

HW 1

Physics 3041  
Mathematical Methods for Physicists

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## 1 Problem 1.6.1

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Expand the function  $f(x) = \frac{\sin(x)}{\cosh(x)+2}$  in Taylor series around the origin going up to  $x^3$ . Calculate  $f(0.1)$  from this series and compare to the exact answer obtained by using a calculator

### Solution

The Taylor series of function  $f(x)$  around origin is given by (1.3.16) ( $\approx$  is used throughout this HW to mean that the left side is the Taylor series approximation of  $f(x)$ ).

$$f(x) \approx \sum_{n=0}^{\infty} \frac{x^n}{n!} f^{(n)}(0)$$

Where  $f^{(n)}(0)$  is the  $n^{\text{th}}$  derivative of  $f(x)$  evaluated at  $x = 0$ .

For  $n = 0$ ,  $f^{(0)}(x) = f(x) = \frac{\sin(x)}{\cosh(x)+2}$ , therefore  $f(0) = 0$ .

For  $n = 1$

$$\begin{aligned} f^{(1)}(x) &= \frac{d}{dx} \left( \frac{\sin(x)}{\cosh(x) + 2} \right) \\ &= \frac{\cos(x)(\cosh(x) + 2) - \sin(x) \sinh(x)}{(\cosh(x) + 2)^2} \\ &= \frac{\cos(x)(\cosh(x) + 2)}{(\cosh(x) + 2)^2} - \frac{\sin(x) \sinh(x)}{(\cosh(x) + 2)^2} \\ &= \frac{\cos(x)}{\cosh(x) + 2} - \frac{\sin(x) \sinh(x)}{(\cosh(x) + 2)^2} \end{aligned}$$

The above evaluated at  $x = 0$  becomes

$$\begin{aligned} f^{(1)}(0) &= \frac{1}{1+2} - \frac{0}{(1+2)^2} \\ &= \frac{1}{3} \end{aligned}$$

For  $n = 2$

$$\begin{aligned}
f^{(2)}(x) &= \frac{d}{dx} \left( \frac{d}{dx} \left( \frac{\sin(x)}{\cosh(x) + 2} \right) \right) \\
&= \frac{d}{dx} \left( \frac{\cos(x)}{\cosh(x) + 2} - \frac{\sin(x) \sinh(x)}{(\cosh(x) + 2)^2} \right) \\
&= \frac{-\sin(x)(\cosh(x) + 2) - \cos(x) \sinh(x)}{(\cosh(x) + 2)^2} \\
&= \frac{(\cos(x) \sinh(x) + \sin(x) \cosh(x))(\cosh(x) + 2)^2 - \sin(x) \sinh(x)(2(\cosh(x) + 2) \sinh(x))}{(\cosh(x) + 2)^4} \\
&= \frac{-\sin(x)(\cosh(x) + 2)}{(\cosh(x) + 2)^2} - \frac{\cos(x) \sinh(x)}{(\cosh(x) + 2)^2} - \frac{\cos(x) \sinh(x)(\cosh(x) + 2)^2}{(\cosh(x) + 2)^4} \\
&= \frac{-\sin(x) \cosh(x)(\cosh(x) + 2)^2}{(\cosh(x) + 2)^4} + \frac{\sin(x) \sinh(x)(2(\cosh(x) + 2) \sinh(x))}{(\cosh(x) + 2)^4} \\
&= \frac{-\sin(x)}{\cosh(x) + 2} - \frac{\cos(x) \sinh(x)}{(\cosh(x) + 2)^2} - \frac{\cos(x) \sinh(x)}{(\cosh(x) + 2)^2} - \frac{\sin(x) \cosh(x)}{(\cosh(x) + 2)^2} + \frac{2 \sin(x) \sinh(x) \sinh(x)}{(\cosh(x) + 2)^3} \\
&= \frac{-\sin(x)}{\cosh(x) + 2} - 2 \frac{\cos(x) \sinh(x)}{(\cosh(x) + 2)^2} - \frac{\sin(x) \cosh(x)}{(\cosh(x) + 2)^2} + \frac{2 \sin(x) \sinh^2(x)}{(\cosh(x) + 2)^3}
\end{aligned}$$

The above evaluated at  $x = 0$  becomes

$$\begin{aligned}
f^{(2)}(0) &= \frac{-0}{1+2} - 2 \frac{0}{(1+2)^2} - \frac{0}{(1+2)^2} + \frac{0}{(1+2)^3} \\
&= 0
\end{aligned}$$

