, du} \\ &= e^{\int 2u \, du} \\ &= e^{u^2}\end{aligned}$$

The ode becomes

$$\frac{d}{du}(\mu v) = \mu p$$

$$\frac{d}{du}(\mu v) = (\mu)(2u)$$

$$\frac{d}{du}(v e^{u^2}) = (e^{u^2})(2u)$$

$$d(v e^{u^2}) = (2u e^{u^2}) du$$

Integrating gives

$$\begin{aligned} v e^{u^2} &= \int 2u e^{u^2} du \\ &= e^{u^2} + c_1 \end{aligned}$$

Dividing throughout by the integrating factor e^{u^2} gives the final solution

$$v = 1 + c_1 e^{-u^2}$$

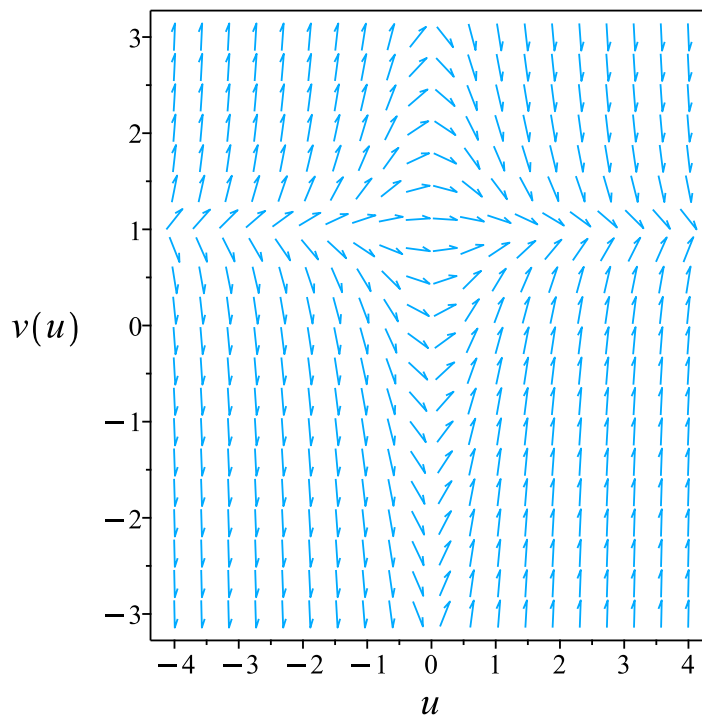


Figure 2.81: Slope field plot
 $v' + 2vu = 2u$

Summary of solutions found

$$v = 1 + c_1 e^{-u^2}$$

Solved as first order separable ode

Time used: 0.112 (sec)

The ode $v' = -2vu + 2u$ is separable as it can be written as

$$\begin{aligned} v' &= -2vu + 2u \\ &= f(u)g(v) \end{aligned}$$

Where

$$\begin{aligned} f(u) &= u \\ g(v) &= -2v + 2 \end{aligned}$$

Integrating gives

$$\begin{aligned} \int \frac{1}{g(v)} dv &= \int f(u) du \\ \int \frac{1}{-2v + 2} dv &= \int u du \\ -\frac{\ln(v - 1)}{2} &= \frac{u^2}{2} + c_1 \end{aligned}$$

We now need to find the singular solutions, these are found by finding for what values $g(v)$ is zero, since we had to divide by this above. Solving $g(v) = 0$ or $-2v + 2 = 0$ for v gives

$$v = 1$$

Now we go over each such singular solution and check if it verifies the ode itself and any initial conditions given. If it does not then the singular solution will not be used.

Therefore the solutions found are

$$\begin{aligned} -\frac{\ln(v - 1)}{2} &= \frac{u^2}{2} + c_1 \\ v &= 1 \end{aligned}$$

Solving for v gives

$$\begin{aligned} v &= 1 \\ v &= e^{-u^2 - 2c_1} + 1 \end{aligned}$$

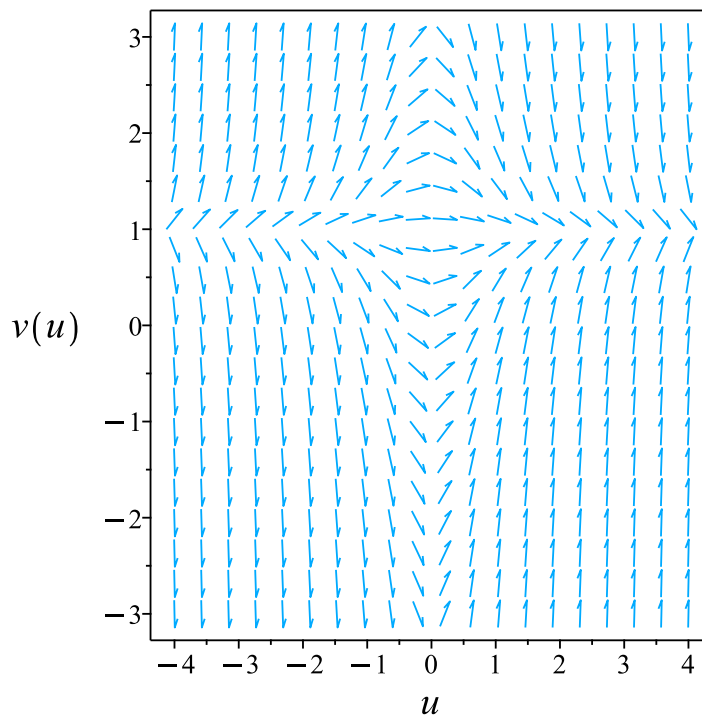


Figure 2.82: Slope field plot
 $v' + 2vu = 2u$

Summary of solutions found

$$v = 1$$

$$v = e^{-u^2 - 2c_1} + 1$$

Solved as first order Exact ode

Time used: 0.369 (sec)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial\phi}{\partial x} &= M \\ \frac{\partial\phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2\phi}{\partial x\partial y} = \frac{\partial^2\phi}{\partial y\partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(u, v) du + N(u, v) dv = 0 \tag{1A}$$

Therefore

$$\begin{aligned}dv &= (-2uv + 2u) du \\ (2uv - 2u) du + dv &= 0\end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(u, v) &= 2uv - 2u \\ N(u, v) &= 1\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial v} = \frac{\partial N}{\partial u}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial v} &= \frac{\partial}{\partial v}(2uv - 2u) \\ &= 2u\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial u} &= \frac{\partial}{\partial u}(1) \\ &= 0\end{aligned}$$

Since $\frac{\partial M}{\partial v} \neq \frac{\partial N}{\partial u}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned} A &= \frac{1}{N} \left(\frac{\partial M}{\partial v} - \frac{\partial N}{\partial u} \right) \\ &= 1((2u) - (0)) \\ &= 2u \end{aligned}$$

Since A does not depend on v , then it can be used to find an integrating factor. The integrating factor μ is

$$\begin{aligned} \mu &= e^{\int A \, du} \\ &= e^{\int 2u \, du} \end{aligned}$$

The result of integrating gives

$$\begin{aligned} \mu &= e^{u^2} \\ &= e^{u^2} \end{aligned}$$

M and N are multiplied by this integrating factor, giving new M and new N which are called \bar{M} and \bar{N} for now so not to confuse them with the original M and N .

$$\begin{aligned} \bar{M} &= \mu M \\ &= e^{u^2}(2uv - 2u) \\ &= 2u(v - 1)e^{u^2} \end{aligned}$$

And

$$\begin{aligned} \bar{N} &= \mu N \\ &= e^{u^2}(1) \\ &= e^{u^2} \end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned} \bar{M} + \bar{N} \frac{dv}{du} &= 0 \\ \left(2u(v - 1)e^{u^2} \right) + \left(e^{u^2} \right) \frac{dv}{du} &= 0 \end{aligned}$$

The following equations are now set up to solve for the function $\phi(u, v)$

$$\frac{\partial \phi}{\partial u} = \bar{M} \tag{1}$$

$$\frac{\partial \phi}{\partial v} = \bar{N} \tag{2}$$

Integrating (2) w.r.t. v gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial v} dv &= \int \bar{N} dv \\ \int \frac{\partial \phi}{\partial v} dv &= \int e^{u^2} dv \\ \phi &= v e^{u^2} + f(u)\end{aligned}\tag{3}$$

Where $f(u)$ is used for the constant of integration since ϕ is a function of both u and v . Taking derivative of equation (3) w.r.t u gives

$$\frac{\partial \phi}{\partial u} = 2vu e^{u^2} + f'(u)\tag{4}$$

But equation (1) says that $\frac{\partial \phi}{\partial u} = 2u(v-1)e^{u^2}$. Therefore equation (4) becomes

$$2u(v-1)e^{u^2} = 2vu e^{u^2} + f'(u)\tag{5}$$

Solving equation (5) for $f'(u)$ gives

$$f'(u) = -2u e^{u^2}$$

Integrating the above w.r.t u gives

$$\begin{aligned}\int f'(u) du &= \int (-2u e^{u^2}) du \\ f(u) &= -e^{u^2} + c_1\end{aligned}$$

Where c_1 is constant of integration. Substituting result found above for $f(u)$ into equation (3) gives ϕ

$$\phi = v e^{u^2} - e^{u^2} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into the constant c_1 gives the solution as

$$c_1 = v e^{u^2} - e^{u^2}$$

Solving for v gives

$$v = e^{-u^2} (e^{u^2} + c_1)$$

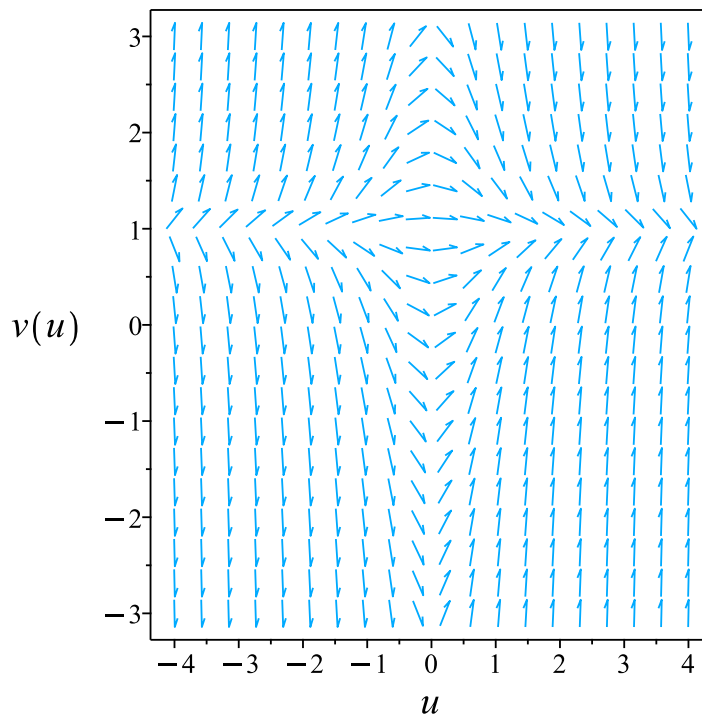


Figure 2.83: Slope field plot
 $v' + 2vu = 2u$

Summary of solutions found

$$v = e^{-u^2} (e^{u^2} + c_1)$$

Solved using Lie symmetry for first order ode

Time used: 0.356 (sec)

Writing the ode as

$$v' = -2uv + 2u$$

$$v' = \omega(u, v)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_u + \omega(\eta_v - \xi_u) - \omega^2 \xi_v - \omega_u \xi - \omega_v \eta = 0 \quad (\text{A})$$

To determine ξ, η then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$\xi = ua_2 + va_3 + a_1 \quad (1E)$$

$$\eta = ub_2 + vb_3 + b_1 \quad (2E)$$

Where the unknown coefficients are

$$\{a_1, a_2, a_3, b_1, b_2, b_3\}$$

Substituting equations (1E,2E) and ω into (A) gives

$$b_2 + (-2uv + 2u)(b_3 - a_2) - (-2uv + 2u)^2 a_3 - (-2v + 2)(ua_2 + va_3 + a_1) + 2u(ub_2 + vb_3 + b_1) = 0 \quad (5E)$$

Putting the above in normal form gives

$$-4u^2v^2a_3 + 8u^2va_3 - 4u^2a_3 + 2u^2b_2 + 4uva_2 + 2v^2a_3 - 4ua_2 + 2ub_1 + 2ub_3 + 2va_1 - 2va_3 - 2a_1 + b_2 = 0$$

Setting the numerator to zero gives

$$-4u^2v^2a_3 + 8u^2va_3 - 4u^2a_3 + 2u^2b_2 + 4uva_2 + 2v^2a_3 - 4ua_2 + 2ub_1 + 2ub_3 + 2va_1 - 2va_3 - 2a_1 + b_2 = 0 \quad (6E)$$

Looking at the above PDE shows the following are all the terms with $\{u, v\}$ in them.

$$\{u, v\}$$

The following substitution is now made to be able to collect on all terms with $\{u, v\}$ in them

$$\{u = v_1, v = v_2\}$$

The above PDE (6E) now becomes

$$-4a_3v_1^2v_2^2 + 8a_3v_1^2v_2 + 4a_2v_1v_2 - 4a_3v_1^2 + 2a_3v_2^2 + 2b_2v_1^2 + 2a_1v_2 - 4a_2v_1 - 2a_3v_2 + 2b_1v_1 + 2b_3v_1 - 2a_1 + b_2 = 0 \quad (7E)$$

Collecting the above on the terms v_i introduced, and these are

$$\{v_1, v_2\}$$

Equation (7E) now becomes

$$\begin{aligned} -4a_3v_1^2v_2^2 + 8a_3v_1^2v_2 + (-4a_3 + 2b_2)v_1^2 + 4a_2v_1v_2 \\ + (-4a_2 + 2b_1 + 2b_3)v_1 + 2a_3v_2^2 + (2a_1 - 2a_3)v_2 - 2a_1 + b_2 = 0 \end{aligned} \quad (8E)$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$\begin{aligned} 4a_2 &= 0 \\ -4a_3 &= 0 \\ 2a_3 &= 0 \\ 8a_3 &= 0 \\ -2a_1 + b_2 &= 0 \\ 2a_1 - 2a_3 &= 0 \\ -4a_3 + 2b_2 &= 0 \\ -4a_2 + 2b_1 + 2b_3 &= 0 \end{aligned}$$

Solving the above equations for the unknowns gives

$$\begin{aligned} a_1 &= 0 \\ a_2 &= 0 \\ a_3 &= 0 \\ b_1 &= -b_3 \\ b_2 &= 0 \\ b_3 &= b_3 \end{aligned}$$

Substituting the above solution in the ansatz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$\begin{aligned} \xi &= 0 \\ \eta &= v - 1 \end{aligned}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(u, v) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{du}{\xi} = \frac{dv}{\eta} = dS \quad (1)$$

The above comes from the requirements that $(\xi \frac{\partial}{\partial u} + \eta \frac{\partial}{\partial v}) S(u, v) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = u$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{v-1} dy \end{aligned}$$

Which results in

$$S = \ln(v-1)$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_u + \omega(u, v)S_v}{R_u + \omega(u, v)R_v} \quad (2)$$

Where in the above R_u, R_v, S_u, S_v are all partial derivatives and $\omega(u, v)$ is the right hand side of the original ode given by

$$\omega(u, v) = -2uv + 2u$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_u &= 1 \\ R_v &= 0 \\ S_u &= 0 \\ S_v &= \frac{1}{v-1} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = -2u \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for u, v in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = -2R$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S .

Since the ode has the form $\frac{d}{dR}S(R) = f(R)$, then we only need to integrate $f(R)$.

$$\int dS = \int -2R dR$$

$$S(R) = -R^2 + c_2$$

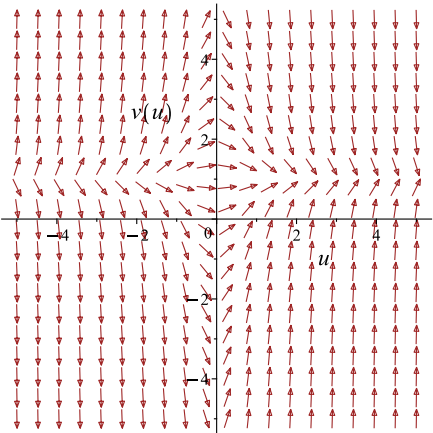
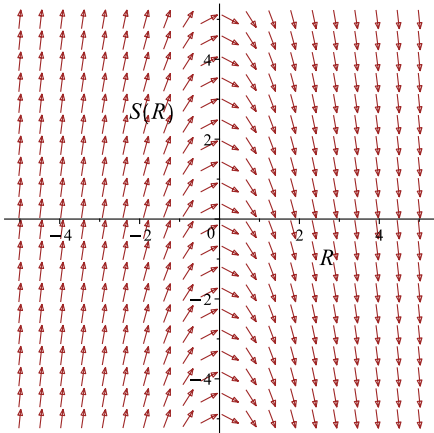
To complete the solution, we just need to transform the above back to u, v coordinates. This results in

$$\ln(v - 1) = -u^2 + c_2$$

Which gives

$$v = e^{-u^2 + c_2} + 1$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in u, v coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dv}{du} = -2uv + 2u$ 	$R = u$ $S = \ln(v - 1)$	$\frac{dS}{dR} = -2R$ 

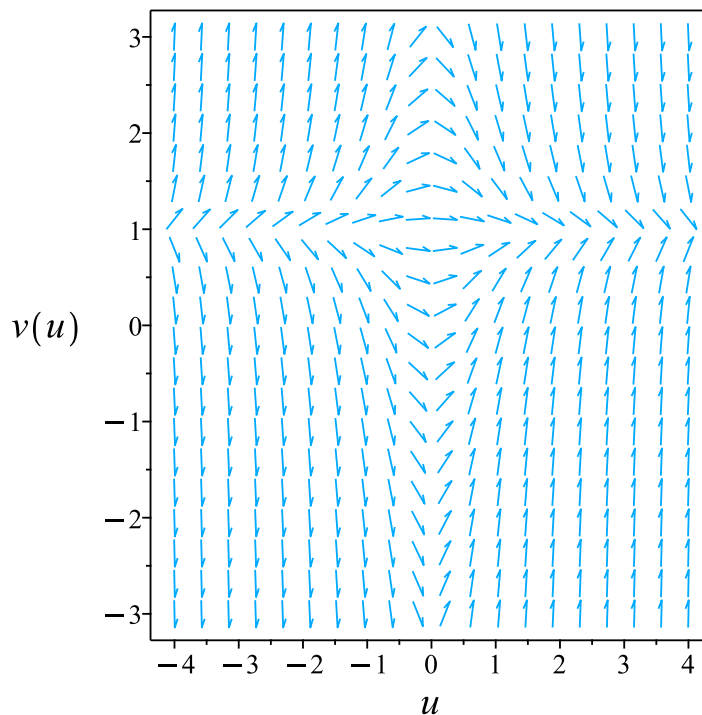


Figure 2.84: Slope field plot
 $v' + 2vu = 2u$

Summary of solutions found

$$v = e^{-u^2+c_2} + 1$$

Maple step by step solution

Let's solve

$$v' + 2vu = 2u$$

- Highest derivative means the order of the ODE is 1

$$v'$$

- Solve for the highest derivative

$$v' = -2vu + 2u$$

- Separate variables

$$\frac{v'}{v-1} = -2u$$

- Integrate both sides with respect to u

$$\int \frac{v'}{v-1} du = \int -2u du + C1$$

- Evaluate integral

$$\ln(v-1) = -u^2 + C1$$

- Solve for v

$$v = e^{-u^2+C_1} + 1$$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`
```

Maple dsolve solution

Solving time : 0.002 (sec)

Leaf size : 14

```
dsolve(diff(v(u),u)+2*v(u)*u = 2*u,
        v(u),singsol=all)
```

$$v = 1 + c_1 e^{-u^2}$$

Mathematica DSolve solution

Solving time : 0.047 (sec)

Leaf size : 22

```
DSolve[{D[v[u],u]+2*u*v[u]==2*u,{}},
        v[u],u,IncludeSingularSolutions->True]
```

$$v(u) \rightarrow 1 + c_1 e^{-u^2}$$

$$v(u) \rightarrow 1$$

2.6.13 problem 10 (b)

Solved as first order separable ode	542
Solved as first order Bernoulli ode	544
Solved as first order Exact ode	548
Solved using Lie symmetry for first order ode	552
Maple step by step solution	559
Maple trace	560
Maple dsolve solution	560
Mathematica DSolve solution	561

Internal problem ID [18253]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter IV. Methods of solution: First order equations. section 33. Problems at page 91

Problem number : 10 (b)

Date solved : Monday, December 23, 2024 at 09:19:25 PM

CAS classification : [_separable]

Solve

$$1 + v^2 + (u^2 + 1)vv' = 0$$

Solved as first order separable ode

Time used: 0.248 (sec)

The ode $v' = -\frac{v^2+1}{(u^2+1)v}$ is separable as it can be written as

$$\begin{aligned} v' &= -\frac{v^2 + 1}{(u^2 + 1)v} \\ &= f(u)g(v) \end{aligned}$$

Where

$$\begin{aligned} f(u) &= -\frac{1}{u^2 + 1} \\ g(v) &= \frac{v^2 + 1}{v} \end{aligned}$$

Integrating gives

$$\int \frac{1}{g(v)} dv = \int f(u) du$$

$$\int \frac{v}{v^2 + 1} dv = \int -\frac{1}{u^2 + 1} du$$

$$\frac{\ln(v^2 + 1)}{2} = -\arctan(u) + c_1$$

We now need to find the singular solutions, these are found by finding for what values $g(v)$ is zero, since we had to divide by this above. Solving $g(v) = 0$ or $\frac{v^2+1}{v} = 0$ for v gives

$$v = -i$$

$$v = i$$

Now we go over each such singular solution and check if it verifies the ode itself and any initial conditions given. If it does not then the singular solution will not be used.

Therefore the solutions found are

$$\frac{\ln(v^2 + 1)}{2} = -\arctan(u) + c_1$$

$$v = -i$$

$$v = i$$

Solving for v gives

$$v = -i$$

$$v = i$$

$$v = \sqrt{-1 + e^{-2 \arctan(u) + 2c_1}}$$

$$v = -\sqrt{-1 + e^{-2 \arctan(u) + 2c_1}}$$

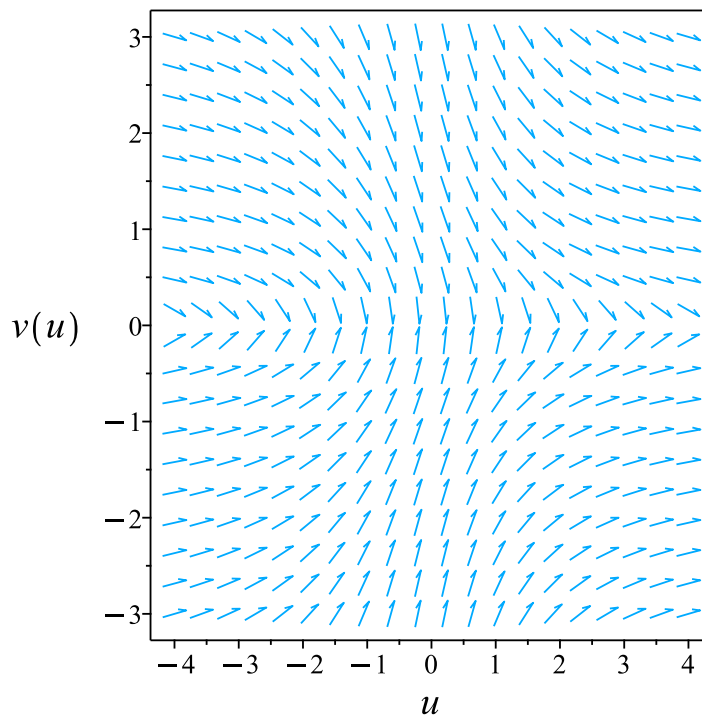


Figure 2.85: Slope field plot
 $1 + v^2 + (u^2 + 1)vv' = 0$

Summary of solutions found

$$v = -i$$

$$v = i$$

$$v = \sqrt{-1 + e^{-2 \arctan(u) + 2c_1}}$$

$$v = -\sqrt{-1 + e^{-2 \arctan(u) + 2c_1}}$$

Solved as first order Bernoulli ode

Time used: 0.276 (sec)

In canonical form, the ODE is

$$\begin{aligned} v' &= F(u, v) \\ &= -\frac{v^2 + 1}{(u^2 + 1)v} \end{aligned}$$

This is a Bernoulli ODE.

$$v' = \left(-\frac{1}{u^2 + 1}\right)v + \left(-\frac{1}{u^2 + 1}\right)\frac{1}{v} \quad (1)$$

The standard Bernoulli ODE has the form

$$v' = f_0(u)v + f_1(u)v^n \quad (2)$$

Comparing this to (1) shows that

$$f_0 = -\frac{1}{u^2 + 1}$$

$$f_1 = -\frac{1}{u^2 + 1}$$

The first step is to divide the above equation by v^n which gives

$$\frac{v'}{v^n} = f_0(u)v^{1-n} + f_1(u) \quad (3)$$

The next step is use the substitution $v = v^{1-n}$ in equation (3) which generates a new ODE in $v(u)$ which will be linear and can be easily solved using an integrating factor. Backsubstitution then gives the solution $v(u)$ which is what we want.

This method is now applied to the ODE at hand. Comparing the ODE (1) With (2) Shows that

$$f_0(u) = -\frac{1}{u^2 + 1}$$

$$f_1(u) = -\frac{1}{u^2 + 1}$$

$$n = -1$$

Dividing both sides of ODE (1) by $v^n = \frac{1}{v}$ gives

$$v'v = -\frac{v^2}{u^2 + 1} - \frac{1}{u^2 + 1} \quad (4)$$

Let

$$v = v^{1-n}$$

$$= v^2 \quad (5)$$

Taking derivative of equation (5) w.r.t u gives

$$v' = 2vv' \quad (6)$$

Substituting equations (5) and (6) into equation (4) gives

$$\frac{v'(u)}{2} = -\frac{v(u)}{u^2 + 1} - \frac{1}{u^2 + 1}$$

$$v' = -\frac{2v}{u^2 + 1} - \frac{2}{u^2 + 1} \quad (7)$$

The above now is a linear ODE in $v(u)$ which is now solved.

In canonical form a linear first order is

$$v'(u) + q(u)v(u) = p(u)$$

Comparing the above to the given ode shows that

$$q(u) = \frac{2}{u^2 + 1}$$

$$p(u) = -\frac{2}{u^2 + 1}$$

The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int q \, du} \\ &= e^{\int \frac{2}{u^2+1} \, du} \\ &= e^{2 \arctan(u)}\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{du}(\mu v) &= \mu p \\ \frac{d}{du}(\mu v) &= (\mu) \left(-\frac{2}{u^2 + 1} \right) \\ \frac{d}{du}(v e^{2 \arctan(u)}) &= (e^{2 \arctan(u)}) \left(-\frac{2}{u^2 + 1} \right) \\ d(v e^{2 \arctan(u)}) &= \left(-\frac{2 e^{2 \arctan(u)}}{u^2 + 1} \right) du\end{aligned}$$

Integrating gives

$$\begin{aligned}v e^{2 \arctan(u)} &= \int -\frac{2 e^{2 \arctan(u)}}{u^2 + 1} \, du \\ &= -e^{2 \arctan(u)} + c_1\end{aligned}$$

Dividing throughout by the integrating factor $e^{2 \arctan(u)}$ gives the final solution

$$v(u) = -1 + c_1 e^{-2 \arctan(u)}$$

The substitution $v = v^{1-n}$ is now used to convert the above solution back to v which results in

$$v^2 = -1 + c_1 e^{-2 \arctan(u)}$$

Solving for v gives

$$v = \sqrt{-1 + c_1 e^{-2 \arctan(u)}}$$

$$v = -\sqrt{-1 + c_1 e^{-2 \arctan(u)}}$$

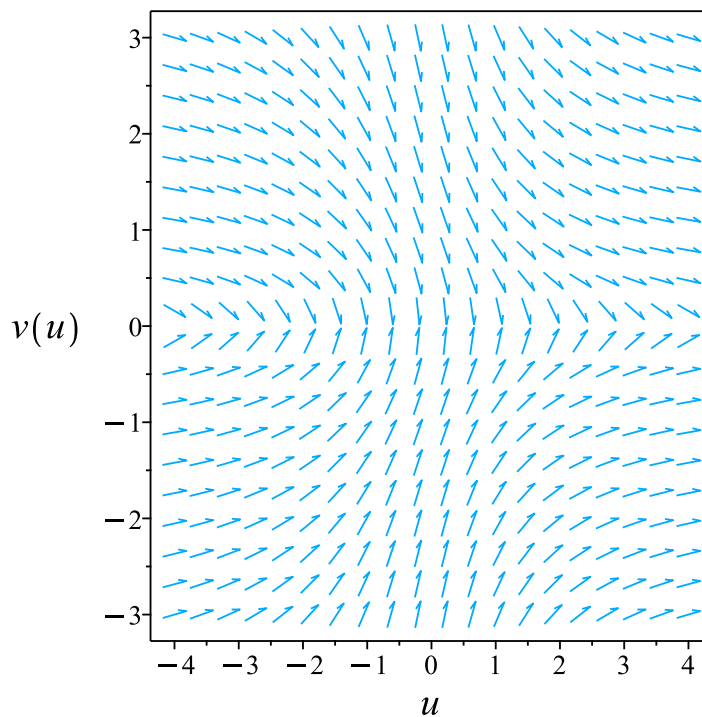


Figure 2.86: Slope field plot
 $1 + v^2 + (u^2 + 1)vv' = 0$

Summary of solutions found

$$v = \sqrt{-1 + c_1 e^{-2 \arctan(u)}}$$

$$v = -\sqrt{-1 + c_1 e^{-2 \arctan(u)}}$$

Solved as first order Exact ode

Time used: 0.514 (sec)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N \end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(u, v) du + N(u, v) dv = 0 \quad (\text{1A})$$

Therefore

$$\begin{aligned} ((u^2 + 1) v) dv &= (-v^2 - 1) du \\ (v^2 + 1) du + ((u^2 + 1) v) dv &= 0 \end{aligned} \quad (\text{2A})$$

Comparing (1A) and (2A) shows that

$$\begin{aligned} M(u, v) &= v^2 + 1 \\ N(u, v) &= (u^2 + 1) v \end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial v} = \frac{\partial N}{\partial u}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial v} &= \frac{\partial}{\partial v}(v^2 + 1) \\ &= 2v\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial u} &= \frac{\partial}{\partial u}((u^2 + 1)v) \\ &= 2uv\end{aligned}$$

Since $\frac{\partial M}{\partial v} \neq \frac{\partial N}{\partial u}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned}A &= \frac{1}{N} \left(\frac{\partial M}{\partial v} - \frac{\partial N}{\partial u} \right) \\ &= \frac{1}{(u^2 + 1)v} ((2v) - (2uv)) \\ &= \frac{-2u + 2}{u^2 + 1}\end{aligned}$$

Since A does not depend on v , then it can be used to find an integrating factor. The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int A \, du} \\ &= e^{\int \frac{-2u+2}{u^2+1} \, du}\end{aligned}$$

The result of integrating gives

$$\begin{aligned}\mu &= e^{-\ln(u^2+1)+2\arctan(u)} \\ &= \frac{e^{2\arctan(u)}}{u^2 + 1}\end{aligned}$$

M and N are multiplied by this integrating factor, giving new M and new N which are called \bar{M} and \bar{N} for now so not to confuse them with the original M and N .

$$\begin{aligned}\bar{M} &= \mu M \\ &= \frac{e^{2\arctan(u)}}{u^2 + 1} (v^2 + 1) \\ &= \frac{(v^2 + 1)e^{2\arctan(u)}}{u^2 + 1}\end{aligned}$$

And

$$\begin{aligned}\bar{N} &= \mu N \\ &= \frac{e^{2\arctan(u)}}{u^2 + 1} ((u^2 + 1)v) \\ &= v e^{2\arctan(u)}\end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned}\bar{M} + \bar{N} \frac{dv}{du} &= 0 \\ \left(\frac{(v^2 + 1)e^{2\arctan(u)}}{u^2 + 1} \right) + (v e^{2\arctan(u)}) \frac{dv}{du} &= 0\end{aligned}$$

The following equations are now set up to solve for the function $\phi(u, v)$

$$\frac{\partial \phi}{\partial u} = \bar{M} \tag{1}$$

$$\frac{\partial \phi}{\partial v} = \bar{N} \tag{2}$$

Integrating (1) w.r.t. u gives

$$\begin{aligned}\int \frac{\partial \phi}{\partial u} du &= \int \bar{M} du \\ \int \frac{\partial \phi}{\partial u} du &= \int \frac{(v^2 + 1)e^{2\arctan(u)}}{u^2 + 1} du \\ \phi &= \frac{(v^2 + 1)e^{2\arctan(u)}}{2} + f(v)\end{aligned} \tag{3}$$

Where $f(v)$ is used for the constant of integration since ϕ is a function of both u and v . Taking derivative of equation (3) w.r.t v gives

$$\frac{\partial \phi}{\partial v} = v e^{2\arctan(u)} + f'(v) \tag{4}$$

But equation (2) says that $\frac{\partial \phi}{\partial v} = v e^{2\arctan(u)}$. Therefore equation (4) becomes

$$v e^{2\arctan(u)} = v e^{2\arctan(u)} + f'(v) \tag{5}$$

Solving equation (5) for $f'(v)$ gives

$$f'(v) = 0$$

Therefore

$$f(v) = c_1$$

Where c_1 is constant of integration. Substituting this result for $f(v)$ into equation (3) gives ϕ

$$\phi = \frac{(v^2 + 1) e^{2 \arctan(u)}}{2} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into the constant c_1 gives the solution as

$$c_1 = \frac{(v^2 + 1) e^{2 \arctan(u)}}{2}$$

Solving for v gives

$$v = e^{-2 \arctan(u)} \sqrt{-e^{2 \arctan(u)} (e^{2 \arctan(u)} - 2c_1)}$$

$$v = -e^{-2 \arctan(u)} \sqrt{-e^{2 \arctan(u)} (e^{2 \arctan(u)} - 2c_1)}$$

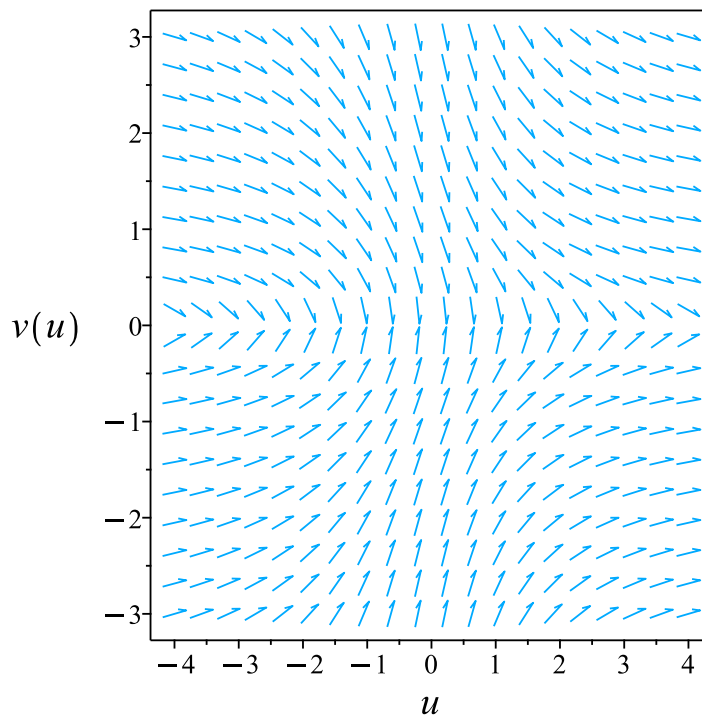


Figure 2.87: Slope field plot
 $1 + v^2 + (u^2 + 1)vv' = 0$

Summary of solutions found

$$v = e^{-2 \arctan(u)} \sqrt{-e^{2 \arctan(u)} (e^{2 \arctan(u)} - 2c_1)}$$

$$v = -e^{-2 \arctan(u)} \sqrt{-e^{2 \arctan(u)} (e^{2 \arctan(u)} - 2c_1)}$$

Solved using Lie symmetry for first order ode

Time used: 1.408 (sec)

Writing the ode as

$$v' = -\frac{v^2 + 1}{(u^2 + 1)v}$$

$$v' = \omega(u, v)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_u + \omega(\eta_v - \xi_u) - \omega^2 \xi_v - \omega_u \xi - \omega_v \eta = 0 \quad (\text{A})$$

To determine ξ, η then (A) is solved using ansatz. Making bivariate polynomials of degree 2 to use as anstaz gives

$$\xi = u^2 a_4 + uva_5 + v^2 a_6 + ua_2 + va_3 + a_1 \quad (1E)$$

$$\eta = u^2 b_4 + uvb_5 + v^2 b_6 + ub_2 + vb_3 + b_1 \quad (2E)$$

Where the unknown coefficients are

$$\{a_1, a_2, a_3, a_4, a_5, a_6, b_1, b_2, b_3, b_4, b_5, b_6\}$$

Substituting equations (1E,2E) and ω into (A) gives

$$\begin{aligned} & 2ub_4 + vb_5 + b_2 - \frac{(v^2 + 1)(-2ua_4 + ub_5 - va_5 + 2vb_6 - a_2 + b_3)}{(u^2 + 1)v} \\ & - \frac{(v^2 + 1)^2 (ua_5 + 2va_6 + a_3)}{(u^2 + 1)^2 v^2} \\ & - \frac{2(v^2 + 1)u(u^2 a_4 + uva_5 + v^2 a_6 + ua_2 + va_3 + a_1)}{(u^2 + 1)^2 v} \\ & - \left(-\frac{2}{u^2 + 1} + \frac{v^2 + 1}{(u^2 + 1)v^2} \right) (u^2 b_4 + uvb_5 + v^2 b_6 + ub_2 + vb_3 + b_1) = 0 \end{aligned} \quad (5E)$$

Putting the above in normal form gives

$$\begin{aligned} & \underline{2u^5 v^2 b_4 + u^4 v^3 b_5 + u^4 v^2 b_2 + u^4 v^2 b_4 - u^2 v^4 a_5 - u^2 v^4 b_6 - 2u v^5 a_6 + u^3 v^2 b_2 + 4u^3 v^2 b_4 - u^2 v^3 a_2 + 2u^2 v^3 b_5} \\ & = 0 \end{aligned}$$

Setting the numerator to zero gives

$$\begin{aligned} & 2u^5 v^2 b_4 + u^4 v^3 b_5 + u^4 v^2 b_2 + u^4 v^2 b_4 - u^2 v^4 a_5 - u^2 v^4 b_6 - 2u v^5 a_6 \\ & + u^3 v^2 b_2 + 4u^3 v^2 b_4 - u^2 v^3 a_2 + 2u^2 v^3 b_5 - 2u v^4 a_3 - u v^4 a_5 - 2v^5 a_6 \\ & - u^4 b_4 - 2u^3 v b_5 - u^2 v^2 a_5 + u^2 v^2 b_1 + 2u^2 v^2 b_2 + u^2 v^2 b_4 - 3u^2 v^2 b_6 \\ & - 2u v^3 a_1 + 2u v^3 a_4 - 2u v^3 a_6 - v^4 a_3 + v^4 a_5 - v^4 b_6 - u^3 b_2 - u^2 v a_2 \\ & - 2u^2 v b_3 - 2u v^2 a_3 - 2u v^2 a_5 + u v^2 b_2 + 2ub_4 v^2 + v^3 a_2 - 4v^3 a_6 \\ & + v^3 b_5 - u^2 b_1 - u^2 b_4 - 2uva_1 + 2uva_4 - 2uvb_5 - 2v^2 a_3 + v^2 a_5 + v^2 b_1 \\ & + b_2 v^2 - 3v^2 b_6 - ua_5 - ub_2 + va_2 - 2va_6 - 2vb_3 - a_3 - b_1 = 0 \end{aligned} \quad (6E)$$

Looking at the above PDE shows the following are all the terms with $\{u, v\}$ in them.

$$\{u, v\}$$

The following substitution is now made to be able to collect on all terms with $\{u, v\}$ in them

$$\{u = v_1, v = v_2\}$$

The above PDE (6E) now becomes

$$\begin{aligned} & 2b_4v_1^5v_2^2 + b_5v_1^4v_2^3 - a_5v_1^2v_2^4 - 2a_6v_1v_2^5 + b_2v_1^4v_2^2 + b_4v_1^4v_2^2 - b_6v_1^2v_2^4 \\ & - a_2v_1^2v_2^3 - 2a_3v_1v_2^4 - a_5v_1v_2^4 - 2a_6v_2^5 + b_2v_1^3v_2^2 + 4b_4v_1^3v_2^2 + 2b_5v_1^2v_2^3 \\ & - 2a_1v_1v_2^3 - a_3v_2^4 + 2a_4v_1v_2^3 - a_5v_1^2v_2^2 + a_5v_2^4 - 2a_6v_1v_2^3 + b_1v_1^2v_2^2 \\ & + 2b_2v_1^2v_2^2 - b_4v_1^4 + b_4v_1^2v_2^2 - 2b_5v_1^3v_2 - 3b_6v_1^2v_2^2 - b_6v_2^4 - a_2v_1^2v_2 + a_2v_2^3 \\ & - 2a_3v_1v_2^2 - 2a_5v_1v_2^2 - 4a_6v_2^3 - b_2v_1^3 + b_2v_1v_2^2 - 2b_3v_1^2v_2 + 2b_4v_1v_2^2 \\ & + b_5v_2^3 - 2a_1v_1v_2 - 2a_3v_2^2 + 2a_4v_1v_2 + a_5v_2^2 - b_1v_1^2 + b_1v_2^2 + b_2v_2^2 - b_4v_1^2 \\ & - 2b_5v_1v_2 - 3b_6v_2^2 + a_2v_2 - a_5v_1 - 2a_6v_2 - b_2v_1 - 2b_3v_2 - a_3 - b_1 = 0 \end{aligned} \quad (7E)$$

Collecting the above on the terms v_i introduced, and these are

$$\{v_1, v_2\}$$

Equation (7E) now becomes

$$\begin{aligned} & 2b_4v_1^5v_2^2 + b_5v_1^4v_2^3 - 2a_6v_1v_2^5 - 2b_5v_1^3v_2 - 2a_6v_2^5 - b_4v_1^4 - b_2v_1^3 \\ & - a_3 - b_1 + (-b_1 - b_4)v_1^2 + (-a_5 - b_2)v_1 + (-a_3 + a_5 - b_6)v_2^4 \\ & + (a_2 - 4a_6 + b_5)v_2^3 + (-2a_3 + a_5 + b_1 + b_2 - 3b_6)v_2^2 \\ & + (a_2 - 2a_6 - 2b_3)v_2 + (-a_5 + b_1 + 2b_2 + b_4 - 3b_6)v_1^2v_2^2 \\ & + (-a_2 - 2b_3)v_1^2v_2 + (-2a_3 - a_5)v_1v_2^4 + (-2a_1 + 2a_4 - 2a_6)v_1v_2^3 \\ & + (-2a_3 - 2a_5 + b_2 + 2b_4)v_1v_2^2 + (-2a_1 + 2a_4 - 2b_5)v_1v_2 + (b_2 + b_4)v_1^4v_2^2 \\ & + (b_2 + 4b_4)v_1^3v_2^2 + (-a_5 - b_6)v_1^2v_2^4 + (-a_2 + 2b_5)v_1^2v_2^3 = 0 \end{aligned} \quad (8E)$$

Setting each coefficients in (8E) to zero gives the following equations to solve

$$\begin{aligned}b_5 &= 0 \\-2a_6 &= 0 \\-b_2 &= 0 \\-b_4 &= 0 \\2b_4 &= 0 \\-2b_5 &= 0 \\-a_2 - 2b_3 &= 0 \\-a_2 + 2b_5 &= 0 \\-2a_3 - a_5 &= 0 \\-a_3 - b_1 &= 0 \\-a_5 - b_2 &= 0 \\-a_5 - b_6 &= 0 \\-b_1 - b_4 &= 0 \\b_2 + b_4 &= 0 \\b_2 + 4b_4 &= 0 \\-2a_1 + 2a_4 - 2a_6 &= 0 \\-2a_1 + 2a_4 - 2b_5 &= 0 \\a_2 - 4a_6 + b_5 &= 0 \\a_2 - 2a_6 - 2b_3 &= 0 \\-a_3 + a_5 - b_6 &= 0 \\-2a_3 - 2a_5 + b_2 + 2b_4 &= 0 \\-2a_3 + a_5 + b_1 + b_2 - 3b_6 &= 0 \\-a_5 + b_1 + 2b_2 + b_4 - 3b_6 &= 0\end{aligned}$$

Solving the above equations for the unknowns gives

$$a_1 = a_4$$

$$a_2 = 0$$

$$a_3 = 0$$

$$a_4 = a_4$$

$$a_5 = 0$$

$$a_6 = 0$$

$$b_1 = 0$$

$$b_2 = 0$$

$$b_3 = 0$$

$$b_4 = 0$$

$$b_5 = 0$$

$$b_6 = 0$$

Substituting the above solution in the ansatz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$\xi = u^2 + 1$$

$$\eta = 0$$

Shifting is now applied to make $\xi = 0$ in order to simplify the rest of the computation

$$\begin{aligned} \eta &= \eta - \omega(u, v) \xi \\ &= 0 - \left(-\frac{v^2 + 1}{(u^2 + 1)v} \right) (u^2 + 1) \\ &= \frac{v^2 + 1}{v} \\ \xi &= 0 \end{aligned}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(u, v) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{du}{\xi} = \frac{dv}{\eta} = dS \quad (1)$$

The above comes from the requirements that $(\xi \frac{\partial}{\partial u} + \eta \frac{\partial}{\partial v}) S(u, v) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Since $\xi = 0$ then in this special case

$$R = u$$

S is found from

$$\begin{aligned} S &= \int \frac{1}{\eta} dy \\ &= \int \frac{1}{\frac{v^2+1}{v}} dy \end{aligned}$$

Which results in

$$S = \frac{\ln(v^2 + 1)}{2}$$

Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_u + \omega(u, v)S_v}{R_u + \omega(u, v)R_v} \quad (2)$$

Where in the above R_u, R_v, S_u, S_v are all partial derivatives and $\omega(u, v)$ is the right hand side of the original ode given by

$$\omega(u, v) = -\frac{v^2 + 1}{(u^2 + 1)v}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_u &= 1 \\ R_v &= 0 \\ S_u &= 0 \\ S_v &= \frac{v}{v^2 + 1} \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = -\frac{1}{u^2 + 1} \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for u, v in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = -\frac{1}{R^2 + 1}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S .

Since the ode has the form $\frac{d}{dR}S(R) = f(R)$, then we only need to integrate $f(R)$.

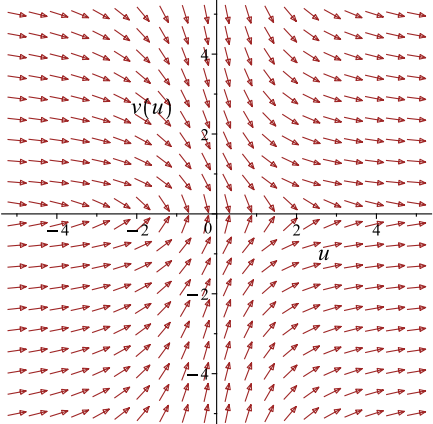
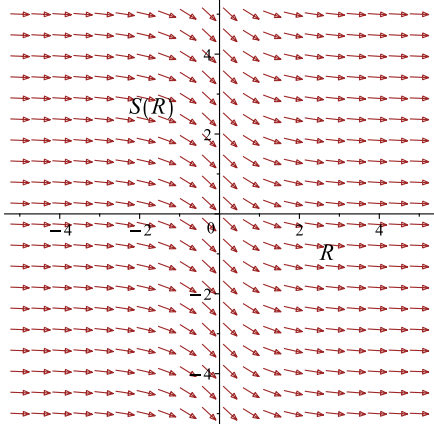
$$\int dS = \int -\frac{1}{R^2 + 1} dR$$

$$S(R) = -\arctan(R) + c_2$$

To complete the solution, we just need to transform the above back to u, v coordinates. This results in

$$\frac{\ln(v^2 + 1)}{2} = -\arctan(u) + c_2$$

The following diagram shows solution curves of the original ode and how they transform in the canonical coordinates space using the mapping shown.

Original ode in u, v coordinates	Canonical coordinates transformation	ODE in canonical coordinates (R, S)
$\frac{dv}{du} = -\frac{v^2+1}{(u^2+1)v}$ 	$R = u$ $S = \frac{\ln(v^2 + 1)}{2}$	$\frac{dS}{dR} = -\frac{1}{R^2+1}$ 

Solving for v gives

$$v = \sqrt{-1 + e^{-2 \arctan(u) + 2c_2}}$$

$$v = -\sqrt{-1 + e^{-2 \arctan(u) + 2c_2}}$$

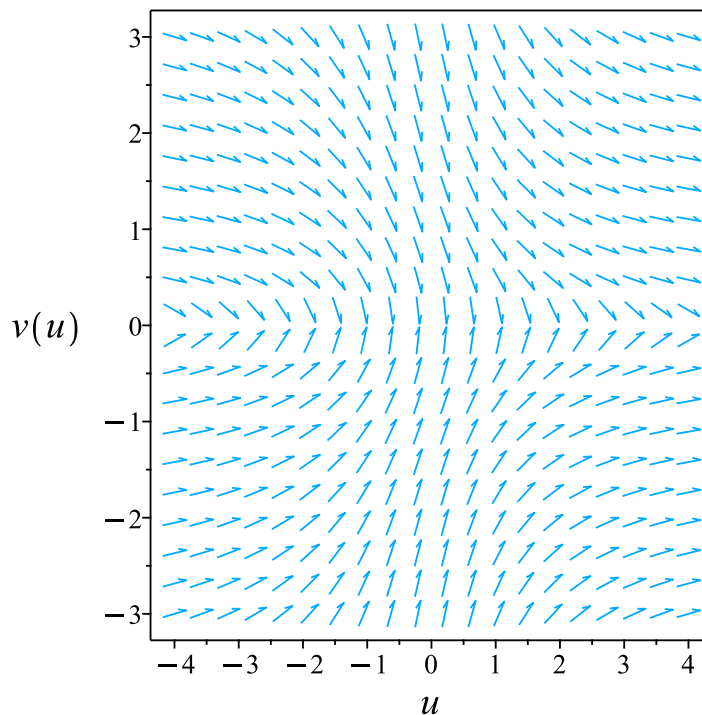


Figure 2.88: Slope field plot
 $1 + v^2 + (u^2 + 1)vv' = 0$

Summary of solutions found

$$v = \sqrt{-1 + e^{-2 \arctan(u) + 2c_2}}$$

$$v = -\sqrt{-1 + e^{-2 \arctan(u) + 2c_2}}$$

Maple step by step solution

Let's solve

$$1 + v^2 + (u^2 + 1)vv' = 0$$

- Highest derivative means the order of the ODE is 1

v'

- Solve for the highest derivative

$$v' = \frac{-v^2 - 1}{(u^2 + 1)v}$$

- Separate variables

$$\frac{v'v}{-v^2 - 1} = \frac{1}{u^2 + 1}$$

- Integrate both sides with respect to u

$$\int \frac{v'v}{-v^2-1} du = \int \frac{1}{u^2+1} du + C1$$

- Evaluate integral

$$-\frac{\ln(v^2+1)}{2} = \arctan(u) + C1$$

- Solve for v

$$\left\{ v = \sqrt{-1 + e^{-2 \arctan(u) - 2C1}}, v = -\sqrt{-1 + e^{-2 \arctan(u) - 2C1}} \right\}$$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
<- Bernoulli successful`
```

Maple dsolve solution

Solving time : 0.009 (sec)

Leaf size : 31

```
dsolve(1+v(u)^2+(u^2+1)*v(u)*diff(v(u),u) = 0,
v(u),singsol=all)
```

$$v = \sqrt{-1 + c_1 e^{-2 \arctan(u)}}$$

$$v = -\sqrt{-1 + c_1 e^{-2 \arctan(u)}}$$

Mathematica DSolve solution

Solving time : 2.406 (sec)

Leaf size : 59

```
DSolve[{(1+v[u]^2)+(1+u^2)*v[u]*D[v[u],u]==0,{}},  
v[u],u,IncludeSingularSolutions->True]
```

$$v(u) \rightarrow -\sqrt{-1 + e^{-2 \arctan(u) + 2c_1}}$$

$$v(u) \rightarrow \sqrt{-1 + e^{-2 \arctan(u) + 2c_1}}$$

$$v(u) \rightarrow -i$$

$$v(u) \rightarrow i$$

2.6.14 problem 10 (c)

Solved as first order separable ode	562
Maple step by step solution	564
Maple trace	565
Maple dsolve solution	565
Mathematica DSolve solution	566

Internal problem ID [18254]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter IV. Methods of solution: First order equations. section 33. Problems at page 91

Problem number : 10 (c)

Date solved : Monday, December 23, 2024 at 09:19:28 PM

CAS classification : [_separable]

Solve

$$u \ln(u) v' + \sin(v)^2 = 1$$

Solved as first order separable ode

Time used: 0.190 (sec)

The ode $v' = -\frac{\sin(v)^2-1}{\ln(u)u}$ is separable as it can be written as

$$\begin{aligned} v' &= -\frac{\sin(v)^2-1}{\ln(u)u} \\ &= f(u)g(v) \end{aligned}$$

Where

$$\begin{aligned} f(u) &= \frac{1}{\ln(u)u} \\ g(v) &= -\sin(v)^2 + 1 \end{aligned}$$

Integrating gives

$$\begin{aligned} \int \frac{1}{g(v)} dv &= \int f(u) du \\ \int \frac{1}{-\sin(v)^2 + 1} dv &= \int \frac{1}{\ln(u)u} du \\ \tan(v) &= \ln(\ln(u)) + c_1 \end{aligned}$$

We now need to find the singular solutions, these are found by finding for what values $g(v)$ is zero, since we had to divide by this above. Solving $g(v) = 0$ or $-\sin(v)^2 + 1 = 0$ for v gives

$$v = -\frac{\pi}{2}$$
$$v = \frac{\pi}{2}$$

Now we go over each such singular solution and check if it verifies the ode itself and any initial conditions given. If it does not then the singular solution will not be used.

Therefore the solutions found are

$$\tan(v) = \ln(\ln(u)) + c_1$$
$$v = -\frac{\pi}{2}$$
$$v = \frac{\pi}{2}$$

Solving for v gives

$$v = -\frac{\pi}{2}$$
$$v = \frac{\pi}{2}$$
$$v = \arctan(\ln(\ln(u)) + c_1)$$

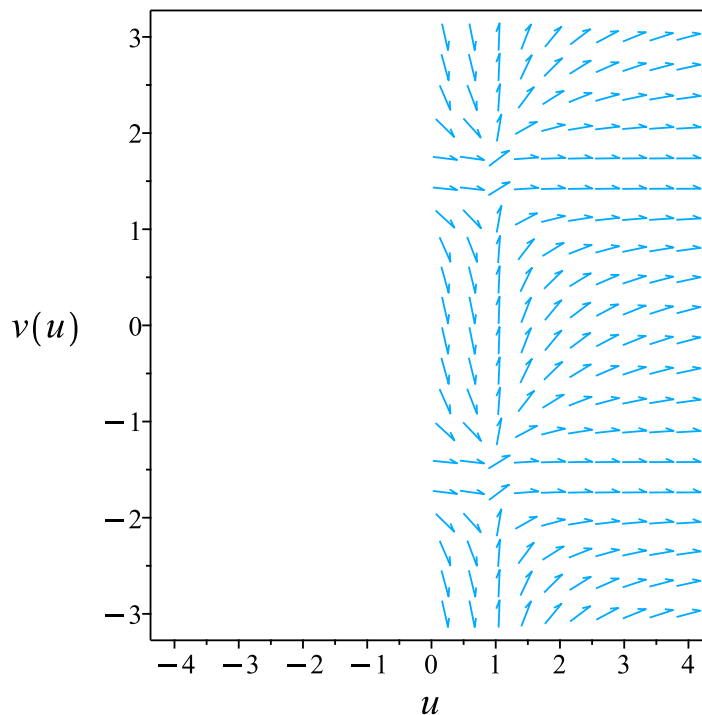


Figure 2.89: Slope field plot
 $u \ln(u) v' + \sin(v)^2 = 1$

Summary of solutions found

$$v = -\frac{\pi}{2}$$

$$v = \frac{\pi}{2}$$

$$v = \arctan(\ln(\ln(u)) + c_1)$$

Maple step by step solution

Let's solve

$$u \ln(u) v' + \sin(v)^2 = 1$$

- Highest derivative means the order of the ODE is 1
 v'
- Solve for the highest derivative

$$v' = \frac{-\sin(v)^2 + 1}{u \ln(u)}$$
- Separate variables

$$\frac{v'}{-\sin(v)^2+1} = \frac{1}{\ln(u)u}$$

- Integrate both sides with respect to u

$$\int \frac{v'}{-\sin(v)^2+1} du = \int \frac{1}{\ln(u)u} du + C1$$

- Evaluate integral

$$\tan(v) = \ln(\ln(u)) + C1$$

- Solve for v

$$v = \arctan(\ln(\ln(u)) + C1)$$

Maple trace

```
`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
trying Bernoulli
trying separable
<- separable successful`
```

Maple dsolve solution

Solving time : 0.006 (sec)

Leaf size : 10

```
dsolve(u*ln(u)*diff(v(u),u)+sin(v(u))^2 = 1,
        v(u),singsol=all)
```

$$v = \arctan(\ln(\ln(u)) + c_1)$$

Mathematica DSolve solution

Solving time : 0.366 (sec)

Leaf size : 52

```
DSolve[{u*Log[u]*D[v[u],u]+Sin[v[u]]==0,{}},  
v[u],u,IncludeSingularSolutions->True]
```

$$v(u) \rightarrow -\arccos(-\tanh(-\log(\log(u)) + c_1))$$

$$v(u) \rightarrow \arccos(-\tanh(-\log(\log(u)) + c_1))$$

$$v(u) \rightarrow 0$$

$$v(u) \rightarrow -\pi$$

$$v(u) \rightarrow \pi$$

2.7 Chapter V. Singular solutions. section 36.

Problems at page 99

2.7.1	problem 1 (eq 98)	568
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2.7.1 problem 1 (eq 98)

Maple step by step solution 585
 Maple dsolve solution 585
 Mathematica DSolve solution 586

Internal problem ID [18255]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter V. Singular solutions. section 36. Problems at page 99

Problem number : 1 (eq 98)

Date solved : Monday, December 23, 2024 at 09:19:33 PM

CAS classification : [[_1st_order, _with_linear_symmetries]]

Solve

$$4yy'^3 - 2y'^2x^2 + 4xyy' + x^3 = 16y^2$$

Solving for the derivative gives these ODE's to solve

$$y' = \frac{(x^6 - 45x^3y^2 + 432y^4 + 3\sqrt{3}\sqrt{-2x^9 + 91y^2x^6 - 1376x^3y^4 + 6912y^6y})^{1/3}}{6y} - \frac{x(-x^3 + 12y^2)}{6y(x^6 - 45x^3y^2 + 432y^4 + 3\sqrt{3}\sqrt{-2x^9 + 91y^2x^6 - 1376x^3y^4 + 6912y^6y})^{1/3}} + \frac{x^2}{6y} \tag{1}$$

$$y' = -\frac{(x^6 - 45x^3y^2 + 432y^4 + 3\sqrt{3}\sqrt{-2x^9 + 91y^2x^6 - 1376x^3y^4 + 6912y^6y})^{1/3}}{12y} + \frac{x(-x^3 + 12y^2)}{12y(x^6 - 45x^3y^2 + 432y^4 + 3\sqrt{3}\sqrt{-2x^9 + 91y^2x^6 - 1376x^3y^4 + 6912y^6y})^{1/3}} + \frac{x^2}{6y} + \frac{i\sqrt{3}\left(\frac{(x^6 - 45x^3y^2 + 432y^4 + 3\sqrt{3}\sqrt{-2x^9 + 91y^2x^6 - 1376x^3y^4 + 6912y^6y})^{1/3}}{6y} + \frac{x(-x^3 + 12y^2)}{6y(x^6 - 45x^3y^2 + 432y^4 + 3\sqrt{3}\sqrt{-2x^9 + 91y^2x^6 - 1376x^3y^4 + 6912y^6y})^{1/3}}\right)}{2} \tag{2}$$

$$y' = -\frac{(x^6 - 45x^3y^2 + 432y^4 + 3\sqrt{3}\sqrt{-2x^9 + 91y^2x^6 - 1376x^3y^4 + 6912y^6y})^{1/3}}{12y} \tag{3}$$

$$+ \frac{x(-x^3 + 12y^2)}{12y(x^6 - 45x^3y^2 + 432y^4 + 3\sqrt{3}\sqrt{-2x^9 + 91y^2x^6 - 1376x^3y^4 + 6912y^6y})^{1/3}} + \frac{x^2}{6y}$$

$$i\sqrt{3} \left(\frac{(x^6 - 45x^3y^2 + 432y^4 + 3\sqrt{3}\sqrt{-2x^9 + 91y^2x^6 - 1376x^3y^4 + 6912y^6y})^{1/3}}{6y} + \frac{x(-x^3 + 12y^2)}{6y(x^6 - 45x^3y^2 + 432y^4 + 3\sqrt{3}\sqrt{-2x^9 + 91y^2x^6 - 1376x^3y^4 + 6912y^6y})^{1/3}} \right)$$

2

Now each of the above is solved separately.

Solving Eq. (1)

Solving for y' gives

$$y' = \frac{-x^4 + 12xy^2 - x^2(x^6 - 45x^3y^2 + 432y^4 + 3\sqrt{3}\sqrt{-2x^9 + 91y^2x^6 - 1376x^3y^4 + 6912y^6y})^{1/3} - (x^6 - 45x^3y^2 + 432y^4 + 3\sqrt{3}\sqrt{-2x^9 + 91y^2x^6 - 1376x^3y^4 + 6912y^6y})^{1/3}}{6y(x^6 - 45x^3y^2 + 432y^4 + 3\sqrt{3}\sqrt{-2x^9 + 91y^2x^6 - 1376x^3y^4 + 6912y^6y})^{1/3}} \tag{1}$$

Each of the above ode's is now solved An ode $y' = f(x, y)$ is isobaric if

$$f(tx, t^m y) = t^{m-1} f(x, y) \tag{1}$$

Where here

$$f(x, y) = -\frac{-x^4 + 12xy^2 - x^2(x^6 - 45x^3y^2 + 432y^4 + 3\sqrt{3}\sqrt{-2x^9 + 91y^2x^6 - 1376x^3y^4 + 6912y^6y})^{1/3} - (x^6 - 45x^3y^2 + 432y^4 + 3\sqrt{3}\sqrt{-2x^9 + 91y^2x^6 - 1376x^3y^4 + 6912y^6y})^{1/3}}{6y(x^6 - 45x^3y^2 + 432y^4 + 3\sqrt{3}\sqrt{-2x^9 + 91y^2x^6 - 1376x^3y^4 + 6912y^6y})^{1/3}} \tag{2}$$

m is the order of isobaric. Substituting (2) into (1) and solving for m gives

$$m = \frac{3}{2}$$

Since the ode is isobaric of order $m = \frac{3}{2}$, then the substitution

$$y = ux^m$$

$$= ux^{3/2}$$

Converts the ODE to a separable in $u(x)$. Performing this substitution gives

$$\frac{3\sqrt{x}u(x)}{2} + x^{3/2}u'(x) = -\frac{-x^4 + 12x^4u(x)^2 - x^2(x^6 - 45x^6u(x)^2 + 432x^6u(x)^4 + 3\sqrt{3}\sqrt{-2x^9 + 91x^9u(x)^2 - 1376x^6u(x)^3 + 6912x^6u(x)^6})^{1/3} - (x^6 - 45x^6u(x)^2 + 432x^6u(x)^4 + 3\sqrt{3}\sqrt{-2x^9 + 91x^9u(x)^2 - 1376x^6u(x)^3 + 6912x^6u(x)^6})^{1/3}}{6x^{3/2}u(x)(x^6 - 45x^6u(x)^2 + 432x^6u(x)^4 + 3\sqrt{3}\sqrt{-2x^9 + 91x^9u(x)^2 - 1376x^6u(x)^3 + 6912x^6u(x)^6})^{1/3}}$$

The ode $u'(x) = \frac{\left(-9 \cdot 3^{1/3} \left(\sqrt{3} \left(432\sqrt{3} u(x)^4 - 45\sqrt{3} u(x)^2 + 9u(x) \sqrt{(27u(x)^2 - 2)(4u(x) - 1)^2(4u(x) + 1)^2 + \sqrt{3}}\right)\right)^{1/3} u(x)^2 - 12u(x)^2 \cdot 3^{2/3}\right)}{18 \left(\sqrt{3} \left(432\sqrt{3} u(x)^4 - 45\sqrt{3} u(x)^2 + 9u(x) \sqrt{(27u(x)^2 - 2)(4u(x) - 1)^2(4u(x) + 1)^2 + \sqrt{3}}\right)\right)^{1/3}}$

is separable as it can be written as

$$u'(x) = \frac{\left(-9 \cdot 3^{1/3} \left(\sqrt{3} \left(432\sqrt{3} u(x)^4 - 45\sqrt{3} u(x)^2 + 9u(x) \sqrt{(27u(x)^2 - 2)(4u(x) - 1)^2(4u(x) + 1)^2 + \sqrt{3}}\right)\right)^{1/3} u(x)^2 - 12u(x)^2 \cdot 3^{2/3}\right)}{18 \left(\sqrt{3} \left(432\sqrt{3} u(x)^4 - 45\sqrt{3} u(x)^2 + 9u(x) \sqrt{(27u(x)^2 - 2)(4u(x) - 1)^2(4u(x) + 1)^2 + \sqrt{3}}\right)\right)^{1/3}}$$

$$= f(x)g(u)$$

Where

$$f(x) = \frac{3^{2/3}}{18x}$$

$$g(u) = \frac{-9 \cdot 3^{1/3} \left(\sqrt{3} \left(432\sqrt{3} u^4 - 45\sqrt{3} u^2 + 9u \sqrt{(27u^2 - 2)(4u - 1)^2(4u + 1)^2 + \sqrt{3}}\right)\right)^{1/3} u^2 - 12u^2 \cdot 3^{2/3}}{18 \left(\sqrt{3} \left(432\sqrt{3} u^4 - 45\sqrt{3} u^2 + 9u \sqrt{(27u^2 - 2)(4u - 1)^2(4u + 1)^2 + \sqrt{3}}\right)\right)^{1/3}}$$

Integrating gives

$$\int \frac{dx}{-9 \cdot 3^{1/3} \left(\sqrt{3} \left(432\sqrt{3} u^4 - 45\sqrt{3} u^2 + 9u \sqrt{(27u^2 - 2)(4u - 1)^2(4u + 1)^2 + \sqrt{3}}\right)\right)^{1/3} u^2 - 12u^2 \cdot 3^{2/3}}$$

$$\int^{u(x)} \frac{d\tau}{-9 \cdot 3^{1/3} \left(\sqrt{3} \left(432\sqrt{3} \tau^4 - 45\sqrt{3} \tau^2 + 9\tau \sqrt{(27\tau^2 - 2)(4\tau - 1)^2(4\tau + 1)^2 + \sqrt{3}}\right)\right)^{1/3} \tau^2 - 12\tau^2 \cdot 3^{2/3}}$$

We now need to find the singular solutions, these are found by finding for what values

$g(u)$ is zero, since we had to divide by this above. Solving $g(u) = 0$ or $\frac{-9 \cdot 3^{1/3} \left(\sqrt{3} \left(432\sqrt{3} u^4 - 45\sqrt{3} u^2 + 9u \sqrt{(27u^2 - 2)(4u - 1)^2(4u + 1)^2 + \sqrt{3}}\right)\right)^{1/3} u^2 - 12u^2 \cdot 3^{2/3}}{18 \left(\sqrt{3} \left(432\sqrt{3} u^4 - 45\sqrt{3} u^2 + 9u \sqrt{(27u^2 - 2)(4u - 1)^2(4u + 1)^2 + \sqrt{3}}\right)\right)^{1/3}} = 0$

0 for $u(x)$ gives

$$u(x) = 1$$

Now we go over each such singular solution and check if it verifies the ode itself and any initial conditions given. If it does not then the singular solution will not be used.

Therefore the solutions found are

$$\int^{u(x)} \frac{1}{-9 \cdot 3^{1/3} \left(\sqrt{3} \left(432\sqrt{3} \tau^4 - 45\sqrt{3} \tau^2 + 9\tau \sqrt{(27\tau^2 - 2)(4\tau - 1)^2(4\tau + 1)^2 + \sqrt{3}} \right) \right)^{1/3} \tau^2 - 12\tau^2 3^{2/3}} d\tau$$

Converting $\int^{u(x)} \frac{1}{-9 \cdot 3^{1/3} \left(\sqrt{3} \left(432\sqrt{3} \tau^4 - 45\sqrt{3} \tau^2 + 9\tau \sqrt{(27\tau^2 - 2)(4\tau - 1)^2(4\tau + 1)^2 + \sqrt{3}} \right) \right)^{1/3} \tau^2 - 12\tau^2 3^{2/3}} d\tau$
 $\frac{3^{2/3} \ln(x)}{18} + c_1$ back to y gives

$$\int^{\frac{y}{x^{3/2}}} \frac{1}{-9 \cdot 3^{1/3} \left(\sqrt{3} \left(432\sqrt{3} \tau^4 - 45\sqrt{3} \tau^2 + 9\tau \sqrt{(27\tau^2 - 2)(4\tau - 1)^2(4\tau + 1)^2 + \sqrt{3}} \right) \right)^{1/3} \tau^2 - 12\tau^2 3^{2/3}} d\tau$$

Converting $u(x) = 1$ back to y gives

$$\frac{y}{x^{3/2}} = 1$$

Solving for y gives

$$\int^{\frac{y}{x^{3/2}}} \frac{1}{-9 \cdot 3^{1/3} \left(\sqrt{3} \left(432\sqrt{3} \tau^4 - 45\sqrt{3} \tau^2 + 9\tau \sqrt{(27\tau^2 - 2)(4\tau - 1)^2(4\tau + 1)^2 + \sqrt{3}} \right) \right)^{1/3} \tau^2 - 12\tau^2 3^{2/3}} d\tau$$

+ c_1
 $y = x^{3/2}$

We now need to find the singular solutions, these are found by finding for what values

$$\left(\frac{(x^6 - 45x^3y^2 + 432y^4 + 3\sqrt{3} \sqrt{-2x^9 + 91y^2x^6 - 1376x^3y^4 + 6912y^6} y)^{1/3}}{6y} - \frac{x(-x^3 + 12y^2)}{6y(x^6 - 45x^3y^2 + 432y^4 + 3\sqrt{3} \sqrt{-2x^9 + 91y^2x^6 - 1376x^3y^4 + 6912y^6} y)} \right)^{1/3}$$

$\frac{x^2}{6y}$) is zero. These give

$$y = \text{RootOf} \left(-x^4 - x^2 \left(x^6 - 45x^3 Z^2 + 432 Z^4 + 3\sqrt{3} \sqrt{-2x^9 + 91 Z^2 x^6 - 1376 Z^4 x^3 + 6912 Z^6 Z} \right)^{1/3} + 12x Z^2 - \left(x^6 - 45x^3 Z^2 + 432 Z^4 + 3\sqrt{3} \sqrt{-2x^9 + 91 Z^2 x^6 - 1376 Z^4 x^3 + 6912 Z^6 Z} \right)^{2/3} \right)$$

Now we go over each such singular solution and check if it verifies the ode itself and any initial conditions given. If it does not then the singular solution will not be used.

The solution $\text{RootOf} \left(-x^4 - x^2 \left(x^6 - 45x^3 Z^2 + 432 Z^4 + 3\sqrt{3} \sqrt{-2x^9 + 91 Z^2 x^6 - 1376 Z^4 x^3 + 6912 Z^6 Z} \right)^{1/3} + 12x Z^2 - \left(x^6 - 45x^3 Z^2 + 432 Z^4 + 3\sqrt{3} \sqrt{-2x^9 + 91 Z^2 x^6 - 1376 Z^4 x^3 + 6912 Z^6 Z} \right)^{2/3} \right)$ will not be used

Solving Eq. (2)

Writing the ode as

$$y' = \frac{-i\sqrt{3}x^4 + 12i\sqrt{3}y^2x + i\sqrt{3}(x^6 - 45x^3y^2 + 432y^4 + 3\sqrt{3}\sqrt{-2x^9 + 91y^2x^6 - 1376x^3y^4 + 6912y^6y})^{1/3}}{12}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2\xi_y - \omega_x\xi - \omega_y\eta = 0 \tag{A}$$

To determine ξ, η then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

$$\xi = xa_2 + ya_3 + a_1 \tag{1E}$$

$$\eta = xb_2 + yb_3 + b_1 \tag{2E}$$

Where the unknown coefficients are

$$\{a_1, a_2, a_3, b_1, b_2, b_3\}$$

Substituting equations (1E,2E) and ω into (A) gives

$$\text{Expression too large to display} \tag{5E}$$

Putting the above in normal form gives

Expression too large to display

Setting the numerator to zero gives

Expression too large to display (6E)

Simplifying the above gives

Expression too large to display (6E)

Since the PDE has radicals, simplifying gives

Expression too large to display

Looking at the above PDE shows the following are all the terms with $\{x, y\}$ in them.

$$\left\{ x, y, \sqrt{-(2x^3 - 27y^2)(x^3 - 16y^2)^2}, \left(x^6 - 45x^3y^2 + 432y^4 \right. \right. \\ \left. \left. + 3\sqrt{3} \sqrt{-(2x^3 - 27y^2)(x^3 - 16y^2)^2} y \right)^{1/3}, \left(x^6 - 45x^3y^2 + 432y^4 + 3\sqrt{3} \sqrt{-(2x^3 - 27y^2)(x^3 - 16y^2)^2} y \right)^{1/3} \right\}$$

The following substitution is now made to be able to collect on all terms with $\{x, y\}$ in them

$$\left\{ x = v_1, y = v_2, \sqrt{-(2x^3 - 27y^2)(x^3 - 16y^2)^2} = v_3, \left(x^6 - 45x^3y^2 + 432y^4 \right. \right. \\ \left. \left. + 3\sqrt{3} \sqrt{-(2x^3 - 27y^2)(x^3 - 16y^2)^2} y \right)^{1/3} = v_4, \left(x^6 - 45x^3y^2 + 432y^4 + 3\sqrt{3} \sqrt{-(2x^3 - 27y^2)(x^3 - 16y^2)^2} y \right)^{1/3} = v_5 \right\}$$

The above PDE (6E) now becomes

Expression too large to display (7E)

Collecting the above on the terms v_i introduced, and these are

$$\{v_1, v_2, v_3, v_4, v_5\}$$

Equation (7E) now becomes

Expression too large to display (8E)

Setting each coefficients in (8E) to zero gives the following equations to solve

$$\begin{aligned} -20736a_1 &= 0 \\ -48a_1 &= 0 \\ 2160a_1 &= 0 \\ -7344a_3 &= 0 \\ -12a_3 &= 0 \\ 588a_3 &= 0 \\ 20736a_3 &= 0 \\ -1080b_1 &= 0 \\ 24b_1 &= 0 \\ 10368b_1 &= 0 \\ -936b_2 &= 0 \\ 24b_2 &= 0 \\ 3888b_2 &= 0 \\ 62208b_2 &= 0 \\ -995328\sqrt{3} a_1 &= 0 \\ -13104\sqrt{3} a_1 &= 0 \\ 288\sqrt{3} a_1 &= 0 \\ 198144\sqrt{3} a_1 &= 0 \\ -446976\sqrt{3} a_3 &= 0 \\ -3564\sqrt{3} a_3 &= 0 \\ 72\sqrt{3} a_3 &= 0 \\ 62640\sqrt{3} a_3 &= 0 \\ 995328\sqrt{3} a_3 &= 0 \\ -99072\sqrt{3} b_1 &= 0 \\ -144\sqrt{3} b_1 &= 0 \\ 6552\sqrt{3} b_1 &= 0 \\ 497664\sqrt{3} b_1 &= 0 \\ -96768\sqrt{3} b_2 &= 0 \\ -59760\sqrt{3} b_2 &= 0 \\ -144\sqrt{3} b_2 &= 0 \\ 5688\sqrt{3} b_2 &= 0 \\ 2985984\sqrt{3} b_2 &= 0 \\ -31104a_2 + 20736b_3 &= 0 \\ -72a_2 + 48b_3 &= 0 \\ 3240a_2 - 2160b_3 &= 0 \end{aligned}$$

Solving the above equations for the unknowns gives

$$\begin{aligned} a_1 &= 0 \\ a_2 &= \frac{2b_3}{3} \\ a_3 &= 0 \\ b_1 &= 0 \\ b_2 &= 0 \\ b_3 &= b_3 \end{aligned}$$

Substituting the above solution in the ansatz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$\begin{aligned} \xi &= \frac{2x}{3} \\ \eta &= y \end{aligned}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

The characteristic pde which is used to find the canonical coordinates is

$$\frac{dx}{\xi} = \frac{dy}{\eta} = dS \quad (1)$$

The above comes from the requirements that $\left(\xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y}\right) S(x, y) = 1$. Starting with the first pair of ode's in (1) gives an ode to solve for the independent variable R in the canonical coordinates, where $S(R)$. Therefore

$$\begin{aligned} \frac{dy}{dx} &= \frac{\eta}{\xi} \\ &= \frac{y}{\frac{2x}{3}} \\ &= \frac{3y}{2x} \end{aligned}$$

This is easily solved to give

$$y = c_1 x^{3/2}$$

Where now the coordinate R is taken as the constant of integration. Hence

$$R = \frac{y}{x^{3/2}}$$

And S is found from

$$\begin{aligned} dS &= \frac{dx}{\xi} \\ &= \frac{dx}{\frac{2x}{3}} \end{aligned}$$

Integrating gives

$$\begin{aligned} S &= \int \frac{dx}{T} \\ &= \frac{3 \ln(x)}{2} \end{aligned}$$

Where the constant of integration is set to zero as we just need one solution. Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = \frac{-i\sqrt{3}x^4 + 12i\sqrt{3}y^2x + i\sqrt{3}(x^6 - 45x^3y^2 + 432y^4 + 3\sqrt{3}\sqrt{-2x^9 + 91y^2x^6 - 1376x^3y^4 + 6912y^6})}{2x^5}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= -\frac{3y}{2x^{5/2}} \\ R_y &= \frac{1}{x^{3/2}} \\ S_x &= \frac{3}{2x} \\ S_y &= 0 \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = -\frac{18x^{3/2} \left(x^6 - 45x^3y^2 + 432y^4 + 3\sqrt{3}\sqrt{-2x^9 + 91y^2x^6 - 1376x^3y^4 + 6912y^6} \right)}{(-i\sqrt{3}x + x) \left(x^6 - 45x^3y^2 + 432y^4 + 3\sqrt{3}\sqrt{-2 \left(x^3 - \frac{27y^2}{2} \right) (x^3 - 16y^2)^2 y} \right)^{2/3} + (-2x^3 + 18y^2)} \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = \frac{18R((48R^3 - 3R)\sqrt{3}\sqrt{27R^2 - 2} + 432R^4 - 45R^2 + 1)}{(-1 + i\sqrt{3})((48R^3 - 3R)\sqrt{3}\sqrt{27R^2 - 2} + 432R^4 - 45R^2 + 1)^{2/3} + (-18R^2 + 2)((48R^3 - 3R)\sqrt{3}\sqrt{27R^2 - 2} + 432R^4 - 45R^2 + 1)^{1/3}}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S .

Since the ode has the form $\frac{d}{dR}S(R) = f(R)$, then we only need to integrate $f(R)$.

$$\int dS = \int \frac{18R((48R^3 - 3R)\sqrt{3}\sqrt{27R^2 - 2} + 432R^4 - 45R^2 + 1)^{2/3}\sqrt{3} + 12i\sqrt{3}R^2 - 18((48R^3 - 3R)\sqrt{3}\sqrt{27R^2 - 2} + 432R^4 - 45R^2 + 1)^{1/3}}{i((48R^3 - 3R)\sqrt{3}\sqrt{27R^2 - 2} + 432R^4 - 45R^2 + 1)^{2/3}\sqrt{3} + 12i\sqrt{3}R^2 - 18((48R^3 - 3R)\sqrt{3}\sqrt{27R^2 - 2} + 432R^4 - 45R^2 + 1)^{1/3}}$$

$$S(R) = \int \frac{18R((48R^3 - 3R)\sqrt{3}\sqrt{27R^2 - 2} + 432R^4 - 45R^2 + 1)^{2/3}\sqrt{3} + 12i\sqrt{3}R^2 - 18((48R^3 - 3R)\sqrt{3}\sqrt{27R^2 - 2} + 432R^4 - 45R^2 + 1)^{1/3}}{i((48R^3 - 3R)\sqrt{3}\sqrt{27R^2 - 2} + 432R^4 - 45R^2 + 1)^{2/3}\sqrt{3} + 12i\sqrt{3}R^2 - 18((48R^3 - 3R)\sqrt{3}\sqrt{27R^2 - 2} + 432R^4 - 45R^2 + 1)^{1/3}}$$

$$S(R) = \int \frac{18R((48R^3 - 3R)\sqrt{3}\sqrt{27R^2 - 2} + 432R^4 - 45R^2 + 1)^{2/3}\sqrt{3} + 12i\sqrt{3}R^2 - 18((48R^3 - 3R)\sqrt{3}\sqrt{27R^2 - 2} + 432R^4 - 45R^2 + 1)^{1/3}}{i((48R^3 - 3R)\sqrt{3}\sqrt{27R^2 - 2} + 432R^4 - 45R^2 + 1)^{2/3}\sqrt{3} + 12i\sqrt{3}R^2 - 18((48R^3 - 3R)\sqrt{3}\sqrt{27R^2 - 2} + 432R^4 - 45R^2 + 1)^{1/3}}$$

To complete the solution, we just need to transform the above back to x, y coordinates.

This results in

$$\frac{3 \ln(x)}{2} = \int \frac{\frac{y}{x^{3/2}}}{i((48a^3 - 3a)\sqrt{3}\sqrt{27a^2 - 2} + 432a^4 - 45a^2 + 1)^{2/3}\sqrt{3} + 12i\sqrt{3}a^2 - 18((48a^3 - 3a)\sqrt{3}\sqrt{27a^2 - 2} + 432a^4 - 45a^2 + 1)^{1/3}}$$

We now need to find the singular solutions, these are found by finding for what values

$$\left(-\frac{(x^6 - 45x^3y^2 + 432y^4 + 3\sqrt{3}\sqrt{-2x^9 + 91y^2x^6 - 1376x^3y^4 + 6912y^6}y)^{1/3}}{12y} + \frac{x(-x^3 + 12y^2)}{12y(x^6 - 45x^3y^2 + 432y^4 + 3\sqrt{3}\sqrt{-2x^9 + 91y^2x^6 - 1376x^3y^4 + 6912y^6}y)^{1/3}} \right) + \frac{i\sqrt{3}\left(\frac{(x^6 - 45x^3y^2 + 432y^4 + 3\sqrt{3}\sqrt{-2x^9 + 91y^2x^6 - 1376x^3y^4 + 6912y^6}y)^{1/3}}{6y} + \frac{x(-x^3 + 12y^2)}{6y(x^6 - 45x^3y^2 + 432y^4 + 3\sqrt{3}\sqrt{-2x^9 + 91y^2x^6 - 1376x^3y^4 + 6912y^6}y)^{1/3}}\right)}{2} + \frac{x^2}{6y}$$

is zero. These give

$$y = \text{RootOf}\left(-i\sqrt{3}x^4 + 12i\sqrt{3}Z^2x + i\sqrt{3}\left(x^6 - 45x^3Z^2 + 432Z^4 + 3\sqrt{3}\sqrt{-2x^9 + 91Z^2x^6 - 1376Z^4x^3 + 6912Z^6Z}\right)^{2/3} - x^4 + 12xZ^2 + 2x^2\left(x^6 - 45x^3Z^2 + 432Z^4 + 3\sqrt{3}\sqrt{-2x^9 + 91Z^2x^6 - 1376Z^4x^3 + 6912Z^6Z}\right)^{1/3} - \left(x^6 - 45x^3Z^2 + 432Z^4 + 3\sqrt{3}\sqrt{-2x^9 + 91Z^2x^6 - 1376Z^4x^3 + 6912Z^6Z}\right)^{2/3}\right)$$

Now we go over each such singular solution and check if it verifies the ode itself and any initial conditions given. If it does not then the singular solution will not be used.

The solution $\text{RootOf} \left(-i\sqrt{3}x^4 + 12i\sqrt{3}Z^2x + i\sqrt{3} \left(x^6 - 45x^3Z^2 + 432Z^4 + 3\sqrt{3} \sqrt{-2x^9 + 91Z^2x} \right) \right)$ will not be used

Solving Eq. (3)

Writing the ode as

$$y' = -\frac{-i\sqrt{3}x^4 + 12i\sqrt{3}y^2x + i\sqrt{3}(x^6 - 45x^3y^2 + 432y^4 + 3\sqrt{3}\sqrt{-2x^9 + 91y^2x^6 - 1376x^3y^4 + 6912y^6})}{1}$$

$$y' = \omega(x, y)$$

The condition of Lie symmetry is the linearized PDE given by

$$\eta_x + \omega(\eta_y - \xi_x) - \omega^2\xi_y - \omega_x\xi - \omega_y\eta = 0 \tag{A}$$

To determine ξ, η then (A) is solved using ansatz. Making bivariate polynomials of degree 1 to use as anstaz gives

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Where the unknown coefficients are

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Substituting equations (1E,2E) and ω into (A) gives

$$\text{Expression too large to display} \tag{5E}$$

Putting the above in normal form gives

$$\text{Expression too large to display}$$

Setting the numerator to zero gives

$$\text{Expression too large to display} \tag{6E}$$

Simplifying the above gives

$$\text{Expression too large to display} \tag{6E}$$

Since the PDE has radicals, simplifying gives

$$\text{Expression too large to display}$$

Looking at the above PDE shows the following are all the terms with $\{x, y\}$ in them.

$$\left\{ x, y, \sqrt{-(2x^3 - 27y^2)(x^3 - 16y^2)^2}, \left(x^6 - 45x^3y^2 + 432y^4 \right. \right. \\ \left. \left. + 3\sqrt{3} \sqrt{-(2x^3 - 27y^2)(x^3 - 16y^2)^2} y \right)^{1/3}, \left(x^6 - 45x^3y^2 + 432y^4 + 3\sqrt{3} \sqrt{-(2x^3 - 27y^2)(x^3 - 16y^2)^2} y \right) \right\}$$

The following substitution is now made to be able to collect on all terms with $\{x, y\}$ in them

$$\left\{ x = v_1, y = v_2, \sqrt{-(2x^3 - 27y^2)(x^3 - 16y^2)^2} = v_3, \left(x^6 - 45x^3y^2 + 432y^4 \right. \right. \\ \left. \left. + 3\sqrt{3} \sqrt{-(2x^3 - 27y^2)(x^3 - 16y^2)^2} y \right)^{1/3} = v_4, \left(x^6 - 45x^3y^2 + 432y^4 + 3\sqrt{3} \sqrt{-(2x^3 - 27y^2)(x^3 - 16y^2)^2} y \right) \right\}$$

The above PDE (6E) now becomes

$$\text{Expression too large to display} \tag{7E}$$

Collecting the above on the terms v_i introduced, and these are

$$\{v_1, v_2, v_3, v_4, v_5\}$$

Equation (7E) now becomes

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Solving the above equations for the unknowns gives

$$\begin{aligned} a_1 &= 0 \\ a_2 &= \frac{2b_3}{3} \\ a_3 &= 0 \\ b_1 &= 0 \\ b_2 &= 0 \\ b_3 &= b_3 \end{aligned}$$

Substituting the above solution in the ansatz (1E,2E) (using 1 as arbitrary value for any unknown in the RHS) gives

$$\begin{aligned} \xi &= \frac{2x}{3} \\ \eta &= y \end{aligned}$$

The next step is to determine the canonical coordinates R, S . The canonical coordinates map $(x, y) \rightarrow (R, S)$ where (R, S) are the canonical coordinates which make the original ode become a quadrature and hence solved by integration.

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$$\begin{aligned} \frac{dy}{dx} &= \frac{\eta}{\xi} \\ &= \frac{y}{\frac{2x}{3}} \\ &= \frac{3y}{2x} \end{aligned}$$

This is easily solved to give

$$y = c_1 x^{3/2}$$

Where now the coordinate R is taken as the constant of integration. Hence

$$R = \frac{y}{x^{3/2}}$$

And S is found from

$$\begin{aligned} dS &= \frac{dx}{\xi} \\ &= \frac{dx}{\frac{2x}{3}} \end{aligned}$$

Integrating gives

$$\begin{aligned} S &= \int \frac{dx}{T} \\ &= \frac{3 \ln(x)}{2} \end{aligned}$$

Where the constant of integration is set to zero as we just need one solution. Now that R, S are found, we need to setup the ode in these coordinates. This is done by evaluating

$$\frac{dS}{dR} = \frac{S_x + \omega(x, y)S_y}{R_x + \omega(x, y)R_y} \quad (2)$$

Where in the above R_x, R_y, S_x, S_y are all partial derivatives and $\omega(x, y)$ is the right hand side of the original ode given by

$$\omega(x, y) = -\frac{-i\sqrt{3}x^4 + 12i\sqrt{3}y^2x + i\sqrt{3}(x^6 - 45x^3y^2 + 432y^4 + 3\sqrt{3}\sqrt{-2x^9 + 91y^2x^6 - 1376x^3y^4 + 6912y^6})}{2x^5}$$

Evaluating all the partial derivatives gives

$$\begin{aligned} R_x &= -\frac{3y}{2x^{5/2}} \\ R_y &= \frac{1}{x^{3/2}} \\ S_x &= \frac{3}{2x} \\ S_y &= 0 \end{aligned}$$

Substituting all the above in (2) and simplifying gives the ode in canonical coordinates.

$$\frac{dS}{dR} = \frac{18x^{3/2} \left(x^6 - 45x^3y^2 + 432y^4 + 3\sqrt{3}\sqrt{-2x^9 + 91y^2x^6 - 1376x^3y^4 + 6912y^6} \right)}{x(-i\sqrt{3} - 1) \left(x^6 - 45x^3y^2 + 432y^4 + 3\sqrt{3}\sqrt{-2 \left(x^3 - \frac{27y^2}{2} \right) (x^3 - 16y^2)^2 y} \right)^{2/3} + 2(x^3 - 9y^2) \left(x^3 - \frac{27y^2}{2} \right)} \quad (2A)$$

We now need to express the RHS as function of R only. This is done by solving for x, y in terms of R, S from the result obtained earlier and simplifying. This gives

$$\frac{dS}{dR} = -\frac{18R((48R^3 - 3R)\sqrt{3}\sqrt{27R^2 - 2} + 432R^4 - 45R^2 + 1)^{2/3} + (18R^2 - 2)((48R^3 - 3R)\sqrt{3}\sqrt{27R^2 - 2} + 432R^4 - 45R^2 + 1)^{2/3}}{(1 + i\sqrt{3})((48R^3 - 3R)\sqrt{3}\sqrt{27R^2 - 2} + 432R^4 - 45R^2 + 1)^{2/3} + (18R^2 - 2)((48R^3 - 3R)\sqrt{3}\sqrt{27R^2 - 2} + 432R^4 - 45R^2 + 1)^{2/3}}$$

The above is a quadrature ode. This is the whole point of Lie symmetry method. It converts an ode, no matter how complicated it is, to one that can be solved by integration when the ode is in the canonical coordinates R, S .

Since the ode has the form $\frac{d}{dR}S(R) = f(R)$, then we only need to integrate $f(R)$.

$$\int dS = \int -\frac{18R((48R^3 - 3R)\sqrt{3}\sqrt{27R^2 - 2} + 432R^4 - 45R^2 + 1)^{2/3} + 12i\sqrt{3}R^2 + 18((48R^3 - 3R)\sqrt{3}\sqrt{27R^2 - 2} + 432R^4 - 45R^2 + 1)^{2/3}}{i\sqrt{3}((48R^3 - 3R)\sqrt{3}\sqrt{27R^2 - 2} + 432R^4 - 45R^2 + 1)^{2/3} + 12i\sqrt{3}R^2 + 18((48R^3 - 3R)\sqrt{3}\sqrt{27R^2 - 2} + 432R^4 - 45R^2 + 1)^{2/3}}$$

$$S(R) = \int -\frac{18R((48R^3 - 3R)\sqrt{3}\sqrt{27R^2 - 2} + 432R^4 - 45R^2 + 1)^{2/3} + 12i\sqrt{3}R^2 + 18((48R^3 - 3R)\sqrt{3}\sqrt{27R^2 - 2} + 432R^4 - 45R^2 + 1)^{2/3}}{i\sqrt{3}((48R^3 - 3R)\sqrt{3}\sqrt{27R^2 - 2} + 432R^4 - 45R^2 + 1)^{2/3} + 12i\sqrt{3}R^2 + 18((48R^3 - 3R)\sqrt{3}\sqrt{27R^2 - 2} + 432R^4 - 45R^2 + 1)^{2/3}}$$

$$S(R) = \int -\frac{18R((48R^3 - 3R)\sqrt{3}\sqrt{27R^2 - 2} + 432R^4 - 45R^2 + 1)^{2/3} + 12i\sqrt{3}R^2 + 18((48R^3 - 3R)\sqrt{3}\sqrt{27R^2 - 2} + 432R^4 - 45R^2 + 1)^{2/3}}{i\sqrt{3}((48R^3 - 3R)\sqrt{3}\sqrt{27R^2 - 2} + 432R^4 - 45R^2 + 1)^{2/3} + 12i\sqrt{3}R^2 + 18((48R^3 - 3R)\sqrt{3}\sqrt{27R^2 - 2} + 432R^4 - 45R^2 + 1)^{2/3}}$$

To complete the solution, we just need to transform the above back to x, y coordinates.