For  $n = 3$

$$\begin{aligned}
f^{(3)}(x) &= \frac{d}{dx} \left( \frac{d^2}{dx^2} \left( \frac{\sin(x)}{\cosh(x) + 2} \right) \right) \\
&= \frac{d}{dx} \left( \frac{-\sin(x)}{\cosh(x) + 2} - 2 \frac{\cos(x) \sinh(x)}{(\cosh(x) + 2)^2} - \frac{\sin(x) \cosh(x)}{(\cosh(x) + 2)^2} + \frac{2 \sin(x) \sinh^2(x)}{(\cosh(x) + 2)^3} \right) \\
&= \frac{-\cos(x)(\cosh(x) + 2) + \sin(x) \sinh(x)}{(\cosh(x) + 2)^2} \\
&= \frac{(-\sin(x) \sinh(x) + \cos(x) \cosh(x))(\cosh(x) + 2)^2 - \cos(x) \sinh(x)(2(\cosh(x) + 2) \sinh(x))}{(\cosh(x) + 2)^4} \\
&= \frac{(\cos(x) \cosh(x) + \sin(x) \sinh(x))(\cosh(x) + 2)^2 - \sin(x) \cosh(x)(2(\cosh(x) + 2) \sinh(x))}{(\cosh(x) + 2)^4} \\
&+ 2 \frac{(\cos(x) \sinh^2(x) + 2 \sin(x) \cosh(x))(\cosh(x) + 2)^3 - (\sin(x) \sinh^2(x))(3(\cosh(x) + 2)^2 \sinh(x))}{(\cosh(x) + 2)^6}
\end{aligned}$$

The above evaluated at  $x = 0$  gives

$$\begin{aligned}
 f^{(3)}(0) &= \frac{-1(1+2)+0}{(1+2)^2} - 2 \frac{(-0+1)(1+2)^2-0}{(1+2)^4} - \frac{(1+0)(1+2)^2-0}{(1+2)^4} + 2 \frac{(0+0)(1+2)^3-(0)(3(1+2)^2 \cdot 0)}{(1+2)^6} \\
 &= \frac{-3}{3^2} - 2 \frac{(1)(3)^2}{(3)^4} - \frac{(1)(3)^2}{(3)^4} + 2 \frac{0}{(3)^6} \\
 &= \frac{-1}{3} - 2 \frac{1}{3^2} - \frac{1}{3^2} \\
 &= -\frac{2}{3}
 \end{aligned}$$

The process stops here, because the problem is asking for  $n = 3$ . Substituting all the derivatives  $f^{(n)}(0)$  values above into

$$f(x) \approx \sum_{n=0}^{\infty} \frac{x^n}{n!} f^{(n)}(0)$$

For up to  $n = 3$  gives the following

$$\begin{aligned}
 f(x) &\approx f(0) + x f^{(1)}(0) + \frac{x^2}{2} f^{(2)}(0) + \frac{x^3}{3!} f^{(3)}(0) + \dots \\
 &\approx 0 + x \frac{1}{3} + \frac{x^2}{2} (0) + \frac{x^3}{3!} \left( -\frac{2}{3} \right) \\
 &\approx x \frac{1}{3} - \frac{2x^3}{3 \cdot 6} \\
 &\approx \frac{x}{3} - \frac{x^3}{9}
 \end{aligned}$$

When  $x = \frac{1}{10}$  the above becomes

$$\begin{aligned}
 f_{n=3}\left(\frac{1}{10}\right) &\approx \frac{1}{30} - \frac{1}{(1000)9} \\
 &\approx \frac{1}{30} - \frac{1}{9000} \\
 &\approx \frac{300-1}{9000} \\
 &\approx \frac{299}{9000}
 \end{aligned}$$

From the calculator

$$\frac{299}{9000} \approx 0.0332222$$

And from the exact expression

$$\begin{aligned}
 \frac{\sin(x)}{\cosh(x)+2} &= \frac{\sin(0.1)}{\cosh(0.1)+2} \\
 &= 0.0332224
 \end{aligned}$$

The error is about  $1.67 \times 10^{-7}$ .

## 2 Problem 2

---

Consider  $f(x) = (1+x)^p$  for (a)  $p = \frac{1}{3}$  and (b)  $p = -2$ , respectively. (1) Find the Taylor series of  $f(x)$  around  $x = 0$ . (2) From the form of the general term, find the interval of convergence of the series. (3) How many terms in the series do you need to estimate  $f(0.1)$  to within 1%? Check that the difference between your estimate and the actual result has approximately the same magnitude as the next term in the series.