This results in

$$\frac{3 \ln(x)}{2} = \int \frac{-\frac{y}{x^{3/2}}}{i((48a^3 - 3a)\sqrt{3}\sqrt{27a^2 - 2} + 432a^4 - 45a^2 + 1)^{2/3}\sqrt{3} + 12i\sqrt{3}a^2 + 18((48a^3 - 3a)\sqrt{3}\sqrt{27a^2 - 2} + 432a^4 - 45a^2 + 1)^{2/3}}$$

We now need to find the singular solutions, these are found by finding for what values

$$\left(-\frac{(x^6 - 45x^3y^2 + 432y^4 + 3\sqrt{3}\sqrt{-2x^9 + 91y^2x^6 - 1376x^3y^4 + 6912y^6}y)^{1/3}}{12y} + \frac{x(-x^3 + 12y^2)}{12y(x^6 - 45x^3y^2 + 432y^4 + 3\sqrt{3}\sqrt{-2x^9 + 91y^2x^6 - 1376x^3y^4 + 6912y^6})^{1/3}}\right) + \frac{i\sqrt{3}\left(\frac{(x^6 - 45x^3y^2 + 432y^4 + 3\sqrt{3}\sqrt{-2x^9 + 91y^2x^6 - 1376x^3y^4 + 6912y^6}y)^{1/3}}{6y} + \frac{x(-x^3 + 12y^2)}{6y(x^6 - 45x^3y^2 + 432y^4 + 3\sqrt{3}\sqrt{-2x^9 + 91y^2x^6 - 1376x^3y^4 + 6912y^6})^{1/3}}\right)}{\frac{x^2}{6y}}$$

is zero. These give

$$y = \text{RootOf}\left(-i\sqrt{3}x^4 + 12i\sqrt{3}Z^2x + i\sqrt{3}\left(x^6 - 45x^3Z^2 + 432Z^4 + 3\sqrt{3}\sqrt{-2x^9 + 91Z^2x^6 - 1376Z^4x^3 + 6912Z^6}Z\right)^{2/3} + x^4 - 2x^2\left(x^6 - 45x^3Z^2 + 432Z^4 + 3\sqrt{3}\sqrt{-2x^9 + 91Z^2x^6 - 1376Z^4x^3 + 6912Z^6}Z\right)^{1/3} - 12xZ^2 + \left(x^6 - 45x^3Z^2 + 432Z^4 + 3\sqrt{3}\sqrt{-2x^9 + 91Z^2x^6 - 1376Z^4x^3 + 6912Z^6}Z\right)^{2/3}\right)$$

Now we go over each such singular solution and check if it verifies the ode itself and any initial conditions given. If it does not then the singular solution will not be used.

The solution $\text{RootOf}\left(-i\sqrt{3}x^4 + 12i\sqrt{3}Z^2x + i\sqrt{3}\left(x^6 - 45x^3Z^2 + 432Z^4 + 3\sqrt{3}\sqrt{-2x^9 + 91Z^2x}\right)\right)$ will not be used

Maple step by step solution

Let's solve

$$4yy'^3 - 2y'^2x^2 + 4xyy' + x^3 = 16y^2$$

- Highest derivative means the order of the ODE is 1

y'

- Solve for the highest derivative

$$\left[y' = \frac{\left(x^6 - 45x^3y^2 + 432y^4 + 3\sqrt{3}\sqrt{-2x^9 + 91y^2x^6 - 1376x^3y^4 + 6912y^6}y\right)^{1/3}}{6y} - \frac{x(-x^3 + 12y^2)}{6y\left(x^6 - 45x^3y^2 + 432y^4 + 3\sqrt{3}\sqrt{-2x^9 + 91y^2x^6 - 1376x^3y^4 + 6912y^6}\right)^{1/3}} \right]$$

- Solve the equation $y' = \frac{\left(x^6 - 45x^3y^2 + 432y^4 + 3\sqrt{3}\sqrt{-2x^9 + 91y^2x^6 - 1376x^3y^4 + 6912y^6}\right)^{1/3}}{6y} - \frac{x(-x^3 + 12y^2)}{6y\left(x^6 - 45x^3y^2 + 432y^4 + 3\sqrt{3}\sqrt{-2x^9 + 91y^2x^6 - 1376x^3y^4 + 6912y^6}\right)^{1/3}}$

- Solve the equation $y' = -\frac{\left(x^6 - 45x^3y^2 + 432y^4 + 3\sqrt{3}\sqrt{-2x^9 + 91y^2x^6 - 1376x^3y^4 + 6912y^6}\right)^{1/3}}{12y} + \frac{x(-x^3 + 12y^2)}{12y\left(x^6 - 45x^3y^2 + 432y^4 + 3\sqrt{3}\sqrt{-2x^9 + 91y^2x^6 - 1376x^3y^4 + 6912y^6}\right)^{1/3}}$

- Solve the equation $y' = -\frac{\left(x^6 - 45x^3y^2 + 432y^4 + 3\sqrt{3}\sqrt{-2x^9 + 91y^2x^6 - 1376x^3y^4 + 6912y^6}\right)^{1/3}}{12y} + \frac{x(-x^3 + 12y^2)}{12y\left(x^6 - 45x^3y^2 + 432y^4 + 3\sqrt{3}\sqrt{-2x^9 + 91y^2x^6 - 1376x^3y^4 + 6912y^6}\right)^{1/3}}$

- Set of solutions

$$\{\text{workingODE}, \text{workingODE}, \text{workingODE}\}$$

Maple dsolve solution

Solving time : 393.754 (sec)

Leaf size : maple_leaf_size

```
dsolve(4*y(x)*diff(y(x),x)^3-2*diff(y(x),x)^2*x^2+4*x*y(x)*diff(y(x),x)+x^3 = 16*y(x)^2,
y(x),singsol=all)
```

No solution found

Mathematica DSolve solution

Solving time : 35.094 (sec)

Leaf size : 49162

```
DSolve[{4*y[x]*D[y[x],x]^3-2*x^2*D[y[x],x]^2+4*x*y[x]*D[y[x],x]+x^3==16*y[x]^2,{}},  
y[x],x,IncludeSingularSolutions->True]
```

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2.8 Chapter VII. Linear equations of order higher than the first. section 56. Problems at page 163

2.8.1	problem 1 (eq 100)	588
2.8.2	problem 2	600
2.8.3	problem 3	614
2.8.4	problem 4	625
2.8.5	problem 5	637
2.8.6	problem 6	639
2.8.7	problem 7	655
2.8.8	problem 8	670
2.8.9	problem 10	686
2.8.10	problem 11	695
2.8.11	problem 14	701
2.8.12	problem 15	708

2.8.1 problem 1 (eq 100)

Solved as second order linear constant coeff ode	588
Solved as second order can be made integrable	589
Solved as second order ode using Kovacic algorithm	592
Solved as second order ode adjoint method	595
Maple step by step solution	598
Maple trace	598
Maple dsolve solution	599
Mathematica DSolve solution	599

Internal problem ID [18256]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter VII. Linear equations of order higher than the first. section 56. Problems at page 163

Problem number : 1 (eq 100)

Date solved : Monday, December 23, 2024 at 09:24:03 PM

CAS classification : [[_2nd_order, _missing_x]]

Solve

$$\theta'' - p^2\theta = 0$$

Solved as second order linear constant coeff ode

Time used: 0.556 (sec)

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$A\theta''(x) + B\theta'(x) + C\theta(x) = 0$$

Where in the above $A = 1, B = 0, C = -p^2$. Let the solution be $\theta = e^{\lambda x}$. Substituting this into the ODE gives

$$\lambda^2 e^{x\lambda} - p^2 e^{x\lambda} = 0 \tag{1}$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$\lambda^2 - p^2 = 0 \tag{2}$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = 0, C = -p^2$ into the above gives

$$\begin{aligned}\lambda_{1,2} &= \frac{0}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{0^2 - (4)(1)(-p^2)} \\ &= \pm \sqrt{p^2}\end{aligned}$$

Hence

$$\lambda_1 = +\sqrt{p^2}$$

$$\lambda_2 = -\sqrt{p^2}$$

Which simplifies to

$$\lambda_1 = \sqrt{p^2}$$

$$\lambda_2 = -\sqrt{p^2}$$

Since roots are real and distinct, then the solution is

$$\theta = c_1 e^{\lambda_1 x} + c_2 e^{\lambda_2 x}$$

$$\theta = c_1 e^{(\sqrt{p^2})x} + c_2 e^{(-\sqrt{p^2})x}$$

Or

$$\theta = c_1 e^{x\sqrt{p^2}} + c_2 e^{-x\sqrt{p^2}}$$

Will add steps showing solving for IC soon.

Summary of solutions found

$$\theta = c_1 e^{x\sqrt{p^2}} + c_2 e^{-x\sqrt{p^2}}$$

Solved as second order can be made integrable

Time used: 5.419 (sec)

Multiplying the ode by θ' gives

$$\theta' \theta'' - p^2 \theta' \theta = 0$$

Integrating the above w.r.t x gives

$$\begin{aligned}\int (\theta' \theta'' - p^2 \theta' \theta) dx &= 0 \\ \frac{\theta'^2}{2} - \frac{p^2 \theta^2}{2} &= c_1\end{aligned}$$

Which is now solved for θ . Solving for the derivative gives these ODE's to solve

$$\theta' = \sqrt{p^2\theta^2 + 2c_1} \quad (1)$$

$$\theta' = -\sqrt{p^2\theta^2 + 2c_1} \quad (2)$$

Now each of the above is solved separately.

Solving Eq. (1)

Integrating gives

$$\int \frac{1}{\sqrt{p^2\theta^2 + 2c_1}} d\theta = dx$$

$$\frac{\ln\left(\frac{p^2\theta}{\sqrt{p^2}} + \sqrt{p^2\theta^2 + 2c_1}\right)}{\sqrt{p^2}} = x + c_2$$

Singular solutions are found by solving

$$\sqrt{p^2\theta^2 + 2c_1} = 0$$

for θ . This is because we had to divide by this in the above step. This gives the following singular solution(s), which also have to satisfy the given ODE.

$$\theta = \frac{\sqrt{-2c_1}}{p}$$

$$\theta = -\frac{\sqrt{-2c_1}}{p}$$

Solving for θ gives

$$\theta = \frac{\sqrt{-2c_1}}{p}$$

$$\theta = -\frac{\sqrt{-2c_1}}{p}$$

$$\theta = -\frac{\sqrt{p^2} \left(-e^{2c_2\sqrt{p^2}+2x\sqrt{p^2}} + 2c_1 \right) e^{-c_2\sqrt{p^2}-x\sqrt{p^2}}}{2p^2}$$

Solving Eq. (2)

Integrating gives

$$\int -\frac{1}{\sqrt{p^2\theta^2 + 2c_1}} d\theta = dx$$

$$-\frac{\ln\left(\frac{p^2\theta}{\sqrt{p^2}} + \sqrt{p^2\theta^2 + 2c_1}\right)}{\sqrt{p^2}} = x + c_3$$

Singular solutions are found by solving

$$-\sqrt{p^2\theta^2 + 2c_1} = 0$$

for θ . This is because we had to divide by this in the above step. This gives the following singular solution(s), which also have to satisfy the given ODE.

$$\theta = \frac{\sqrt{-2c_1}}{p}$$

$$\theta = -\frac{\sqrt{-2c_1}}{p}$$

Solving for θ gives

$$\theta = \frac{\sqrt{-2c_1}}{p}$$

$$\theta = -\frac{\sqrt{-2c_1}}{p}$$

$$\theta = -\frac{\sqrt{p^2} \left(-e^{-2c_3\sqrt{p^2}-2x\sqrt{p^2}} + 2c_1 \right) e^{c_3\sqrt{p^2}+x\sqrt{p^2}}}{2p^2}$$

Will add steps showing solving for IC soon.

The solution

$$\theta = \frac{\sqrt{-2c_1}}{p}$$

was found not to satisfy the ode or the IC. Hence it is removed. The solution

$$\theta = -\frac{\sqrt{-2c_1}}{p}$$

was found not to satisfy the ode or the IC. Hence it is removed.

Summary of solutions found

$$\theta = -\frac{\sqrt{p^2} \left(-e^{2c_2\sqrt{p^2}+2x\sqrt{p^2}} + 2c_1 \right) e^{-c_2\sqrt{p^2}-x\sqrt{p^2}}}{2p^2}$$

$$\theta = -\frac{\sqrt{p^2} \left(-e^{-2c_3\sqrt{p^2}-2x\sqrt{p^2}} + 2c_1 \right) e^{c_3\sqrt{p^2}+x\sqrt{p^2}}}{2p^2}$$

Solved as second order ode using Kovacic algorithm

Time used: 0.081 (sec)

Writing the ode as

$$\theta'' - p^2\theta = 0 \quad (1)$$

$$A\theta'' + B\theta' + C\theta = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$A = 1$$

$$B = 0 \quad (3)$$

$$C = -p^2$$

Applying the Liouville transformation on the dependent variable gives

$$z(x) = \theta e^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \quad (4)$$

Where r is given by

$$r = \frac{s}{t} \quad (5)$$

$$= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2}$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{p^2}{1} \quad (6)$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= p^2 \\ t &= 1 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(x) = (p^2) z(x) \quad (7)$$

Equation (7) is now solved. After finding $z(x)$ then θ is found using the inverse transformation

$$\theta = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 2.47: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - 0 \\ &= 0 \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = p^2$ is not a function of x , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(x) = e^{x\sqrt{p^2}}$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in θ is found from

$$\theta_1 = z_1 e^{\int -\frac{1}{2} \frac{B}{A} dx}$$

Since $B = 0$ then the above reduces to

$$\begin{aligned} \theta_1 &= z_1 \\ &= e^{x\sqrt{p^2}} \end{aligned}$$

Which simplifies to

$$\theta_1 = e^{x\sqrt{p^2}}$$

The second solution θ_2 to the original ode is found using reduction of order

$$\theta_2 = \theta_1 \int \frac{e^{\int -\frac{B}{A} dx}}{\theta_1^2} dx$$

Since $B = 0$ then the above becomes

$$\begin{aligned} \theta_2 &= \theta_1 \int \frac{1}{\theta_1^2} dx \\ &= e^{x\sqrt{p^2}} \int \frac{1}{e^{2x\sqrt{p^2}}} dx \\ &= e^{x\sqrt{p^2}} \left(-\frac{\sqrt{p^2} e^{-2x\sqrt{p^2}}}{2p^2} \right) \end{aligned}$$

Therefore the solution is

$$\begin{aligned} \theta &= c_1 \theta_1 + c_2 \theta_2 \\ &= c_1 \left(e^{x\sqrt{p^2}} \right) + c_2 \left(e^{x\sqrt{p^2}} \left(-\frac{\sqrt{p^2} e^{-2x\sqrt{p^2}}}{2p^2} \right) \right) \end{aligned}$$

Will add steps showing solving for IC soon.

Summary of solutions found

$$\theta = c_1 e^{x\sqrt{p^2}} - \frac{c_2 \operatorname{csgn}(p) e^{-x \operatorname{csgn}(p)p}}{2p}$$

Solved as second order ode adjoint method

Time used: 0.247 (sec)

In normal form the ode

$$\theta'' - p^2\theta = 0 \quad (1)$$

Becomes

$$\theta'' + p(x)\theta' + q(x)\theta = r(x) \quad (2)$$

Where

$$\begin{aligned} p(x) &= 0 \\ q(x) &= -p^2 \\ r(x) &= 0 \end{aligned}$$

The Lagrange adjoint ode is given by

$$\begin{aligned} \xi'' - (\xi p)' + \xi q &= 0 \\ \xi'' - (0)' + (-p^2\xi(x)) &= 0 \\ \xi''(x) - p^2\xi(x) &= 0 \end{aligned}$$

Which is solved for $\xi(x)$. This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$A\xi''(x) + B\xi'(x) + C\xi(x) = 0$$

Where in the above $A = 1, B = 0, C = -p^2$. Let the solution be $\xi = e^{\lambda x}$. Substituting this into the ODE gives

$$\lambda^2 e^{x\lambda} - p^2 e^{x\lambda} = 0 \quad (1)$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$\lambda^2 - p^2 = 0 \quad (2)$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = 0, C = -p^2$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{0}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{0^2 - (4)(1)(-p^2)} \\ &= \pm \sqrt{p^2} \end{aligned}$$

Hence

$$\begin{aligned} \lambda_1 &= +\sqrt{p^2} \\ \lambda_2 &= -\sqrt{p^2} \end{aligned}$$

Which simplifies to

$$\begin{aligned} \lambda_1 &= \sqrt{p^2} \\ \lambda_2 &= -\sqrt{p^2} \end{aligned}$$

Since roots are real and distinct, then the solution is

$$\begin{aligned} \xi &= c_1 e^{\lambda_1 x} + c_2 e^{\lambda_2 x} \\ \xi &= c_1 e^{(\sqrt{p^2})x} + c_2 e^{(-\sqrt{p^2})x} \end{aligned}$$

Or

$$\xi = c_1 e^{x\sqrt{p^2}} + c_2 e^{-x\sqrt{p^2}}$$

Will add steps showing solving for IC soon.

The original ode now reduces to first order ode

$$\begin{aligned} \xi(x) \theta' - \theta \xi'(x) + \xi(x) p(x) \theta &= \int \xi(x) r(x) dx \\ \theta' + \theta \left(p(x) - \frac{\xi'(x)}{\xi(x)} \right) &= \frac{\int \xi(x) r(x) dx}{\xi(x)} \end{aligned}$$

Or

$$\theta' - \frac{\theta \left(c_1 \sqrt{p^2} e^{x\sqrt{p^2}} - c_2 \sqrt{p^2} e^{-x\sqrt{p^2}} \right)}{c_1 e^{x\sqrt{p^2}} + c_2 e^{-x\sqrt{p^2}}} = 0$$

Which is now a first order ode. This is now solved for θ . In canonical form a linear first order is

$$\theta' + q(x)\theta = p(x)$$

Comparing the above to the given ode shows that

$$q(x) = -\frac{\operatorname{csgn}(p)p(c_1e^{2x\operatorname{csgn}(p)p} - c_2)}{c_1e^{2x\operatorname{csgn}(p)p} + c_2}$$

$$p(x) = 0$$

The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int q dx} \\ &= e^{\int -\frac{\operatorname{csgn}(p)p(c_1e^{2x\operatorname{csgn}(p)p} - c_2)}{c_1e^{2x\operatorname{csgn}(p)p} + c_2} dx} \\ &= e^{\frac{\ln(e^{2x\operatorname{csgn}(p)p})}{2} - \ln(c_1e^{2x\operatorname{csgn}(p)p} + c_2)}\end{aligned}$$

The ode becomes

$$\frac{d}{dx}\left(\theta e^{\frac{\ln(e^{2x\operatorname{csgn}(p)p})}{2} - \ln(c_1e^{2x\operatorname{csgn}(p)p} + c_2)}\right) = 0$$

Integrating gives

$$\begin{aligned}\theta e^{\frac{\ln(e^{2x\operatorname{csgn}(p)p})}{2} - \ln(c_1e^{2x\operatorname{csgn}(p)p} + c_2)} &= \int 0 dx + c_3 \\ &= c_3\end{aligned}$$

Dividing throughout by the integrating factor $e^{\frac{\ln(e^{2x\operatorname{csgn}(p)p})}{2} - \ln(c_1e^{2x\operatorname{csgn}(p)p} + c_2)}$ gives the final solution

$$\theta = \frac{(c_1e^{2x\operatorname{csgn}(p)p} + c_2) c_3}{\sqrt{e^{2x\operatorname{csgn}(p)p}}}$$

Hence, the solution found using Lagrange adjoint equation method is

$$\theta = \frac{(c_1e^{2x\operatorname{csgn}(p)p} + c_2) c_3}{\sqrt{e^{2x\operatorname{csgn}(p)p}}}$$

The constants can be merged to give

$$\theta = \frac{c_1e^{2x\operatorname{csgn}(p)p} + c_2}{\sqrt{e^{2x\operatorname{csgn}(p)p}}}$$

Will add steps showing solving for IC soon.

Summary of solutions found

$$\theta = \frac{c_1 e^{2x \operatorname{sgn}(p)p} + c_2}{\sqrt{e^{2x \operatorname{sgn}(p)p}}}$$

Maple step by step solution

Let's solve

$$\theta'' - p^2\theta = 0$$

- Highest derivative means the order of the ODE is 2
 θ''
- Characteristic polynomial of ODE
 $-p^2 + r^2 = 0$
- Factor the characteristic polynomial
 $-(p - r)(p + r) = 0$
- Roots of the characteristic polynomial
 $r = (p, -p)$
- 1st solution of the ODE
 $\theta_1(x) = e^{xp}$
- 2nd solution of the ODE
 $\theta_2(x) = e^{-xp}$
- General solution of the ODE
 $\theta = C1\theta_1(x) + C2\theta_2(x)$
- Substitute in solutions
 $\theta = C1 e^{xp} + C2 e^{-xp}$

Maple trace

```
`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
<- constant coefficients successful`
```

Maple dsolve solution

Solving time : 0.004 (sec)

Leaf size : 18

```
dsolve(diff(diff(theta(x),x),x)-p^2*theta(x) = 0,  
        theta(x),singsol=all)
```

$$\theta = c_1 e^{-xp} + c_2 e^{xp}$$

Mathematica DSolve solution

Solving time : 0.013 (sec)

Leaf size : 23

```
DSolve[{D[theta[x],{x,2}]-p^2*theta[x]==0,{}},  
        theta[x],x,IncludeSingularSolutions->True]
```

$$\theta(x) \rightarrow c_1 e^{px} + c_2 e^{-px}$$

2.8.2 problem 2

Solved as second order linear constant coeff ode	600
Solved as second order can be made integrable	602
Solved as second order ode using Kovacic algorithm	605
Solved as second order ode adjoint method	609
Maple step by step solution	612
Maple trace	613
Maple dsolve solution	613
Mathematica DSolve solution	613

Internal problem ID [18257]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter VII. Linear equations of order higher than the first. section 56. Problems at page 163

Problem number : 2

Date solved : Monday, December 23, 2024 at 09:24:10 PM

CAS classification : [[_2nd_order, _missing_x]]

Solve

$$y'' + y = 0$$

Solved as second order linear constant coeff ode

Time used: 0.117 (sec)

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = 0$$

Where in the above $A = 1, B = 0, C = 1$. Let the solution be $y = e^{\lambda x}$. Substituting this into the ODE gives

$$\lambda^2 e^{x\lambda} + e^{x\lambda} = 0 \tag{1}$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$\lambda^2 + 1 = 0 \tag{2}$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = 0, C = 1$ into the above gives

$$\begin{aligned}\lambda_{1,2} &= \frac{0}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{0^2 - (4)(1)(1)} \\ &= \pm i\end{aligned}$$

Hence

$$\lambda_1 = +i$$

$$\lambda_2 = -i$$

Which simplifies to

$$\lambda_1 = i$$

$$\lambda_2 = -i$$

Since roots are complex conjugate of each others, then let the roots be

$$\lambda_{1,2} = \alpha \pm i\beta$$

Where $\alpha = 0$ and $\beta = 1$. Therefore the final solution, when using Euler relation, can be written as

$$y = e^{\alpha x}(c_1 \cos(\beta x) + c_2 \sin(\beta x))$$

Which becomes

$$y = e^0(c_1 \cos(x) + c_2 \sin(x))$$

Or

$$y = c_1 \cos(x) + c_2 \sin(x)$$

Will add steps showing solving for IC soon.

Summary of solutions found

$$y = c_1 \cos(x) + c_2 \sin(x)$$

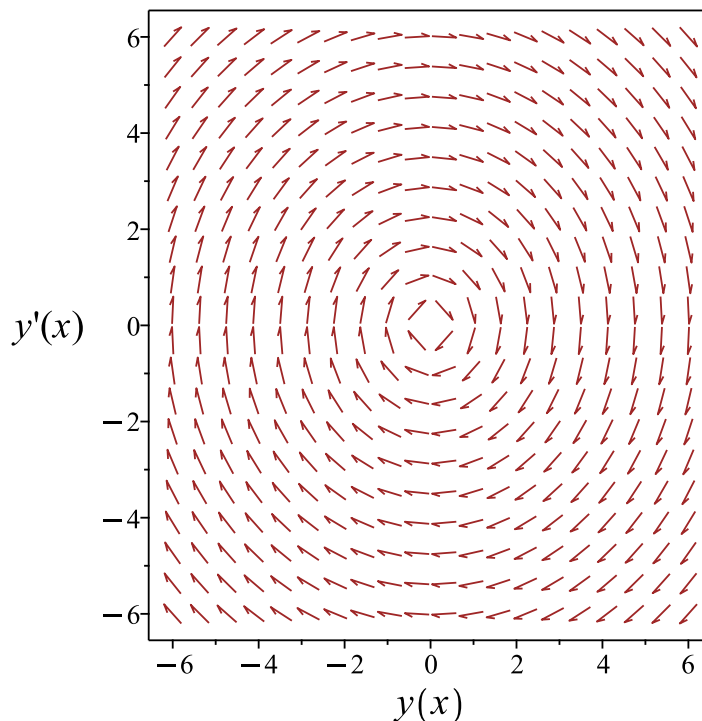


Figure 2.90: Slope field plot
 $y'' + y = 0$

Solved as second order can be made integrable

Time used: 3.275 (sec)

Multiplying the ode by y' gives

$$y'y'' + yy' = 0$$

Integrating the above w.r.t x gives

$$\int (y'y'' + yy') dx = 0$$

$$\frac{y'^2}{2} + \frac{y^2}{2} = c_1$$

Which is now solved for y . Solving for the derivative gives these ODE's to solve

$$y' = \sqrt{-y^2 + 2c_1} \quad (1)$$

$$y' = -\sqrt{-y^2 + 2c_1} \quad (2)$$

Now each of the above is solved separately.

Solving Eq. (1)

Integrating gives

$$\int \frac{1}{\sqrt{-y^2 + 2c_1}} dy = dx$$

$$\arctan\left(\frac{y}{\sqrt{-y^2 + 2c_1}}\right) = x + c_2$$

Singular solutions are found by solving

$$\sqrt{-y^2 + 2c_1} = 0$$

for y . This is because we had to divide by this in the above step. This gives the following singular solution(s), which also have to satisfy the given ODE.

$$y = \sqrt{2} \sqrt{c_1}$$

$$y = -\sqrt{2} \sqrt{c_1}$$

Solving for y gives

$$y = \sqrt{2} \sqrt{c_1}$$

$$y = \tan(x + c_2) \sqrt{2} \sqrt{\frac{c_1}{\tan^2(x + c_2) + 1}}$$

$$y = -\sqrt{2} \sqrt{c_1}$$

Solving Eq. (2)

Integrating gives

$$\int -\frac{1}{\sqrt{-y^2 + 2c_1}} dy = dx$$

$$-\arctan\left(\frac{y}{\sqrt{-y^2 + 2c_1}}\right) = x + c_3$$

Singular solutions are found by solving

$$-\sqrt{-y^2 + 2c_1} = 0$$

for y . This is because we had to divide by this in the above step. This gives the following singular solution(s), which also have to satisfy the given ODE.

$$y = \sqrt{2} \sqrt{c_1}$$

$$y = -\sqrt{2} \sqrt{c_1}$$

Solving for y gives

$$y = \sqrt{2} \sqrt{c_1}$$

$$y = -\sqrt{2} \sqrt{c_1}$$

$$y = -\tan(x + c_3) \sqrt{2} \sqrt{\frac{c_1}{\tan(x + c_3)^2 + 1}}$$

Will add steps showing solving for IC soon.

The solution

$$y = \sqrt{2} \sqrt{c_1}$$

was found not to satisfy the ode or the IC. Hence it is removed. The solution

$$y = -\sqrt{2} \sqrt{c_1}$$

was found not to satisfy the ode or the IC. Hence it is removed.

Summary of solutions found

$$y = \tan(x + c_2) \sqrt{2} \sqrt{\frac{c_1}{\tan(x + c_2)^2 + 1}}$$

$$y = -\tan(x + c_3) \sqrt{2} \sqrt{\frac{c_1}{\tan(x + c_3)^2 + 1}}$$

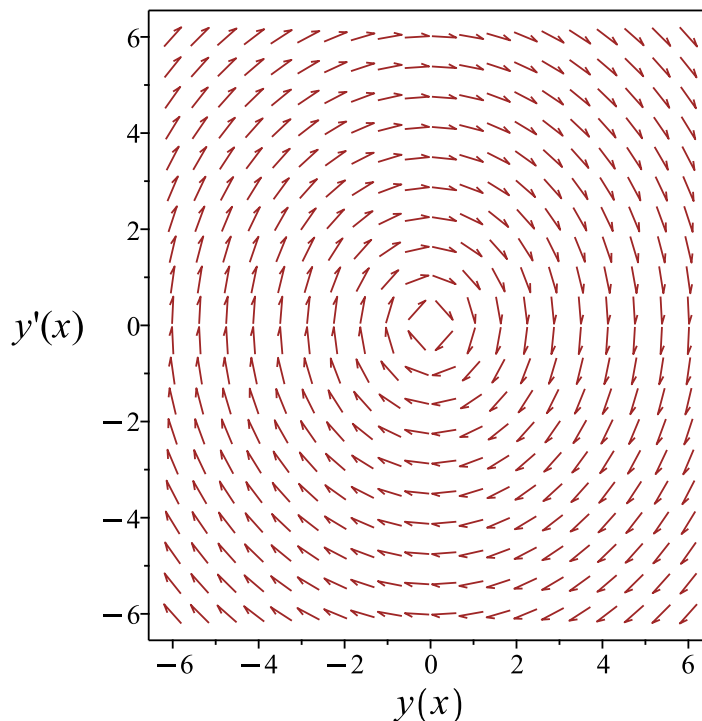


Figure 2.91: Slope field plot
 $y'' + y = 0$

Solved as second order ode using Kovacic algorithm

Time used: 0.105 (sec)

Writing the ode as

$$y'' + y = 0 \quad (1)$$

$$Ay'' + By' + Cy = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$A = 1 \quad (3)$$

$$B = 0$$

$$C = 1$$

Applying the Liouville transformation on the dependent variable gives

$$z(x) = ye^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \quad (4)$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \quad (5)$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{-1}{1} \quad (6)$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= -1 \\ t &= 1 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(x) = -z(x) \quad (7)$$

Equation (7) is now solved. After finding $z(x)$ then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 2.49: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - 0 \\ &= 0 \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = -1$ is not a function of x , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(x) = \cos(x)$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$y_1 = z_1 e^{\int -\frac{1}{2} \frac{B}{A} dx}$$

Since $B = 0$ then the above reduces to

$$\begin{aligned} y_1 &= z_1 \\ &= \cos(x) \end{aligned}$$

Which simplifies to

$$y_1 = \cos(x)$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dx}}{y_1^2} dx$$

Since $B = 0$ then the above becomes

$$\begin{aligned} y_2 &= y_1 \int \frac{1}{y_1^2} dx \\ &= \cos(x) \int \frac{1}{\cos(x)^2} dx \\ &= \cos(x) (\tan(x)) \end{aligned}$$

Therefore the solution is

$$\begin{aligned}y &= c_1 y_1 + c_2 y_2 \\ &= c_1(\cos(x)) + c_2(\cos(x)(\tan(x)))\end{aligned}$$

Will add steps showing solving for IC soon.

Summary of solutions found

$$y = c_1 \cos(x) + c_2 \sin(x)$$

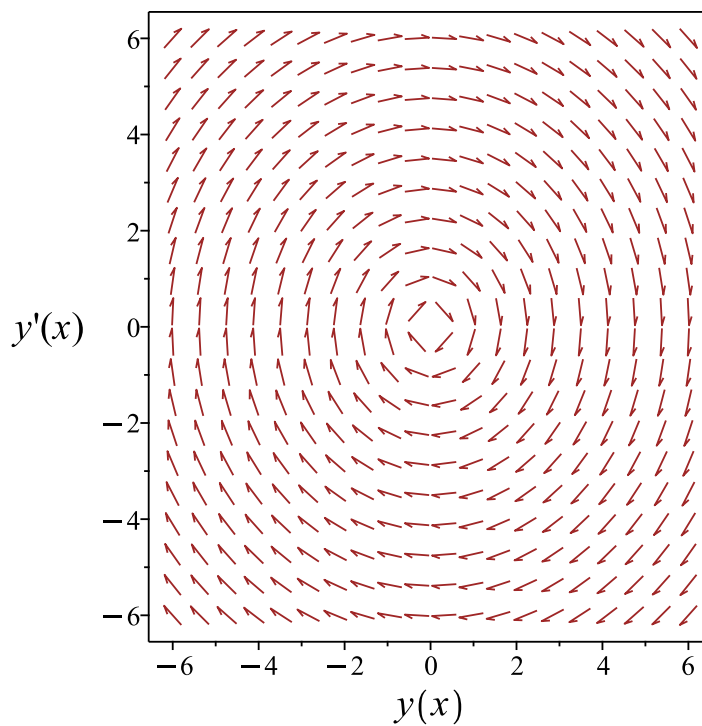


Figure 2.92: Slope field plot
 $y'' + y = 0$

Solved as second order ode adjoint method

Time used: 0.288 (sec)

In normal form the ode

$$y'' + y = 0 \quad (1)$$

Becomes

$$y'' + p(x)y' + q(x)y = r(x) \quad (2)$$

Where

$$p(x) = 0$$

$$q(x) = 1$$

$$r(x) = 0$$

The Lagrange adjoint ode is given by

$$\xi'' - (\xi p)' + \xi q = 0$$

$$\xi'' - (0)' + (\xi(x)) = 0$$

$$\xi''(x) + \xi(x) = 0$$

Which is solved for $\xi(x)$. This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$A\xi''(x) + B\xi'(x) + C\xi(x) = 0$$

Where in the above $A = 1, B = 0, C = 1$. Let the solution be $\xi = e^{\lambda x}$. Substituting this into the ODE gives

$$\lambda^2 e^{x\lambda} + e^{x\lambda} = 0 \quad (1)$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$\lambda^2 + 1 = 0 \quad (2)$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = 0, C = 1$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{0}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{0^2 - (4)(1)(1)} \\ &= \pm i \end{aligned}$$

Hence

$$\lambda_1 = +i$$

$$\lambda_2 = -i$$

Which simplifies to

$$\lambda_1 = i$$

$$\lambda_2 = -i$$

Since roots are complex conjugate of each others, then let the roots be

$$\lambda_{1,2} = \alpha \pm i\beta$$

Where $\alpha = 0$ and $\beta = 1$. Therefore the final solution, when using Euler relation, can be written as

$$\xi = e^{\alpha x}(c_1 \cos(\beta x) + c_2 \sin(\beta x))$$

Which becomes

$$\xi = e^0(c_1 \cos(x) + c_2 \sin(x))$$

Or

$$\xi = c_1 \cos(x) + c_2 \sin(x)$$

Will add steps showing solving for IC soon.

The original ode now reduces to first order ode

$$\begin{aligned} \xi(x) y' - y \xi'(x) + \xi(x) p(x) y &= \int \xi(x) r(x) dx \\ y' + y \left(p(x) - \frac{\xi'(x)}{\xi(x)} \right) &= \frac{\int \xi(x) r(x) dx}{\xi(x)} \end{aligned}$$

Or

$$y' - \frac{y(-c_1 \sin(x) + c_2 \cos(x))}{c_1 \cos(x) + c_2 \sin(x)} = 0$$

Which is now a first order ode. This is now solved for y . In canonical form a linear first order is

$$y' + q(x)y = p(x)$$

Comparing the above to the given ode shows that

$$q(x) = -\frac{-c_1 \sin(x) + c_2 \cos(x)}{c_1 \cos(x) + c_2 \sin(x)}$$

$$p(x) = 0$$

The integrating factor μ is

$$\begin{aligned} \mu &= e^{\int q dx} \\ &= e^{\int -\frac{-c_1 \sin(x) + c_2 \cos(x)}{c_1 \cos(x) + c_2 \sin(x)} dx} \\ &= \frac{1}{c_1 \cos(x) + c_2 \sin(x)} \end{aligned}$$

The ode becomes

$$\frac{d}{dx} \mu y = 0$$

$$\frac{d}{dx} \left(\frac{y}{c_1 \cos(x) + c_2 \sin(x)} \right) = 0$$

Integrating gives

$$\begin{aligned} \frac{y}{c_1 \cos(x) + c_2 \sin(x)} &= \int 0 dx + c_3 \\ &= c_3 \end{aligned}$$

Dividing throughout by the integrating factor $\frac{1}{c_1 \cos(x) + c_2 \sin(x)}$ gives the final solution

$$y = (c_1 \cos(x) + c_2 \sin(x)) c_3$$

Hence, the solution found using Lagrange adjoint equation method is

$$y = (c_1 \cos(x) + c_2 \sin(x)) c_3$$

The constants can be merged to give

$$y = c_1 \cos(x) + c_2 \sin(x)$$

Will add steps showing solving for IC soon.

Summary of solutions found

$$y = c_1 \cos(x) + c_2 \sin(x)$$

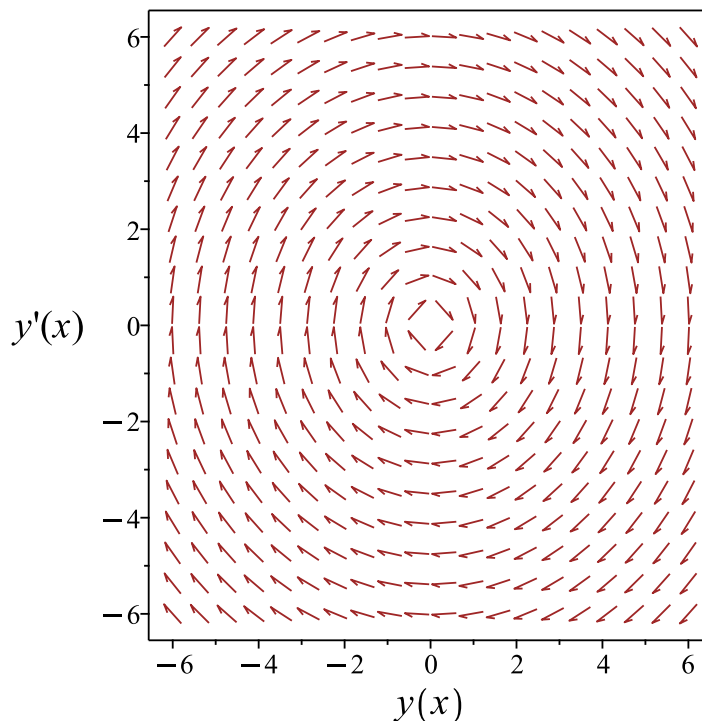


Figure 2.93: Slope field plot
 $y'' + y = 0$

Maple step by step solution

Let's solve

$$y'' + y = 0$$

- Highest derivative means the order of the ODE is 2
- y''
- Characteristic polynomial of ODE
- $r^2 + 1 = 0$
- Use quadratic formula to solve for r
- $r = \frac{0 \pm (\sqrt{-4})}{2}$
- Roots of the characteristic polynomial
- $r = (-i, i)$
- 1st solution of the ODE
- $y_1(x) = \cos(x)$
- 2nd solution of the ODE
- $y_2(x) = \sin(x)$
- General solution of the ODE

- $y = C1y_1(x) + C2y_2(x)$
Substitute in solutions
 $y = C1 \cos(x) + C2 \sin(x)$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
<- constant coefficients successful`

```

Maple dsolve solution

Solving time : 0.001 (sec)

Leaf size : 13

```

dsolve(diff(diff(y(x),x),x)+y(x) = 0,
        y(x),singsol=all)

```

$$y = c_1 \sin(x) + c_2 \cos(x)$$

Mathematica DSolve solution

Solving time : 0.011 (sec)

Leaf size : 16

```

DSolve[{D[y[x],{x,2}]+y[x]==0,{}}],
        y[x],x,IncludeSingularSolutions->True]

```

$$y(x) \rightarrow c_1 \cos(x) + c_2 \sin(x)$$

2.8.3 problem 3

Solved as second order linear constant coeff ode	614
Solved as second order ode using Kovacic algorithm	616
Solved as second order ode adjoint method	620
Maple step by step solution	623
Maple trace	624
Maple dsolve solution	624
Mathematica DSolve solution	624

Internal problem ID [18258]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter VII. Linear equations of order higher than the first. section 56. Problems at page 163

Problem number : 3

Date solved : Monday, December 23, 2024 at 09:24:14 PM

CAS classification : [[_2nd_order, _missing_x]]

Solve

$$y'' + 12y = 7y'$$

Solved as second order linear constant coeff ode

Time used: 0.108 (sec)

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = 0$$

Where in the above $A = 1, B = -7, C = 12$. Let the solution be $y = e^{\lambda x}$. Substituting this into the ODE gives

$$\lambda^2 e^{x\lambda} - 7\lambda e^{x\lambda} + 12e^{x\lambda} = 0 \quad (1)$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$\lambda^2 - 7\lambda + 12 = 0 \quad (2)$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = -7, C = 12$ into the above gives

$$\begin{aligned}\lambda_{1,2} &= \frac{7}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{-7^2 - (4)(1)(12)} \\ &= \frac{7}{2} \pm \frac{1}{2}\end{aligned}$$

Hence

$$\begin{aligned}\lambda_1 &= \frac{7}{2} + \frac{1}{2} \\ \lambda_2 &= \frac{7}{2} - \frac{1}{2}\end{aligned}$$

Which simplifies to

$$\begin{aligned}\lambda_1 &= 4 \\ \lambda_2 &= 3\end{aligned}$$

Since roots are real and distinct, then the solution is

$$\begin{aligned}y &= c_1 e^{\lambda_1 x} + c_2 e^{\lambda_2 x} \\ y &= c_1 e^{(4)x} + c_2 e^{(3)x}\end{aligned}$$

Or

$$y = c_1 e^{4x} + c_2 e^{3x}$$

Will add steps showing solving for IC soon.

Summary of solutions found

$$y = c_1 e^{4x} + c_2 e^{3x}$$

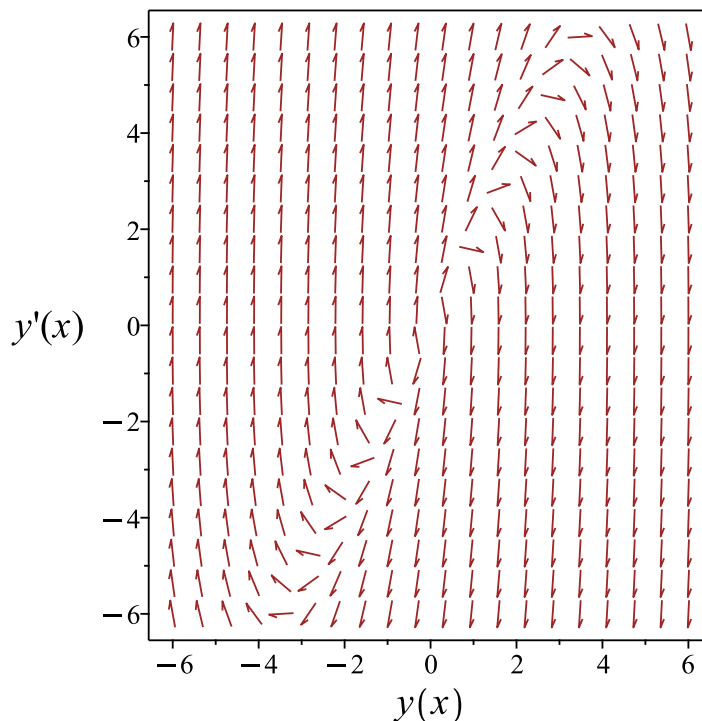


Figure 2.94: Slope field plot
 $y'' + 12y = 7y'$

Solved as second order ode using Kovacic algorithm

Time used: 0.516 (sec)

Writing the ode as

$$y'' + 12y - 7y' = 0 \quad (1)$$

$$Ay'' + By' + Cy = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$\begin{aligned} A &= 1 \\ B &= -7 \\ C &= 12 \end{aligned} \quad (3)$$

Applying the Liouville transformation on the dependent variable gives

$$z(x) = ye^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \quad (4)$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \quad (5)$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{1}{4} \quad (6)$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= 1 \\ t &= 4 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(x) = \frac{z(x)}{4} \quad (7)$$

Equation (7) is now solved. After finding $z(x)$ then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 2.51: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - 0 \\ &= 0 \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = \frac{1}{4}$ is not a function of x , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(x) = e^{-\frac{x}{2}}$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned} y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dx} \\ &= z_1 e^{-\int \frac{1}{2} \frac{-7}{1} dx} \\ &= z_1 e^{\frac{7x}{2}} \\ &= z_1 \left(e^{\frac{7x}{2}} \right) \end{aligned}$$

Which simplifies to

$$y_1 = e^{3x}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dx}}{y_1^2} dx$$

Substituting gives

$$\begin{aligned} y_2 &= y_1 \int \frac{e^{\int -\frac{-7}{1} dx}}{(y_1)^2} dx \\ &= y_1 \int \frac{e^{7x}}{(y_1)^2} dx \\ &= y_1 (e^{7x} e^{-6x}) \end{aligned}$$

Therefore the solution is

$$\begin{aligned} y &= c_1 y_1 + c_2 y_2 \\ &= c_1 (e^{3x}) + c_2 (e^{3x} (e^{7x} e^{-6x})) \end{aligned}$$

Will add steps showing solving for IC soon.

Summary of solutions found

$$y = c_1 e^{3x} + c_2 e^{4x}$$

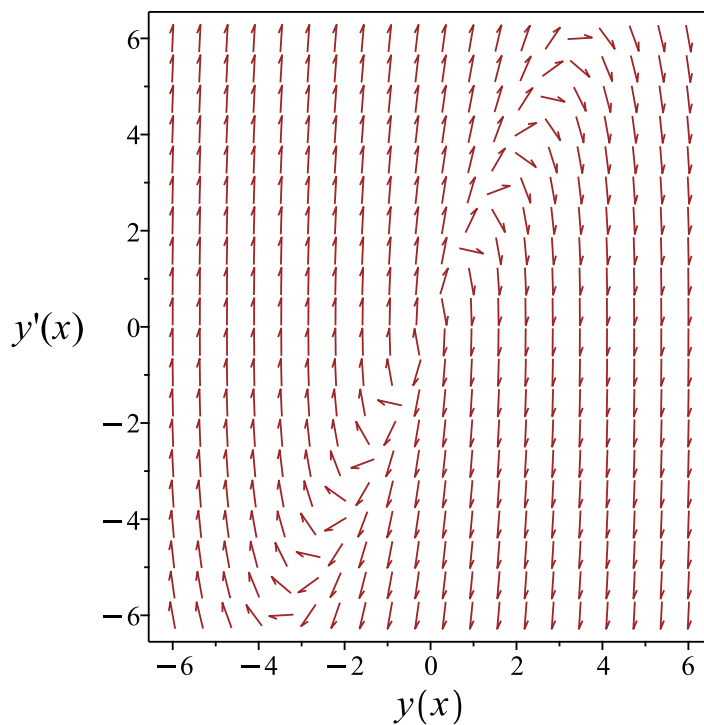


Figure 2.95: Slope field plot
 $y'' + 12y = 7y'$

Solved as second order ode adjoint method

Time used: 0.181 (sec)

In normal form the ode

$$y'' + 12y = 7y' \quad (1)$$

Becomes

$$y'' + p(x)y' + q(x)y = r(x) \quad (2)$$

Where

$$p(x) = -7$$

$$q(x) = 12$$

$$r(x) = 0$$

The Lagrange adjoint ode is given by

$$\xi'' - (\xi p)' + \xi q = 0$$

$$\xi'' - (-7\xi(x))' + (12\xi(x)) = 0$$

$$\xi''(x) + 7\xi'(x) + 12\xi(x) = 0$$

Which is solved for $\xi(x)$. This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$A\xi''(x) + B\xi'(x) + C\xi(x) = 0$$

Where in the above $A = 1, B = 7, C = 12$. Let the solution be $\xi = e^{\lambda x}$. Substituting this into the ODE gives

$$\lambda^2 e^{x\lambda} + 7\lambda e^{x\lambda} + 12e^{x\lambda} = 0 \quad (1)$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$\lambda^2 + 7\lambda + 12 = 0 \quad (2)$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = 7, C = 12$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{-7}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{7^2 - (4)(1)(12)} \\ &= -\frac{7}{2} \pm \frac{1}{2} \end{aligned}$$

Hence

$$\lambda_1 = -\frac{7}{2} + \frac{1}{2}$$

$$\lambda_2 = -\frac{7}{2} - \frac{1}{2}$$

Which simplifies to

$$\lambda_1 = -3$$

$$\lambda_2 = -4$$

Since roots are real and distinct, then the solution is

$$\xi = c_1 e^{\lambda_1 x} + c_2 e^{\lambda_2 x}$$

$$\xi = c_1 e^{(-3)x} + c_2 e^{(-4)x}$$

Or

$$\xi = c_1 e^{-3x} + c_2 e^{-4x}$$

Will add steps showing solving for IC soon.

The original ode now reduces to first order ode

$$\xi(x) y' - y \xi'(x) + \xi(x) p(x) y = \int \xi(x) r(x) dx$$

$$y' + y \left(p(x) - \frac{\xi'(x)}{\xi(x)} \right) = \frac{\int \xi(x) r(x) dx}{\xi(x)}$$

Or

$$y' + y \left(-7 - \frac{-3c_1 e^{-3x} - 4c_2 e^{-4x}}{c_1 e^{-3x} + c_2 e^{-4x}} \right) = 0$$

Which is now a first order ode. This is now solved for y . In canonical form a linear first order is

$$y' + q(x)y = p(x)$$

Comparing the above to the given ode shows that

$$q(x) = -\frac{3e^{-x}c_2 + 4c_1}{e^{-x}c_2 + c_1}$$

$$p(x) = 0$$

The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int q dx} \\ &= e^{\int -\frac{3e^{-x}c_2+4c_1}{e^{-x}c_2+c_1} dx} \\ &= \frac{e^{-4x}}{e^{-x}c_2+c_1}\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dx}\mu y &= 0 \\ \frac{d}{dx}\left(\frac{y e^{-4x}}{e^{-x}c_2+c_1}\right) &= 0\end{aligned}$$

Integrating gives

$$\begin{aligned}\frac{y e^{-4x}}{e^{-x}c_2+c_1} &= \int 0 dx + c_3 \\ &= c_3\end{aligned}$$

Dividing throughout by the integrating factor $\frac{e^{-4x}}{e^{-x}c_2+c_1}$ gives the final solution

$$y = e^{3x}(c_1e^x + c_2) c_3$$

Hence, the solution found using Lagrange adjoint equation method is

$$y = e^{3x}(c_1e^x + c_2) c_3$$

The constants can be merged to give

$$y = e^{3x}(c_1e^x + c_2)$$

Will add steps showing solving for IC soon.

Summary of solutions found

$$y = e^{3x}(c_1e^x + c_2)$$

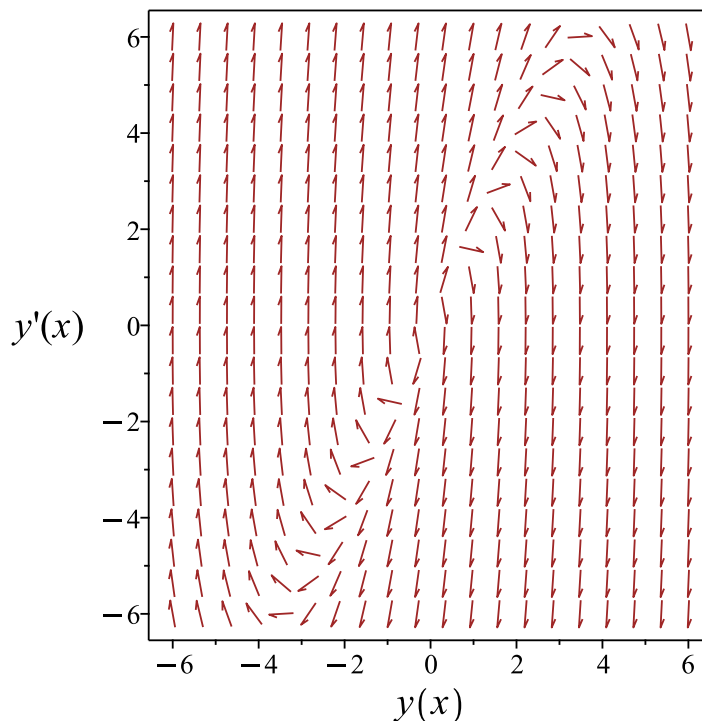


Figure 2.96: Slope field plot
 $y'' + 12y = 7y'$

Maple step by step solution

Let's solve

$$y'' + 12y = 7y'$$

- Highest derivative means the order of the ODE is 2

$$y''$$

- Isolate 2nd derivative

$$y'' = -12y + 7y'$$

- Group terms with y on the lhs of the ODE and the rest on the rhs of the ODE; ODE is linear

$$y'' + 12y - 7y' = 0$$

- Characteristic polynomial of ODE

$$r^2 - 7r + 12 = 0$$

- Factor the characteristic polynomial

$$(r - 3)(r - 4) = 0$$

- Roots of the characteristic polynomial

$$r = (3, 4)$$

- 1st solution of the ODE

- $y_1(x) = e^{3x}$
- 2nd solution of the ODE
 $y_2(x) = e^{4x}$
- General solution of the ODE
 $y = C1y_1(x) + C2y_2(x)$
- Substitute in solutions
 $y = C1 e^{3x} + C2 e^{4x}$

Maple trace

```
`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
<- constant coefficients successful`
```

Maple dsolve solution

Solving time : 0.003 (sec)

Leaf size : 17

```
dsolve(diff(diff(y(x),x),x)+12*y(x) = 7*diff(y(x),x),
        y(x),singsol=all)
```

$$y = c_1 e^{3x} + c_2 e^{4x}$$

Mathematica DSolve solution

Solving time : 0.014 (sec)

Leaf size : 20

```
DSolve[{D[y[x],{x,2}]+12*y[x]==7*D[y[x],x],{}}],
        y[x],x,IncludeSingularSolutions->True]
```

$$y(x) \rightarrow e^{3x}(c_2 e^x + c_1)$$

2.8.4 problem 4

Solved as second order linear constant coeff ode	625
Solved as second order can be made integrable	626
Solved as second order ode using Kovacic algorithm	629
Solved as second order ode adjoint method	632
Maple step by step solution	635
Maple trace	635
Maple dsolve solution	636
Mathematica DSolve solution	636

Internal problem ID [18259]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter VII. Linear equations of order higher than the first. section 56. Problems at page 163

Problem number : 4

Date solved : Monday, December 23, 2024 at 09:24:15 PM

CAS classification : [[_2nd_order, _missing_x]]

Solve

$$r'' - a^2r = 0$$

Solved as second order linear constant coeff ode

Time used: 0.129 (sec)

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ar''(\phi) + Br'(\phi) + Cr(\phi) = 0$$

Where in the above $A = 1, B = 0, C = -a^2$. Let the solution be $r = e^{\lambda\phi}$. Substituting this into the ODE gives

$$\lambda^2 e^{\phi\lambda} - a^2 e^{\phi\lambda} = 0 \tag{1}$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda\phi}$ gives

$$-a^2 + \lambda^2 = 0 \tag{2}$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = 0, C = -a^2$ into the above gives

$$\begin{aligned}\lambda_{1,2} &= \frac{0}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{0^2 - (4)(1)(-a^2)} \\ &= \pm \sqrt{a^2}\end{aligned}$$

Hence

$$\lambda_1 = +\sqrt{a^2}$$

$$\lambda_2 = -\sqrt{a^2}$$

Which simplifies to

$$\lambda_1 = \sqrt{a^2}$$

$$\lambda_2 = -\sqrt{a^2}$$

Since roots are real and distinct, then the solution is

$$r = c_1 e^{\lambda_1 \phi} + c_2 e^{\lambda_2 \phi}$$

$$r = c_1 e^{(\sqrt{a^2})\phi} + c_2 e^{(-\sqrt{a^2})\phi}$$

Or

$$r = c_1 e^{\phi\sqrt{a^2}} + c_2 e^{-\phi\sqrt{a^2}}$$

Will add steps showing solving for IC soon.

Summary of solutions found

$$r = c_1 e^{\phi\sqrt{a^2}} + c_2 e^{-\phi\sqrt{a^2}}$$

Solved as second order can be made integrable

Time used: 5.807 (sec)

Multiplying the ode by r' gives

$$r' r'' - a^2 r' r = 0$$

Integrating the above w.r.t ϕ gives

$$\begin{aligned}\int (r' r'' - a^2 r' r) d\phi &= 0 \\ \frac{r'^2}{2} - \frac{a^2 r^2}{2} &= c_1\end{aligned}$$

Which is now solved for r . Solving for the derivative gives these ODE's to solve

$$r' = \sqrt{a^2 r^2 + 2c_1} \quad (1)$$

$$r' = -\sqrt{a^2 r^2 + 2c_1} \quad (2)$$

Now each of the above is solved separately.

Solving Eq. (1)

Integrating gives

$$\int \frac{1}{\sqrt{a^2 r^2 + 2c_1}} dr = d\phi$$

$$\frac{\ln\left(\frac{a^2 r}{\sqrt{a^2}} + \sqrt{a^2 r^2 + 2c_1}\right)}{\sqrt{a^2}} = \phi + c_2$$

Singular solutions are found by solving

$$\sqrt{a^2 r^2 + 2c_1} = 0$$

for r . This is because we had to divide by this in the above step. This gives the following singular solution(s), which also have to satisfy the given ODE.

$$r = \frac{\sqrt{-2c_1}}{a}$$

$$r = -\frac{\sqrt{-2c_1}}{a}$$

Solving for r gives

$$r = \frac{\sqrt{-2c_1}}{a}$$

$$r = -\frac{\sqrt{-2c_1}}{a}$$

$$r = -\frac{\sqrt{a^2} \left(-e^{2c_2\sqrt{a^2}+2\phi\sqrt{a^2}} + 2c_1 \right) e^{-c_2\sqrt{a^2}-\phi\sqrt{a^2}}}{2a^2}$$

Solving Eq. (2)

Integrating gives

$$\int -\frac{1}{\sqrt{a^2r^2 + 2c_1}} dr = d\phi$$

$$-\frac{\ln\left(\frac{a^2r}{\sqrt{a^2}} + \sqrt{a^2r^2 + 2c_1}\right)}{\sqrt{a^2}} = \phi + c_3$$

Singular solutions are found by solving

$$-\sqrt{a^2r^2 + 2c_1} = 0$$

for r . This is because we had to divide by this in the above step. This gives the following singular solution(s), which also have to satisfy the given ODE.

$$r = \frac{\sqrt{-2c_1}}{a}$$

$$r = -\frac{\sqrt{-2c_1}}{a}$$

Solving for r gives

$$r = \frac{\sqrt{-2c_1}}{a}$$

$$r = -\frac{\sqrt{-2c_1}}{a}$$

$$r = -\frac{\sqrt{a^2} \left(-e^{-2c_3\sqrt{a^2} - 2\phi\sqrt{a^2}} + 2c_1 \right) e^{c_3\sqrt{a^2} + \phi\sqrt{a^2}}}{2a^2}$$

Will add steps showing solving for IC soon.

The solution

$$r = \frac{\sqrt{-2c_1}}{a}$$

was found not to satisfy the ode or the IC. Hence it is removed. The solution

$$r = -\frac{\sqrt{-2c_1}}{a}$$

was found not to satisfy the ode or the IC. Hence it is removed.

Summary of solutions found

$$r = -\frac{\sqrt{a^2} \left(-e^{2c_2\sqrt{a^2}+2\phi\sqrt{a^2}} + 2c_1 \right) e^{-c_2\sqrt{a^2}-\phi\sqrt{a^2}}}{2a^2}$$

$$r = -\frac{\sqrt{a^2} \left(-e^{-2c_3\sqrt{a^2}-2\phi\sqrt{a^2}} + 2c_1 \right) e^{c_3\sqrt{a^2}+\phi\sqrt{a^2}}}{2a^2}$$

Solved as second order ode using Kovacic algorithm

Time used: 0.084 (sec)

Writing the ode as

$$r'' - a^2r = 0 \tag{1}$$

$$Ar'' + Br' + Cr = 0 \tag{2}$$

Comparing (1) and (2) shows that

$$\begin{aligned} A &= 1 \\ B &= 0 \\ C &= -a^2 \end{aligned} \tag{3}$$

Applying the Liouville transformation on the dependent variable gives

$$z(\phi) = r e^{\int \frac{B}{2A} d\phi}$$

Then (2) becomes

$$z''(\phi) = rz(\phi) \tag{4}$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \tag{5}$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{a^2}{1} \tag{6}$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= a^2 \\ t &= 1 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(\phi) = (a^2) z(\phi) \quad (7)$$

Equation (7) is now solved. After finding $z(\phi)$ then r is found using the inverse transformation

$$r = z(\phi) e^{-\int \frac{B}{2A} d\phi}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 2.53: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - 0 \\ &= 0 \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = a^2$ is not a function of ϕ , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(\phi) = e^{\phi\sqrt{a^2}}$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in r is found from

$$r_1 = z_1 e^{\int -\frac{1}{2} \frac{B}{A} d\phi}$$

Since $B = 0$ then the above reduces to

$$\begin{aligned} r_1 &= z_1 \\ &= e^{\phi\sqrt{a^2}} \end{aligned}$$

Which simplifies to

$$r_1 = e^{\phi\sqrt{a^2}}$$

The second solution r_2 to the original ode is found using reduction of order

$$r_2 = r_1 \int \frac{e^{\int -\frac{B}{A} d\phi}}{r_1^2} d\phi$$

Since $B = 0$ then the above becomes

$$\begin{aligned} r_2 &= r_1 \int \frac{1}{r_1^2} d\phi \\ &= e^{\phi\sqrt{a^2}} \int \frac{1}{e^{2\phi\sqrt{a^2}}} d\phi \\ &= e^{\phi\sqrt{a^2}} \left(-\frac{\sqrt{a^2} e^{-2\phi\sqrt{a^2}}}{2a^2} \right) \end{aligned}$$

Therefore the solution is

$$\begin{aligned} r &= c_1 r_1 + c_2 r_2 \\ &= c_1 \left(e^{\phi\sqrt{a^2}} \right) + c_2 \left(e^{\phi\sqrt{a^2}} \left(-\frac{\sqrt{a^2} e^{-2\phi\sqrt{a^2}}}{2a^2} \right) \right) \end{aligned}$$

Will add steps showing solving for IC soon.

Summary of solutions found

$$r = c_1 e^{\phi\sqrt{a^2}} - \frac{c_2 \operatorname{csgn}(a) e^{-\phi \operatorname{csgn}(a)a}}{2a}$$

Solved as second order ode adjoint method

Time used: 0.707 (sec)

In normal form the ode

$$r'' - a^2 r = 0 \quad (1)$$

Becomes

$$r'' + p(\phi) r' + q(\phi) r = r(\phi) \quad (2)$$

Where

$$\begin{aligned} p(\phi) &= 0 \\ q(\phi) &= -a^2 \\ r(\phi) &= 0 \end{aligned}$$

The Lagrange adjoint ode is given by

$$\begin{aligned} \xi'' - (\xi p)' + \xi q &= 0 \\ \xi'' - (0)' + (-a^2 \xi(\phi)) &= 0 \\ \xi''(\phi) - a^2 \xi(\phi) &= 0 \end{aligned}$$

Which is solved for $\xi(\phi)$. This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$A\xi''(\phi) + B\xi'(\phi) + C\xi(\phi) = 0$$

Where in the above $A = 1, B = 0, C = -a^2$. Let the solution be $\xi = e^{\lambda\phi}$. Substituting this into the ODE gives

$$\lambda^2 e^{\lambda\phi} - a^2 e^{\lambda\phi} = 0 \quad (1)$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda\phi}$ gives

$$-a^2 + \lambda^2 = 0 \quad (2)$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = 0, C = -a^2$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{0}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{0^2 - (4)(1)(-a^2)} \\ &= \pm \sqrt{a^2} \end{aligned}$$

Hence

$$\lambda_1 = +\sqrt{a^2}$$

$$\lambda_2 = -\sqrt{a^2}$$

Which simplifies to

$$\lambda_1 = \sqrt{a^2}$$

$$\lambda_2 = -\sqrt{a^2}$$

Since roots are real and distinct, then the solution is

$$\xi = c_1 e^{\lambda_1 \phi} + c_2 e^{\lambda_2 \phi}$$

$$\xi = c_1 e^{(\sqrt{a^2})\phi} + c_2 e^{(-\sqrt{a^2})\phi}$$

Or

$$\xi = c_1 e^{\phi\sqrt{a^2}} + c_2 e^{-\phi\sqrt{a^2}}$$

Will add steps showing solving for IC soon.

The original ode now reduces to first order ode

$$\begin{aligned} \xi(\phi) r' - r \xi'(\phi) + \xi(\phi) p(\phi) r &= \int \xi(\phi) r(\phi) d\phi \\ r' + r \left(p(\phi) - \frac{\xi'(\phi)}{\xi(\phi)} \right) &= \frac{\int \xi(\phi) r(\phi) d\phi}{\xi(\phi)} \end{aligned}$$

Or

$$r' - \frac{r \left(c_1 \sqrt{a^2} e^{\phi\sqrt{a^2}} - c_2 \sqrt{a^2} e^{-\phi\sqrt{a^2}} \right)}{c_1 e^{\phi\sqrt{a^2}} + c_2 e^{-\phi\sqrt{a^2}}} = 0$$

Which is now a first order ode. This is now solved for r . In canonical form a linear first order is

$$r' + q(\phi)r = p(\phi)$$

Comparing the above to the given ode shows that

$$q(\phi) = -\frac{\operatorname{csgn}(a) a (c_1 e^{2\phi \operatorname{csgn}(a)a} - c_2)}{c_1 e^{2\phi \operatorname{csgn}(a)a} + c_2}$$

$$p(\phi) = 0$$

The integrating factor μ is

$$\begin{aligned} \mu &= e^{\int q d\phi} \\ &= e^{\int -\frac{\operatorname{csgn}(a)a(c_1 e^{2\phi \operatorname{csgn}(a)a} - c_2)}{c_1 e^{2\phi \operatorname{csgn}(a)a} + c_2} d\phi} \\ &= e^{\frac{\ln(e^{2\phi \operatorname{csgn}(a)a})}{2} - \ln(c_1 e^{2\phi \operatorname{csgn}(a)a} + c_2)} \end{aligned}$$

The ode becomes

$$\frac{d}{d\phi} \mu r = 0$$

$$\frac{d}{d\phi} \left(r e^{\frac{\ln(e^{2\phi \operatorname{csgn}(a)a})}{2} - \ln(c_1 e^{2\phi \operatorname{csgn}(a)a} + c_2)} \right) = 0$$

Integrating gives

$$\begin{aligned} r e^{\frac{\ln(e^{2\phi \operatorname{csgn}(a)a})}{2} - \ln(c_1 e^{2\phi \operatorname{csgn}(a)a} + c_2)} &= \int 0 d\phi + c_3 \\ &= c_3 \end{aligned}$$

Dividing throughout by the integrating factor $e^{\frac{\ln(e^{2\phi \operatorname{csgn}(a)a})}{2} - \ln(c_1 e^{2\phi \operatorname{csgn}(a)a} + c_2)}$ gives the final solution

$$r = \frac{(c_1 e^{2\phi \operatorname{csgn}(a)a} + c_2) c_3}{\sqrt{e^{2\phi \operatorname{csgn}(a)a}}}$$

Hence, the solution found using Lagrange adjoint equation method is

$$r = \frac{(c_1 e^{2\phi \operatorname{csgn}(a)a} + c_2) c_3}{\sqrt{e^{2\phi \operatorname{csgn}(a)a}}}$$

The constants can be merged to give

$$r = \frac{c_1 e^{2\phi \operatorname{csgn}(a)a} + c_2}{\sqrt{e^{2\phi \operatorname{csgn}(a)a}}}$$

Will add steps showing solving for IC soon.

Summary of solutions found

$$r = \frac{c_1 e^{2\phi \operatorname{csgn}(a)a} + c_2}{\sqrt{e^{2\phi \operatorname{csgn}(a)a}}}$$

Maple step by step solution

Let's solve

$$r'' - a^2 r = 0$$

- Highest derivative means the order of the ODE is 2
- r''
- Characteristic polynomial of ODE
 $-a^2 + s^2 = 0$
- Factor the characteristic polynomial
 $-(a - s)(a + s) = 0$
- Roots of the characteristic polynomial
 $s = (a, -a)$
- 1st solution of the ODE
 $r_1(\phi) = e^{\phi a}$
- 2nd solution of the ODE
 $r_2(\phi) = e^{-\phi a}$
- General solution of the ODE
 $r = C1 r_1(\phi) + C2 r_2(\phi)$
- Substitute in solutions
 $r = C1 e^{\phi a} + C2 e^{-\phi a}$

Maple trace

```
`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
<- constant coefficients successful`
```

Maple dsolve solution

Solving time : 0.003 (sec)

Leaf size : 18

```
dsolve(diff(diff(r(phi),phi),phi)-a^2*r(phi) = 0,  
        r(phi),singsol=all)
```

$$r = c_1 e^{\phi a} + c_2 e^{-\phi a}$$

Mathematica DSolve solution

Solving time : 0.013 (sec)

Leaf size : 23

```
DSolve[{D[r[phi],{phi,2}]-a^2*r[phi]==0,{}},  
        r[phi],phi,IncludeSingularSolutions->True]
```

$$r(\phi) \rightarrow c_1 e^{a\phi} + c_2 e^{-a\phi}$$

2.8.5 problem 5

Solved as higher order constant coeff ode	637
Maple step by step solution	638
Maple trace	638
Maple dsolve solution	638
Mathematica DSolve solution	638

Internal problem ID [18260]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter VII. Linear equations of order higher than the first. section 56. Problems at page 163

Problem number : 5

Date solved : Monday, December 23, 2024 at 09:24:23 PM

CAS classification : [[_high_order, _missing_x]]

Solve

$$y'''' - a^4 y = 0$$

Solved as higher order constant coeff ode

Time used: 0.139 (sec)

The characteristic equation is

$$-a^4 + \lambda^4 = 0$$

The roots of the above equation are

$$\lambda_1 = a$$

$$\lambda_2 = -a$$

$$\lambda_3 = ia$$

$$\lambda_4 = -ia$$

Therefore the homogeneous solution is

$$y_h(x) = e^{ax} c_1 + e^{-ax} c_2 + e^{iax} c_3 + e^{-iax} c_4$$

The fundamental set of solutions for the homogeneous solution are the following

$$\begin{aligned}y_1 &= e^{ax} \\y_2 &= e^{-ax} \\y_3 &= e^{iax} \\y_4 &= e^{-iax}\end{aligned}$$

Maple step by step solution

Maple trace

```
`Methods for high order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
<- constant coefficients successful`
```

Maple dsolve solution

Solving time : 0.007 (sec)
Leaf size : 30

```
dsolve(diff(diff(diff(diff(y(x),x),x),x),x)-a^4*y(x) = 0,
        y(x),singsol=all)
```

$$y = e^{ax}c_1 + e^{-ax}c_2 + c_3 \sin(ax) + c_4 \cos(ax)$$

Mathematica DSolve solution

Solving time : 0.012 (sec)
Leaf size : 53

```
DSolve[{D[y[x],{x,4}]-a^2*y[x]==0,{}},
        y[x],x,IncludeSingularSolutions->True]
```

$$y(x) \rightarrow c_2 e^{-\sqrt{ax}} + c_4 e^{\sqrt{ax}} + c_1 \cos(\sqrt{ax}) + c_3 \sin(\sqrt{ax})$$

2.8.6 problem 6

Solved as second order linear constant coeff ode	639
Solved as second order ode using Kovacic algorithm	642
Solved as second order ode adjoint method	647
Maple step by step solution	653
Maple trace	654
Maple dsolve solution	654
Mathematica DSolve solution	654

Internal problem ID [18261]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter VII. Linear equations of order higher than the first. section 56. Problems at page 163

Problem number : 6

Date solved : Monday, December 23, 2024 at 09:46:35 PM

CAS classification : [[_2nd_order, _with_linear_symmetries]]

Solve

$$v'' - 6v' + 13v = e^{-2u}$$

Solved as second order linear constant coeff ode

Time used: 0.191 (sec)

This is second order non-homogeneous ODE. In standard form the ODE is

$$Av''(u) + Bv'(u) + Cv(u) = f(u)$$

Where $A = 1, B = -6, C = 13, f(u) = e^{-2u}$. Let the solution be

$$v = v_h + v_p$$

Where v_h is the solution to the homogeneous ODE $Av''(u) + Bv'(u) + Cv(u) = 0$, and v_p is a particular solution to the non-homogeneous ODE $Av''(u) + Bv'(u) + Cv(u) = f(u)$. v_h is the solution to

$$v'' - 6v' + 13v = 0$$

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Av''(u) + Bv'(u) + Cv(u) = 0$$

Where in the above $A = 1, B = -6, C = 13$. Let the solution be $v = e^{\lambda u}$. Substituting this into the ODE gives

$$\lambda^2 e^{u\lambda} - 6\lambda e^{u\lambda} + 13 e^{u\lambda} = 0 \quad (1)$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda u}$ gives

$$\lambda^2 - 6\lambda + 13 = 0 \quad (2)$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = -6, C = 13$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{6}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{-6^2 - (4)(1)(13)} \\ &= 3 \pm 2i \end{aligned}$$

Hence

$$\lambda_1 = 3 + 2i$$

$$\lambda_2 = 3 - 2i$$

Which simplifies to

$$\lambda_1 = 3 + 2i$$

$$\lambda_2 = 3 - 2i$$

Since roots are complex conjugate of each others, then let the roots be

$$\lambda_{1,2} = \alpha \pm i\beta$$

Where $\alpha = 3$ and $\beta = 2$. Therefore the final solution, when using Euler relation, can be written as

$$v = e^{\alpha u} (c_1 \cos(\beta u) + c_2 \sin(\beta u))$$

Which becomes

$$v = e^{3u} (c_1 \cos(2u) + c_2 \sin(2u))$$

Therefore the homogeneous solution v_h is

$$v_h = e^{3u} (c_1 \cos(2u) + c_2 \sin(2u))$$

The particular solution is now found using the method of undetermined coefficients. Looking at the RHS of the ode, which is

$$e^{-2u}$$

Shows that the corresponding undetermined set of the basis functions (UC_set) for the trial solution is

$$[e^{-2u}]$$

While the set of the basis functions for the homogeneous solution found earlier is

$$\{e^{3u} \cos(2u), e^{3u} \sin(2u)\}$$

Since there is no duplication between the basis function in the UC_set and the basis functions of the homogeneous solution, the trial solution is a linear combination of all the basis in the UC_set.

$$v_p = A_1 e^{-2u}$$

The unknowns $\{A_1\}$ are found by substituting the above trial solution v_p into the ODE and comparing coefficients. Substituting the trial solution into the ODE and simplifying gives

$$29A_1 e^{-2u} = e^{-2u}$$

Solving for the unknowns by comparing coefficients results in

$$\left[A_1 = \frac{1}{29} \right]$$

Substituting the above back in the above trial solution v_p , gives the particular solution

$$v_p = \frac{e^{-2u}}{29}$$

Therefore the general solution is

$$\begin{aligned} v &= v_h + v_p \\ &= (e^{3u}(c_1 \cos(2u) + c_2 \sin(2u))) + \left(\frac{e^{-2u}}{29} \right) \end{aligned}$$

Will add steps showing solving for IC soon.

Summary of solutions found

$$v = \frac{e^{-2u}}{29} + e^{3u}(c_1 \cos(2u) + c_2 \sin(2u))$$

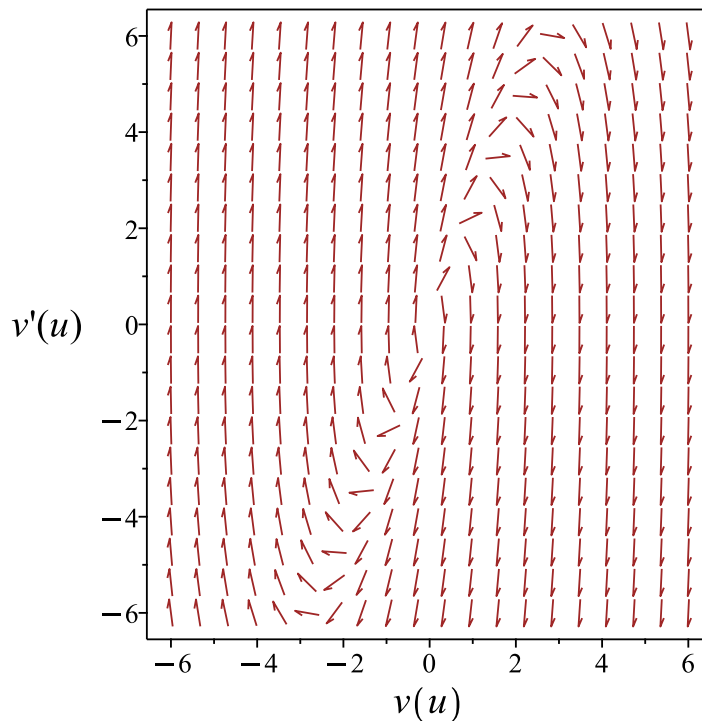


Figure 2.97: Slope field plot

$$v'' - 6v' + 13v = e^{-2u}$$

Solved as second order ode using Kovacic algorithm

Time used: 0.201 (sec)

Writing the ode as

$$v'' - 6v' + 13v = 0 \tag{1}$$

$$Av'' + Bv' + Cv = 0 \tag{2}$$

Comparing (1) and (2) shows that

$$A = 1$$

$$B = -6 \tag{3}$$

$$C = 13$$

Applying the Liouville transformation on the dependent variable gives

$$z(u) = ve^{\int \frac{B}{2A} du}$$

Then (2) becomes

$$z''(u) = rz(u) \quad (4)$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \quad (5)$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{-4}{1} \quad (6)$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= -4 \\ t &= 1 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(u) = -4z(u) \quad (7)$$

Equation (7) is now solved. After finding $z(u)$ then v is found using the inverse transformation

$$v = z(u) e^{-\int \frac{B}{2A} du}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 2.55: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - 0 \\ &= 0 \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = -4$ is not a function of u , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(u) = \cos(2u)$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in v is found from

$$\begin{aligned} v_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} du} \\ &= z_1 e^{-\int \frac{1}{2} \frac{-6}{1} du} \\ &= z_1 e^{3u} \\ &= z_1 (e^{3u}) \end{aligned}$$

Which simplifies to

$$v_1 = e^{3u} \cos(2u)$$

The second solution v_2 to the original ode is found using reduction of order

$$v_2 = v_1 \int \frac{e^{\int -\frac{B}{A} du}}{v_1^2} du$$

Substituting gives

$$\begin{aligned} v_2 &= v_1 \int \frac{e^{\int -\frac{-6}{1} du}}{(v_1)^2} du \\ &= v_1 \int \frac{e^{6u}}{(v_1)^2} du \\ &= v_1 \left(\frac{\tan(2u)}{2} \right) \end{aligned}$$

Therefore the solution is

$$\begin{aligned} v &= c_1 v_1 + c_2 v_2 \\ &= c_1 (e^{3u} \cos(2u)) + c_2 \left(e^{3u} \cos(2u) \left(\frac{\tan(2u)}{2} \right) \right) \end{aligned}$$

This is second order nonhomogeneous ODE. Let the solution be

$$v = v_h + v_p$$

Where v_h is the solution to the homogeneous ODE $Av''(u) + Bv'(u) + Cv(u) = 0$, and v_p is a particular solution to the nonhomogeneous ODE $Av''(u) + Bv'(u) + Cv(u) = f(u)$. v_h is the solution to

$$v'' - 6v' + 13v = 0$$

The homogeneous solution is found using the Kovacic algorithm which results in

$$v_h = c_1 e^{3u} \cos(2u) + \frac{c_2 e^{3u} \sin(2u)}{2}$$

The particular solution is now found using the method of undetermined coefficients. Looking at the RHS of the ode, which is

$$e^{-2u}$$

Shows that the corresponding undetermined set of the basis functions (UC_set) for the trial solution is

$$[\{e^{-2u}\}]$$

While the set of the basis functions for the homogeneous solution found earlier is

$$\left\{ e^{3u} \cos(2u), \frac{e^{3u} \sin(2u)}{2} \right\}$$

Since there is no duplication between the basis function in the UC_set and the basis functions of the homogeneous solution, the trial solution is a linear combination of all the basis in the UC_set.

$$v_p = A_1 e^{-2u}$$

The unknowns $\{A_1\}$ are found by substituting the above trial solution v_p into the ODE and comparing coefficients. Substituting the trial solution into the ODE and simplifying gives

$$29A_1e^{-2u} = e^{-2u}$$

Solving for the unknowns by comparing coefficients results in

$$\left[A_1 = \frac{1}{29} \right]$$

Substituting the above back in the above trial solution v_p , gives the particular solution

$$v_p = \frac{e^{-2u}}{29}$$

Therefore the general solution is

$$\begin{aligned} v &= v_h + v_p \\ &= \left(c_1 e^{3u} \cos(2u) + \frac{c_2 e^{3u} \sin(2u)}{2} \right) + \left(\frac{e^{-2u}}{29} \right) \end{aligned}$$

Will add steps showing solving for IC soon.

Summary of solutions found

$$v = c_1 e^{3u} \cos(2u) + \frac{c_2 e^{3u} \sin(2u)}{2} + \frac{e^{-2u}}{29}$$

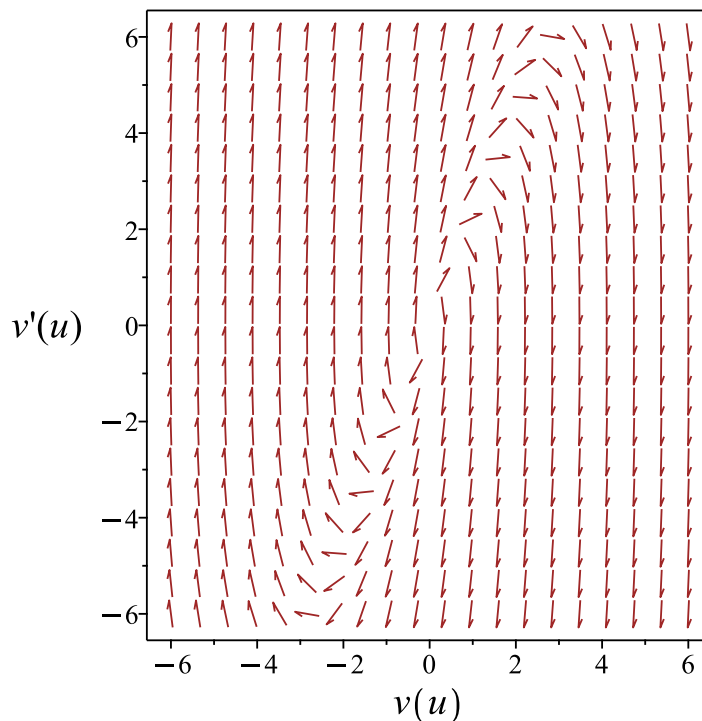


Figure 2.98: Slope field plot
 $v'' - 6v' + 13v = e^{-2u}$

Solved as second order ode adjoint method

Time used: 11.164 (sec)

In normal form the ode

$$v'' - 6v' + 13v = e^{-2u} \quad (1)$$

Becomes

$$v'' + p(u)v' + q(u)v = r(u) \quad (2)$$

Where

$$\begin{aligned} p(u) &= -6 \\ q(u) &= 13 \\ r(u) &= e^{-2u} \end{aligned}$$

The Lagrange adjoint ode is given by

$$\begin{aligned} \xi'' - (\xi p)' + \xi q &= 0 \\ \xi'' - (-6\xi(u))' + (13\xi(u)) &= 0 \\ \xi''(u) + 6\xi'(u) + 13\xi(u) &= 0 \end{aligned}$$

Which is solved for $\xi(u)$. This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$A\xi''(u) + B\xi'(u) + C\xi(u) = 0$$

Where in the above $A = 1, B = 6, C = 13$. Let the solution be $\xi = e^{\lambda u}$. Substituting this into the ODE gives

$$\lambda^2 e^{u\lambda} + 6\lambda e^{u\lambda} + 13e^{u\lambda} = 0 \quad (1)$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda u}$ gives

$$\lambda^2 + 6\lambda + 13 = 0 \quad (2)$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = 6, C = 13$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{-6}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{6^2 - (4)(1)(13)} \\ &= -3 \pm 2i \end{aligned}$$

Hence

$$\lambda_1 = -3 + 2i$$

$$\lambda_2 = -3 - 2i$$

Which simplifies to

$$\lambda_1 = -3 + 2i$$

$$\lambda_2 = -3 - 2i$$

Since roots are complex conjugate of each others, then let the roots be

$$\lambda_{1,2} = \alpha \pm i\beta$$

Where $\alpha = -3$ and $\beta = 2$. Therefore the final solution, when using Euler relation, can be written as

$$\xi = e^{\alpha u} (c_1 \cos(\beta u) + c_2 \sin(\beta u))$$

Which becomes

$$\xi = e^{-3u} (c_1 \cos(2u) + c_2 \sin(2u))$$

Will add steps showing solving for IC soon.

The original ode now reduces to first order ode

$$\begin{aligned}\xi(u) v' - v\xi'(u) + \xi(u) p(u) v &= \int \xi(u) r(u) du \\ v' + v \left(p(u) - \frac{\xi'(u)}{\xi(u)} \right) &= \frac{\int \xi(u) r(u) du}{\xi(u)}\end{aligned}$$

Or

$$v' + v \left(-6 - \frac{(-3e^{-3u}(c_1 \cos(2u) + c_2 \sin(2u)) + e^{-3u}(-2c_1 \sin(2u) + 2c_2 \cos(2u))) e^{3u}}{c_1 \cos(2u) + c_2 \sin(2u)} \right) = \frac{e^{3u} \left(2c_1 \left(\frac{-5}{2} \right) \right)}{c_1 \cos(2u) + c_2 \sin(2u)}$$

Which is now a first order ode. This is now solved for v . In canonical form a linear first order is

$$v' + q(u)v = p(u)$$

Comparing the above to the given ode shows that

$$\begin{aligned}q(u) &= -\frac{(3c_1 + 2c_2) \cos(2u) - 2 \sin(2u) \left(c_1 - \frac{3c_2}{2} \right)}{c_1 \cos(2u) + c_2 \sin(2u)} \\ p(u) &= -\frac{5e^{-2u} \left(\left(c_1 + \frac{2c_2}{5} \right) \cos(2u) - \frac{2 \sin(2u) \left(c_1 - \frac{5c_2}{2} \right)}{5} \right)}{29c_1 \cos(2u) + 29c_2 \sin(2u)}\end{aligned}$$

The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int q du} \\ &= e^{\int -\frac{(3c_1 + 2c_2) \cos(2u) - 2 \sin(2u) \left(c_1 - \frac{3c_2}{2} \right)}{c_1 \cos(2u) + c_2 \sin(2u)} du} \\ &= e^{-\ln(c_1 + c_2 \tan(2u)) + \frac{\ln(\tan(2u)^2 + 1)}{2} - 3u}\end{aligned}$$

The ode becomes

$$\frac{d}{du}(\mu v) = \mu p$$

$$\frac{d}{du}(\mu v) = (\mu) \left(-\frac{5e^{-2u} \left(\left(c_1 + \frac{2c_2}{5} \right) \cos(2u) - \frac{2 \sin(2u) \left(c_1 - \frac{5c_2}{2} \right)}{5} \right)}{29c_1 \cos(2u) + 29c_2 \sin(2u)} \right)$$

$$\begin{aligned} & \frac{d}{du} \left(v e^{-\ln(c_1 + c_2 \tan(2u)) + \frac{\ln(\tan(2u)^2 + 1)}{2} - 3u} \right) \\ &= \left(e^{-\ln(c_1 + c_2 \tan(2u)) + \frac{\ln(\tan(2u)^2 + 1)}{2} - 3u} \right) \left(-\frac{5 e^{-2u} \left(\left(c_1 + \frac{2c_2}{5} \right) \cos(2u) - \frac{2 \sin(2u) \left(c_1 - \frac{5c_2}{2} \right)}{5} \right)}{29c_1 \cos(2u) + 29c_2 \sin(2u)} \right) \end{aligned}$$

$$\begin{aligned} & d \left(v e^{-\ln(c_1 + c_2 \tan(2u)) + \frac{\ln(\tan(2u)^2 + 1)}{2} - 3u} \right) \\ &= \left(-\frac{5 e^{-2u} \left(\left(c_1 + \frac{2c_2}{5} \right) \cos(2u) - \frac{2 \sin(2u) \left(c_1 - \frac{5c_2}{2} \right)}{5} \right) e^{-\ln(c_1 + c_2 \tan(2u)) + \frac{\ln(\tan(2u)^2 + 1)}{2} - 3u}}{29c_1 \cos(2u) + 29c_2 \sin(2u)} \right) du \end{aligned}$$

Integrating gives

$$\begin{aligned} v e^{-\ln(c_1 + c_2 \tan(2u)) + \frac{\ln(\tan(2u)^2 + 1)}{2} - 3u} &= \int -\frac{5 e^{-2u} \left(\left(c_1 + \frac{2c_2}{5} \right) \cos(2u) - \frac{2 \sin(2u) \left(c_1 - \frac{5c_2}{2} \right)}{5} \right) e^{-\ln(c_1 + c_2 \tan(2u)) + \frac{\ln(\tan(2u)^2 + 1)}{2} - 3u}}{29c_1 \cos(2u) + 29c_2 \sin(2u)} \\ &= \int -\frac{5 e^{-2u} \left(\left(c_1 + \frac{2c_2}{5} \right) \cos(2u) - \frac{2 \sin(2u) \left(c_1 - \frac{5c_2}{2} \right)}{5} \right) e^{-\ln(c_1 + c_2 \tan(2u)) + \frac{\ln(\tan(2u)^2 + 1)}{2} - 3u}}{29c_1 \cos(2u) + 29c_2 \sin(2u)} \end{aligned}$$

Dividing throughout by the integrating factor $e^{-\ln(c_1 + c_2 \tan(2u)) + \frac{\ln(\tan(2u)^2 + 1)}{2} - 3u}$ gives the final solution

$$v = (c_1 + c_2 \tan(2u)) e^{\ln\left(\frac{1}{\sqrt{\tan(2u)^2 + 1}}\right) + 3u} \left(\int -\frac{5 e^{-2u} \left(\left(c_1 + \frac{2c_2}{5} \right) \cos(2u) - \frac{2 \sin(2u) \left(c_1 - \frac{5c_2}{2} \right)}{5} \right) e^{-\ln(c_1 + c_2 \tan(2u)) + \frac{\ln(\tan(2u)^2 + 1)}{2} - 3u}}{29c_1 \cos(2u) + 29c_2 \sin(2u)} \right)$$

Hence, the solution found using Lagrange adjoint equation method is

$$v = (c_1 + c_2 \tan(2u)) e^{\ln\left(\frac{1}{\sqrt{\tan(2u)^2+1}}\right)+3u} \left(\int \frac{5 e^{-2u} \left((c_1 + \frac{2c_2}{5}) \cos(2u) - \frac{2 \sin(2u) (c_1 - \frac{5c_2}{2})}{5} \right) e^{-\ln(c_1+c_2 \tan(2u))+\frac{\ln(\tan(2u)^2+1)}{2}-3u}}{29c_1 \cos(2u) + 29c_2 \sin(2u)} du + c_3 \right)$$

The constants can be merged to give

$$v = (c_1 + c_2 \tan(2u)) e^{\ln\left(\frac{1}{\sqrt{\tan(2u)^2+1}}\right)+3u} \left(\int \frac{5 e^{-2u} \left((c_1 + \frac{2c_2}{5}) \cos(2u) - \frac{2 \sin(2u) (c_1 - \frac{5c_2}{2})}{5} \right) e^{-\ln(c_1+c_2 \tan(2u))+\frac{\ln(\tan(2u)^2+1)}{2}-3u}}{29c_1 \cos(2u) + 29c_2 \sin(2u)} du + 1 \right)$$

Will add steps showing solving for IC soon.

Summary of solutions found

$$v = (c_1 + c_2 \tan(2u)) e^{\ln\left(\frac{1}{\sqrt{\tan(2u)^2+1}}\right) + 3u} \left(\int \frac{5 e^{-2u} \left((c_1 + \frac{2c_2}{5}) \cos(2u) - \frac{2 \sin(2u)(c_1 - \frac{5c_2}{2})}{5} \right) e^{-\ln(c_1 + c_2 \tan(2u)) + \frac{\ln(\tan(2u)^2+1)}{2} - 3u}}{29c_1 \cos(2u) + 29c_2 \sin(2u)} du + 1 \right)$$

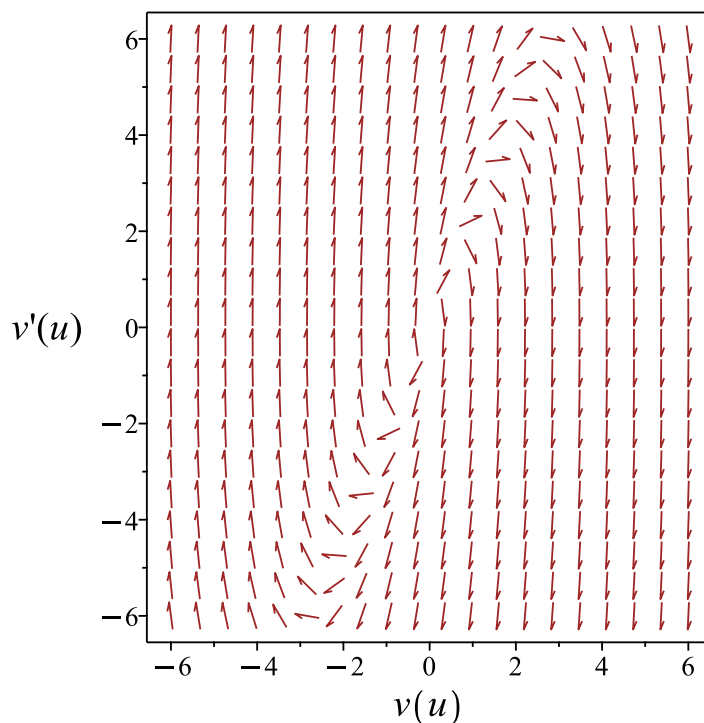


Figure 2.99: Slope field plot
 $v'' - 6v' + 13v = e^{-2u}$

Maple step by step solution

Let's solve

$$v'' - 6v' + 13v = e^{-2u}$$

- Highest derivative means the order of the ODE is 2

$$v''$$

- Characteristic polynomial of homogeneous ODE

$$r^2 - 6r + 13 = 0$$

- Use quadratic formula to solve for r

$$r = \frac{6 \pm (\sqrt{-16})}{2}$$

- Roots of the characteristic polynomial

$$r = (3 - 2I, 3 + 2I)$$

- 1st solution of the homogeneous ODE

$$v_1(u) = e^{3u} \cos(2u)$$

- 2nd solution of the homogeneous ODE

$$v_2(u) = e^{3u} \sin(2u)$$

- General solution of the ODE

$$v = C1v_1(u) + C2v_2(u) + v_p(u)$$

- Substitute in solutions of the homogeneous ODE

$$v = C1 e^{3u} \cos(2u) + C2 e^{3u} \sin(2u) + v_p(u)$$

- Find a particular solution $v_p(u)$ of the ODE

- Use variation of parameters to find v_p here $f(u)$ is the forcing function

$$\left[v_p(u) = -v_1(u) \left(\int \frac{v_2(u)f(u)}{W(v_1(u),v_2(u))} du \right) + v_2(u) \left(\int \frac{v_1(u)f(u)}{W(v_1(u),v_2(u))} du \right), f(u) = e^{-2u} \right]$$

- Wronskian of solutions of the homogeneous equation

$$W(v_1(u), v_2(u)) = \begin{bmatrix} e^{3u} \cos(2u) & e^{3u} \sin(2u) \\ 3e^{3u} \cos(2u) - 2e^{3u} \sin(2u) & 3e^{3u} \sin(2u) + 2e^{3u} \cos(2u) \end{bmatrix}$$

- Compute Wronskian

$$W(v_1(u), v_2(u)) = 2e^{6u}$$

- Substitute functions into equation for $v_p(u)$

$$v_p(u) = \frac{e^{3u} (-\cos(2u) (\int \sin(2u)e^{-5u} du) + \sin(2u) (\int \cos(2u)e^{-5u} du))}{2}$$

- Compute integrals

$$v_p(u) = \frac{e^{-2u}}{29}$$

- Substitute particular solution into general solution to ODE

$$v = C2 e^{3u} \sin(2u) + C1 e^{3u} \cos(2u) + \frac{e^{-2u}}{29}$$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
trying high order exact linear fully integrable
trying differential order: 2; linear nonhomogeneous with symmetry [0,1]
trying a double symmetry of the form [xi=0, eta=F(x)]
-> Try solving first the homogeneous part of the ODE
    checking if the LODE has constant coefficients
    <- constant coefficients successful
<- solving first the homogeneous part of the ODE successful`

```

Maple dsolve solution

Solving time : 0.003 (sec)

Leaf size : 33

```

dsolve(diff(diff(v(u),u),u)-6*diff(v(u),u)+13*v(u) = exp(-2*u),
        v(u),singsol=all)

```

$$v = (c_1 \cos(2u) + c_2 \sin(2u)) e^{-2u} e^{5u} + \frac{e^{-2u}}{29}$$

Mathematica DSolve solution

Solving time : 0.116 (sec)

Leaf size : 39

```

DSolve[{D[v[u],{u,2}]-6*D[v[u],u]+13*v[u]==Exp[-2*u],{}},
        v[u],u,IncludeSingularSolutions->True]

```

$$v(u) \rightarrow \frac{e^{-2u}}{29} + c_2 e^{3u} \cos(2u) + c_1 e^{3u} \sin(2u)$$

2.8.7 problem 7

Solved as second order linear constant coeff ode	655
Solved as second order ode using Kovacic algorithm	658
Solved as second order ode adjoint method	663
Maple step by step solution	667
Maple trace	669
Maple dsolve solution	669
Mathematica DSolve solution	669

Internal problem ID [18262]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter VII. Linear equations of order higher than the first. section 56. Problems at page 163

Problem number : 7

Date solved : Monday, December 23, 2024 at 09:47:04 PM

CAS classification : [[_2nd_order, _linear, _nonhomogeneous]]

Solve

$$y'' + 4y' - y = \sin(t)$$

Solved as second order linear constant coeff ode

Time used: 0.151 (sec)

This is second order non-homogeneous ODE. In standard form the ODE is

$$Ay''(t) + By'(t) + Cy(t) = f(t)$$

Where $A = 1, B = 4, C = -1, f(t) = \sin(t)$. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = 0$, and y_p is a particular solution to the non-homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = f(t)$. y_h is the solution to

$$y'' + 4y' - y = 0$$

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(t) + By'(t) + Cy(t) = 0$$

Where in the above $A = 1, B = 4, C = -1$. Let the solution be $y = e^{\lambda t}$. Substituting this into the ODE gives

$$\lambda^2 e^{t\lambda} + 4\lambda e^{t\lambda} - e^{t\lambda} = 0 \quad (1)$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda t}$ gives

$$\lambda^2 + 4\lambda - 1 = 0 \quad (2)$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = 4, C = -1$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{-4}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{4^2 - (4)(1)(-1)} \\ &= -2 \pm \sqrt{5} \end{aligned}$$

Hence

$$\lambda_1 = -2 + \sqrt{5}$$

$$\lambda_2 = -2 - \sqrt{5}$$

Which simplifies to

$$\lambda_1 = -2 + \sqrt{5}$$

$$\lambda_2 = -2 - \sqrt{5}$$

Since roots are real and distinct, then the solution is

$$y = c_1 e^{\lambda_1 t} + c_2 e^{\lambda_2 t}$$

$$y = c_1 e^{(-2+\sqrt{5})t} + c_2 e^{(-2-\sqrt{5})t}$$

Or

$$y = c_1 e^{t(-2+\sqrt{5})} + c_2 e^{t(-2-\sqrt{5})}$$

Therefore the homogeneous solution y_h is

$$y_h = c_1 e^{t(-2+\sqrt{5})} + c_2 e^{t(-2-\sqrt{5})}$$

The particular solution is now found using the method of undetermined coefficients. Looking at the RHS of the ode, which is

$$\sin(t)$$

Shows that the corresponding undetermined set of the basis functions (UC_set) for the trial solution is

$$[\{\cos(t), \sin(t)\}]$$

While the set of the basis functions for the homogeneous solution found earlier is

$$\left\{ e^{t(-2-\sqrt{5})}, e^{t(-2+\sqrt{5})} \right\}$$

Since there is no duplication between the basis function in the UC_set and the basis functions of the homogeneous solution, the trial solution is a linear combination of all the basis in the UC_set.

$$y_p = A_1 \cos(t) + A_2 \sin(t)$$

The unknowns $\{A_1, A_2\}$ are found by substituting the above trial solution y_p into the ODE and comparing coefficients. Substituting the trial solution into the ODE and simplifying gives

$$-2A_1 \cos(t) - 2A_2 \sin(t) - 4A_1 \sin(t) + 4A_2 \cos(t) = \sin(t)$$

Solving for the unknowns by comparing coefficients results in

$$\left[A_1 = -\frac{1}{5}, A_2 = -\frac{1}{10} \right]$$

Substituting the above back in the above trial solution y_p , gives the particular solution

$$y_p = -\frac{\cos(t)}{5} - \frac{\sin(t)}{10}$$

Therefore the general solution is

$$\begin{aligned} y &= y_h + y_p \\ &= \left(c_1 e^{t(-2+\sqrt{5})} + c_2 e^{t(-2-\sqrt{5})} \right) + \left(-\frac{\cos(t)}{5} - \frac{\sin(t)}{10} \right) \end{aligned}$$

Will add steps showing solving for IC soon.