Solution

### 2.1 Case $p = \frac{1}{3}$

$$f(x) = (1+x)^{\frac{1}{3}}$$

Part (1) The Taylor series is given by

$$f(x) \approx f(0) + xf'(0) + \frac{x^2}{2!}f''(0) + \frac{x^3}{3!}f'''(0) + \dots \quad (1)$$

Where  $f(0) = 1$  and  $f'(x) = \frac{1}{3}(1+x)^{-\frac{2}{3}}$ . Hence  $f'(0) = \frac{1}{3}$  and  $f''(x) = \frac{1}{3}\left(-\frac{2}{3}\right)(1+x)^{-\frac{5}{3}}$ . Hence  $f''(0) = -\frac{(2)}{3^2}$ , and  $f'''(x) = \frac{1}{3}\left(-\frac{2}{3}\right)\left(-\frac{5}{3}\right)(1+x)^{-\frac{8}{3}}$ , hence  $f'''(0) = \frac{1}{3}\left(-\frac{2}{3}\right)\left(-\frac{5}{3}\right) = \frac{(2)(5)}{3^3}$ , and  $f^{(4)}(x) = \frac{1}{3}\left(-\frac{2}{3}\right)\left(-\frac{5}{3}\right)\left(-\frac{8}{3}\right)(1+x)^{-\frac{11}{3}}$ , hence  $f^{(4)}(0) = \frac{1}{3}\left(-\frac{2}{3}\right)\left(-\frac{5}{3}\right)\left(-\frac{8}{3}\right) = -\frac{1}{3^4}((2)(5)(8))$  and on. The series in (1) becomes

$$\begin{aligned} f(x) &\approx 1 + \frac{1}{3}x - \frac{(2)x^2}{3^2 2!} + \frac{(2)(5)x^3}{3^3 3!} - \frac{(2)(5)(8)x^4}{3^4 4!} + \frac{(2)(5)(8)(11)x^5}{3^5 5!} - \frac{(2)(5)(8)(11)(14)x^6}{3^6 6!} - \dots \\ &\approx 1 + \frac{1}{3}x - \frac{1}{3^2}x^2 + \frac{5}{3^3}x^3 - \frac{10}{3^4}x^4 + \frac{22}{3^5}x^5 - \frac{154}{3^6}x^6 + \dots \end{aligned} \quad (2)$$

The general term is found by comparing the above to the general term obtained from binomial expansion. Since

$$(1+x)^p = \binom{p}{0}x^0 + \binom{p}{1}x + \binom{p}{2}x^2 + \dots \quad (3)$$

Comparing (2,3) shows that the general term is the binomial coefficient  $\binom{\frac{1}{3}}{n}$ . Therefore

the Taylor series for  $(1+x)^{\frac{1}{3}}$  can be written as

$$f(x) \approx \sum_{n=0}^{\infty} \binom{\frac{1}{3}}{n} x^n$$

For  $p = \frac{1}{3}$  the above becomes

$$f(x) \approx \sum_{n=0}^{\infty} \binom{\frac{1}{3}}{n} x^n$$

Part(2)

$$\begin{aligned} R &= \lim_{n \rightarrow \infty} \left| \frac{a_n}{a_{n+1}} \right| \\ &= \lim_{n \rightarrow \infty} \left| \frac{\binom{p}{n}}{\binom{p}{n+1}} \right| \end{aligned}$$

The Binomial coefficient  $\binom{p}{n} = \frac{p!}{n!(p-n)!}$ , for when  $p$  is integer. This is not the case here.

For non-integer  $p$  The Binomial coefficient becomes  $\binom{p}{n} = \frac{\Gamma(p+1)}{\Gamma(n+1)\Gamma(p-n+1)}$  where  $\Gamma(p)$  is the Gamma function. The above ratio now becomes

$$\begin{aligned} R &= \lim_{n \rightarrow \infty} \left| \frac{\frac{\Gamma(p+1)}{\Gamma(n+1)\Gamma(p-n+1)}}{\frac{\Gamma(p+1)}{\Gamma(n+2)\Gamma(p-n)}} \right| \\ &= \lim_{n \rightarrow \infty} \left| \frac{\Gamma(n+2)\Gamma(p-n)}{\Gamma(n+1)\Gamma(p-n+1)} \right| \\ &= \lim_{n \rightarrow \infty} \left| \frac{n\Gamma(p-n)}{\Gamma(p-n+1)} \right| \\ &= \lim_{n \rightarrow \infty} \left| \frac{n}{p-n} \right| \end{aligned}$$