Summary of solutions found

$$y = -\frac{\cos(t)}{5} - \frac{\sin(t)}{10} + c_1 e^{t(-2+\sqrt{5})} + c_2 e^{t(-2-\sqrt{5})}$$

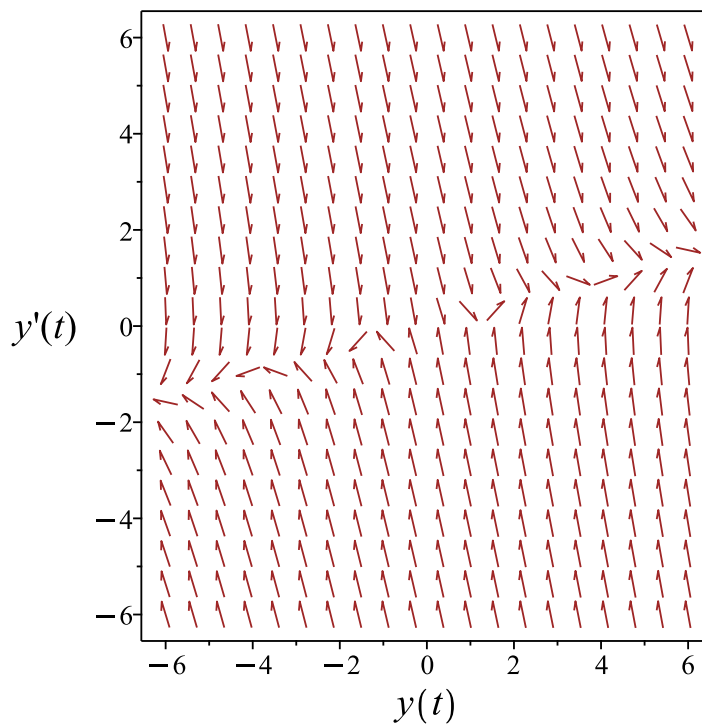


Figure 2.100: Slope field plot
 $y'' + 4y' - y = \sin(t)$

Solved as second order ode using Kovacic algorithm

Time used: 0.233 (sec)

Writing the ode as

$$y'' + 4y' - y = 0 \tag{1}$$

$$Ay'' + By' + Cy = 0 \tag{2}$$

Comparing (1) and (2) shows that

$$A = 1$$

$$B = 4 \tag{3}$$

$$C = -1$$

Applying the Liouville transformation on the dependent variable gives

$$z(t) = ye^{\int \frac{B}{2A} dt}$$

Then (2) becomes

$$z''(t) = rz(t) \tag{4}$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \tag{5}$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{5}{1} \tag{6}$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= 5 \\ t &= 1 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(t) = 5z(t) \tag{7}$$

Equation (7) is now solved. After finding $z(t)$ then y is found using the inverse transformation

$$y = z(t) e^{-\int \frac{B}{2A} dt}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 2.57: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - 0 \\ &= 0 \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = 5$ is not a function of t , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(t) = e^{-\sqrt{5}t}$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$\begin{aligned} y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dt} \\ &= z_1 e^{-\int \frac{1}{2} \frac{4}{1} dt} \\ &= z_1 e^{-2t} \\ &= z_1 (e^{-2t}) \end{aligned}$$

Which simplifies to

$$y_1 = e^{-t(2+\sqrt{5})}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dt}}{y_1^2} dt$$

Substituting gives

$$\begin{aligned} y_2 &= y_1 \int \frac{e^{\int -\frac{4}{1} dt}}{(y_1)^2} dt \\ &= y_1 \int \frac{e^{-4t}}{(y_1)^2} dt \\ &= y_1 \left(\frac{\sqrt{5} e^{-4t} e^{2t(2+\sqrt{5})}}{10} \right) \end{aligned}$$

Therefore the solution is

$$\begin{aligned} y &= c_1 y_1 + c_2 y_2 \\ &= c_1 \left(e^{-t(2+\sqrt{5})} \right) + c_2 \left(e^{-t(2+\sqrt{5})} \left(\frac{\sqrt{5} e^{-4t} e^{2t(2+\sqrt{5})}}{10} \right) \right) \end{aligned}$$

This is second order nonhomogeneous ODE. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE $Ay''(t) + By'(t) + Cy(t) = 0$, and y_p is a particular solution to the nonhomogeneous ODE $Ay''(t) + By'(t) + Cy(t) = f(t)$. y_h is the solution to

$$y'' + 4y' - y = 0$$

The homogeneous solution is found using the Kovacic algorithm which results in

$$y_h = c_1 e^{-t(2+\sqrt{5})} + \frac{c_2 \sqrt{5} e^{t(-2+\sqrt{5})}}{10}$$

The particular solution is now found using the method of undetermined coefficients. Looking at the RHS of the ode, which is

$$\sin(t)$$

Shows that the corresponding undetermined set of the basis functions (UC_set) for the trial solution is

$$[\{\cos(t), \sin(t)\}]$$

While the set of the basis functions for the homogeneous solution found earlier is

$$\left\{ \frac{\sqrt{5} e^{t(-2+\sqrt{5})}}{10}, e^{-t(2+\sqrt{5})} \right\}$$

Since there is no duplication between the basis function in the UC_set and the basis functions of the homogeneous solution, the trial solution is a linear combination of all the basis in the UC_set.

$$y_p = A_1 \cos(t) + A_2 \sin(t)$$

The unknowns $\{A_1, A_2\}$ are found by substituting the above trial solution y_p into the ODE and comparing coefficients. Substituting the trial solution into the ODE and simplifying gives

$$-2A_1 \cos(t) - 2A_2 \sin(t) - 4A_1 \sin(t) + 4A_2 \cos(t) = \sin(t)$$

Solving for the unknowns by comparing coefficients results in

$$\left[A_1 = -\frac{1}{5}, A_2 = -\frac{1}{10} \right]$$

Substituting the above back in the above trial solution y_p , gives the particular solution

$$y_p = -\frac{\cos(t)}{5} - \frac{\sin(t)}{10}$$

Therefore the general solution is

$$\begin{aligned} y &= y_h + y_p \\ &= \left(c_1 e^{-t(2+\sqrt{5})} + \frac{c_2 \sqrt{5} e^{t(-2+\sqrt{5})}}{10} \right) + \left(-\frac{\cos(t)}{5} - \frac{\sin(t)}{10} \right) \end{aligned}$$

Will add steps showing solving for IC soon.

Summary of solutions found

$$y = c_1 e^{-t(2+\sqrt{5})} + \frac{c_2 \sqrt{5} e^{t(-2+\sqrt{5})}}{10} - \frac{\cos(t)}{5} - \frac{\sin(t)}{10}$$

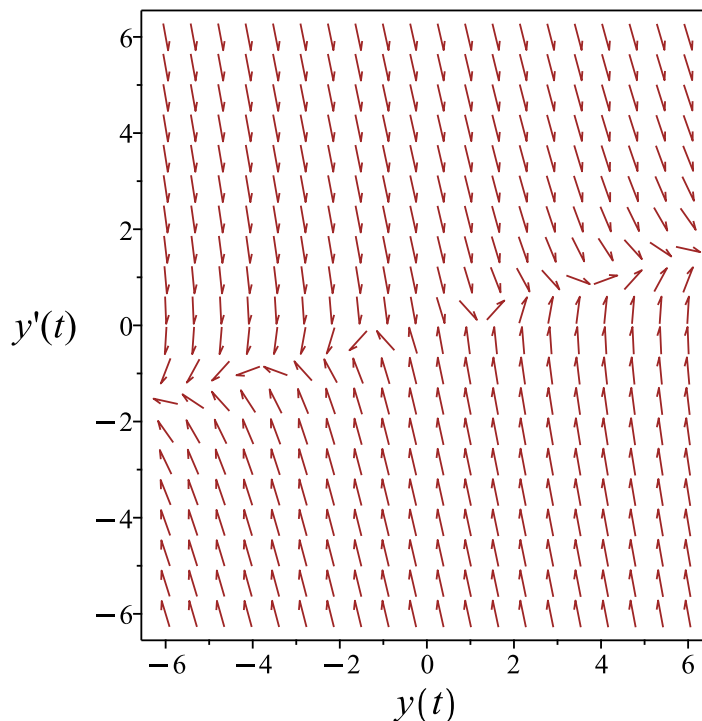


Figure 2.101: Slope field plot
 $y'' + 4y' - y = \sin(t)$

Solved as second order ode adjoint method

Time used: 1.243 (sec)

In normal form the ode

$$y'' + 4y' - y = \sin(t) \quad (1)$$

Becomes

$$y'' + p(t)y' + q(t)y = r(t) \quad (2)$$

Where

$$p(t) = 4$$

$$q(t) = -1$$

$$r(t) = \sin(t)$$

The Lagrange adjoint ode is given by

$$\xi'' - (\xi p)' + \xi q = 0$$

$$\xi'' - (4\xi(t))' + (-\xi(t)) = 0$$

$$\xi''(t) - 4\xi'(t) - \xi(t) = 0$$

Which is solved for $\xi(t)$. This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$A\xi''(t) + B\xi'(t) + C\xi(t) = 0$$

Where in the above $A = 1, B = -4, C = -1$. Let the solution be $\xi = e^{\lambda t}$. Substituting this into the ODE gives

$$\lambda^2 e^{t\lambda} - 4\lambda e^{t\lambda} - e^{t\lambda} = 0 \quad (1)$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda t}$ gives

$$\lambda^2 - 4\lambda - 1 = 0 \quad (2)$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = -4, C = -1$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{4}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{-4^2 - (4)(1)(-1)} \\ &= 2 \pm \sqrt{5} \end{aligned}$$

Hence

$$\lambda_1 = 2 + \sqrt{5}$$

$$\lambda_2 = 2 - \sqrt{5}$$

Which simplifies to

$$\lambda_1 = 2 + \sqrt{5}$$

$$\lambda_2 = 2 - \sqrt{5}$$

Since roots are real and distinct, then the solution is

$$\xi = c_1 e^{\lambda_1 t} + c_2 e^{\lambda_2 t}$$

$$\xi = c_1 e^{(2+\sqrt{5})t} + c_2 e^{(2-\sqrt{5})t}$$

Or

$$\xi = c_1 e^{t(2+\sqrt{5})} + c_2 e^{t(2-\sqrt{5})}$$

Will add steps showing solving for IC soon.

The original ode now reduces to first order ode

$$\begin{aligned}\xi(t) y' - y \xi'(t) + \xi(t) p(t) y &= \int \xi(t) r(t) dt \\ y' + y \left(p(t) - \frac{\xi'(t)}{\xi(t)} \right) &= \frac{\int \xi(t) r(t) dt}{\xi(t)}\end{aligned}$$

Or

$$y' + y \left(4 - \frac{c_1(2 + \sqrt{5}) e^{t(2+\sqrt{5})} + c_2(2 - \sqrt{5}) e^{t(2-\sqrt{5})}}{c_1 e^{t(2+\sqrt{5})} + c_2 e^{t(2-\sqrt{5})}} \right) = \frac{c_1 \left(-\frac{e^{t(2+\sqrt{5})} \cos(t)}{(2+\sqrt{5})^2+1} + \frac{(2+\sqrt{5}) e^{t(2+\sqrt{5})} \sin(t)}{(2+\sqrt{5})^2+1} \right) + c_2 \left(-\frac{e^{t(2-\sqrt{5})} \cos(t)}{(2-\sqrt{5})^2+1} + \frac{(2-\sqrt{5}) e^{t(2-\sqrt{5})} \sin(t)}{(2-\sqrt{5})^2+1} \right)}{c_1 e^{t(2+\sqrt{5})} + c_2 e^{t(2-\sqrt{5})}}$$

Which is now a first order ode. This is now solved for y . In canonical form a linear first order is

$$y' + q(t)y = p(t)$$

Comparing the above to the given ode shows that

$$\begin{aligned}q(t) &= -\frac{-c_2(2 + \sqrt{5}) e^{-t(-2+\sqrt{5})} + c_1 e^{t(2+\sqrt{5})} (-2 + \sqrt{5})}{c_1 e^{t(2+\sqrt{5})} + c_2 e^{-t(-2+\sqrt{5})}} \\ p(t) &= \frac{-2 \left(\left(\cos(t) + \frac{\sin(t)}{2} \right) \sqrt{5} + \frac{5 \cos(t)}{2} \right) c_2 e^{-t(-2+\sqrt{5})} + 2c_1 \left(\left(\cos(t) + \frac{\sin(t)}{2} \right) \sqrt{5} - \frac{5 \cos(t)}{2} \right) e^{t(2+\sqrt{5})}}{10c_1 e^{t(2+\sqrt{5})} + 10c_2 e^{-t(-2+\sqrt{5})}}\end{aligned}$$

The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int q dt} \\ &= e^{\int -\frac{-c_2(2+\sqrt{5})e^{-t(-2+\sqrt{5})} + c_1 e^{t(2+\sqrt{5})}(-2+\sqrt{5})}{c_1 e^{t(2+\sqrt{5})} + c_2 e^{-t(-2+\sqrt{5})}} dt} \\ &= \frac{e^{t(2+\sqrt{5})}}{e^{2\sqrt{5}t} c_1 + c_2}\end{aligned}$$

The ode becomes

$$\frac{d}{dt}(\mu y) = \mu p$$

$$\begin{aligned}&\frac{d}{dt}(\mu y) \\ &= (\mu) \left(\frac{-2 \left(\left(\cos(t) + \frac{\sin(t)}{2} \right) \sqrt{5} + \frac{5 \cos(t)}{2} \right) c_2 e^{-t(-2+\sqrt{5})} + 2c_1 \left(\left(\cos(t) + \frac{\sin(t)}{2} \right) \sqrt{5} - \frac{5 \cos(t)}{2} \right) e^{t(2+\sqrt{5})}}{10c_1 e^{t(2+\sqrt{5})} + 10c_2 e^{-t(-2+\sqrt{5})}} \right)\end{aligned}$$

$$\begin{aligned} & \frac{d}{dt} \left(\frac{y e^{t(2+\sqrt{5})}}{e^{2\sqrt{5}t} c_1 + c_2} \right) \\ &= \left(\frac{e^{t(2+\sqrt{5})}}{e^{2\sqrt{5}t} c_1 + c_2} \right) \left(\frac{-2 \left(\left(\cos(t) + \frac{\sin(t)}{2} \right) \sqrt{5} + \frac{5 \cos(t)}{2} \right) c_2 e^{-t(-2+\sqrt{5})} + 2c_1 \left(\left(\cos(t) + \frac{\sin(t)}{2} \right) \sqrt{5} - \frac{5 \cos(t)}{2} \right) e^{t(2+\sqrt{5})}}{10c_1 e^{t(2+\sqrt{5})} + 10c_2 e^{-t(-2+\sqrt{5})}} \right) \\ & \frac{d}{dt} \left(\frac{y e^{t(2+\sqrt{5})}}{e^{2\sqrt{5}t} c_1 + c_2} \right) \\ &= \left(\frac{\left(-2 \left(\left(\cos(t) + \frac{\sin(t)}{2} \right) \sqrt{5} + \frac{5 \cos(t)}{2} \right) c_2 e^{-t(-2+\sqrt{5})} + 2c_1 \left(\left(\cos(t) + \frac{\sin(t)}{2} \right) \sqrt{5} - \frac{5 \cos(t)}{2} \right) e^{t(2+\sqrt{5})} \right) e^{t(2+\sqrt{5})}}{\left(10c_1 e^{t(2+\sqrt{5})} + 10c_2 e^{-t(-2+\sqrt{5})} \right) (e^{2\sqrt{5}t} c_1 + c_2)} \right) e^{t(2+\sqrt{5})} \end{aligned}$$

Integrating gives

$$\begin{aligned} \frac{y e^{t(2+\sqrt{5})}}{e^{2\sqrt{5}t} c_1 + c_2} &= \int \frac{\left(-2 \left(\left(\cos(t) + \frac{\sin(t)}{2} \right) \sqrt{5} + \frac{5 \cos(t)}{2} \right) c_2 e^{-t(-2+\sqrt{5})} + 2c_1 \left(\left(\cos(t) + \frac{\sin(t)}{2} \right) \sqrt{5} - \frac{5 \cos(t)}{2} \right) e^{t(2+\sqrt{5})} \right) e^{t(2+\sqrt{5})}}{\left(10c_1 e^{t(2+\sqrt{5})} + 10c_2 e^{-t(-2+\sqrt{5})} \right) (e^{2\sqrt{5}t} c_1 + c_2)} dt \\ &= \frac{(i\sqrt{5} - 2\sqrt{5}) \sqrt{5} e^{t(2+i+\sqrt{5})}}{100 e^{2\sqrt{5}t} c_1 + 100c_2} - \frac{(i\sqrt{5} + 2\sqrt{5}) \sqrt{5} e^{t(2-i+\sqrt{5})}}{100 (e^{2\sqrt{5}t} c_1 + c_2)} + c_3 \end{aligned}$$

Dividing throughout by the integrating factor $\frac{e^{t(2+\sqrt{5})}}{e^{2\sqrt{5}t} c_1 + c_2}$ gives the final solution

$$y = - \frac{\left((2+i) e^{t(2-i+\sqrt{5})} + (2-i) e^{t(2+i+\sqrt{5})} - 20 e^{2\sqrt{5}t} c_1 c_3 - 20 c_2 c_3 \right) e^{-t(2+\sqrt{5})}}{20}$$

Hence, the solution found using Lagrange adjoint equation method is

$$y = - \frac{\left((2+i) e^{t(2-i+\sqrt{5})} + (2-i) e^{t(2+i+\sqrt{5})} - 20 e^{2\sqrt{5}t} c_1 c_3 - 20 c_2 c_3 \right) e^{-t(2+\sqrt{5})}}{20}$$

The constants can be merged to give

$$y = - \frac{\left((2+i) e^{t(2-i+\sqrt{5})} + (2-i) e^{t(2+i+\sqrt{5})} - 20 e^{2\sqrt{5}t} c_1 - 20 c_2 \right) e^{-t(2+\sqrt{5})}}{20}$$

Will add steps showing solving for IC soon.

Summary of solutions found

$$y = -\frac{\left((2+i)e^{t(2-i+\sqrt{5})} + (2-i)e^{t(2+i+\sqrt{5})} - 20e^{2\sqrt{5}t}c_1 - 20c_2 \right) e^{-t(2+\sqrt{5})}}{20}$$

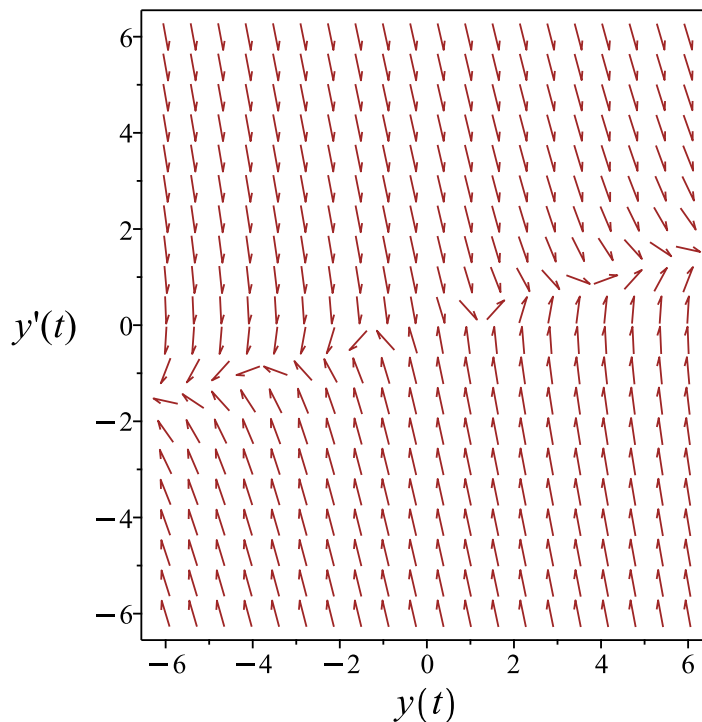


Figure 2.102: Slope field plot
 $y'' + 4y' - y = \sin(t)$

Maple step by step solution

Let's solve

$$y'' + 4y' - y = \sin(t)$$

- Highest derivative means the order of the ODE is 2
- y''
- Characteristic polynomial of homogeneous ODE
- $r^2 + 4r - 1 = 0$
- Use quadratic formula to solve for r

$$r = \frac{(-4) \pm (\sqrt{20})}{2}$$

- Roots of the characteristic polynomial

$$r = (-2 - \sqrt{5}, -2 + \sqrt{5})$$

- 1st solution of the homogeneous ODE

$$y_1(t) = e^{t(-2-\sqrt{5})}$$

- 2nd solution of the homogeneous ODE

$$y_2(t) = e^{t(-2+\sqrt{5})}$$

- General solution of the ODE

$$y = C1y_1(t) + C2y_2(t) + y_p(t)$$

- Substitute in solutions of the homogeneous ODE

$$y = C1 e^{t(-2-\sqrt{5})} + e^{t(-2+\sqrt{5})} C2 + y_p(t)$$

- Find a particular solution $y_p(t)$ of the ODE

- Use variation of parameters to find y_p here $f(t)$ is the forcing function

$$\left[y_p(t) = -y_1(t) \left(\int \frac{y_2(t)f(t)}{W(y_1(t),y_2(t))} dt \right) + y_2(t) \left(\int \frac{y_1(t)f(t)}{W(y_1(t),y_2(t))} dt \right), f(t) = \sin(t) \right]$$

- Wronskian of solutions of the homogeneous equation

$$W(y_1(t), y_2(t)) = \begin{bmatrix} e^{t(-2-\sqrt{5})} & e^{t(-2+\sqrt{5})} \\ (-2 - \sqrt{5}) e^{t(-2-\sqrt{5})} & (-2 + \sqrt{5}) e^{t(-2+\sqrt{5})} \end{bmatrix}$$

- Compute Wronskian

$$W(y_1(t), y_2(t)) = 2\sqrt{5} e^{-4t}$$

- Substitute functions into equation for $y_p(t)$

$$y_p(t) = -\frac{\sqrt{5} \left(e^{-t(2+\sqrt{5})} \left(\int \sin(t) e^{t(2+\sqrt{5})} dt \right) - e^{t(-2+\sqrt{5})} \left(\int \sin(t) e^{-t(-2+\sqrt{5})} dt \right) \right)}{10}$$

- Compute integrals

$$y_p(t) = -\frac{\cos(t)}{5} - \frac{\sin(t)}{10}$$

- Substitute particular solution into general solution to ODE

$$y = C1 e^{t(-2-\sqrt{5})} + e^{t(-2+\sqrt{5})} C2 - \frac{\cos(t)}{5} - \frac{\sin(t)}{10}$$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
trying high order exact linear fully integrable
trying differential order: 2; linear nonhomogeneous with symmetry [0,1]
trying a double symmetry of the form [xi=0, eta=F(x)]
-> Try solving first the homogeneous part of the ODE
    checking if the LODE has constant coefficients
    <- constant coefficients successful
<- solving first the homogeneous part of the ODE successful`

```

Maple dsolve solution

Solving time : 0.004 (sec)

Leaf size : 34

```

dsolve(diff(diff(y(t),t),t)+4*diff(y(t),t)-y(t) = sin(t),
        y(t),singsol=all)

```

$$y = e^{t(-2+\sqrt{5})}c_2 + c_1 e^{-t(2+\sqrt{5})} - \frac{\cos(t)}{5} - \frac{\sin(t)}{10}$$

Mathematica DSolve solution

Solving time : 0.245 (sec)

Leaf size : 47

```

DSolve[{D[y[t],{t,2}]+4*D[y[t],t]-y[t]==Sin[t],{}},
        y[t],t,IncludeSingularSolutions->True]

```

$$y(t) \rightarrow -\frac{\sin(t)}{10} - \frac{\cos(t)}{5} + e^{-((2+\sqrt{5})t)}(c_2 e^{2\sqrt{5}t} + c_1)$$

2.8.8 problem 8

Solved as second order linear constant coeff ode	670
Solved as second order ode using Kovacic algorithm	674
Solved as second order ode adjoint method	678
Maple step by step solution	683
Maple trace	684
Maple dsolve solution	685
Mathematica DSolve solution	685

Internal problem ID [18263]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter VII. Linear equations of order higher than the first. section 56. Problems at page 163

Problem number : 8

Date solved : Monday, December 23, 2024 at 09:47:06 PM

CAS classification : [[_2nd_order, _linear, _nonhomogeneous]]

Solve

$$y'' + 3y = \sin(x) + \frac{\sin(3x)}{3}$$

Solved as second order linear constant coeff ode

Time used: 0.208 (sec)

This is second order non-homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = f(x)$$

Where $A = 1, B = 0, C = 3, f(x) = \sin(x) + \frac{\sin(3x)}{3}$. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE $Ay''(x) + By'(x) + Cy(x) = 0$, and y_p is a particular solution to the non-homogeneous ODE $Ay''(x) + By'(x) + Cy(x) = f(x)$. y_h is the solution to

$$y'' + 3y = 0$$

This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = 0$$

Where in the above $A = 1, B = 0, C = 3$. Let the solution be $y = e^{\lambda x}$. Substituting this into the ODE gives

$$\lambda^2 e^{x\lambda} + 3e^{x\lambda} = 0 \quad (1)$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$\lambda^2 + 3 = 0 \quad (2)$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = 0, C = 3$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{0}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{0^2 - (4)(1)(3)} \\ &= \pm i\sqrt{3} \end{aligned}$$

Hence

$$\begin{aligned} \lambda_1 &= +i\sqrt{3} \\ \lambda_2 &= -i\sqrt{3} \end{aligned}$$

Which simplifies to

$$\begin{aligned} \lambda_1 &= i\sqrt{3} \\ \lambda_2 &= -i\sqrt{3} \end{aligned}$$

Since roots are complex conjugate of each others, then let the roots be

$$\lambda_{1,2} = \alpha \pm i\beta$$

Where $\alpha = 0$ and $\beta = \sqrt{3}$. Therefore the final solution, when using Euler relation, can be written as

$$y = e^{\alpha x} (c_1 \cos(\beta x) + c_2 \sin(\beta x))$$

Which becomes

$$y = e^0 \left(c_1 \cos(\sqrt{3}x) + c_2 \sin(\sqrt{3}x) \right)$$

Or

$$y = c_1 \cos(\sqrt{3}x) + c_2 \sin(\sqrt{3}x)$$

Therefore the homogeneous solution y_h is

$$y_h = c_1 \cos(\sqrt{3}x) + c_2 \sin(\sqrt{3}x)$$

The particular solution is now found using the method of undetermined coefficients. Looking at the RHS of the ode, which is

$$\sin(x) + \frac{\sin(3x)}{3}$$

Shows that the corresponding undetermined set of the basis functions (UC_set) for the trial solution is

$$[\{\cos(x), \sin(x)\}, \{\cos(3x), \sin(3x)\}]$$

While the set of the basis functions for the homogeneous solution found earlier is

$$\{\cos(\sqrt{3}x), \sin(\sqrt{3}x)\}$$

Since there is no duplication between the basis function in the UC_set and the basis functions of the homogeneous solution, the trial solution is a linear combination of all the basis in the UC_set.

$$y_p = A_1 \cos(x) + A_2 \sin(x) + A_3 \cos(3x) + A_4 \sin(3x)$$

The unknowns $\{A_1, A_2, A_3, A_4\}$ are found by substituting the above trial solution y_p into the ODE and comparing coefficients. Substituting the trial solution into the ODE and simplifying gives

$$2A_1 \cos(x) + 2A_2 \sin(x) - 6A_3 \cos(3x) - 6A_4 \sin(3x) = \sin(x) + \frac{\sin(3x)}{3}$$

Solving for the unknowns by comparing coefficients results in

$$\left[A_1 = 0, A_2 = \frac{1}{2}, A_3 = 0, A_4 = -\frac{1}{18} \right]$$

Substituting the above back in the above trial solution y_p , gives the particular solution

$$y_p = \frac{\sin(x)}{2} - \frac{\sin(3x)}{18}$$

Therefore the general solution is

$$y = y_h + y_p$$

$$= \left(c_1 \cos(\sqrt{3}x) + c_2 \sin(\sqrt{3}x) \right) + \left(\frac{\sin(x)}{2} - \frac{\sin(3x)}{18} \right)$$

Will add steps showing solving for IC soon.

Summary of solutions found

$$y = \frac{\sin(x)}{2} - \frac{\sin(3x)}{18} + c_1 \cos(\sqrt{3}x) + c_2 \sin(\sqrt{3}x)$$

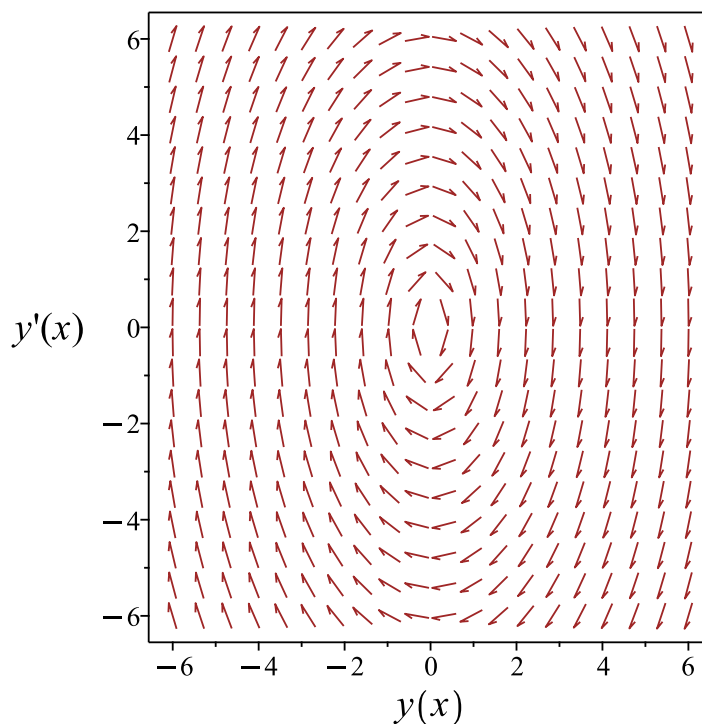


Figure 2.103: Slope field plot
 $y'' + 3y = \sin(x) + \frac{\sin(3x)}{3}$

Solved as second order ode using Kovacic algorithm

Time used: 0.216 (sec)

Writing the ode as

$$y'' + 3y = 0 \quad (1)$$

$$Ay'' + By' + Cy = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$\begin{aligned} A &= 1 \\ B &= 0 \\ C &= 3 \end{aligned} \quad (3)$$

Applying the Liouville transformation on the dependent variable gives

$$z(x) = ye^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \quad (4)$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \quad (5)$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{-3}{1} \quad (6)$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= -3 \\ t &= 1 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(x) = -3z(x) \quad (7)$$

Equation (7) is now solved. After finding $z(x)$ then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 2.59: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 0 - 0 \\ &= 0 \end{aligned}$$

There are no poles in r . Therefore the set of poles Γ is empty. Since there is no odd order pole larger than 2 and the order at ∞ is 0 then the necessary conditions for case one are met. Therefore

$$L = [1]$$

Since $r = -3$ is not a function of x , then there is no need run Kovacic algorithm to obtain a solution for transformed ode $z'' = rz$ as one solution is

$$z_1(x) = \cos(\sqrt{3}x)$$

Using the above, the solution for the original ode can now be found. The first solution to the original ode in y is found from

$$y_1 = z_1 e^{\int -\frac{1}{2} \frac{B}{A} dx}$$

Since $B = 0$ then the above reduces to

$$\begin{aligned} y_1 &= z_1 \\ &= \cos(\sqrt{3}x) \end{aligned}$$

Which simplifies to

$$y_1 = \cos(\sqrt{3}x)$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dx}}{y_1^2} dx$$

Since $B = 0$ then the above becomes

$$\begin{aligned} y_2 &= y_1 \int \frac{1}{y_1^2} dx \\ &= \cos(\sqrt{3}x) \int \frac{1}{\cos^2(\sqrt{3}x)} dx \\ &= \cos(\sqrt{3}x) \left(\frac{\sqrt{3} \tan(\sqrt{3}x)}{3} \right) \end{aligned}$$

Therefore the solution is

$$\begin{aligned} y &= c_1 y_1 + c_2 y_2 \\ &= c_1 \left(\cos(\sqrt{3}x) \right) + c_2 \left(\cos(\sqrt{3}x) \left(\frac{\sqrt{3} \tan(\sqrt{3}x)}{3} \right) \right) \end{aligned}$$

This is second order nonhomogeneous ODE. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE $Ay''(x) + By'(x) + Cy(x) = 0$, and y_p is a particular solution to the nonhomogeneous ODE $Ay''(x) + By'(x) + Cy(x) = f(x)$. y_h is the solution to

$$y'' + 3y = 0$$

The homogeneous solution is found using the Kovacic algorithm which results in

$$y_h = c_1 \cos(\sqrt{3}x) + \frac{c_2 \sqrt{3} \sin(\sqrt{3}x)}{3}$$

The particular solution is now found using the method of undetermined coefficients. Looking at the RHS of the ode, which is

$$\sin(x) + \frac{\sin(3x)}{3}$$

Shows that the corresponding undetermined set of the basis functions (UC_set) for the trial solution is

$$[\{\cos(x), \sin(x)\}, \{\cos(3x), \sin(3x)\}]$$

While the set of the basis functions for the homogeneous solution found earlier is

$$\left\{ \frac{\sqrt{3} \sin(\sqrt{3}x)}{3}, \cos(\sqrt{3}x) \right\}$$

Since there is no duplication between the basis function in the UC_set and the basis functions of the homogeneous solution, the trial solution is a linear combination of all the basis in the UC_set.

$$y_p = A_1 \cos(x) + A_2 \sin(x) + A_3 \cos(3x) + A_4 \sin(3x)$$

The unknowns $\{A_1, A_2, A_3, A_4\}$ are found by substituting the above trial solution y_p into the ODE and comparing coefficients. Substituting the trial solution into the ODE and simplifying gives

$$2A_1 \cos(x) + 2A_2 \sin(x) - 6A_3 \cos(3x) - 6A_4 \sin(3x) = \sin(x) + \frac{\sin(3x)}{3}$$

Solving for the unknowns by comparing coefficients results in

$$\left[A_1 = 0, A_2 = \frac{1}{2}, A_3 = 0, A_4 = -\frac{1}{18} \right]$$

Substituting the above back in the above trial solution y_p , gives the particular solution

$$y_p = \frac{\sin(x)}{2} - \frac{\sin(3x)}{18}$$

Therefore the general solution is

$$\begin{aligned} y &= y_h + y_p \\ &= \left(c_1 \cos(\sqrt{3}x) + \frac{c_2 \sqrt{3} \sin(\sqrt{3}x)}{3} \right) + \left(\frac{\sin(x)}{2} - \frac{\sin(3x)}{18} \right) \end{aligned}$$

Will add steps showing solving for IC soon.

Summary of solutions found

$$y = c_1 \cos(\sqrt{3}x) + \frac{c_2 \sqrt{3} \sin(\sqrt{3}x)}{3} + \frac{\sin(x)}{2} - \frac{\sin(3x)}{18}$$

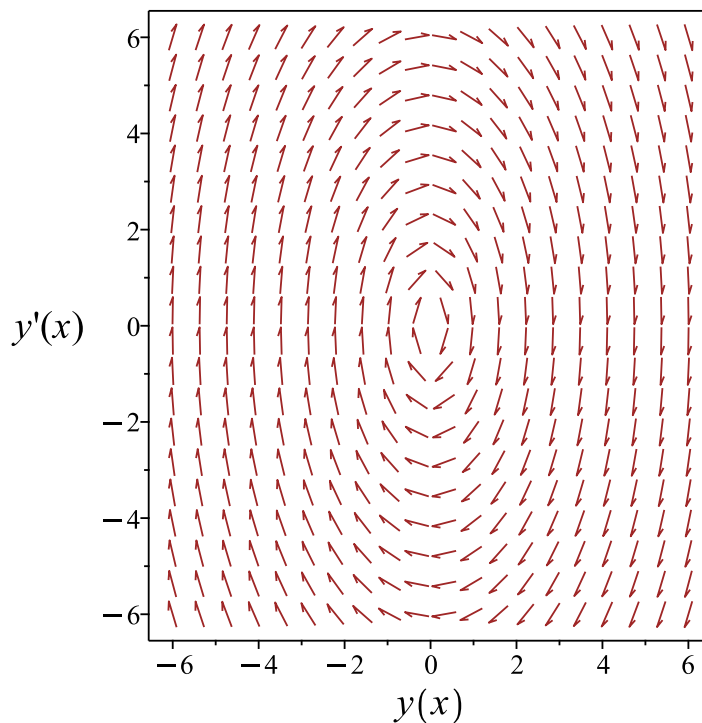


Figure 2.104: Slope field plot
 $y'' + 3y = \sin(x) + \frac{\sin(3x)}{3}$

Solved as second order ode adjoint method

Time used: 3.849 (sec)

In normal form the ode

$$y'' + 3y = \sin(x) + \frac{\sin(3x)}{3} \quad (1)$$

Becomes

$$y'' + p(x)y' + q(x)y = r(x) \quad (2)$$

Where

$$\begin{aligned} p(x) &= 0 \\ q(x) &= 3 \\ r(x) &= \sin(x) + \frac{\sin(3x)}{3} \end{aligned}$$

The Lagrange adjoint ode is given by

$$\begin{aligned} \xi'' - (\xi p)' + \xi q &= 0 \\ \xi'' - (0)' + (3\xi(x)) &= 0 \\ \xi''(x) + 3\xi(x) &= 0 \end{aligned}$$

Which is solved for $\xi(x)$. This is second order with constant coefficients homogeneous ODE. In standard form the ODE is

$$A\xi''(x) + B\xi'(x) + C\xi(x) = 0$$

Where in the above $A = 1, B = 0, C = 3$. Let the solution be $\xi = e^{\lambda x}$. Substituting this into the ODE gives

$$\lambda^2 e^{x\lambda} + 3e^{x\lambda} = 0 \quad (1)$$

Since exponential function is never zero, then dividing Eq(2) throughout by $e^{\lambda x}$ gives

$$\lambda^2 + 3 = 0 \quad (2)$$

Equation (2) is the characteristic equation of the ODE. Its roots determine the general solution form. Using the quadratic formula

$$\lambda_{1,2} = \frac{-B}{2A} \pm \frac{1}{2A} \sqrt{B^2 - 4AC}$$

Substituting $A = 1, B = 0, C = 3$ into the above gives

$$\begin{aligned} \lambda_{1,2} &= \frac{0}{(2)(1)} \pm \frac{1}{(2)(1)} \sqrt{0^2 - (4)(1)(3)} \\ &= \pm i\sqrt{3} \end{aligned}$$

Hence

$$\begin{aligned} \lambda_1 &= +i\sqrt{3} \\ \lambda_2 &= -i\sqrt{3} \end{aligned}$$

Which simplifies to

$$\begin{aligned}\lambda_1 &= i\sqrt{3} \\ \lambda_2 &= -i\sqrt{3}\end{aligned}$$

Since roots are complex conjugate of each others, then let the roots be

$$\lambda_{1,2} = \alpha \pm i\beta$$

Where $\alpha = 0$ and $\beta = \sqrt{3}$. Therefore the final solution, when using Euler relation, can be written as

$$\xi = e^{\alpha x}(c_1 \cos(\beta x) + c_2 \sin(\beta x))$$

Which becomes

$$\xi = e^0(c_1 \cos(\sqrt{3}x) + c_2 \sin(\sqrt{3}x))$$

Or

$$\xi = c_1 \cos(\sqrt{3}x) + c_2 \sin(\sqrt{3}x)$$

Will add steps showing solving for IC soon.

The original ode now reduces to first order ode

$$\begin{aligned}\xi(x)y' - y\xi'(x) + \xi(x)p(x)y &= \int \xi(x)r(x)dx \\ y' + y\left(p(x) - \frac{\xi'(x)}{\xi(x)}\right) &= \frac{\int \xi(x)r(x)dx}{\xi(x)}\end{aligned}$$

Or

$$y' - \frac{y(-c_1\sqrt{3}\sin(\sqrt{3}x) + c_2\sqrt{3}\cos(\sqrt{3}x))}{c_1\cos(\sqrt{3}x) + c_2\sin(\sqrt{3}x)} = \frac{\frac{c_1\cos(x(-3+\sqrt{3}))}{-18+6\sqrt{3}} - \frac{c_1\cos(x(1+\sqrt{3}))}{2(1+\sqrt{3})} - \frac{c_1\cos(x(3+\sqrt{3}))}{6(3+\sqrt{3})} + \frac{c_1\cos(x(-3-\sqrt{3}))}{-18-6\sqrt{3}}}{c_1\cos(\sqrt{3}x) + c_2\sin(\sqrt{3}x)}$$

Which is now a first order ode. This is now solved for y . In canonical form a linear first order is

$$y' + q(x)y = p(x)$$

Comparing the above to the given ode shows that

$$q(x) = -\frac{\sqrt{3}(c_2 \cos(\sqrt{3}x) - c_1 \sin(\sqrt{3}x))}{c_1 \cos(\sqrt{3}x) + c_2 \sin(\sqrt{3}x)}$$

$$p(x) = \frac{(2(\cos(x)^2 - \frac{5}{2}) \sin(x) c_2 \sqrt{3} - 6c_1 \cos(x)^3 + 9c_1 \cos(x)) \cos(\sqrt{3}x) - 2 \sin(\sqrt{3}x) (c_1 \sin(x) (\cos(x) \cos(\sqrt{3}x) + \sin(x) \sin(\sqrt{3}x)))}{9c_1 \cos(\sqrt{3}x) + 9c_2 \sin(\sqrt{3}x)}$$

The integrating factor μ is

$$\begin{aligned} \mu &= e^{\int q dx} \\ &= e^{\int -\frac{\sqrt{3}(c_2 \cos(\sqrt{3}x) - c_1 \sin(\sqrt{3}x))}{c_1 \cos(\sqrt{3}x) + c_2 \sin(\sqrt{3}x)} dx} \\ &= \frac{1}{c_1 \cos(\sqrt{3}x) + c_2 \sin(\sqrt{3}x)} \end{aligned}$$

The ode becomes

$$\begin{aligned} \frac{d}{dx}(\mu y) &= \mu p \\ \frac{d}{dx}(\mu y) &= (\mu) \left(\frac{(2(\cos(x)^2 - \frac{5}{2}) \sin(x) c_2 \sqrt{3} - 6c_1 \cos(x)^3 + 9c_1 \cos(x)) \cos(\sqrt{3}x) - 2 \sin(\sqrt{3}x) (c_1 \sin(x) (\cos(x) \cos(\sqrt{3}x) + \sin(x) \sin(\sqrt{3}x)))}{9c_1 \cos(\sqrt{3}x) + 9c_2 \sin(\sqrt{3}x)} \right) \\ \frac{d}{dx} \left(\frac{y}{c_1 \cos(\sqrt{3}x) + c_2 \sin(\sqrt{3}x)} \right) &= \left(\frac{1}{c_1 \cos(\sqrt{3}x) + c_2 \sin(\sqrt{3}x)} \right) \left(\frac{(2(\cos(x)^2 - \frac{5}{2}) \sin(x) c_2 \sqrt{3} - 6c_1 \cos(x)^3 + 9c_1 \cos(x)) \cos(\sqrt{3}x) - 2 \sin(\sqrt{3}x) (c_1 \sin(x) (\cos(x) \cos(\sqrt{3}x) + \sin(x) \sin(\sqrt{3}x)))}{9c_1 \cos(\sqrt{3}x) + 9c_2 \sin(\sqrt{3}x)} \right) \\ d \left(\frac{y}{c_1 \cos(\sqrt{3}x) + c_2 \sin(\sqrt{3}x)} \right) &= \left(\frac{(2(\cos(x)^2 - \frac{5}{2}) \sin(x) c_2 \sqrt{3} - 6c_1 \cos(x)^3 + 9c_1 \cos(x)) \cos(\sqrt{3}x) - 2 \sin(\sqrt{3}x) (c_1 \sin(x) (\cos(x) \cos(\sqrt{3}x) + \sin(x) \sin(\sqrt{3}x)))}{(9c_1 \cos(\sqrt{3}x) + 9c_2 \sin(\sqrt{3}x)) (c_1 \cos(\sqrt{3}x) + c_2 \sin(\sqrt{3}x))} \right) dx \end{aligned}$$

Integrating gives

$$\begin{aligned} \frac{y}{c_1 \cos(\sqrt{3}x) + c_2 \sin(\sqrt{3}x)} &= \int \frac{(2(\cos(x)^2 - \frac{5}{2}) \sin(x) c_2 \sqrt{3} - 6c_1 \cos(x)^3 + 9c_1 \cos(x)) \cos(\sqrt{3}x) - 2 \sin(\sqrt{3}x) (c_1 \sin(x) (\cos(x) \cos(\sqrt{3}x) + \sin(x) \sin(\sqrt{3}x)))}{(9c_1 \cos(\sqrt{3}x) + 9c_2 \sin(\sqrt{3}x)) (c_1 \cos(\sqrt{3}x) + c_2 \sin(\sqrt{3}x))} dx \\ &= \frac{i \left(e^{ix(3+\sqrt{3})} - 9e^{ix(1+\sqrt{3})} + 9e^{ix(\sqrt{3}-1)} - e^{ix(-3+\sqrt{3})} \right)}{-18ic_2 e^{2i\sqrt{3}x} + 18c_1 e^{2i\sqrt{3}x} + 18ic_2 + 18c_1} + c_3 \end{aligned}$$

Dividing throughout by the integrating factor $\frac{1}{c_1 \cos(\sqrt{3}x) + c_2 \sin(\sqrt{3}x)}$ gives the final solution

$$y = \frac{\left(18ie^{2i\sqrt{3}x}c_1c_3 + 18e^{2i\sqrt{3}x}c_2c_3 + 18ic_1c_3 - 18c_2c_3 - e^{ix(3+\sqrt{3})} + 9e^{ix(1+\sqrt{3})} - 9e^{ix(\sqrt{3}-1)} + e^{ix(-3+\sqrt{3})} \right)}{(18ic_1 + 18c_2)e^{2i\sqrt{3}x} + 18ic_1 - 18c_2}$$

Hence, the solution found using Lagrange adjoint equation method is

$$y = \frac{\left(18ie^{2i\sqrt{3}x}c_1c_3 + 18e^{2i\sqrt{3}x}c_2c_3 + 18ic_1c_3 - 18c_2c_3 - e^{ix(3+\sqrt{3})} + 9e^{ix(1+\sqrt{3})} - 9e^{ix(\sqrt{3}-1)} + e^{ix(-3+\sqrt{3})} \right)}{(18ic_1 + 18c_2)e^{2i\sqrt{3}x} + 18ic_1 - 18c_2}$$

The constants can be merged to give

$$y = \frac{\left(18ie^{2i\sqrt{3}x}c_1 + 18c_2e^{2i\sqrt{3}x} + 18ic_1 - 18c_2 - e^{ix(3+\sqrt{3})} + 9e^{ix(1+\sqrt{3})} - 9e^{ix(\sqrt{3}-1)} + e^{ix(-3+\sqrt{3})} \right) (c_1 \cos(\sqrt{3}x) + c_2 \sin(\sqrt{3}x))}{(18ic_1 + 18c_2)e^{2i\sqrt{3}x} + 18ic_1 - 18c_2}$$

Will add steps showing solving for IC soon.

Summary of solutions found

$$y = \frac{\left(18ie^{2i\sqrt{3}x}c_1 + 18c_2e^{2i\sqrt{3}x} + 18ic_1 - 18c_2 - e^{ix(3+\sqrt{3})} + 9e^{ix(1+\sqrt{3})} - 9e^{ix(\sqrt{3}-1)} + e^{ix(-3+\sqrt{3})} \right) (c_1 \cos(\sqrt{3}x) + c_2 \sin(\sqrt{3}x))}{(18ic_1 + 18c_2)e^{2i\sqrt{3}x} + 18ic_1 - 18c_2}$$

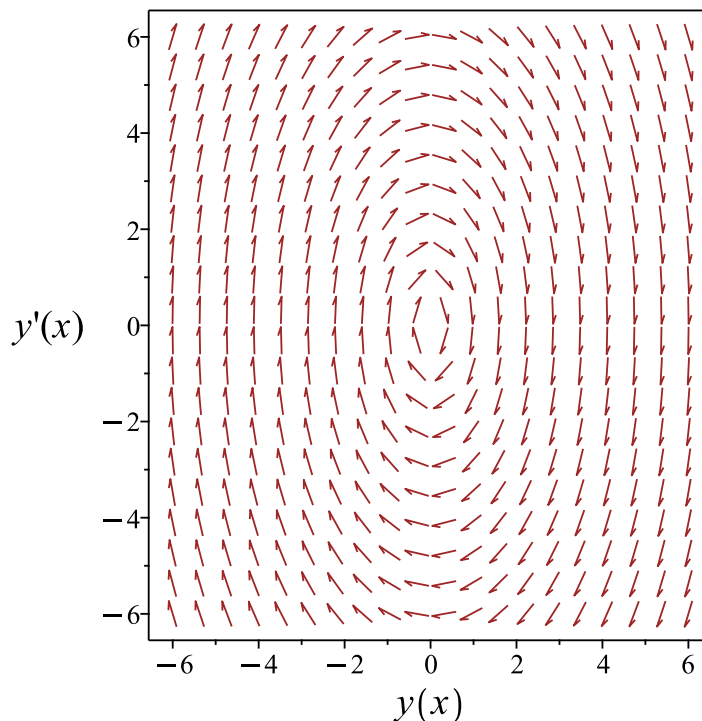


Figure 2.105: Slope field plot
 $y'' + 3y = \sin(x) + \frac{\sin(3x)}{3}$

Maple step by step solution

Let's solve

$$y'' + 3y = \sin(x) + \frac{\sin(3x)}{3}$$

- Highest derivative means the order of the ODE is 2
- y''
- Characteristic polynomial of homogeneous ODE
 $r^2 + 3 = 0$
- Use quadratic formula to solve for r
 $r = \frac{0 \pm \sqrt{-12}}{2}$
- Roots of the characteristic polynomial
 $r = (-I\sqrt{3}, I\sqrt{3})$
- 1st solution of the homogeneous ODE
 $y_1(x) = \cos(\sqrt{3}x)$
- 2nd solution of the homogeneous ODE
 $y_2(x) = \sin(\sqrt{3}x)$

- General solution of the ODE

$$y = C1y_1(x) + C2y_2(x) + y_p(x)$$
- Substitute in solutions of the homogeneous ODE

$$y = C1 \cos(\sqrt{3}x) + C2 \sin(\sqrt{3}x) + y_p(x)$$
- Find a particular solution $y_p(x)$ of the ODE
 - Use variation of parameters to find y_p here $f(x)$ is the forcing function

$$\left[y_p(x) = -y_1(x) \left(\int \frac{y_2(x)f(x)}{W(y_1(x),y_2(x))} dx \right) + y_2(x) \left(\int \frac{y_1(x)f(x)}{W(y_1(x),y_2(x))} dx \right), f(x) = \sin(x) + \frac{\sin(3x)}{3} \right]$$
 - Wronskian of solutions of the homogeneous equation

$$W(y_1(x), y_2(x)) = \begin{bmatrix} \cos(\sqrt{3}x) & \sin(\sqrt{3}x) \\ -\sqrt{3} \sin(\sqrt{3}x) & \sqrt{3} \cos(\sqrt{3}x) \end{bmatrix}$$
 - Compute Wronskian

$$W(y_1(x), y_2(x)) = \sqrt{3}$$
 - Substitute functions into equation for $y_p(x)$

$$y_p(x) = -\frac{\sqrt{3} \left(\cos(\sqrt{3}x) \left(\int \sin(\sqrt{3}x) (\sin(3x) + 3 \sin(x)) dx \right) - \sin(\sqrt{3}x) \left(\int \cos(\sqrt{3}x) (\sin(3x) + 3 \sin(x)) dx \right) \right)}{9}$$
 - Compute integrals

$$y_p(x) = \frac{\sin(x)}{2} - \frac{\sin(3x)}{18}$$
- Substitute particular solution into general solution to ODE

$$y = \frac{\sin(x)}{2} - \frac{\sin(3x)}{18} + C1 \cos(\sqrt{3}x) + C2 \sin(\sqrt{3}x)$$

Maple trace

```

`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
trying high order exact linear fully integrable
trying differential order: 2; linear nonhomogeneous with symmetry [0,1]
trying a double symmetry of the form [xi=0, eta=F(x)]
-> Try solving first the homogeneous part of the ODE
    checking if the LODE has constant coefficients
    <- constant coefficients successful
<- solving first the homogeneous part of the ODE successful`

```


Maple dsolve solution

Solving time : 0.004 (sec)

Leaf size : 31

```
dsolve(diff(diff(y(x),x),x)+3*y(x) = sin(x)+1/3*sin(3*x),
        y(x),singsol=all)
```

$$y = \frac{\sin(x)}{2} - \frac{\sin(3x)}{18} + c_1 \cos(\sqrt{3}x) + c_2 \sin(\sqrt{3}x)$$

Mathematica DSolve solution

Solving time : 0.536 (sec)

Leaf size : 42

```
DSolve[{D[y[x],{x,2}]+3*y[x]==Sin[x]+1/3*Sin[3*x],{}},
        y[x],x,IncludeSingularSolutions->True]
```

$$y(x) \rightarrow \frac{\sin(x)}{2} - \frac{1}{18} \sin(3x) + c_1 \cos(\sqrt{3}x) + c_2 \sin(\sqrt{3}x)$$

2.8.9 problem 10

Solved as first order linear ode	686
Solved as first order Exact ode	688
Maple step by step solution	693
Maple trace	694
Maple dsolve solution	694
Mathematica DSolve solution	694

Internal problem ID [18264]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter VII. Linear equations of order higher than the first. section 56. Problems at page 163

Problem number : 10

Date solved : Monday, December 23, 2024 at 09:47:12 PM

CAS classification : [[_linear, 'class A']]

Solve

$$5x' + x = \sin(3t)$$

Solved as first order linear ode

Time used: 0.227 (sec)

In canonical form a linear first order is

$$x' + q(t)x = p(t)$$

Comparing the above to the given ode shows that

$$q(t) = \frac{1}{5}$$

$$p(t) = \frac{\sin(3t)}{5}$$

The integrating factor μ is

$$\begin{aligned} \mu &= e^{\int q dt} \\ &= e^{\int \frac{1}{5} dt} \\ &= e^{\frac{t}{5}} \end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dt}(\mu x) &= \mu p \\ \frac{d}{dt}(\mu x) &= (\mu) \left(\frac{\sin(3t)}{5} \right) \\ \frac{d}{dt} \left(x e^{\frac{t}{5}} \right) &= \left(e^{\frac{t}{5}} \right) \left(\frac{\sin(3t)}{5} \right) \\ d \left(x e^{\frac{t}{5}} \right) &= \left(\frac{\sin(3t) e^{\frac{t}{5}}}{5} \right) dt\end{aligned}$$

Integrating gives

$$\begin{aligned}x e^{\frac{t}{5}} &= \int \frac{\sin(3t) e^{\frac{t}{5}}}{5} dt \\ &= -\frac{15 \cos(3t) e^{\frac{t}{5}}}{226} + \frac{\sin(3t) e^{\frac{t}{5}}}{226} + c_1\end{aligned}$$

Dividing throughout by the integrating factor $e^{\frac{t}{5}}$ gives the final solution

$$x = \frac{\sin(3t)}{226} - \frac{15 \cos(3t)}{226} + c_1 e^{-\frac{t}{5}}$$

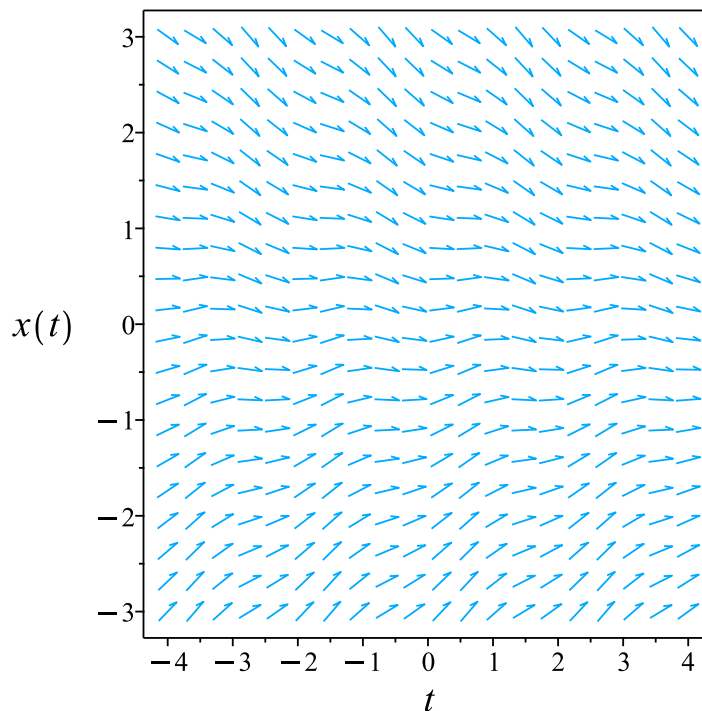


Figure 2.106: Slope field plot
 $5x' + x = \sin(3t)$

Summary of solutions found

$$x = \frac{\sin(3t)}{226} - \frac{15 \cos(3t)}{226} + c_1 e^{-\frac{t}{5}}$$

Solved as first order Exact ode

Time used: 0.142 (sec)

To solve an ode of the form

$$M(x, y) + N(x, y) \frac{dy}{dx} = 0 \quad (\text{A})$$

We assume there exists a function $\phi(x, y) = c$ where c is constant, that satisfies the ode. Taking derivative of ϕ w.r.t. x gives

$$\frac{d}{dx} \phi(x, y) = 0$$

Hence

$$\frac{\partial \phi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{dy}{dx} = 0 \quad (\text{B})$$

Comparing (A,B) shows that

$$\begin{aligned}\frac{\partial \phi}{\partial x} &= M \\ \frac{\partial \phi}{\partial y} &= N\end{aligned}$$

But since $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ then for the above to be valid, we require that

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

If the above condition is satisfied, then the original ode is called exact. We still need to determine $\phi(x, y)$ but at least we know now that we can do that since the condition $\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$ is satisfied. If this condition is not satisfied then this method will not work and we have to now look for an integrating factor to force this condition, which might or might not exist. The first step is to write the ODE in standard form to check for exactness, which is

$$M(t, x) dt + N(t, x) dx = 0 \tag{1A}$$

Therefore

$$\begin{aligned}(5) dx &= (-x + \sin(3t)) dt \\ (x - \sin(3t)) dt + (5) dx &= 0\end{aligned} \tag{2A}$$

Comparing (1A) and (2A) shows that

$$\begin{aligned}M(t, x) &= x - \sin(3t) \\ N(t, x) &= 5\end{aligned}$$

The next step is to determine if the ODE is exact or not. The ODE is exact when the following condition is satisfied

$$\frac{\partial M}{\partial x} = \frac{\partial N}{\partial t}$$

Using result found above gives

$$\begin{aligned}\frac{\partial M}{\partial x} &= \frac{\partial}{\partial x}(x - \sin(3t)) \\ &= 1\end{aligned}$$

And

$$\begin{aligned}\frac{\partial N}{\partial t} &= \frac{\partial}{\partial t}(5) \\ &= 0\end{aligned}$$

Since $\frac{\partial M}{\partial x} \neq \frac{\partial N}{\partial t}$, then the ODE is not exact. Since the ODE is not exact, we will try to find an integrating factor to make it exact. Let

$$\begin{aligned} A &= \frac{1}{N} \left(\frac{\partial M}{\partial x} - \frac{\partial N}{\partial t} \right) \\ &= \frac{1}{5} ((1) - (0)) \\ &= \frac{1}{5} \end{aligned}$$

Since A does not depend on x , then it can be used to find an integrating factor. The integrating factor μ is

$$\begin{aligned} \mu &= e^{\int A dt} \\ &= e^{\int \frac{1}{5} dt} \end{aligned}$$

The result of integrating gives

$$\begin{aligned} \mu &= e^{\frac{t}{5}} \\ &= e^{\frac{t}{5}} \end{aligned}$$

M and N are multiplied by this integrating factor, giving new M and new N which are called \overline{M} and \overline{N} for now so not to confuse them with the original M and N .

$$\begin{aligned} \overline{M} &= \mu M \\ &= e^{\frac{t}{5}} (x - \sin(3t)) \\ &= (x - \sin(3t)) e^{\frac{t}{5}} \end{aligned}$$

And

$$\begin{aligned} \overline{N} &= \mu N \\ &= e^{\frac{t}{5}} (5) \\ &= 5 e^{\frac{t}{5}} \end{aligned}$$

Now a modified ODE is obtained from the original ODE, which is exact and can be solved. The modified ODE is

$$\begin{aligned} \overline{M} + \overline{N} \frac{dx}{dt} &= 0 \\ \left((x - \sin(3t)) e^{\frac{t}{5}} \right) + \left(5 e^{\frac{t}{5}} \right) \frac{dx}{dt} &= 0 \end{aligned}$$

The following equations are now set up to solve for the function $\phi(t, x)$

$$\frac{\partial \phi}{\partial t} = \overline{M} \quad (1)$$

$$\frac{\partial \phi}{\partial x} = \overline{N} \quad (2)$$

Integrating (2) w.r.t. x gives

$$\begin{aligned} \int \frac{\partial \phi}{\partial x} dx &= \int \overline{N} dx \\ \int \frac{\partial \phi}{\partial x} dx &= \int 5 e^{\frac{t}{5}} dx \\ \phi &= 5x e^{\frac{t}{5}} + f(t) \end{aligned} \quad (3)$$

Where $f(t)$ is used for the constant of integration since ϕ is a function of both t and x . Taking derivative of equation (3) w.r.t t gives

$$\frac{\partial \phi}{\partial t} = x e^{\frac{t}{5}} + f'(t) \quad (4)$$

But equation (1) says that $\frac{\partial \phi}{\partial t} = (x - \sin(3t)) e^{\frac{t}{5}}$. Therefore equation (4) becomes

$$(x - \sin(3t)) e^{\frac{t}{5}} = x e^{\frac{t}{5}} + f'(t) \quad (5)$$

Solving equation (5) for $f'(t)$ gives

$$f'(t) = -\sin(3t) e^{\frac{t}{5}}$$

Integrating the above w.r.t t gives

$$\begin{aligned} \int f'(t) dt &= \int \left(-\sin(3t) e^{\frac{t}{5}} \right) dt \\ f(t) &= \frac{75 \cos(3t) e^{\frac{t}{5}}}{226} - \frac{5 \sin(3t) e^{\frac{t}{5}}}{226} + c_1 \end{aligned}$$

Where c_1 is constant of integration. Substituting result found above for $f(t)$ into equation (3) gives ϕ

$$\phi = 5x e^{\frac{t}{5}} + \frac{75 \cos(3t) e^{\frac{t}{5}}}{226} - \frac{5 \sin(3t) e^{\frac{t}{5}}}{226} + c_1$$

But since ϕ itself is a constant function, then let $\phi = c_2$ where c_2 is new constant and combining c_1 and c_2 constants into the constant c_1 gives the solution as

$$c_1 = 5x e^{\frac{t}{5}} + \frac{75 \cos(3t) e^{\frac{t}{5}}}{226} - \frac{5 \sin(3t) e^{\frac{t}{5}}}{226}$$

Solving for x gives

$$x = -\frac{\left(75 \cos(3t) e^{\frac{t}{5}} - 5 \sin(3t) e^{\frac{t}{5}} - 226c_1\right) e^{-\frac{t}{5}}}{1130}$$

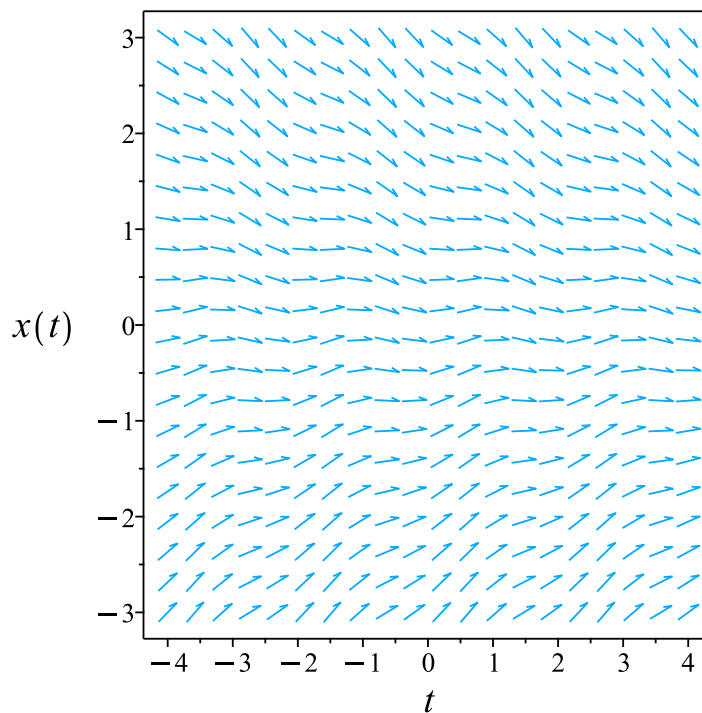


Figure 2.107: Slope field plot
 $5x' + x = \sin(3t)$

Summary of solutions found

$$x = -\frac{\left(75 \cos(3t) e^{\frac{t}{5}} - 5 \sin(3t) e^{\frac{t}{5}} - 226c_1\right) e^{-\frac{t}{5}}}{1130}$$

Maple step by step solution

Let's solve

$$5x' + x = \sin(3t)$$

- Highest derivative means the order of the ODE is 1

$$x'$$

- Solve for the highest derivative

$$x' = -\frac{x}{5} + \frac{\sin(3t)}{5}$$

- Group terms with x on the lhs of the ODE and the rest on the rhs of the ODE

$$x' + \frac{x}{5} = \frac{\sin(3t)}{5}$$

- The ODE is linear; multiply by an integrating factor $\mu(t)$

$$\mu(t) \left(x' + \frac{x}{5} \right) = \frac{\mu(t) \sin(3t)}{5}$$

- Assume the lhs of the ODE is the total derivative $\frac{d}{dt}(x\mu(t))$

$$\mu(t) \left(x' + \frac{x}{5} \right) = x'\mu(t) + x\mu'(t)$$

- Isolate $\mu'(t)$

$$\mu'(t) = \frac{\mu(t)}{5}$$

- Solve to find the integrating factor

$$\mu(t) = e^{\frac{t}{5}}$$

- Integrate both sides with respect to t

$$\int \left(\frac{d}{dt}(x\mu(t)) \right) dt = \int \frac{\mu(t) \sin(3t)}{5} dt + C1$$

- Evaluate the integral on the lhs

$$x\mu(t) = \int \frac{\mu(t) \sin(3t)}{5} dt + C1$$

- Solve for x

$$x = \frac{\int \frac{\mu(t) \sin(3t)}{5} dt + C1}{\mu(t)}$$

- Substitute $\mu(t) = e^{\frac{t}{5}}$

$$x = \frac{\int \frac{\sin(3t)e^{\frac{t}{5}}}{5} dt + C1}{e^{\frac{t}{5}}}$$

- Evaluate the integrals on the rhs

$$x = \frac{-\frac{15 \cos(3t)e^{\frac{t}{5}}}{226} + \frac{\sin(3t)e^{\frac{t}{5}}}{226} + C1}{e^{\frac{t}{5}}}$$

- Simplify

$$x = \frac{\sin(3t)}{226} - \frac{15 \cos(3t)}{226} + C1 e^{-\frac{t}{5}}$$

Maple trace

```

`Methods for first order ODEs:
--- Trying classification methods ---
trying a quadrature
trying 1st order linear
<- 1st order linear successful`

```

Maple dsolve solution

Solving time : 0.003 (sec)

Leaf size : 23

```

dsolve(5*diff(x(t),t)+x(t) = sin(3*t),
       x(t),singsol=all)

```

$$x = \frac{\sin(3t)}{226} - \frac{15 \cos(3t)}{226} + c_1 e^{-\frac{t}{5}}$$

Mathematica DSolve solution

Solving time : 0.084 (sec)

Leaf size : 31

```

DSolve[{5*D[x[t],t]+x[t]==Sin[3*t],{}},
       x[t],t,IncludeSingularSolutions->True]

```

$$x(t) \rightarrow \frac{1}{226}(\sin(3t) - 15 \cos(3t)) + c_1 e^{-t/5}$$

2.8.10 problem 11

Solved as higher order constant coeff ode	695
Maple step by step solution	697
Maple trace	699
Maple dsolve solution	700
Mathematica DSolve solution	700

Internal problem ID [18265]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter VII. Linear equations of order higher than the first. section 56. Problems at page 163

Problem number : 11

Date solved : Monday, December 23, 2024 at 09:47:13 PM

CAS classification : [[_high_order, _missing_y]]

Solve

$$x'''' - 6x''' + 11x'' - 6x' = e^{-3t}$$

Solved as higher order constant coeff ode

Time used: 0.104 (sec)

The characteristic equation is

$$\lambda^4 - 6\lambda^3 + 11\lambda^2 - 6\lambda = 0$$

The roots of the above equation are

$$\lambda_1 = 0$$

$$\lambda_2 = 1$$

$$\lambda_3 = 2$$

$$\lambda_4 = 3$$

Therefore the homogeneous solution is

$$x_h(t) = c_1 + e^t c_2 + e^{2t} c_3 + e^{3t} c_4$$

The fundamental set of solutions for the homogeneous solution are the following

$$\begin{aligned}x_1 &= 1 \\x_2 &= e^t \\x_3 &= e^{2t} \\x_4 &= e^{3t}\end{aligned}$$

This is higher order nonhomogeneous ODE. Let the solution be

$$x = x_h + x_p$$

Where x_h is the solution to the homogeneous ODE And x_p is a particular solution to the nonhomogeneous ODE. x_h is the solution to

$$x'''' - 6x''' + 11x'' - 6x' = 0$$

Now the particular solution to the given ODE is found

$$x'''' - 6x''' + 11x'' - 6x' = e^{-3t}$$

The particular solution is now found using the method of undetermined coefficients.

Looking at the RHS of the ode, which is

$$e^{-3t}$$

Shows that the corresponding undetermined set of the basis functions (UC_set) for the trial solution is

$$[\{e^{-3t}\}]$$

While the set of the basis functions for the homogeneous solution found earlier is

$$\{1, e^t, e^{2t}, e^{3t}\}$$

Since there is no duplication between the basis function in the UC_set and the basis functions of the homogeneous solution, the trial solution is a linear combination of all the basis in the UC_set.

$$x_p = A_1 e^{-3t}$$

The unknowns $\{A_1\}$ are found by substituting the above trial solution x_p into the ODE and comparing coefficients. Substituting the trial solution into the ODE and simplifying gives

$$360A_1 e^{-3t} = e^{-3t}$$

Solving for the unknowns by comparing coefficients results in

$$\left[A_1 = \frac{1}{360} \right]$$

Substituting the above back in the above trial solution x_p , gives the particular solution

$$x_p = \frac{e^{-3t}}{360}$$

Therefore the general solution is

$$\begin{aligned} x &= x_h + x_p \\ &= (c_1 + e^t c_2 + e^{2t} c_3 + e^{3t} c_4) + \left(\frac{e^{-3t}}{360} \right) \end{aligned}$$

Maple step by step solution

Let's solve

$$x'''' - 6x''' + 11x'' - 6x' = e^{-3t}$$

- Highest derivative means the order of the ODE is 4
 x''''
- Characteristic polynomial of homogeneous ODE
 $r^4 - 6r^3 + 11r^2 - 6r = 0$
- Roots of the characteristic polynomial
 $r = [0, 1, 2, 3]$
- Homogeneous solution from $r = 0$
 $x_1(t) = 1$
- Homogeneous solution from $r = 1$
 $x_2(t) = e^t$
- Homogeneous solution from $r = 2$
 $x_3(t) = e^{2t}$
- Homogeneous solution from $r = 3$
 $x_4(t) = e^{3t}$
- General solution of the ODE
 $x = C_1 x_1(t) + C_2 x_2(t) + C_3 x_3(t) + C_4 x_4(t) + x_p(t)$
- Substitute in solutions of the homogeneous ODE
 $x = C_1 + e^t C_2 + e^{2t} C_3 + e^{3t} C_4 + x_p(t)$

□ Find a particular solution $x_p(t)$ of the ODE

- Define the forcing function of the ODE

$$f(t) = e^{-3t}$$

- Form of the particular solution to the ODE where the $u_i(t)$ are to be found

$$x_p(t) = \sum_{i=1}^4 u_i(t) x_i(t)$$

- Calculate the 1st derivative of $x_p(t)$

$$x_p'(t) = \sum_{i=1}^4 (u_i'(t) x_i(t) + u_i(t) x_i'(t))$$

- Choose equation to add to a system of equations in $u_i'(t)$

$$\sum_{i=1}^4 u_i'(t) x_i(t) = 0$$

- Calculate the 2nd derivative of $x_p(t)$

$$x_p''(t) = \sum_{i=1}^4 (u_i'(t) x_i'(t) + u_i(t) x_i''(t))$$

- Choose equation to add to a system of equations in $u_i'(t)$

$$\sum_{i=1}^4 u_i'(t) x_i'(t) = 0$$

- Calculate the 3rd derivative of $x_p(t)$

$$x_p'''(t) = \sum_{i=1}^4 (u_i'(t) x_i''(t) + u_i(t) x_i'''(t))$$

- Choose equation to add to a system of equations in $u_i'(t)$

$$\sum_{i=1}^4 u_i'(t) x_i''(t) = 0$$

- The ODE is of the following form where the $P_i(t)$ in this situation are the coefficients of the ODE

$$x'''' + \left(\sum_{i=0}^3 P_i(t) x^{(i)} \right) = f(t)$$

- Substitute $x_p(t) = \sum_{i=1}^4 u_i(t) x_i(t)$ into the ODE

$$\left(\sum_{j=0}^3 P_j(t) \left(\sum_{i=1}^4 u_i(t) x_i^{(j)}(t) \right) \right) + \sum_{i=1}^4 (u_i'(t) x_i'''(t) + u_i(t) x_i''''(t)) = f(t)$$

- Rearrange the ODE

$$\sum_{i=1}^4 \left(u_i(t) \cdot \left(\left(\sum_{j=0}^3 P_j(t) x_i^{(j)}(t) \right) + x_i''''(t) \right) + u_i'(t) x_i'''(t) \right) = f(t)$$

- Notice that $x_i(t)$ are solutions to the homogeneous equation so the first term in the sum is 0

$$\sum_{i=1}^4 u_i'(t) x_i'''(t) = f(t)$$

- We have now made a system of 4 equations in 4 unknowns ($u_i'(t)$)

$$\left[\sum_{i=1}^4 u_i'(t) x_i(t) = 0, \sum_{i=1}^4 u_i'(t) x_i'(t) = 0, \sum_{i=1}^4 u_i'(t) x_i''(t) = 0, \sum_{i=1}^4 u_i'(t) x_i'''(t) = f(t) \right]$$

- Convert the system to linear algebra format, notice that the matrix is the wronskian W

$$\begin{bmatrix} x_1(t) & x_2(t) & x_3(t) & x_4(t) \\ x_1'(t) & x_2'(t) & x_3'(t) & x_4'(t) \\ x_1''(t) & x_2''(t) & x_3''(t) & x_4''(t) \\ x_1'''(t) & x_2'''(t) & x_3'''(t) & x_4'''(t) \end{bmatrix} \cdot \begin{bmatrix} u_1'(t) \\ u_2'(t) \\ u_3'(t) \\ u_4'(t) \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \\ f(t) \end{bmatrix}$$

- Solve for the varied parameters

$$\begin{bmatrix} u_1(t) \\ u_2(t) \\ u_3(t) \\ u_4(t) \end{bmatrix} = \int \frac{1}{W} \cdot \begin{bmatrix} 0 \\ 0 \\ 0 \\ f(t) \end{bmatrix} dt$$

- Substitute in the homogeneous solutions and forcing function and solve

$$\begin{bmatrix} u_1(t) \\ u_2(t) \\ u_3(t) \\ u_4(t) \end{bmatrix} = \begin{bmatrix} \frac{e^{-3t}}{18} \\ -\frac{e^{-3t}}{8e^t} \\ \frac{e^{-3t}}{10e^{2t}} \\ -\frac{e^{-3t}}{36e^{3t}} \end{bmatrix}$$

Find a particular solution $x_p(t)$ of the ODE

$$x_p(t) = \frac{e^{-3t}}{360}$$

- Substitute particular solution into general solution to ODE

$$x = C_1 + e^t C_2 + e^{2t} C_3 + e^{3t} C_4 + \frac{e^{-3t}}{360}$$

Maple trace

```

`Methods for high order ODEs:
--- Trying classification methods ---
trying a quadrature
trying high order exact linear fully integrable
trying differential order: 4; linear nonhomogeneous with symmetry [0,1]
-> Calling odsolve with the ODE`, diff(diff(diff(_b(_a), _a), _a), _a) = 6*(diff(diff

```

```
Methods for third order ODEs:  
--- Trying classification methods ---  
trying a quadrature  
trying high order exact linear fully integrable  
trying differential order: 3; linear nonhomogeneous with symmetry [0,1]  
trying high order linear exact nonhomogeneous  
trying differential order: 3; missing the dependent variable  
checking if the LODE has constant coefficients  
<- constant coefficients successful  
<- differential order: 4; linear nonhomogeneous with symmetry [0,1] successful`
```

Maple dsolve solution

Solving time : 0.003 (sec)

Leaf size : 39

```
dsolve(diff(diff(diff(diff(x(t),t),t),t),t)-6*diff(diff(diff(x(t),t),t),t)+11*diff(diff(x(t),t),t),t),t),t),t),t),t)+11*diff(diff(x(t),t),t),t)+6*diff(diff(x(t),t),t),t)-6*diff(x(t),t),t)+11*x(t),singsol=all)
```

$$x = \frac{\left(c_3 e^{6t} + 3c_1 e^{4t} + \frac{3c_2 e^{5t}}{2} + 3e^{3t}c_4 + \frac{1}{120}\right) e^{-3t}}{3}$$

Mathematica DSolve solution

Solving time : 0.06 (sec)

Leaf size : 45

```
DSolve[{D[x[t],{t,4}]-6*D[x[t],{t,3}]+11*D[x[t],{t,2}]-6*D[x[t],t]==Exp[-3*t],{}}],x[t],t,IncludeSingularSolutions->True]
```

$$x(t) \rightarrow \frac{e^{-3t}}{360} + c_1 e^t + \frac{1}{2} c_2 e^{2t} + \frac{1}{3} c_3 e^{3t} + c_4$$

2.8.11 problem 14

Solved as higher order Euler type ode	701
Maple step by step solution	706
Maple trace	706
Maple dsolve solution	707
Mathematica DSolve solution	707

Internal problem ID [18266]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter VII. Linear equations of order higher than the first. section 56. Problems at page 163

Problem number : 14

Date solved : Monday, December 23, 2024 at 09:47:14 PM

CAS classification : [[_high_order, _missing_y]]

Solve

$$x^4 y'''' + x^3 y''' - 20x^2 y'' + 20xy' = 17x^6$$

Solved as higher order Euler type ode

Time used: 0.562 (sec)

This is Euler ODE of higher order. Let $y = x^\lambda$. Hence

$$\begin{aligned} y' &= \lambda x^{\lambda-1} \\ y'' &= \lambda(\lambda-1) x^{\lambda-2} \\ y''' &= \lambda(\lambda-1)(\lambda-2) x^{\lambda-3} \\ y'''' &= \lambda(\lambda-1)(\lambda-2)(\lambda-3) x^{\lambda-4} \end{aligned}$$

Substituting these back into

$$x^4 y'''' + x^3 y''' - 20x^2 y'' + 20xy' = 17x^6$$

gives

$$\begin{aligned} 20x\lambda x^{\lambda-1} - 20x^2\lambda(\lambda-1) x^{\lambda-2} + x^3\lambda(\lambda-1)(\lambda-2) x^{\lambda-3} \\ + x^4\lambda(\lambda-1)(\lambda-2)(\lambda-3) x^{\lambda-4} = 0 \end{aligned}$$

Which simplifies to

$$20\lambda x^\lambda - 20\lambda(\lambda-1) x^\lambda + \lambda(\lambda-1)(\lambda-2) x^\lambda + \lambda(\lambda-1)(\lambda-2)(\lambda-3) x^\lambda = 0$$

And since $x^\lambda \neq 0$ then dividing through by x^λ , the above becomes

$$20\lambda - 20\lambda(\lambda - 1) + \lambda(\lambda - 1)(\lambda - 2) + \lambda(\lambda - 1)(\lambda - 2)(\lambda - 3) = 0$$

Simplifying gives the characteristic equation as

$$\lambda^4 - 5\lambda^3 - 12\lambda^2 + 36\lambda = 0$$

Solving the above gives the following roots

$$\lambda_1 = 0$$

$$\lambda_2 = 2$$

$$\lambda_3 = 6$$

$$\lambda_4 = -3$$

This table summarises the result

root	multiplicity	type of root
0	1	real root
-3	1	real root
2	1	real root
6	1	real root

The solution is generated by going over the above table. For each real root λ of multiplicity one generates a $c_1 x^\lambda$ basis solution. Each real root of multiplicity two, generates $c_1 x^\lambda$ and $c_2 x^\lambda \ln(x)$ basis solutions. Each real root of multiplicity three, generates $c_1 x^\lambda$ and $c_2 x^\lambda \ln(x)$ and $c_3 x^\lambda \ln(x)^2$ basis solutions, and so on. Each complex root $\alpha \pm i\beta$ of multiplicity one generates $x^\alpha(c_1 \cos(\beta \ln(x)) + c_2 \sin(\beta \ln(x)))$ basis solutions. And each complex root $\alpha \pm i\beta$ of multiplicity two generates $\ln(x) x^\alpha(c_1 \cos(\beta \ln(x)) + c_2 \sin(\beta \ln(x)))$ basis solutions. And each complex root $\alpha \pm i\beta$ of multiplicity three generates $\ln(x)^2 x^\alpha(c_1 \cos(\beta \ln(x)) + c_2 \sin(\beta \ln(x)))$ basis solutions. And so on. Using the above show that the solution is

$$y = c_1 + \frac{c_2}{x^3} + c_3 x^2 + c_4 x^6$$

The fundamental set of solutions for the homogeneous solution are the following

$$y_1 = 1$$

$$y_2 = \frac{1}{x^3}$$

$$y_3 = x^2$$

$$y_4 = x^6$$

This is higher order nonhomogeneous Euler type ODE. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous Euler ODE And y_p is a particular solution to the nonhomogeneous Euler ODE. y_h is the solution to

$$x^4 y'''' + x^3 y''' - 20x^2 y'' + 20xy' = 0$$

Now the particular solution to the given ODE is found

$$x^4 y'''' + x^3 y''' - 20x^2 y'' + 20xy' = 17x^6$$

Let the particular solution be

$$y_p = U_1 y_1 + U_2 y_2 + U_3 y_3 + U_4 y_4$$

Where y_i are the basis solutions found above for the homogeneous solution y_h and $U_i(x)$ are functions to be determined as follows

$$U_i = (-1)^{n-i} \int \frac{F(x)W_i(x)}{aW(x)} dx$$

Where $W(x)$ is the Wronskian and $W_i(x)$ is the Wronskian that results after deleting the last row and the i -th column of the determinant and n is the order of the ODE or equivalently, the number of basis solutions, and a is the coefficient of the leading derivative in the ODE, and $F(x)$ is the RHS of the ODE. Therefore, the first step is to find the Wronskian $W(x)$. This is given by

$$W(x) = \begin{vmatrix} y_1 & y_2 & y_3 & y_4 \\ y_1' & y_2' & y_3' & y_4' \\ y_1'' & y_2'' & y_3'' & y_4'' \\ y_1''' & y_2''' & y_3''' & y_4''' \end{vmatrix}$$

Substituting the fundamental set of solutions y_i found above in the Wronskian gives

$$W = \begin{bmatrix} 1 & \frac{1}{x^3} & x^2 & x^6 \\ 0 & -\frac{3}{x^4} & 2x & 6x^5 \\ 0 & \frac{12}{x^5} & 2 & 30x^4 \\ 0 & -\frac{60}{x^6} & 0 & 120x^3 \end{bmatrix}$$

$$|W| = -\frac{6480}{x}$$

The determinant simplifies to

$$|W| = -\frac{6480}{x}$$

Now we determine W_i for each U_i .

$$\begin{aligned} W_1(x) &= \det \begin{bmatrix} \frac{1}{x^3} & x^2 & x^6 \\ -\frac{3}{x^4} & 2x & 6x^5 \\ \frac{12}{x^5} & 2 & 30x^4 \end{bmatrix} \\ &= 180x^2 \end{aligned}$$

$$\begin{aligned} W_2(x) &= \det \begin{bmatrix} 1 & x^2 & x^6 \\ 0 & 2x & 6x^5 \\ 0 & 2 & 30x^4 \end{bmatrix} \\ &= 48x^5 \end{aligned}$$

$$\begin{aligned} W_3(x) &= \det \begin{bmatrix} 1 & \frac{1}{x^3} & x^6 \\ 0 & -\frac{3}{x^4} & 6x^5 \\ 0 & \frac{12}{x^5} & 30x^4 \end{bmatrix} \\ &= -162 \end{aligned}$$

$$\begin{aligned} W_4(x) &= \det \begin{bmatrix} 1 & \frac{1}{x^3} & x^2 \\ 0 & -\frac{3}{x^4} & 2x \\ 0 & \frac{12}{x^5} & 2 \end{bmatrix} \\ &= -\frac{30}{x^4} \end{aligned}$$

Now we are ready to evaluate each $U_i(x)$.

$$\begin{aligned} U_1 &= (-1)^{4-1} \int \frac{F(x)W_1(x)}{aW(x)} dx \\ &= (-1)^3 \int \frac{(17x^6)(180x^2)}{(x^4)\left(-\frac{6480}{x}\right)} dx \\ &= - \int \frac{3060x^8}{-6480x^3} dx \\ &= - \int \left(-\frac{17x^5}{36}\right) dx \\ &= \frac{17x^6}{216} \end{aligned}$$

$$\begin{aligned}
U_2 &= (-1)^{4-2} \int \frac{F(x)W_2(x)}{aW(x)} dx \\
&= (-1)^2 \int \frac{(17x^6)(48x^5)}{(x^4)\left(-\frac{6480}{x}\right)} dx \\
&= \int \frac{816x^{11}}{-6480x^3} dx \\
&= \int \left(-\frac{17x^8}{135}\right) dx \\
&= -\frac{17x^9}{1215}
\end{aligned}$$

$$\begin{aligned}
U_3 &= (-1)^{4-3} \int \frac{F(x)W_3(x)}{aW(x)} dx \\
&= (-1)^1 \int \frac{(17x^6)\left(-\frac{162}{x}\right)}{(x^4)\left(-\frac{6480}{x}\right)} dx \\
&= - \int \frac{-2754x^6}{-6480x^3} dx \\
&= - \int \left(\frac{17x^3}{40}\right) dx \\
&= -\frac{17x^4}{160}
\end{aligned}$$

$$\begin{aligned}
U_4 &= (-1)^{4-4} \int \frac{F(x)W_4(x)}{aW(x)} dx \\
&= (-1)^0 \int \frac{(17x^6)\left(-\frac{30}{x^4}\right)}{(x^4)\left(-\frac{6480}{x}\right)} dx \\
&= \int \frac{-510x^2}{-6480x^3} dx \\
&= \int \left(\frac{17}{216x}\right) dx \\
&= \frac{17 \ln(x)}{216}
\end{aligned}$$

Now that all the U_i functions have been determined, the particular solution is found from

$$y_p = U_1y_1 + U_2y_2 + U_3y_3 + U_4y_4$$

Hence

$$\begin{aligned}
 y_p &= \left(\frac{17x^6}{216} \right) \\
 &+ \left(-\frac{17x^9}{1215} \right) \left(\frac{1}{x^3} \right) \\
 &+ \left(-\frac{17x^4}{160} \right) (x^2) \\
 &+ \left(\frac{17 \ln(x)}{216} \right) (x^6)
 \end{aligned}$$

Therefore the particular solution is

$$y_p = \frac{17x^6(-19 + 36 \ln(x))}{7776}$$

Therefore the general solution is

$$\begin{aligned}
 y &= y_h + y_p \\
 &= \left(c_1 + \frac{c_2}{x^3} + c_3 x^2 + c_4 x^6 \right) + \left(\frac{17x^6(-19 + 36 \ln(x))}{7776} \right)
 \end{aligned}$$

Maple step by step solution

Let's solve

$$x^4 y'''' + x^3 y''' - 20x^2 y'' + 20xy' = 17x^6$$

- Highest derivative means the order of the ODE is 4
 y''''

Maple trace

```

`Methods for high order ODEs:
--- Trying classification methods ---
trying a quadrature
trying high order exact linear fully integrable
trying differential order: 4; linear nonhomogeneous with symmetry [0,1]
-> Calling odsolve with the ODE`, diff(diff(diff(_b(_a), _a), _a), _a) = -(-17*_a^5+(c
Methods for third order ODEs:
--- Trying classification methods ---
trying a quadrature
trying high order exact linear fully integrable

```

```

trying differential order: 3; linear nonhomogeneous with symmetry [0,1]
trying high order linear exact nonhomogeneous
trying differential order: 3; missing the dependent variable
checking if the LODE is of Euler type
<- LODE of Euler type successful
Euler equation successful
<- differential order: 4; linear nonhomogeneous with symmetry [0,1] successful`

```

Maple dsolve solution

Solving time : 0.003 (sec)

Leaf size : 41

```

dsolve(x^4*diff(diff(diff(diff(y(x),x),x),x),x)+x^3*diff(diff(diff(y(x),x),x),x)-20*x^2*diff(diff(y(x),x),x)-20*x*diff(y(x),x)+17*x^6,y(x),singsol=all)

```

$$y = \frac{612 \ln(x) x^9 + (1296c_3 - 323) x^9 + 3888c_1 x^5 + 7776c_4 x^3 - 2592c_2}{7776x^3}$$

Mathematica DSolve solution

Solving time : 0.014 (sec)

Leaf size : 49

```

DSolve[{x^4*D[y[x],{x,4}]+x^3*D[y[x],{x,3}]-20*x^2*D[y[x],{x,2}]+20*x*D[y[x],x]==17*x^6,{}}],y[x],x,IncludeSingularSolutions->True]

```

$$y(x) \rightarrow \frac{17}{216} x^6 \log(x) + \left(-\frac{323}{7776} + \frac{c_3}{6} \right) x^6 - \frac{c_1}{3x^3} + \frac{c_2 x^2}{2} + c_4$$

2.8.12 problem 15

Solved as higher order Euler type ode	708
Maple step by step solution	714
Maple trace	714
Maple dsolve solution	715
Mathematica DSolve solution	715

Internal problem ID [18267]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter VII. Linear equations of order higher than the first. section 56. Problems at page 163

Problem number : 15

Date solved : Monday, December 23, 2024 at 09:47:15 PM

CAS classification : [[_high_order, _exact, _linear, _nonhomogeneous]]

Solve

$$t^4 x'''' - 2t^3 x''' - 20t^2 x'' + 12x't + 16x = \cos(3 \ln(t))$$

Solved as higher order Euler type ode

Time used: 0.786 (sec)

This is Euler ODE of higher order. Let $x = t^\lambda$. Hence

$$\begin{aligned} x' &= \lambda t^{\lambda-1} \\ x'' &= \lambda(\lambda-1) t^{\lambda-2} \\ x''' &= \lambda(\lambda-1)(\lambda-2) t^{\lambda-3} \\ x'''' &= \lambda(\lambda-1)(\lambda-2)(\lambda-3) t^{\lambda-4} \end{aligned}$$

Substituting these back into

$$t^4 x'''' - 2t^3 x''' - 20t^2 x'' + 12x't + 16x = \cos(3 \ln(t))$$

gives

$$\begin{aligned} 12t\lambda t^{\lambda-1} - 20t^2\lambda(\lambda-1) t^{\lambda-2} - 2t^3\lambda(\lambda-1)(\lambda-2) t^{\lambda-3} \\ + t^4\lambda(\lambda-1)(\lambda-2)(\lambda-3) t^{\lambda-4} + 16t^\lambda = 0 \end{aligned}$$

Which simplifies to

$$12\lambda t^\lambda - 20\lambda(\lambda-1) t^\lambda - 2\lambda(\lambda-1)(\lambda-2) t^\lambda + \lambda(\lambda-1)(\lambda-2)(\lambda-3) t^\lambda + 16t^\lambda = 0$$

And since $t^\lambda \neq 0$ then dividing through by t^λ , the above becomes

$$12\lambda - 20\lambda(\lambda - 1) - 2\lambda(\lambda - 1)(\lambda - 2) + \lambda(\lambda - 1)(\lambda - 2)(\lambda - 3) + 16 = 0$$

Simplifying gives the characteristic equation as

$$(\lambda - 2)(\lambda - 8)(\lambda + 1)^2 = 0$$

Solving the above gives the following roots

$$\lambda_1 = 2$$

$$\lambda_2 = 8$$

$$\lambda_3 = -1$$

$$\lambda_4 = -1$$

This table summarises the result

root	multiplicity	type of root
-1	2	real root
2	1	real root
8	1	real root

The solution is generated by going over the above table. For each real root λ of multiplicity one generates a $c_1 t^\lambda$ basis solution. Each real root of multiplicity two, generates $c_1 t^\lambda$ and $c_2 t^\lambda \ln(t)$ basis solutions. Each real root of multiplicity three, generates $c_1 t^\lambda$ and $c_2 t^\lambda \ln(t)$ and $c_3 t^\lambda \ln(t)^2$ basis solutions, and so on. Each complex root $\alpha \pm i\beta$ of multiplicity one generates $t^\alpha (c_1 \cos(\beta \ln(t)) + c_2 \sin(\beta \ln(t)))$ basis solutions. And each complex root $\alpha \pm i\beta$ of multiplicity two generates $\ln(t) t^\alpha (c_1 \cos(\beta \ln(t)) + c_2 \sin(\beta \ln(t)))$ basis solutions. And each complex root $\alpha \pm i\beta$ of multiplicity three generates $\ln(t)^2 t^\alpha (c_1 \cos(\beta \ln(t)) + c_2 \sin(\beta \ln(t)))$ basis solutions. And so on. Using the above show that the solution is

$$x = \frac{c_1}{t} + \frac{c_2 \ln(t)}{t} + c_3 t^2 + c_4 t^8$$

The fundamental set of solutions for the homogeneous solution are the following

$$x_1 = \frac{1}{t}$$

$$x_2 = \frac{\ln(t)}{t}$$

$$x_3 = t^2$$

$$x_4 = t^8$$

This is higher order nonhomogeneous Euler type ODE. Let the solution be

$$x = x_h + x_p$$

Where x_h is the solution to the homogeneous Euler ODE And x_p is a particular solution to the nonhomogeneous Euler ODE. x_h is the solution to

$$t^4 x'''' - 2t^3 x''' - 20t^2 x'' + 12x't + 16x = 0$$

Now the particular solution to the given ODE is found

$$t^4 x'''' - 2t^3 x''' - 20t^2 x'' + 12x't + 16x = \cos(3 \ln(t))$$

Let the particular solution be

$$x_p = U_1 x_1 + U_2 x_2 + U_3 x_3 + U_4 x_4$$

Where x_i are the basis solutions found above for the homogeneous solution x_h and $U_i(t)$ are functions to be determined as follows

$$U_i = (-1)^{n-i} \int \frac{F(t)W_i(t)}{aW(t)} dt$$

Where $W(t)$ is the Wronskian and $W_i(t)$ is the Wronskian that results after deleting the last row and the i -th column of the determinant and n is the order of the ODE or equivalently, the number of basis solutions, and a is the coefficient of the leading derivative in the ODE, and $F(t)$ is the RHS of the ODE. Therefore, the first step is to find the Wronskian $W(t)$. This is given by

$$W(t) = \begin{vmatrix} x_1 & x_2 & x_3 & x_4 \\ x_1' & x_2' & x_3' & x_4' \\ x_1'' & x_2'' & x_3'' & x_4'' \\ x_1''' & x_2''' & x_3''' & x_4''' \end{vmatrix}$$

Substituting the fundamental set of solutions x_i found above in the Wronskian gives

$$W = \begin{bmatrix} \frac{1}{t} & \frac{\ln(t)}{t} & t^2 & t^8 \\ -\frac{1}{t^2} & \frac{1-\ln(t)}{t^2} & 2t & 8t^7 \\ \frac{2}{t^3} & \frac{-3+2\ln(t)}{t^3} & 2 & 56t^6 \\ -\frac{6}{t^4} & \frac{11-6\ln(t)}{t^4} & 0 & 336t^5 \end{bmatrix}$$

$$|W| = 4374t^2$$

The determinant simplifies to

$$|W| = 4374t^2$$

Now we determine W_i for each U_i .

$$\begin{aligned} W_1(t) &= \det \begin{bmatrix} \frac{\ln(t)}{t} & t^2 & t^8 \\ \frac{1-\ln(t)}{t^2} & 2t & 8t^7 \\ \frac{-3+2\ln(t)}{t^3} & 2 & 56t^6 \end{bmatrix} \\ &= 18t^6(-4 + 9 \ln(t)) \end{aligned}$$

$$\begin{aligned} W_2(t) &= \det \begin{bmatrix} \frac{1}{t} & t^2 & t^8 \\ -\frac{1}{t^2} & 2t & 8t^7 \\ \frac{2}{t^3} & 2 & 56t^6 \end{bmatrix} \\ &= 162t^6 \end{aligned}$$

$$\begin{aligned} W_3(t) &= \det \begin{bmatrix} \frac{1}{t} & \frac{\ln(t)}{t} & t^8 \\ -\frac{1}{t^2} & \frac{1-\ln(t)}{t^2} & 8t^7 \\ \frac{2}{t^3} & \frac{-3+2\ln(t)}{t^3} & 56t^6 \end{bmatrix} \\ &= 81t^3 \end{aligned}$$

$$\begin{aligned} W_4(t) &= \det \begin{bmatrix} \frac{1}{t} & \frac{\ln(t)}{t} & t^2 \\ -\frac{1}{t^2} & \frac{1-\ln(t)}{t^2} & 2t \\ \frac{2}{t^3} & \frac{-3+2\ln(t)}{t^3} & 2 \end{bmatrix} \\ &= \frac{9}{t^3} \end{aligned}$$

Now we are ready to evaluate each $U_i(t)$.

$$\begin{aligned}
 U_1 &= (-1)^{4-1} \int \frac{F(t)W_1(t)}{aW(t)} dt \\
 &= (-1)^3 \int \frac{(\cos(3 \ln(t))) (18t^6(-4 + 9 \ln(t)))}{(t^4)(4374t^2)} dt \\
 &= - \int \frac{18 \cos(3 \ln(t)) t^6(-4 + 9 \ln(t))}{4374t^6} dt \\
 &= - \int \left(\frac{\cos(3 \ln(t))(-4 + 9 \ln(t))}{243} \right) dt \\
 &= - \frac{\left(\frac{8}{25} + \frac{9 \ln(t)}{10}\right) t \cos(3 \ln(t))}{243} + \frac{\left(\frac{87}{50} - \frac{27 \ln(t)}{10}\right) t \sin(3 \ln(t))}{243}
 \end{aligned}$$

$$\begin{aligned}
 U_2 &= (-1)^{4-2} \int \frac{F(t)W_2(t)}{aW(t)} dt \\
 &= (-1)^2 \int \frac{(\cos(3 \ln(t))) (162t^6)}{(t^4)(4374t^2)} dt \\
 &= \int \frac{162 \cos(3 \ln(t)) t^6}{4374t^6} dt \\
 &= \int \left(\frac{\cos(3 \ln(t))}{27} \right) dt \\
 &= \frac{\cos(3 \ln(t)) t}{270} + \frac{t \sin(3 \ln(t))}{90}
 \end{aligned}$$

$$\begin{aligned}
 U_3 &= (-1)^{4-3} \int \frac{F(t)W_3(t)}{aW(t)} dt \\
 &= (-1)^1 \int \frac{(\cos(3 \ln(t))) (81t^3)}{(t^4)(4374t^2)} dt \\
 &= - \int \frac{81 \cos(3 \ln(t)) t^3}{4374t^6} dt \\
 &= - \int \left(\frac{\cos(3 \ln(t))}{54t^3} \right) dt \\
 &= - \frac{\frac{1}{351} + \frac{\tan\left(\frac{3 \ln(t)}{2}\right)^2}{351} + \frac{\tan\left(\frac{3 \ln(t)}{2}\right)}{117}}{\left(1 + \tan\left(\frac{3 \ln(t)}{2}\right)\right)^2} t^2
 \end{aligned}$$

$$\begin{aligned}
U_4 &= (-1)^{4-4} \int \frac{F(t)W_4(t)}{aW(t)} dt \\
&= (-1)^0 \int \frac{(\cos(3 \ln(t))) \left(\frac{9}{t^3}\right)}{(t^4)(4374t^2)} dt \\
&= \int \frac{\frac{9 \cos(3 \ln(t))}{t^3}}{4374t^6} dt \\
&= \int \left(\frac{\cos(3 \ln(t))}{486t^9} \right) dt \\
&= \frac{\left(-\frac{2}{17739} - \frac{i}{23652}\right) t^{3i}}{t^8} + \frac{\left(-\frac{2}{17739} + \frac{i}{23652}\right) t^{-3i}}{t^8}
\end{aligned}$$

Now that all the U_i functions have been determined, the particular solution is found from

$$x_p = U_1x_1 + U_2x_2 + U_3x_3 + U_4x_4$$

Hence

$$\begin{aligned}
x_p &= \left(-\frac{\left(\frac{8}{25} + \frac{9 \ln(t)}{10}\right) t \cos(3 \ln(t))}{243} + \frac{\left(\frac{87}{50} - \frac{27 \ln(t)}{10}\right) t \sin(3 \ln(t))}{243} \right) \left(\frac{1}{t} \right) \\
&+ \left(\frac{\cos(3 \ln(t)) t}{270} + \frac{t \sin(3 \ln(t))}{90} \right) \left(\frac{\ln(t)}{t} \right) \\
&+ \left(\frac{-\frac{1}{351} + \frac{\tan\left(\frac{3 \ln(t)}{2}\right)^2}{351} + \frac{\tan\left(\frac{3 \ln(t)}{2}\right)}{117}}{\left(1 + \tan\left(\frac{3 \ln(t)}{2}\right)^2\right) t^2} \right) (t^2) \\
&+ \left(\frac{\left(-\frac{2}{17739} - \frac{i}{23652}\right) t^{3i}}{t^8} + \frac{\left(-\frac{2}{17739} + \frac{i}{23652}\right) t^{-3i}}{t^8} \right) (t^8)
\end{aligned}$$

Therefore the particular solution is

$$x_p = \left(\frac{31}{47450} - \frac{141i}{94900} \right) t^{-3i} t^{6i} + \left(\frac{31}{47450} + \frac{141i}{94900} \right) t^{-3i}$$

Therefore the general solution is

$$x = x_h + x_p$$

$$= \left(\frac{c_1}{t} + \frac{c_2 \ln(t)}{t} + c_3 t^2 + c_4 t^8 \right) + \left(\left(\frac{31}{47450} - \frac{141i}{94900} \right) t^{-3i} t^{6i} + \left(\frac{31}{47450} + \frac{141i}{94900} \right) t^{-3i} \right)$$

Maple step by step solution

Let's solve

$$t^4 x'''' - 2t^3 x''' - 20t^2 x'' + 12x't + 16x = \cos(3 \ln(t))$$

- Highest derivative means the order of the ODE is 4
 x''''

Maple trace

```

`Methods for high order ODEs:
--- Trying classification methods ---
trying a quadrature
trying high order exact linear fully integrable
trying differential order: 4; linear nonhomogeneous with symmetry [0,1]
trying high order linear exact nonhomogeneous
-> Calling odsolve with the ODE`, diff(diff(diff(_b(_a), _a), _a), _a) = (_C1-16*_b(_a)
  Methods for third order ODEs:
  --- Trying classification methods ---
  trying a quadrature
  trying high order exact linear fully integrable
  trying differential order: 3; linear nonhomogeneous with symmetry [0,1]
  trying high order linear exact nonhomogeneous
  -> Calling odsolve with the ODE`, diff(diff(_g(_f), _f), _f) = _C2+8*_g(_f)/_f^2+6*_
    Methods for second order ODEs:
    --- Trying classification methods ---
    trying a quadrature
    trying high order exact linear fully integrable
    trying differential order: 2; linear nonhomogeneous with symmetry [0,1]
    trying a double symmetry of the form [xi=0, eta=F(x)]
    -> Try solving first the homogeneous part of the ODE
        checking if the LODE has constant coefficients
        checking if the LODE is of Euler type
        <- LODE of Euler type successful
        <- solving first the homogeneous part of the ODE successful
        <- high order exact_linear_nonhomogeneous successful
    <- high order exact_linear_nonhomogeneous successful`

```

Maple dsolve solution

Solving time : 0.002 (sec)

Leaf size : 43

```
dsolve(t^4*diff(diff(diff(diff(x(t),t),t),t),t)-2*t^3*diff(diff(diff(x(t),t),t),t)-20*
x(t),singsol=all)
```

$$x = \frac{(15066 + 34263i)t^{1-3i} + (15066 - 34263i)t^{1+3i} + 23060700t^9c_3 - 1281150c_2t^3 + 854100c_1 \ln(t) + 9490}{23060700t}$$

Mathematica DSolve solution

Solving time : 0.079 (sec)

Leaf size : 48

```
DSolve[{t^4*D[x[t],{t,4}]-2*t^3*D[x[t],{t,3}]-20*t^2*D[x[t],{t,2}]+12*t*D[x[t],t]+16*x[t]==0,
x[t],t,IncludeSingularSolutions->True]
```

$$x(t) \rightarrow \frac{c_4 t^9 + c_3 t^3 + c_2 \log(t) + c_1}{t} + \frac{141 \sin(3 \log(t))}{47450} + \frac{31 \cos(3 \log(t))}{23725}$$

2.9 Chapter VII. Linear equations of order higher than the first. section 63. Problems at page 196

2.9.1	problem 1	717
2.9.2	problem 2	720
2.9.3	problem 3	726
2.9.4	problem 8	733

2.9.1 problem 1

Solved as higher order constant coeff ode	717
Maple step by step solution	718
Maple trace	718
Maple dsolve solution	718
Mathematica DSolve solution	719

Internal problem ID [18268]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter VII. Linear equations of order higher than the first. section 63. Problems at page 196

Problem number : 1

Date solved : Monday, December 23, 2024 at 09:47:17 PM

CAS classification : [[_3rd_order, _missing_x]]

Solve

$$y''' - y'' - y' + y = 0$$

Solved as higher order constant coeff ode

Time used: 0.186 (sec)

The characteristic equation is

$$\lambda^3 - \lambda^2 - \lambda + 1 = 0$$

The roots of the above equation are

$$\lambda_1 = -1$$

$$\lambda_2 = 1$$

$$\lambda_3 = 1$$

Therefore the homogeneous solution is

$$y_h(x) = e^{-x} c_1 + e^x c_2 + x e^x c_3$$

The fundamental set of solutions for the homogeneous solution are the following

$$y_1 = e^{-x}$$

$$y_2 = e^x$$

$$y_3 = x e^x$$

Maple step by step solution

Let's solve

$$y''' - y'' - y' + y = 0$$

- Highest derivative means the order of the ODE is 3
 y'''
- Characteristic polynomial of ODE
 $r^3 - r^2 - r + 1 = 0$
- Roots of the characteristic polynomial and corresponding multiplicities
 $r = [[-1, 1], [1, 2]]$
- Solution from $r = -1$
 $y_1(x) = e^{-x}$
- 1st solution from $r = 1$
 $y_2(x) = e^x$
- 2nd solution from $r = 1$
 $y_3(x) = x e^x$
- General solution of the ODE
 $y = C1y_1(x) + C2y_2(x) + C3y_3(x)$
- Substitute in solutions and simplify
 $y = e^{-x} C1 + e^x(C3x + C2)$

Maple trace

```
`Methods for third order ODEs:
--- Trying classification methods ---
trying a quadrature
checking if the LODE has constant coefficients
<- constant coefficients successful`
```

Maple dsolve solution

Solving time : 0.003 (sec)

Leaf size : 19

```
dsolve(diff(diff(diff(y(x),x),x),x)-diff(diff(y(x),x),x)-diff(y(x),x)+y(x) = 0,
      y(x),singsol=all)
```

$$y = e^{-x}c_1 + e^x(c_3x + c_2)$$

Mathematica DSolve solution

Solving time : 0.003 (sec)

Leaf size : 25

```
DSolve[{D[y[x],{x,3}]-D[y[x],{x,2}]-D[y[x],x]+y[x]==0,{}},  
y[x],x,IncludeSingularSolutions->True]
```

$$y(x) \rightarrow c_1 e^{-x} + e^x (c_3 x + c_2)$$

2.9.2 problem 2

Solved as higher order constant coeff ode	720
Maple step by step solution	722
Maple trace	724
Maple dsolve solution	725
Mathematica DSolve solution	725

Internal problem ID [18269]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter VII. Linear equations of order higher than the first. section 63. Problems at page 196

Problem number : 2

Date solved : Monday, December 23, 2024 at 09:47:17 PM

CAS classification : [[_high_order, _missing_y]]

Solve

$$y'''' - 3y''' + 3y'' - y' = e^{2x}$$

Solved as higher order constant coeff ode

Time used: 0.109 (sec)

The characteristic equation is

$$\lambda^4 - 3\lambda^3 + 3\lambda^2 - \lambda = 0$$

The roots of the above equation are

$$\lambda_1 = 0$$

$$\lambda_2 = 1$$

$$\lambda_3 = 1$$

$$\lambda_4 = 1$$

Therefore the homogeneous solution is

$$y_h(x) = c_1 + e^x c_2 + x e^x c_3 + x^2 e^x c_4$$

The fundamental set of solutions for the homogeneous solution are the following

$$\begin{aligned}y_1 &= 1 \\y_2 &= e^x \\y_3 &= x e^x \\y_4 &= x^2 e^x\end{aligned}$$

This is higher order nonhomogeneous ODE. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE And y_p is a particular solution to the nonhomogeneous ODE. y_h is the solution to

$$y'''' - 3y''' + 3y'' - y' = 0$$

Now the particular solution to the given ODE is found

$$y'''' - 3y''' + 3y'' - y' = e^{2x}$$

The particular solution is now found using the method of undetermined coefficients.

Looking at the RHS of the ode, which is

$$e^{2x}$$

Shows that the corresponding undetermined set of the basis functions (UC_set) for the trial solution is

$$[\{e^{2x}\}]$$

While the set of the basis functions for the homogeneous solution found earlier is

$$\{1, x e^x, x^2 e^x, e^x\}$$

Since there is no duplication between the basis function in the UC_set and the basis functions of the homogeneous solution, the trial solution is a linear combination of all the basis in the UC_set.

$$y_p = A_1 e^{2x}$$

The unknowns $\{A_1\}$ are found by substituting the above trial solution y_p into the ODE and comparing coefficients. Substituting the trial solution into the ODE and simplifying gives

$$2A_1 e^{2x} = e^{2x}$$

Solving for the unknowns by comparing coefficients results in

$$\left[A_1 = \frac{1}{2} \right]$$

Substituting the above back in the above trial solution y_p , gives the particular solution

$$y_p = \frac{e^{2x}}{2}$$

Therefore the general solution is

$$\begin{aligned} y &= y_h + y_p \\ &= (c_1 + e^x c_2 + x e^x c_3 + x^2 e^x c_4) + \left(\frac{e^{2x}}{2} \right) \end{aligned}$$

Maple step by step solution

Let's solve

$$y'''' - 3y''' + 3y'' - y' = e^{2x}$$

- Highest derivative means the order of the ODE is 4
 y''''
- Characteristic polynomial of homogeneous ODE
 $r^4 - 3r^3 + 3r^2 - r = 0$
- Roots of the characteristic polynomial and corresponding multiplicities
 $r = [[0, 1], [1, 3]]$
- Homogeneous solution from $r = 0$
 $y_1(x) = 1$
- 1st homogeneous solution from $r = 1$
 $y_2(x) = e^x$
- 2nd homogeneous solution from $r = 1$
 $y_3(x) = x e^x$
- 3rd homogeneous solution from $r = 1$
 $y_4(x) = x^2 e^x$
- General solution of the ODE
 $y = C1 y_1(x) + C2 y_2(x) + C3 y_3(x) + C4 y_4(x) + y_p(x)$
- Substitute in solutions of the homogeneous ODE
 $y = C1 + e^x C2 + x e^x C3 + x^2 e^x C4 + y_p(x)$

□ Find a particular solution $y_p(x)$ of the ODE

- Define the forcing function of the ODE

$$f(x) = e^{2x}$$

- Form of the particular solution to the ODE where the $u_i(x)$ are to be found

$$y_p(x) = \sum_{i=1}^4 u_i(x) y_i(x)$$

- Calculate the 1st derivative of $y_p(x)$

$$y_p'(x) = \sum_{i=1}^4 (u_i'(x) y_i(x) + u_i(x) y_i'(x))$$

- Choose equation to add to a system of equations in $u_i'(x)$

$$\sum_{i=1}^4 u_i'(x) y_i(x) = 0$$

- Calculate the 2nd derivative of $y_p(x)$

$$y_p''(x) = \sum_{i=1}^4 (u_i'(x) y_i'(x) + u_i(x) y_i''(x))$$

- Choose equation to add to a system of equations in $u_i'(x)$

$$\sum_{i=1}^4 u_i'(x) y_i'(x) = 0$$

- Calculate the 3rd derivative of $y_p(x)$

$$y_p'''(x) = \sum_{i=1}^4 (u_i'(x) y_i''(x) + u_i(x) y_i'''(x))$$

- Choose equation to add to a system of equations in $u_i'(x)$

$$\sum_{i=1}^4 u_i'(x) y_i''(x) = 0$$

- The ODE is of the following form where the $P_i(x)$ in this situation are the coefficients of the

$$y'''' + \left(\sum_{i=0}^3 P_i(x) y^{(i)} \right) = f(x)$$

- Substitute $y_p(x) = \sum_{i=1}^4 u_i(x) y_i(x)$ into the ODE

$$\left(\sum_{j=0}^3 P_j(x) \left(\sum_{i=1}^4 u_i(x) y_i^{(j)}(x) \right) \right) + \sum_{i=1}^4 (u_i'(x) y_i'''(x) + u_i(x) y_i''''(x)) = f(x)$$

- Rearrange the ODE

$$\sum_{i=1}^4 \left(u_i(x) \cdot \left(\left(\sum_{j=0}^3 P_j(x) y_i^{(j)}(x) \right) + y_i''''(x) \right) + u_i'(x) y_i'''(x) \right) = f(x)$$

- Notice that $y_i(x)$ are solutions to the homogeneous equation so the first term in the sum is 0

$$\sum_{i=1}^4 u_i'(x) y_i'''(x) = f(x)$$

- We have now made a system of 4 equations in 4 unknowns ($u_i'(x)$)

$$\left[\sum_{i=1}^4 u_i'(x) y_i(x) = 0, \sum_{i=1}^4 u_i'(x) y_i'(x) = 0, \sum_{i=1}^4 u_i'(x) y_i''(x) = 0, \sum_{i=1}^4 u_i'(x) y_i'''(x) = f(x) \right]$$

- Convert the system to linear algebra format, notice that the matrix is the wronskian W

$$\begin{bmatrix} y_1(x) & y_2(x) & y_3(x) & y_4(x) \\ y_1'(x) & y_2'(x) & y_3'(x) & y_4'(x) \\ y_1''(x) & y_2''(x) & y_3''(x) & y_4''(x) \\ y_1'''(x) & y_2'''(x) & y_3'''(x) & y_4'''(x) \end{bmatrix} \cdot \begin{bmatrix} u_1'(x) \\ u_2'(x) \\ u_3'(x) \\ u_4'(x) \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \\ f(x) \end{bmatrix}$$

- Solve for the varied parameters

$$\begin{bmatrix} u_1(x) \\ u_2(x) \\ u_3(x) \\ u_4(x) \end{bmatrix} = \int \frac{1}{W} \cdot \begin{bmatrix} 0 \\ 0 \\ 0 \\ f(x) \end{bmatrix} dx$$

- Substitute in the homogeneous solutions and forcing function and solve

$$\begin{bmatrix} u_1(x) \\ u_2(x) \\ u_3(x) \\ u_4(x) \end{bmatrix} = \begin{bmatrix} -\frac{e^{2x}}{2} \\ \frac{(x^2+2)e^{2x}}{2e^x} \\ -\frac{x e^{2x}}{e^x} \\ \frac{e^{2x}}{2e^x} \end{bmatrix}$$

Find a particular solution $y_p(x)$ of the ODE

$$y_p(x) = \frac{e^{2x}}{2}$$

- Substitute particular solution into general solution to ODE

$$y = C_1 + e^x C_2 + x e^x C_3 + x^2 e^x C_4 + \frac{e^{2x}}{2}$$

Maple trace

```

`Methods for high order ODEs:
--- Trying classification methods ---
trying a quadrature
trying high order exact linear fully integrable
trying differential order: 4; linear nonhomogeneous with symmetry [0,1]
-> Calling odsolve with the ODE`, diff(diff(diff(_b(_a), _a), _a), _a) = 3*(diff(diff

```



```

Methods for third order ODEs:
--- Trying classification methods ---
trying a quadrature
trying high order exact linear fully integrable
trying differential order: 3; linear nonhomogeneous with symmetry [0,1]
trying high order linear exact nonhomogeneous
trying differential order: 3; missing the dependent variable
checking if the LODE has constant coefficients
<- constant coefficients successful
<- differential order: 4; linear nonhomogeneous with symmetry [0,1] successful`

```

Maple dsolve solution

Solving time : 0.004 (sec)

Leaf size : 33

```

dsolve(diff(diff(diff(diff(y(x),x),x),x),x)-3*diff(diff(diff(y(x),x),x),x)+3*diff(diff(y(x),x),x),x),
        y(x),singsol=all)

```

$$y = \frac{e^{2x}}{2} + ((x^2 - 2x + 2)c_3 + c_2x + c_1 - c_2)e^x + c_4$$

Mathematica DSolve solution

Solving time : 0.063 (sec)

Leaf size : 41

```

DSolve[{D[y[x],{x,4}]-3*D[y[x],{x,3}]+3*D[y[x],{x,2}]-D[y[x],x]==Exp[2*x],{}}],
        y[x],x,IncludeSingularSolutions->True]

```

$$y(x) \rightarrow \frac{1}{2}e^x(e^x + 2(c_3(x^2 - 2x + 2) + c_2(x - 1) + c_1)) + c_4$$

2.9.3 problem 3

Solved as higher order constant coeff ode	726
Maple step by step solution	730
Maple trace	732
Maple dsolve solution	732
Mathematica DSolve solution	732

Internal problem ID [18270]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter VII. Linear equations of order higher than the first. section 63. Problems at page 196

Problem number : 3

Date solved : Monday, December 23, 2024 at 09:47:18 PM

CAS classification : [[_3rd_order, _linear, _nonhomogeneous]]

Solve

$$y''' - y'' + y' - y = \cos(x)$$

Solved as higher order constant coeff ode

Time used: 0.620 (sec)

The characteristic equation is

$$\lambda^3 - \lambda^2 + \lambda - 1 = 0$$

The roots of the above equation are

$$\lambda_1 = 1$$

$$\lambda_2 = i$$

$$\lambda_3 = -i$$

Therefore the homogeneous solution is

$$y_h(x) = c_1 e^x + e^{-ix} c_2 + e^{ix} c_3$$

The fundamental set of solutions for the homogeneous solution are the following

$$y_1 = e^x$$

$$y_2 = e^{-ix}$$

$$y_3 = e^{ix}$$

This is higher order nonhomogeneous ODE. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE And y_p is a particular solution to the nonhomogeneous ODE. y_h is the solution to

$$y''' - y'' + y' - y = 0$$

Now the particular solution to the given ODE is found

$$y''' - y'' + y' - y = \cos(x)$$

Let the particular solution be

$$y_p = U_1 y_1 + U_2 y_2 + U_3 y_3$$

Where y_i are the basis solutions found above for the homogeneous solution y_h and $U_i(x)$ are functions to be determined as follows

$$U_i = (-1)^{n-i} \int \frac{F(x)W_i(x)}{aW(x)} dx$$

Where $W(x)$ is the Wronskian and $W_i(x)$ is the Wronskian that results after deleting the last row and the i -th column of the determinant and n is the order of the ODE or equivalently, the number of basis solutions, and a is the coefficient of the leading derivative in the ODE, and $F(x)$ is the RHS of the ODE. Therefore, the first step is to find the Wronskian $W(x)$. This is given by

$$W(x) = \begin{vmatrix} y_1 & y_2 & y_3 \\ y_1' & y_2' & y_3' \\ y_1'' & y_2'' & y_3'' \end{vmatrix}$$

Substituting the fundamental set of solutions y_i found above in the Wronskian gives

$$W = \begin{bmatrix} e^x & e^{-ix} & e^{ix} \\ e^x & -ie^{-ix} & ie^{ix} \\ e^x & -e^{-ix} & -e^{ix} \end{bmatrix}$$

$$|W| = 4ie^x e^{-ix} e^{ix}$$

The determinant simplifies to

$$|W| = 4ie^x$$

Now we determine W_i for each U_i .

$$\begin{aligned} W_1(x) &= \det \begin{bmatrix} e^{-ix} & e^{ix} \\ -ie^{-ix} & ie^{ix} \end{bmatrix} \\ &= 2i \end{aligned}$$

$$\begin{aligned} W_2(x) &= \det \begin{bmatrix} e^x & e^{ix} \\ e^x & ie^{ix} \end{bmatrix} \\ &= (-1 + i) e^{(1+i)x} \end{aligned}$$

$$\begin{aligned} W_3(x) &= \det \begin{bmatrix} e^x & e^{-ix} \\ e^x & -ie^{-ix} \end{bmatrix} \\ &= (-1 - i) e^{(1-i)x} \end{aligned}$$

Now we are ready to evaluate each $U_i(x)$.

$$\begin{aligned} U_1 &= (-1)^{3-1} \int \frac{F(x)W_1(x)}{aW(x)} dx \\ &= (-1)^2 \int \frac{(\cos(x))(2i)}{(1)(4ie^x)} dx \\ &= \int \frac{2i \cos(x)}{4ie^x} dx \\ &= \int \left(\frac{\cos(x) e^{-x}}{2} \right) dx \\ &= -\frac{\cos(x) e^{-x}}{4} + \frac{e^{-x} \sin(x)}{4} \end{aligned}$$

$$\begin{aligned} U_2 &= (-1)^{3-2} \int \frac{F(x)W_2(x)}{aW(x)} dx \\ &= (-1)^1 \int \frac{(\cos(x))((-1+i)e^{(1+i)x})}{(1)(4ie^x)} dx \\ &= - \int \frac{(-1+i) \cos(x) e^{(1+i)x}}{4ie^x} dx \\ &= - \int \left(\left(\frac{1}{4} + \frac{i}{4} \right) \cos(x) e^{ix} \right) dx \\ &= -\frac{x}{8} - \frac{ix}{8} - \frac{e^{2ix}}{16} + \frac{ie^{2ix}}{16} \\ &= -\frac{x}{8} - \frac{ix}{8} - \frac{e^{2ix}}{16} + \frac{ie^{2ix}}{16} \end{aligned}$$

$$\begin{aligned}
U_3 &= (-1)^{3-3} \int \frac{F(x)W_3(x)}{aW(x)} dx \\
&= (-1)^0 \int \frac{(\cos(x))((-1-i)e^{(1-i)x})}{(1)(4ie^x)} dx \\
&= \int \frac{(-1-i)\cos(x)e^{(1-i)x}}{4ie^x} dx \\
&= \int \left(\left(-\frac{1}{4} + \frac{i}{4} \right) \cos(x)e^{-ix} \right) dx \\
&= \int \left(-\frac{1}{4} + \frac{i}{4} \right) \cos(x)e^{-ix} dx
\end{aligned}$$

Now that all the U_i functions have been determined, the particular solution is found from

$$y_p = U_1y_1 + U_2y_2 + U_3y_3$$

Hence

$$\begin{aligned}
y_p &= \left(-\frac{\cos(x)e^{-x}}{4} + \frac{e^{-x}\sin(x)}{4} \right) (e^x) \\
&\quad + \left(-\frac{x}{8} - \frac{ix}{8} - \frac{e^{2ix}}{16} + \frac{ie^{2ix}}{16} \right) (e^{-ix}) \\
&\quad + \left(\int \left(-\frac{1}{4} + \frac{i}{4} \right) \cos(x)e^{-ix} dx \right) (e^{ix})
\end{aligned}$$

Therefore the particular solution is

$$y_p = \frac{(-5+i-4x)\cos(x)}{16} + \frac{(1+i-4x)\sin(x)}{16}$$

Which simplifies to

$$y_p = \frac{(-5+i-4x)\cos(x)}{16} + \frac{(1+i-4x)\sin(x)}{16}$$

Therefore the general solution is

$$\begin{aligned}
y &= y_h + y_p \\
&= (c_1 e^x + e^{-ix}c_2 + e^{ix}c_3) + \left(\frac{(-5+i-4x)\cos(x)}{16} + \frac{(1+i-4x)\sin(x)}{16} \right)
\end{aligned}$$

Maple step by step solution

Let's solve

$$y''' - y'' + y' - y = \cos(x)$$

- Highest derivative means the order of the ODE is 3

$$y'''$$

- Characteristic polynomial of homogeneous ODE

$$r^3 - r^2 + r - 1 = 0$$

- Roots of the characteristic polynomial

$$r = [1, I, -I]$$

- Homogeneous solution from $r = 1$

$$y_1(x) = e^x$$

- Homogeneous solutions from $r = I$ and $r = -I$

$$[y_2(x) = \sin(x), y_3(x) = \cos(x)]$$

- General solution of the ODE

$$y = C1 y_1(x) + C2 y_2(x) + C3 y_3(x) + y_p(x)$$

- Substitute in solutions of the homogeneous ODE

$$y = C1 e^x + C2 \sin(x) + C3 \cos(x) + y_p(x)$$

- Find a particular solution $y_p(x)$ of the ODE

- Define the forcing function of the ODE

$$f(x) = \cos(x)$$

- Form of the particular solution to the ODE where the $u_i(x)$ are to be found

$$y_p(x) = \sum_{i=1}^3 u_i(x) y_i(x)$$

- Calculate the 1st derivative of $y_p(x)$

$$y_p'(x) = \sum_{i=1}^3 (u_i'(x) y_i(x) + u_i(x) y_i'(x))$$

- Choose equation to add to a system of equations in $u_i'(x)$

$$\sum_{i=1}^3 u_i'(x) y_i(x) = 0$$

- Calculate the 2nd derivative of $y_p(x)$

$$y_p''(x) = \sum_{i=1}^3 (u_i'(x) y_i'(x) + u_i(x) y_i''(x))$$

- Choose equation to add to a system of equations in $u_i'(x)$

$$\sum_{i=1}^3 u_i'(x) y_i'(x) = 0$$

- The ODE is of the following form where the $P_i(x)$ in this situation are the coefficients of the

$$y''' + \left(\sum_{i=0}^2 P_i(x) y^{(i)} \right) = f(x)$$

- Substitute $y_p(x) = \sum_{i=1}^3 u_i(x) y_i(x)$ into the ODE

$$\left(\sum_{j=0}^2 P_j(x) \left(\sum_{i=1}^3 u_i(x) y_i^{(j)}(x) \right) \right) + \sum_{i=1}^3 (u_i'(x) y_i''(x) + u_i(x) y_i'''(x)) = f(x)$$

- Rearrange the ODE

$$\sum_{i=1}^3 \left(u_i(x) \cdot \left(\left(\sum_{j=0}^2 P_j(x) y_i^{(j)}(x) \right) + y_i'''(x) \right) + u_i'(x) y_i''(x) \right) = f(x)$$

- Notice that $y_i(x)$ are solutions to the homogeneous equation so the first term in the sum is 0

$$\sum_{i=1}^3 u_i'(x) y_i''(x) = f(x)$$

- We have now made a system of 3 equations in 3 unknowns ($u_i'(x)$)

$$\left[\sum_{i=1}^3 u_i'(x) y_i(x) = 0, \sum_{i=1}^3 u_i'(x) y_i'(x) = 0, \sum_{i=1}^3 u_i'(x) y_i''(x) = f(x) \right]$$

- Convert the system to linear algebra format, notice that the matrix is the wronskian W

$$\begin{bmatrix} y_1(x) & y_2(x) & y_3(x) \\ y_1'(x) & y_2'(x) & y_3'(x) \\ y_1''(x) & y_2''(x) & y_3''(x) \end{bmatrix} \cdot \begin{bmatrix} u_1'(x) \\ u_2'(x) \\ u_3'(x) \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ f(x) \end{bmatrix}$$

- Solve for the varied parameters

$$\begin{bmatrix} u_1(x) \\ u_2(x) \\ u_3(x) \end{bmatrix} = \int \frac{1}{W} \cdot \begin{bmatrix} 0 \\ 0 \\ f(x) \end{bmatrix} dx$$

- Substitute in the homogeneous solutions and forcing function and solve

$$\begin{bmatrix} u_1(x) \\ u_2(x) \\ u_3(x) \end{bmatrix} = \begin{bmatrix} -\frac{\cos(x)e^{-x}}{4} + \frac{e^{-x} \sin(x)}{4} \\ -\frac{x}{4} - \frac{\sin(2x)}{8} + \frac{\cos(2x)}{8} \\ -\frac{x}{4} - \frac{\sin(2x)}{8} - \frac{\cos(2x)}{8} \end{bmatrix}$$

Find a particular solution $y_p(x)$ of the ODE

$$y_p(x) = \frac{(-2x-3)\cos(x)}{8} + \frac{(-2x+1)\sin(x)}{8}$$

- Substitute particular solution into general solution to ODE

$$y = C1 e^x + C2 \sin(x) + C3 \cos(x) + \frac{(-2x-3)\cos(x)}{8} + \frac{(-2x+1)\sin(x)}{8}$$

Maple trace

```

`Methods for third order ODEs:
--- Trying classification methods ---
trying a quadrature
trying high order exact linear fully integrable
trying differential order: 3; linear nonhomogeneous with symmetry [0,1]
trying high order linear exact nonhomogeneous
trying differential order: 3; missing the dependent variable
checking if the LODE has constant coefficients
<- constant coefficients successful`

```

Maple dsolve solution

Solving time : 0.003 (sec)

Leaf size : 33

```

dsolve(diff(diff(diff(y(x),x),x),x)-diff(diff(y(x),x),x)+diff(y(x),x)-y(x) = cos(x),
        y(x),singsol=all)

```

$$y = \frac{(4c_1 - x - 2) \cos(x)}{4} + \frac{(-x + 4c_3 + 1) \sin(x)}{4} + e^x c_2$$

Mathematica DSolve solution

Solving time : 0.038 (sec)

Leaf size : 40

```

DSolve[{D[y[x],{x,3}]-D[y[x],{x,2}]+D[y[x],x]-y[x]==Cos[x],{}}],
        y[x],x,IncludeSingularSolutions->True]

```

$$y(x) \rightarrow \frac{1}{4}(4c_3 e^x - (x + 2 - 4c_1) \cos(x) + (-x + 1 + 4c_2) \sin(x))$$

2.9.4 problem 8

Solved as second order Euler type ode	733
Solved as second order linear exact ode	737
Solved as second order integrable as is ode	739
Solved as second order integrable as is ode (ABC method) . . .	740
Solved as second order ode using change of variable on x method 2	741
Solved as second order ode using change of variable on y method 2	751
Solved as second order ode using Kovacic algorithm	756
Solved as second order ode adjoint method	763
Maple step by step solution	766
Maple trace	766
Maple dsolve solution	766
Mathematica DSolve solution	767

Internal problem ID [18271]

Book : Elementary Differential Equations. By Thornton C. Fry. D Van Nostrand. NY. First Edition (1929)

Section : Chapter VII. Linear equations of order higher than the first. section 63. Problems at page 196

Problem number : 8

Date solved : Monday, December 23, 2024 at 09:47:19 PM

CAS classification : [[_2nd_order, _exact, _linear, _nonhomogeneous]]

Solve

$$x^2y'' + 3xy' + y = \frac{1}{x}$$

Solved as second order Euler type ode

Time used: 0.298 (sec)

This is second order non-homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = f(x)$$

Where $A = x^2$, $B = 3x$, $C = 1$, $f(x) = \frac{1}{x}$. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE $Ay''(x) + By'(x) + Cy(x) = 0$, and y_p is a particular solution to the non-homogeneous ODE $Ay''(x) + By'(x) + Cy(x) = f(x)$. Solving for y_h from

$$x^2y'' + 3xy' + y = 0$$

This is Euler second order ODE. Let the solution be $y = x^r$, then $y' = rx^{r-1}$ and $y'' = r(r-1)x^{r-2}$. Substituting these back into the given ODE gives

$$x^2(r(r-1))x^{r-2} + 3rx^{r-1} + x^r = 0$$

Simplifying gives

$$r(r-1)x^r + 3rx^r + x^r = 0$$

Since $x^r \neq 0$ then dividing throughout by x^r gives

$$r(r-1) + 3r + 1 = 0$$

Or

$$r^2 + 2r + 1 = 0 \tag{1}$$

Equation (1) is the characteristic equation. Its roots determine the form of the general solution. Using the quadratic equation the roots are

$$r_1 = -1$$

$$r_2 = -1$$

Since the roots are equal, then the general solution is

$$y = c_1y_1 + c_2y_2$$

Where $y_1 = x^r$ and $y_2 = x^r \ln(x)$. Hence

$$y = \frac{c_1}{x} + \frac{c_2 \ln(x)}{x}$$

Next, we find the particular solution to the ODE

$$x^2y'' + 3xy' + y = \frac{1}{x}$$

The particular solution y_p can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of parameters will be used as it is more general and can be used when the coefficients of the ODE depend on x as well. Let

$$y_p(x) = u_1y_1 + u_2y_2 \tag{1}$$

Where u_1, u_2 to be determined, and y_1, y_2 are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as

$$y_1 = \frac{1}{x}$$

$$y_2 = \frac{\ln(x)}{x}$$

In the Variation of parameters u_1, u_2 are found using

$$u_1 = - \int \frac{y_2 f(x)}{aW(x)} \quad (2)$$

$$u_2 = \int \frac{y_1 f(x)}{aW(x)} \quad (3)$$

Where $W(x)$ is the Wronskian and a is the coefficient in front of y'' in the given ODE.

The Wronskian is given by $W = \begin{vmatrix} y_1 & y_2 \\ y_1' & y_2' \end{vmatrix}$. Hence

$$W = \begin{vmatrix} \frac{1}{x} & \frac{\ln(x)}{x} \\ \frac{d}{dx} \left(\frac{1}{x} \right) & \frac{d}{dx} \left(\frac{\ln(x)}{x} \right) \end{vmatrix}$$

Which gives

$$W = \begin{vmatrix} \frac{1}{x} & \frac{\ln(x)}{x} \\ -\frac{1}{x^2} & -\frac{\ln(x)}{x^2} + \frac{1}{x^2} \end{vmatrix}$$

Therefore

$$W = \left(\frac{1}{x} \right) \left(-\frac{\ln(x)}{x^2} + \frac{1}{x^2} \right) - \left(\frac{\ln(x)}{x} \right) \left(-\frac{1}{x^2} \right)$$

Which simplifies to

$$W = \frac{1}{x^3}$$

Which simplifies to

$$W = \frac{1}{x^3}$$

Therefore Eq. (2) becomes

$$u_1 = - \int \frac{\frac{\ln(x)}{x^2}}{\frac{1}{x}} dx$$

Which simplifies to

$$u_1 = - \int \frac{\ln(x)}{x} dx$$

Hence

$$u_1 = - \frac{\ln(x)^2}{2}$$

And Eq. (3) becomes

$$u_2 = \int \frac{\frac{1}{x^2}}{\frac{1}{x}} dx$$

Which simplifies to

$$u_2 = \int \frac{1}{x} dx$$

Hence

$$u_2 = \ln(x)$$

Therefore the particular solution, from equation (1) is

$$y_p(x) = \frac{\ln(x)^2}{2x}$$

Therefore the general solution is

$$\begin{aligned} y &= y_h + y_p \\ &= \frac{\ln(x)^2 + 2c_2 \ln(x) + 2c_1}{2x} \end{aligned}$$

Will add steps showing solving for IC soon.

Summary of solutions found

$$y = \frac{\ln(x)^2 + 2c_2 \ln(x) + 2c_1}{2x}$$

Solved as second order linear exact ode

Time used: 0.100 (sec)

An ode of the form

$$p(x)y'' + q(x)y' + r(x)y = s(x)$$

is exact if

$$p''(x) - q'(x) + r(x) = 0 \tag{1}$$

For the given ode we have

$$p(x) = x^2$$

$$q(x) = 3x$$

$$r(x) = 1$$

$$s(x) = \frac{1}{x}$$

Hence

$$p''(x) = 2$$

$$q'(x) = 3$$

Therefore (1) becomes

$$2 - (3) + (1) = 0$$

Hence the ode is exact. Since we now know the ode is exact, it can be written as

$$(p(x)y' + (q(x) - p'(x))y)' = s(x)$$

Integrating gives

$$p(x)y' + (q(x) - p'(x))y = \int s(x) dx$$

Substituting the above values for p, q, r, s gives

$$x^2y' + yx = \int \frac{1}{x} dx$$

We now have a first order ode to solve which is

$$x^2y' + yx = \ln(x) + c_1$$

In canonical form a linear first order is

$$y' + q(x)y = p(x)$$

Comparing the above to the given ode shows that

$$q(x) = \frac{1}{x}$$

$$p(x) = \frac{\ln(x) + c_1}{x^2}$$

The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int q dx} \\ &= e^{\int \frac{1}{x} dx} \\ &= x\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dx}(\mu y) &= \mu p \\ \frac{d}{dx}(\mu y) &= (\mu) \left(\frac{\ln(x) + c_1}{x^2} \right) \\ \frac{d}{dx}(yx) &= (x) \left(\frac{\ln(x) + c_1}{x^2} \right) \\ d(yx) &= \left(\frac{\ln(x) + c_1}{x} \right) dx\end{aligned}$$

Integrating gives

$$\begin{aligned}yx &= \int \frac{\ln(x) + c_1}{x} dx \\ &= \frac{\ln(x)^2}{2} + c_1 \ln(x) + c_2\end{aligned}$$

Dividing throughout by the integrating factor x gives the final solution

$$y = \frac{\frac{\ln(x)^2}{2} + c_1 \ln(x) + c_2}{x}$$

Will add steps showing solving for IC soon.

Summary of solutions found

$$y = \frac{\frac{\ln(x)^2}{2} + c_1 \ln(x) + c_2}{x}$$

Solved as second order integrable as is ode

Time used: 0.185 (sec)

Integrating both sides of the ODE w.r.t x gives

$$\int (x^2 y'' + 3xy' + y) dx = \int \frac{1}{x} dx$$

$$x^2 y' + yx = \ln(x) + c_1$$

Which is now solved for y . In canonical form a linear first order is

$$y' + q(x)y = p(x)$$

Comparing the above to the given ode shows that

$$q(x) = \frac{1}{x}$$

$$p(x) = \frac{\ln(x) + c_1}{x^2}$$

The integrating factor μ is

$$\begin{aligned} \mu &= e^{\int q dx} \\ &= e^{\int \frac{1}{x} dx} \\ &= x \end{aligned}$$

The ode becomes

$$\frac{d}{dx}(\mu y) = \mu p$$

$$\frac{d}{dx}(\mu y) = (\mu) \left(\frac{\ln(x) + c_1}{x^2} \right)$$

$$\frac{d}{dx}(yx) = (x) \left(\frac{\ln(x) + c_1}{x^2} \right)$$

$$d(yx) = \left(\frac{\ln(x) + c_1}{x} \right) dx$$

Integrating gives

$$\begin{aligned} yx &= \int \frac{\ln(x) + c_1}{x} dx \\ &= \frac{\ln(x)^2}{2} + c_1 \ln(x) + c_2 \end{aligned}$$

Dividing throughout by the integrating factor x gives the final solution

$$y = \frac{\frac{\ln(x)^2}{2} + c_1 \ln(x) + c_2}{x}$$

Will add steps showing solving for IC soon.

Summary of solutions found

$$y = \frac{\frac{\ln(x)^2}{2} + c_1 \ln(x) + c_2}{x}$$

Solved as second order integrable as is ode (ABC method)

Time used: 0.066 (sec)

Writing the ode as

$$x^2 y'' + 3xy' + y = \frac{1}{x}$$

Integrating both sides of the ODE w.r.t x gives

$$\int (x^2 y'' + 3xy' + y) dx = \int \frac{1}{x} dx$$

$$x^2 y' + yx = \ln(x) + c_1$$

Which is now solved for y . In canonical form a linear first order is

$$y' + q(x)y = p(x)$$

Comparing the above to the given ode shows that

$$q(x) = \frac{1}{x}$$

$$p(x) = \frac{\ln(x) + c_1}{x^2}$$

The integrating factor μ is

$$\begin{aligned} \mu &= e^{\int q dx} \\ &= e^{\int \frac{1}{x} dx} \\ &= x \end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dx}(\mu y) &= \mu p \\ \frac{d}{dx}(\mu y) &= (\mu) \left(\frac{\ln(x) + c_1}{x^2} \right) \\ \frac{d}{dx}(yx) &= (x) \left(\frac{\ln(x) + c_1}{x^2} \right) \\ d(yx) &= \left(\frac{\ln(x) + c_1}{x} \right) dx\end{aligned}$$

Integrating gives

$$\begin{aligned}yx &= \int \frac{\ln(x) + c_1}{x} dx \\ &= \frac{\ln(x)^2}{2} + c_1 \ln(x) + c_2\end{aligned}$$

Dividing throughout by the integrating factor x gives the final solution

$$y = \frac{\frac{\ln(x)^2}{2} + c_1 \ln(x) + c_2}{x}$$

Will add steps showing solving for IC soon.

Solved as second order ode using change of variable on x method 2

Time used: 0.495 (sec)

This is second order non-homogeneous ODE. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE $Ay''(x) + By'(x) + Cy(x) = 0$, and y_p is a particular solution to the non-homogeneous ODE $Ay''(x) + By'(x) + Cy(x) = f(x)$. y_h is the solution to

$$x^2 y'' + 3xy' + y = 0$$

In normal form the ode

$$x^2 y'' + 3xy' + y = 0 \tag{1}$$

Becomes

$$y'' + p(x)y' + q(x)y = 0 \quad (2)$$

Where

$$p(x) = \frac{3}{x}$$

$$q(x) = \frac{1}{x^2}$$

Applying change of variables $\tau = g(x)$ to (2) gives

$$\frac{d^2}{d\tau^2}y(\tau) + p_1\left(\frac{d}{d\tau}y(\tau)\right) + q_1y(\tau) = 0 \quad (3)$$

Where τ is the new independent variable, and

$$p_1(\tau) = \frac{\tau''(x) + p(x)\tau'(x)}{\tau'(x)^2} \quad (4)$$

$$q_1(\tau) = \frac{q(x)}{\tau'(x)^2} \quad (5)$$

Let $p_1 = 0$. Eq (4) simplifies to

$$\tau''(x) + p(x)\tau'(x) = 0$$

This ode is solved resulting in

$$\begin{aligned} \tau &= \int e^{-\int p(x)dx} dx \\ &= \int e^{-\int \frac{3}{x} dx} dx \\ &= \int e^{-3\ln(x)} dx \\ &= \int \frac{1}{x^3} dx \\ &= -\frac{1}{2x^2} \end{aligned} \quad (6)$$

Using (6) to evaluate q_1 from (5) gives

$$\begin{aligned} q_1(\tau) &= \frac{q(x)}{\tau'(x)^2} \\ &= \frac{\frac{1}{x^2}}{\frac{1}{x^6}} \\ &= x^4 \end{aligned} \quad (7)$$

Substituting the above in (3) and noting that now $p_1 = 0$ results in

$$\begin{aligned}\frac{d^2}{d\tau^2}y(\tau) + q_1y(\tau) &= 0 \\ \frac{d^2}{d\tau^2}y(\tau) + x^4y(\tau) &= 0\end{aligned}$$

But in terms of τ

$$x^4 = \frac{1}{4\tau^2}$$

Hence the above ode becomes

$$\frac{d^2}{d\tau^2}y(\tau) + \frac{y(\tau)}{4\tau^2} = 0$$

The above ode is now solved for $y(\tau)$. Writing the ode as

$$\frac{d^2}{d\tau^2}y(\tau) + \frac{y(\tau)}{4\tau^2} = 0 \quad (1)$$

$$A\frac{d^2}{d\tau^2}y(\tau) + B\frac{d}{d\tau}y(\tau) + Cy(\tau) = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$\begin{aligned}A &= 1 \\ B &= 0 \\ C &= \frac{1}{4\tau^2}\end{aligned} \quad (3)$$

Applying the Liouville transformation on the dependent variable gives

$$z(\tau) = y(\tau) e^{\int \frac{B}{2A} d\tau}$$

Then (2) becomes

$$z''(\tau) = rz(\tau) \quad (4)$$

Where r is given by

$$\begin{aligned}r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2}\end{aligned} \quad (5)$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{-1}{4\tau^2} \quad (6)$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= -1 \\ t &= 4\tau^2 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(\tau) = \left(-\frac{1}{4\tau^2}\right) z(\tau) \quad (7)$$

Equation (7) is now solved. After finding $z(\tau)$ then $y(\tau)$ is found using the inverse transformation

$$y(\tau) = z(\tau) e^{-\int \frac{B}{2A} d\tau}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 2.68: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} O(\infty) &= \deg(t) - \deg(s) \\ &= 2 - 0 \\ &= 2 \end{aligned}$$

The poles of r in eq. (7) and the order of each pole are determined by solving for the roots of $t = 4\tau^2$. There is a pole at $\tau = 0$ of order 2. Since there is no odd order pole larger than 2 and the order at ∞ is 2 then the necessary conditions for case one are met. Since there is a pole of order 2 then necessary conditions for case two are met. Since pole order is not larger than 2 and the order at ∞ is 2 then the necessary conditions for case three are met. Therefore

$$L = [1, 2, 4, 6, 12]$$

Attempting to find a solution using case $n = 1$.

Looking at poles of order 2. The partial fractions decomposition of r is

$$r = -\frac{1}{4\tau^2}$$

For the pole at $\tau = 0$ let b be the coefficient of $\frac{1}{\tau^2}$ in the partial fractions decomposition of r given above. Therefore $b = -\frac{1}{4}$. Hence

$$\begin{aligned} [\sqrt{r}]_c &= 0 \\ \alpha_c^+ &= \frac{1}{2} + \sqrt{1 + 4b} = \frac{1}{2} \\ \alpha_c^- &= \frac{1}{2} - \sqrt{1 + 4b} = \frac{1}{2} \end{aligned}$$

Since the order of r at ∞ is 2 then $[\sqrt{r}]_\infty = 0$. Let b be the coefficient of $\frac{1}{\tau^2}$ in the Laurent series expansion of r at ∞ , which can be found by dividing the leading coefficient of s by the leading coefficient of t from

$$r = \frac{s}{t} = -\frac{1}{4\tau^2}$$

Since the $\gcd(s, t) = 1$. This gives $b = -\frac{1}{4}$. Hence

$$\begin{aligned} [\sqrt{r}]_\infty &= 0 \\ \alpha_\infty^+ &= \frac{1}{2} + \sqrt{1 + 4b} = \frac{1}{2} \\ \alpha_\infty^- &= \frac{1}{2} - \sqrt{1 + 4b} = \frac{1}{2} \end{aligned}$$

The following table summarizes the findings so far for poles and for the order of r at ∞ where r is

$$r = -\frac{1}{4\tau^2}$$

pole c location	pole order	$[\sqrt{r}]_c$	α_c^+	α_c^-
0	2	0	$\frac{1}{2}$	$\frac{1}{2}$

Order of r at ∞	$[\sqrt{r}]_\infty$	α_∞^+	α_∞^-
2	0	$\frac{1}{2}$	$\frac{1}{2}$

Now that the all $[\sqrt{r}]_c$ and its associated α_c^\pm have been determined for all the poles in the set Γ and $[\sqrt{r}]_\infty$ and its associated α_∞^\pm have also been found, the next step is to determine possible non negative integer d from these using

$$d = \alpha_\infty^{s(\infty)} - \sum_{c \in \Gamma} \alpha_c^{s(c)}$$

Where $s(c)$ is either $+$ or $-$ and $s(\infty)$ is the sign of α_∞^\pm . This is done by trial over all set of families $s = (s(c))_{c \in \Gamma \cup \infty}$ until such d is found to work in finding candidate ω . Trying $\alpha_\infty^- = \frac{1}{2}$ then

$$\begin{aligned} d &= \alpha_\infty^- - (\alpha_{c_1}^+) \\ &= \frac{1}{2} - \left(\frac{1}{2}\right) \\ &= 0 \end{aligned}$$

Since d an integer and $d \geq 0$ then it can be used to find ω using

$$\omega = \sum_{c \in \Gamma} \left(s(c) [\sqrt{r}]_c + \frac{\alpha_c^{s(c)}}{\tau - c} \right) + s(\infty) [\sqrt{r}]_\infty$$

The above gives

$$\begin{aligned} \omega &= \left((+) [\sqrt{r}]_{c_1} + \frac{\alpha_{c_1}^+}{\tau - c_1} \right) + (-) [\sqrt{r}]_\infty \\ &= \frac{1}{2\tau} + (-) (0) \\ &= \frac{1}{2\tau} \\ &= \frac{1}{2\tau} \end{aligned}$$

Now that ω is determined, the next step is find a corresponding minimal polynomial $p(\tau)$ of degree $d = 0$ to solve the ode. The polynomial $p(\tau)$ needs to satisfy the equation

$$p'' + 2\omega p' + (\omega' + \omega^2 - r) p = 0 \quad (1A)$$

Let

$$p(\tau) = 1 \quad (2A)$$

Substituting the above in eq. (1A) gives

$$(0) + 2\left(\frac{1}{2\tau}\right)(0) + \left(\left(-\frac{1}{2\tau^2}\right) + \left(\frac{1}{2\tau}\right)^2 - \left(-\frac{1}{4\tau^2}\right)\right) = 0$$

$$0 = 0$$

The equation is satisfied since both sides are zero. Therefore the first solution to the ode $z'' = rz$ is

$$\begin{aligned} z_1(\tau) &= pe^{\int \omega d\tau} \\ &= e^{\int \frac{1}{2\tau} d\tau} \\ &= \sqrt{\tau} \end{aligned}$$

The first solution to the original ode in $y(\tau)$ is found from

$$y_1 = z_1 e^{\int -\frac{1}{2} \frac{B}{A} d\tau}$$

Since $B = 0$ then the above reduces to

$$\begin{aligned} y_1 &= z_1 \\ &= \sqrt{\tau} \end{aligned}$$

Which simplifies to

$$y_1 = \sqrt{\tau}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} d\tau}}{y_1^2} d\tau$$

Since $B = 0$ then the above becomes

$$\begin{aligned} y_2 &= y_1 \int \frac{1}{y_1^2} d\tau \\ &= \sqrt{\tau} \int \frac{1}{\tau} d\tau \\ &= \sqrt{\tau} (\ln(\tau)) \end{aligned}$$

Therefore the solution is

$$\begin{aligned} y(\tau) &= c_1 y_1 + c_2 y_2 \\ &= c_1(\sqrt{\tau}) + c_2(\sqrt{\tau}(\ln(\tau))) \end{aligned}$$

Will add steps showing solving for IC soon.

The above solution is now transformed back to y using (6) which results in

$$y = c_1 \sqrt{-\frac{1}{2x^2}} + c_2 \sqrt{-\frac{1}{2x^2}} \ln\left(-\frac{1}{2x^2}\right)$$

Therefore the homogeneous solution y_h is

$$y_h = c_1 \sqrt{-\frac{1}{2x^2}} + c_2 \sqrt{-\frac{1}{2x^2}} \ln\left(-\frac{1}{2x^2}\right)$$

The particular solution y_p can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of parameters will be used as it is more general and can be used when the coefficients of the ODE depend on x as well. Let

$$y_p(x) = u_1 y_1 + u_2 y_2 \tag{1}$$

Where u_1, u_2 to be determined, and y_1, y_2 are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as

$$\begin{aligned} y_1 &= \sqrt{-\frac{1}{2x^2}} \\ y_2 &= -\sqrt{-\frac{1}{2x^2}} \ln(2) + \sqrt{-\frac{1}{2x^2}} \ln\left(-\frac{1}{x^2}\right) \end{aligned}$$

In the Variation of parameters u_1, u_2 are found using

$$u_1 = -\int \frac{y_2 f(x)}{aW(x)} \tag{2}$$

$$u_2 = \int \frac{y_1 f(x)}{aW(x)} \tag{3}$$

Where $W(x)$ is the Wronskian and a is the coefficient in front of y'' in the given ODE.

The Wronskian is given by $W = \begin{vmatrix} y_1 & y_2 \\ y_1' & y_2' \end{vmatrix}$. Hence

$$W = \begin{vmatrix} \sqrt{-\frac{1}{2x^2}} & -\sqrt{-\frac{1}{2x^2}} \ln(2) + \sqrt{-\frac{1}{2x^2}} \ln\left(-\frac{1}{x^2}\right) \\ \frac{d}{dx}\left(\sqrt{-\frac{1}{2x^2}}\right) & \frac{d}{dx}\left(-\sqrt{-\frac{1}{2x^2}} \ln(2) + \sqrt{-\frac{1}{2x^2}} \ln\left(-\frac{1}{x^2}\right)\right) \end{vmatrix}$$

Which gives

$$W = \begin{vmatrix} \sqrt{-\frac{1}{2x^2}} & -\sqrt{-\frac{1}{2x^2}} \ln(2) + \sqrt{-\frac{1}{2x^2}} \ln\left(-\frac{1}{x^2}\right) \\ \frac{1}{2\sqrt{-\frac{1}{2x^2}} x^3} & -\frac{\ln(2)}{2\sqrt{-\frac{1}{2x^2}} x^3} + \frac{\ln\left(-\frac{1}{x^2}\right)}{2\sqrt{-\frac{1}{2x^2}} x^3} - \frac{2\sqrt{-\frac{1}{2x^2}}}{x} \end{vmatrix}$$

Therefore

$$W = \left(\sqrt{-\frac{1}{2x^2}}\right) \left(-\frac{\ln(2)}{2\sqrt{-\frac{1}{2x^2}} x^3} + \frac{\ln\left(-\frac{1}{x^2}\right)}{2\sqrt{-\frac{1}{2x^2}} x^3} - \frac{2\sqrt{-\frac{1}{2x^2}}}{x}\right) - \left(-\sqrt{-\frac{1}{2x^2}} \ln(2) + \sqrt{-\frac{1}{2x^2}} \ln\left(-\frac{1}{x^2}\right)\right) \left(\frac{1}{2\sqrt{-\frac{1}{2x^2}} x^3}\right)$$

Which simplifies to

$$W = \frac{1}{x^3}$$

Which simplifies to

$$W = \frac{1}{x^3}$$

Therefore Eq. (2) becomes

$$u_1 = - \int \frac{-\sqrt{-\frac{1}{2x^2}} \ln(2) + \sqrt{-\frac{1}{2x^2}} \ln\left(-\frac{1}{x^2}\right)}{\frac{x}{\frac{1}{x}}} dx$$

Which simplifies to

$$u_1 = - \int -\frac{\sqrt{2} \sqrt{-\frac{1}{x^2}} (\ln(2) - \ln\left(-\frac{1}{x^2}\right))}{2} dx$$

Hence

$$u_1 = \frac{\sqrt{2} \sqrt{-\frac{1}{x^2}} x \ln(2) \ln(x)}{2} + \frac{\sqrt{2} \sqrt{-\frac{1}{x^2}} x \ln\left(-\frac{1}{x^2}\right)^2}{8}$$

And Eq. (3) becomes

$$u_2 = \int \frac{\frac{\sqrt{-\frac{1}{2x^2}}}{x}}{\frac{1}{x}} dx$$

Which simplifies to

$$u_2 = \int \frac{\sqrt{2} \sqrt{-\frac{1}{x^2}}}{2} dx$$

Hence

$$u_2 = \frac{\sqrt{2} \sqrt{-\frac{1}{x^2}} x \ln(x)}{2}$$

Therefore the particular solution, from equation (1) is

$$y_p(x) = \left(\frac{\sqrt{2} \sqrt{-\frac{1}{x^2}} x \ln(2) \ln(x)}{2} + \frac{\sqrt{2} \sqrt{-\frac{1}{x^2}} x \ln\left(-\frac{1}{x^2}\right)^2}{8} \right) \sqrt{-\frac{1}{2x^2}} \\ + \frac{\left(-\sqrt{-\frac{1}{2x^2}} \ln(2) + \sqrt{-\frac{1}{2x^2}} \ln\left(-\frac{1}{x^2}\right) \right) \sqrt{2} \sqrt{-\frac{1}{x^2}} x \ln(x)}{2}$$

Which simplifies to

$$y_p(x) = -\frac{\ln\left(-\frac{1}{x^2}\right) \left(\ln\left(-\frac{1}{x^2}\right) + 4 \ln(x) \right)}{8x}$$

Therefore the general solution is

$$y = y_h + y_p \\ = \left(c_1 \sqrt{-\frac{1}{2x^2}} + c_2 \sqrt{-\frac{1}{2x^2}} \ln\left(-\frac{1}{2x^2}\right) \right) + \left(-\frac{\ln\left(-\frac{1}{x^2}\right) \left(\ln\left(-\frac{1}{x^2}\right) + 4 \ln(x) \right)}{8x} \right)$$

Will add steps showing solving for IC soon.

Summary of solutions found

$$y = -\frac{\ln\left(-\frac{1}{x^2}\right)\left(\ln\left(-\frac{1}{x^2}\right) + 4\ln(x)\right)}{8x} + c_1\sqrt{-\frac{1}{2x^2}} + c_2\sqrt{-\frac{1}{2x^2}}\ln\left(-\frac{1}{2x^2}\right)$$

Solved as second order ode using change of variable on y method 2

Time used: 0.318 (sec)

This is second order non-homogeneous ODE. In standard form the ODE is

$$Ay''(x) + By'(x) + Cy(x) = f(x)$$

Where $A = x^2$, $B = 3x$, $C = 1$, $f(x) = \frac{1}{x}$. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE $Ay''(x) + By'(x) + Cy(x) = 0$, and y_p is a particular solution to the non-homogeneous ODE $Ay''(x) + By'(x) + Cy(x) = f(x)$. Solving for y_h from

$$x^2y'' + 3xy' + y = 0$$

In normal form the ode

$$x^2y'' + 3xy' + y = 0 \tag{1}$$

Becomes

$$y'' + p(x)y' + q(x)y = 0 \tag{2}$$

Where

$$p(x) = \frac{3}{x}$$

$$q(x) = \frac{1}{x^2}$$

Applying change of variables on the dependent variable $y = v(x)x^n$ to (2) gives the following ode where the dependent variables is $v(x)$ and not y .

$$v''(x) + \left(\frac{2n}{x} + p\right)v'(x) + \left(\frac{n(n-1)}{x^2} + \frac{np}{x} + q\right)v(x) = 0 \tag{3}$$

Let the coefficient of $v(x)$ above be zero. Hence

$$\frac{n(n-1)}{x^2} + \frac{np}{x} + q = 0 \quad (4)$$

Substituting the earlier values found for $p(x)$ and $q(x)$ into (4) gives

$$\frac{n(n-1)}{x^2} + \frac{3n}{x^2} + \frac{1}{x^2} = 0 \quad (5)$$

Solving (5) for n gives

$$n = -1 \quad (6)$$

Substituting this value in (3) gives

$$\begin{aligned} v''(x) + \frac{v'(x)}{x} &= 0 \\ v''(x) + \frac{v'(x)}{x} &= 0 \end{aligned} \quad (7)$$

Using the substitution

$$u(x) = v'(x)$$

Then (7) becomes

$$u'(x) + \frac{u(x)}{x} = 0 \quad (8)$$

The above is now solved for $u(x)$. In canonical form a linear first order is

$$u'(x) + q(x)u(x) = p(x)$$

Comparing the above to the given ode shows that

$$\begin{aligned} q(x) &= \frac{1}{x} \\ p(x) &= 0 \end{aligned}$$

The integrating factor μ is

$$\begin{aligned} \mu &= e^{\int q dx} \\ &= e^{\int \frac{1}{x} dx} \\ &= x \end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dx}\mu u &= 0 \\ \frac{d}{dx}(ux) &= 0\end{aligned}$$

Integrating gives

$$\begin{aligned}ux &= \int 0 dx + c_3 \\ &= c_3\end{aligned}$$

Dividing throughout by the integrating factor x gives the final solution

$$u(x) = \frac{c_3}{x}$$

Now that $u(x)$ is known, then

$$\begin{aligned}v'(x) &= u(x) \\ v(x) &= \int u(x) dx + c_4 \\ &= c_3 \ln(x) + c_4\end{aligned}$$

Hence

$$\begin{aligned}y &= v(x)x^n \\ &= \frac{c_3 \ln(x) + c_4}{x} \\ &= \frac{c_3 \ln(x) + c_4}{x}\end{aligned}$$

Now the particular solution to this ODE is found

$$x^2 y'' + 3xy' + y = \frac{1}{x}$$

The particular solution y_p can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of parameters will be used as it is more general and can be used when the coefficients of the ODE depend on x as well. Let

$$y_p(x) = u_1 y_1 + u_2 y_2 \tag{1}$$

Where u_1, u_2 to be determined, and y_1, y_2 are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as

$$y_1 = \frac{1}{x}$$

$$y_2 = \frac{\ln(x)}{x}$$

In the Variation of parameters u_1, u_2 are found using

$$u_1 = - \int \frac{y_2 f(x)}{aW(x)} \quad (2)$$

$$u_2 = \int \frac{y_1 f(x)}{aW(x)} \quad (3)$$

Where $W(x)$ is the Wronskian and a is the coefficient in front of y'' in the given ODE.

The Wronskian is given by $W = \begin{vmatrix} y_1 & y_2 \\ y_1' & y_2' \end{vmatrix}$. Hence

$$W = \begin{vmatrix} \frac{1}{x} & \frac{\ln(x)}{x} \\ \frac{d}{dx} \left(\frac{1}{x} \right) & \frac{d}{dx} \left(\frac{\ln(x)}{x} \right) \end{vmatrix}$$

Which gives

$$W = \begin{vmatrix} \frac{1}{x} & \frac{\ln(x)}{x} \\ -\frac{1}{x^2} & -\frac{\ln(x)}{x^2} + \frac{1}{x^2} \end{vmatrix}$$

Therefore

$$W = \left(\frac{1}{x} \right) \left(-\frac{\ln(x)}{x^2} + \frac{1}{x^2} \right) - \left(\frac{\ln(x)}{x} \right) \left(-\frac{1}{x^2} \right)$$

Which simplifies to

$$W = \frac{1}{x^3}$$

Which simplifies to

$$W = \frac{1}{x^3}$$

Therefore Eq. (2) becomes

$$u_1 = - \int \frac{\frac{\ln(x)}{x^2}}{\frac{1}{x}} dx$$

Which simplifies to

$$u_1 = - \int \frac{\ln(x)}{x} dx$$

Hence

$$u_1 = - \frac{\ln(x)^2}{2}$$

And Eq. (3) becomes

$$u_2 = \int \frac{\frac{1}{x^2}}{\frac{1}{x}} dx$$

Which simplifies to

$$u_2 = \int \frac{1}{x} dx$$

Hence

$$u_2 = \ln(x)$$

Therefore the particular solution, from equation (1) is

$$y_p(x) = \frac{\ln(x)^2}{2x}$$

Therefore the general solution is

$$\begin{aligned} y &= y_h + y_p \\ &= \left(\frac{c_3 \ln(x) + c_4}{x} \right) + \left(\frac{\ln(x)^2}{2x} \right) \end{aligned}$$

Will add steps showing solving for IC soon.

Summary of solutions found

$$y = \frac{c_3 \ln(x) + c_4}{x} + \frac{\ln(x)^2}{2x}$$

Solved as second order ode using Kovacic algorithm

Time used: 0.293 (sec)

Writing the ode as

$$x^2y'' + 3xy' + y = 0 \quad (1)$$

$$Ay'' + By' + Cy = 0 \quad (2)$$

Comparing (1) and (2) shows that

$$\begin{aligned} A &= x^2 \\ B &= 3x \\ C &= 1 \end{aligned} \quad (3)$$

Applying the Liouville transformation on the dependent variable gives

$$z(x) = ye^{\int \frac{B}{2A} dx}$$

Then (2) becomes

$$z''(x) = rz(x) \quad (4)$$

Where r is given by

$$\begin{aligned} r &= \frac{s}{t} \\ &= \frac{2AB' - 2BA' + B^2 - 4AC}{4A^2} \end{aligned} \quad (5)$$

Substituting the values of A, B, C from (3) in the above and simplifying gives

$$r = \frac{-1}{4x^2} \quad (6)$$

Comparing the above to (5) shows that

$$\begin{aligned} s &= -1 \\ t &= 4x^2 \end{aligned}$$

Therefore eq. (4) becomes

$$z''(x) = \left(-\frac{1}{4x^2}\right) z(x) \quad (7)$$

Equation (7) is now solved. After finding $z(x)$ then y is found using the inverse transformation

$$y = z(x) e^{-\int \frac{B}{2A} dx}$$

The first step is to determine the case of Kovacic algorithm this ode belongs to. There are 3 cases depending on the order of poles of r and the order of r at ∞ . The following table summarizes these cases.

Case	Allowed pole order for r	Allowed value for $\mathcal{O}(\infty)$
1	$\{0, 1, 2, 4, 6, 8, \dots\}$	$\{\dots, -6, -4, -2, 0, 2, 3, 4, 5, 6, \dots\}$
2	Need to have at least one pole that is either order 2 or odd order greater than 2. Any other pole order is allowed as long as the above condition is satisfied. Hence the following set of pole orders are all allowed. $\{1, 2\}, \{1, 3\}, \{2\}, \{3\}, \{3, 4\}, \{1, 2, 5\}$.	no condition
3	$\{1, 2\}$	$\{2, 3, 4, 5, 6, 7, \dots\}$

Table 2.69: Necessary conditions for each Kovacic case

The order of r at ∞ is the degree of t minus the degree of s . Therefore

$$\begin{aligned} \mathcal{O}(\infty) &= \deg(t) - \deg(s) \\ &= 2 - 0 \\ &= 2 \end{aligned}$$

The poles of r in eq. (7) and the order of each pole are determined by solving for the roots of $t = 4x^2$. There is a pole at $x = 0$ of order 2. Since there is no odd order pole larger than 2 and the order at ∞ is 2 then the necessary conditions for case one are met. Since there is a pole of order 2 then necessary conditions for case two are met. Since pole order is not larger than 2 and the order at ∞ is 2 then the necessary conditions for case three are met. Therefore

$$L = [1, 2, 4, 6, 12]$$

Attempting to find a solution using case $n = 1$.

Looking at poles of order 2. The partial fractions decomposition of r is

$$r = -\frac{1}{4x^2}$$

For the pole at $x = 0$ let b be the coefficient of $\frac{1}{x^2}$ in the partial fractions decomposition of r given above. Therefore $b = -\frac{1}{4}$. Hence

$$\begin{aligned} [\sqrt{r}]_c &= 0 \\ \alpha_c^+ &= \frac{1}{2} + \sqrt{1 + 4b} = \frac{1}{2} \\ \alpha_c^- &= \frac{1}{2} - \sqrt{1 + 4b} = \frac{1}{2} \end{aligned}$$

Since the order of r at ∞ is 2 then $[\sqrt{r}]_\infty = 0$. Let b be the coefficient of $\frac{1}{x^2}$ in the Laurent series expansion of r at ∞ , which can be found by dividing the leading coefficient of s by the leading coefficient of t from

$$r = \frac{s}{t} = -\frac{1}{4x^2}$$

Since the $\gcd(s, t) = 1$. This gives $b = -\frac{1}{4}$. Hence

$$\begin{aligned} [\sqrt{r}]_\infty &= 0 \\ \alpha_\infty^+ &= \frac{1}{2} + \sqrt{1 + 4b} = \frac{1}{2} \\ \alpha_\infty^- &= \frac{1}{2} - \sqrt{1 + 4b} = \frac{1}{2} \end{aligned}$$

The following table summarizes the findings so far for poles and for the order of r at ∞ where r is

$$r = -\frac{1}{4x^2}$$

pole c location	pole order	$[\sqrt{r}]_c$	α_c^+	α_c^-
0	2	0	$\frac{1}{2}$	$\frac{1}{2}$

Order of r at ∞	$[\sqrt{r}]_\infty$	α_∞^+	α_∞^-
2	0	$\frac{1}{2}$	$\frac{1}{2}$

Now that the all $[\sqrt{r}]_c$ and its associated α_c^\pm have been determined for all the poles in the set Γ and $[\sqrt{r}]_\infty$ and its associated α_∞^\pm have also been found, the next step is to determine possible non negative integer d from these using

$$d = \alpha_\infty^{s(\infty)} - \sum_{c \in \Gamma} \alpha_c^{s(c)}$$

Where $s(c)$ is either $+$ or $-$ and $s(\infty)$ is the sign of α_∞^\pm . This is done by trial over all set of families $s = (s(c))_{c \in \Gamma \cup \infty}$ until such d is found to work in finding candidate ω .

Trying $\alpha_{\infty}^{-} = \frac{1}{2}$ then

$$\begin{aligned} d &= \alpha_{\infty}^{-} - (\alpha_{c_1}^{-}) \\ &= \frac{1}{2} - \left(\frac{1}{2}\right) \\ &= 0 \end{aligned}$$

Since d an integer and $d \geq 0$ then it can be used to find ω using

$$\omega = \sum_{c \in \Gamma} \left(s(c)[\sqrt{r}]_c + \frac{\alpha_c^{s(c)}}{x-c} \right) + s(\infty)[\sqrt{r}]_{\infty}$$

The above gives

$$\begin{aligned} \omega &= \left((-)[\sqrt{r}]_{c_1} + \frac{\alpha_{c_1}^{-}}{x-c_1} \right) + (-)[\sqrt{r}]_{\infty} \\ &= \frac{1}{2x} + (-)(0) \\ &= \frac{1}{2x} \\ &= \frac{1}{2x} \end{aligned}$$

Now that ω is determined, the next step is find a corresponding minimal polynomial $p(x)$ of degree $d = 0$ to solve the ode. The polynomial $p(x)$ needs to satisfy the equation

$$p'' + 2\omega p' + (\omega' + \omega^2 - r)p = 0 \quad (1A)$$

Let

$$p(x) = 1 \quad (2A)$$

Substituting the above in eq. (1A) gives

$$\begin{aligned} (0) + 2\left(\frac{1}{2x}\right)(0) + \left(\left(-\frac{1}{2x^2}\right) + \left(\frac{1}{2x}\right)^2 - \left(-\frac{1}{4x^2}\right) \right) &= 0 \\ 0 &= 0 \end{aligned}$$

The equation is satisfied since both sides are zero. Therefore the first solution to the ode $z'' = rz$ is

$$\begin{aligned} z_1(x) &= p e^{\int \omega dx} \\ &= e^{\int \frac{1}{2x} dx} \\ &= \sqrt{x} \end{aligned}$$

The first solution to the original ode in y is found from

$$\begin{aligned} y_1 &= z_1 e^{\int -\frac{1}{2} \frac{B}{A} dx} \\ &= z_1 e^{-\int \frac{1}{2} \frac{3x}{x^2} dx} \\ &= z_1 e^{-\frac{3 \ln(x)}{2}} \\ &= z_1 \left(\frac{1}{x^{3/2}} \right) \end{aligned}$$

Which simplifies to

$$y_1 = \frac{1}{x}$$

The second solution y_2 to the original ode is found using reduction of order

$$y_2 = y_1 \int \frac{e^{\int -\frac{B}{A} dx}}{y_1^2} dx$$

Substituting gives

$$\begin{aligned} y_2 &= y_1 \int \frac{e^{\int -\frac{3x}{x^2} dx}}{(y_1)^2} dx \\ &= y_1 \int \frac{e^{-3 \ln(x)}}{(y_1)^2} dx \\ &= y_1 (\ln(x)) \end{aligned}$$

Therefore the solution is

$$\begin{aligned} y &= c_1 y_1 + c_2 y_2 \\ &= c_1 \left(\frac{1}{x} \right) + c_2 \left(\frac{1}{x} (\ln(x)) \right) \end{aligned}$$

This is second order nonhomogeneous ODE. Let the solution be

$$y = y_h + y_p$$

Where y_h is the solution to the homogeneous ODE $Ay''(x) + By'(x) + Cy(x) = 0$, and y_p is a particular solution to the nonhomogeneous ODE $Ay''(x) + By'(x) + Cy(x) = f(x)$. y_h is the solution to

$$x^2 y'' + 3xy' + y = 0$$

The homogeneous solution is found using the Kovacic algorithm which results in

$$y_h = \frac{c_1}{x} + \frac{c_2 \ln(x)}{x}$$

The particular solution y_p can be found using either the method of undetermined coefficients, or the method of variation of parameters. The method of variation of parameters will be used as it is more general and can be used when the coefficients of the ODE depend on x as well. Let

$$y_p(x) = u_1 y_1 + u_2 y_2 \quad (1)$$

Where u_1, u_2 to be determined, and y_1, y_2 are the two basis solutions (the two linearly independent solutions of the homogeneous ODE) found earlier when solving the homogeneous ODE as

$$y_1 = \frac{1}{x}$$

$$y_2 = \frac{\ln(x)}{x}$$

In the Variation of parameters u_1, u_2 are found using

$$u_1 = - \int \frac{y_2 f(x)}{aW(x)} \quad (2)$$

$$u_2 = \int \frac{y_1 f(x)}{aW(x)} \quad (3)$$

Where $W(x)$ is the Wronskian and a is the coefficient in front of y'' in the given ODE.

The Wronskian is given by $W = \begin{vmatrix} y_1 & y_2 \\ y_1' & y_2' \end{vmatrix}$. Hence

$$W = \begin{vmatrix} \frac{1}{x} & \frac{\ln(x)}{x} \\ \frac{d}{dx} \left(\frac{1}{x} \right) & \frac{d}{dx} \left(\frac{\ln(x)}{x} \right) \end{vmatrix}$$

Which gives

$$W = \begin{vmatrix} \frac{1}{x} & \frac{\ln(x)}{x} \\ -\frac{1}{x^2} & -\frac{\ln(x)}{x^2} + \frac{1}{x^2} \end{vmatrix}$$

Therefore

$$W = \left(\frac{1}{x}\right) \left(-\frac{\ln(x)}{x^2} + \frac{1}{x^2}\right) - \left(\frac{\ln(x)}{x}\right) \left(-\frac{1}{x^2}\right)$$

Which simplifies to

$$W = \frac{1}{x^3}$$

Which simplifies to

$$W = \frac{1}{x^3}$$

Therefore Eq. (2) becomes

$$u_1 = - \int \frac{\frac{\ln(x)}{x^2}}{\frac{1}{x}} dx$$

Which simplifies to

$$u_1 = - \int \frac{\ln(x)}{x} dx$$

Hence

$$u_1 = -\frac{\ln(x)^2}{2}$$

And Eq. (3) becomes

$$u_2 = \int \frac{\frac{1}{x^2}}{\frac{1}{x}} dx$$

Which simplifies to

$$u_2 = \int \frac{1}{x} dx$$

Hence

$$u_2 = \ln(x)$$

Therefore the particular solution, from equation (1) is

$$y_p(x) = \frac{\ln(x)^2}{2x}$$

Therefore the general solution is

$$\begin{aligned} y &= y_h + y_p \\ &= \left(\frac{c_1}{x} + \frac{c_2 \ln(x)}{x} \right) + \left(\frac{\ln(x)^2}{2x} \right) \end{aligned}$$

Will add steps showing solving for IC soon.

Summary of solutions found

$$y = \frac{c_1}{x} + \frac{c_2 \ln(x)}{x} + \frac{\ln(x)^2}{2x}$$

Solved as second order ode adjoint method

Time used: 0.344 (sec)

In normal form the ode

$$x^2 y'' + 3xy' + y = \frac{1}{x} \tag{1}$$

Becomes

$$y'' + p(x)y' + q(x)y = r(x) \tag{2}$$

Where

$$\begin{aligned} p(x) &= \frac{3}{x} \\ q(x) &= \frac{1}{x^2} \\ r(x) &= \frac{1}{x^3} \end{aligned}$$

The Lagrange adjoint ode is given by

$$\begin{aligned} \xi'' - (\xi p)' + \xi q &= 0 \\ \xi'' - \left(\frac{3\xi(x)}{x} \right)' + \left(\frac{\xi(x)}{x^2} \right) &= 0 \\ \xi''(x) + \frac{4\xi(x)}{x^2} - \frac{3\xi'(x)}{x} &= 0 \end{aligned}$$

Which is solved for $\xi(x)$. This is Euler second order ODE. Let the solution be $\xi = x^r$, then $\xi' = rx^{r-1}$ and $\xi'' = r(r-1)x^{r-2}$. Substituting these back into the given ODE gives

$$x^2(r(r-1))x^{r-2} - 3rxr^{r-1} + 4x^r = 0$$

Simplifying gives

$$r(r-1)x^r - 3rx^r + 4x^r = 0$$

Since $x^r \neq 0$ then dividing throughout by x^r gives

$$r(r-1) - 3r + 4 = 0$$

Or

$$r^2 - 4r + 4 = 0 \tag{1}$$

Equation (1) is the characteristic equation. Its roots determine the form of the general solution. Using the quadratic equation the roots are

$$r_1 = 2$$

$$r_2 = 2$$

Since the roots are equal, then the general solution is

$$\xi = c_1\xi_1 + c_2\xi_2$$

Where $\xi_1 = x^r$ and $\xi_2 = x^r \ln(x)$. Hence

$$\xi = c_1 x^2 + c_2 x^2 \ln(x)$$

Will add steps showing solving for IC soon.

The original ode now reduces to first order ode

$$\begin{aligned} \xi(x)y' - y\xi'(x) + \xi(x)p(x)y &= \int \xi(x)r(x) dx \\ y' + y\left(p(x) - \frac{\xi'(x)}{\xi(x)}\right) &= \frac{\int \xi(x)r(x) dx}{\xi(x)} \end{aligned}$$

Or

$$y' + y\left(\frac{3}{x} - \frac{2c_1x + 2c_2x \ln(x) + c_2x}{c_1x^2 + c_2x^2 \ln(x)}\right) = \frac{\frac{c_2 \ln(x)^2}{2} + c_1 \ln(x)}{c_1x^2 + c_2x^2 \ln(x)}$$

Which is now a first order ode. This is now solved for y . In canonical form a linear first order is

$$y' + q(x)y = p(x)$$

Comparing the above to the given ode shows that

$$q(x) = -\frac{-c_2 \ln(x) - c_1 + c_2}{x(c_2 \ln(x) + c_1)}$$

$$p(x) = \frac{\ln(x)(c_2 \ln(x) + 2c_1)}{2x^2(c_2 \ln(x) + c_1)}$$

The integrating factor μ is

$$\begin{aligned}\mu &= e^{\int q dx} \\ &= e^{\int -\frac{-c_2 \ln(x) - c_1 + c_2}{x(c_2 \ln(x) + c_1)} dx} \\ &= \frac{x}{c_2 \ln(x) + c_1}\end{aligned}$$

The ode becomes

$$\begin{aligned}\frac{d}{dx}(\mu y) &= \mu p \\ \frac{d}{dx}(\mu y) &= (\mu) \left(\frac{\ln(x)(c_2 \ln(x) + 2c_1)}{2x^2(c_2 \ln(x) + c_1)} \right) \\ \frac{d}{dx} \left(\frac{yx}{c_2 \ln(x) + c_1} \right) &= \left(\frac{x}{c_2 \ln(x) + c_1} \right) \left(\frac{\ln(x)(c_2 \ln(x) + 2c_1)}{2x^2(c_2 \ln(x) + c_1)} \right) \\ d \left(\frac{yx}{c_2 \ln(x) + c_1} \right) &= \left(\frac{\ln(x)(c_2 \ln(x) + 2c_1)}{2x(c_2 \ln(x) + c_1)^2} \right) dx\end{aligned}$$

Integrating gives

$$\begin{aligned}\frac{yx}{c_2 \ln(x) + c_1} &= \int \frac{\ln(x)(c_2 \ln(x) + 2c_1)}{2x(c_2 \ln(x) + c_1)^2} dx \\ &= \frac{\ln(x)}{2c_2} + \frac{c_1^2}{2c_2^2(c_2 \ln(x) + c_1)} + c_3\end{aligned}$$

Dividing throughout by the integrating factor $\frac{x}{c_2 \ln(x) + c_1}$ gives the final solution

$$y = \frac{\ln(x)^2 c_2^2 + c_2(2c_2^2 c_3 + c_1) \ln(x) + c_1(2c_2^2 c_3 + c_1)}{2x c_2^2}$$

Hence, the solution found using Lagrange adjoint equation method is

$$y = \frac{\ln(x)^2 c_2^2 + c_2(2c_2^2 c_3 + c_1) \ln(x) + c_1(2c_2^2 c_3 + c_1)}{2x c_2^2}$$

The constants can be merged to give

$$y = \frac{\ln(x)^2 c_2^2 + c_2(2c_2^2 + c_1) \ln(x) + c_1(2c_2^2 + c_1)}{2x c_2^2}$$

Will add steps showing solving for IC soon.

Summary of solutions found

$$y = \frac{\ln(x)^2 c_2^2 + c_2(2c_2^2 + c_1) \ln(x) + c_1(2c_2^2 + c_1)}{2x c_2^2}$$

Maple step by step solution

Maple trace

```
`Methods for second order ODEs:
--- Trying classification methods ---
trying a quadrature
trying high order exact linear fully integrable
<- high order exact linear fully integrable successful`
```

Maple dsolve solution

Solving time : 0.003 (sec)

Leaf size : 20

```
dsolve(x^2*diff(diff(y(x),x),x)+3*x*diff(y(x),x)+y(x) = 1/x,
y(x),singsol=all)
```

$$y = \frac{\frac{\ln(x)^2}{2} + c_1 \ln(x) + c_2}{x}$$

Mathematica DSolve solution

Solving time : 0.024 (sec)

Leaf size : 27

```
DSolve[{x^2*D[y[x],{x,2}]+3*x*D[y[x],x]+y[x]==1/x,{}},  
y[x],x,IncludeSingularSolutions->True]
```

$$y(x) \rightarrow \frac{\log^2(x) + 2c_2 \log(x) + 2c_1}{2x}$$