But  $p = \frac{1}{3}$ , hence the above becomes

$$\begin{aligned} R &= \lim_{n \rightarrow \infty} \left| \frac{n}{\frac{1}{3} - n} \right| \\ &= \lim_{n \rightarrow \infty} \left| \frac{n}{n - \frac{1}{3}} \right| \\ &= 1 \end{aligned}$$

Therefore the radius of convergence is 1. This means the Taylor series found above converges to  $f(x)$  for  $|x| < 1$ .

### Part 3

$$f(x) = (1 + x)^{\frac{1}{3}}$$

When  $x = 0.1$

$$\begin{aligned} f(0.1) &= (1.1)^{\frac{1}{3}} \\ &= 1.032280115 \end{aligned}$$

one percent of the above is

$$\frac{1}{100}(1.032280115) = 0.01032280115$$

The value  $n$  is now found such that

$$|R_n(x)| \leq M \frac{(0.1)^{n+1}}{(n+1)!} \leq 0.01032280115$$

Where  $R_n(x)$  is the Taylor series remainder using  $n$  terms.  $M$  is the upper bound for the  $n + 1$  derivative of  $f(x)$  any where between  $[0, 0.1]$ . Instead of trying to find  $M$ , few calculations are used to find how many terms are needed.

For  $n = 0$ ,  $\tilde{f}(0.1) = 1$  and the error is  $1.032280115 - 1 = 0.032280115$ .

For  $n = 1$ ,  $\tilde{f}(0.1) = 1 + \left( \frac{1}{3} \right) 0.1 = 1.0333333$ , and the error is  $|1.032280115 - 1.0333333| =$

0.001053218. Because this is smaller than  $R_n(x)$  then only two terms are needed in the Taylor series to obtained the required accuracy. Therefore

$$f(x) \approx 1 + \frac{1}{3}x$$

## 2.2 Case $p = -2$

$$f(x) = (1 + x)^{-2}$$

Part (1) The Taylor series is

$$f(x) \approx f(0) + xf'(0) + \frac{x^2}{2!}f''(0) + \frac{x^3}{3!}f'''(0) + \dots$$

But  $f(0) = 1$  and  $f'(x) = (-2)(1 + x)^{-3}$ . Hence  $f'(0) = -2$  and  $f''(x) = (-2)(-3)(1 + x)^{-4}$ . Hence  $f''(0) = (-2)(-3)$ , and  $f'''(x) = -2(-3)(-4)(1 + x)^{-5}$ , hence  $f'''(0) = (-2)(-3)(-4)(-5)$  and so on. The above becomes

$$\begin{aligned} f(x) &\approx 1 + (-2)x - (-2)(-3)\frac{x^2}{2!} + (-2)(-3)(-4)\frac{x^3}{3!} + \dots \\ &\approx 1 - 2x + (2)(3)\frac{x^2}{2!} - (2)(3)(4)\frac{x^3}{3!} + \dots \\ &\approx 1 - 2x + 3x^2 - 4x^3 + \dots \end{aligned}$$

The general term is therefore

$$f(x) \approx \sum_{n=0}^{\infty} (-1)^n (n+1)x^n$$

Part(2)

$$\begin{aligned} R &= \lim_{n \rightarrow \infty} \left| \frac{a_n}{a_{n+1}} \right| \\ &= \lim_{n \rightarrow \infty} \left| \frac{(n+1)}{(n+2)} \right| \\ &= 1 \end{aligned}$$

Hence the series converges to  $f(x)$  for  $|x| < 1$ .

Part 3

$$f(x) = (1 + x)^{-2}$$

For  $x = 0.1$

$$\begin{aligned} f(0.1) &= (1.1)^{-2} \\ &= \frac{1}{1.1^2} \\ &= 0.82644628 \end{aligned}$$

One percent of the above is

$$\frac{1}{100}(0.8264462810) = 0.0082644628$$

The value  $n$  is now found such that

$$|R_n(x)| \leq M \frac{(0.1)^{n+1}}{(n+1)!} \leq 0.0082644628$$

Where  $R_n(x)$  is the Taylor series remainder using  $n$  terms.  $M$  is the upper bound for the  $n+1$  derivative of  $f(x)$  any where between  $[0, 0.1]$ . Doing few calculations gives

For  $n = 0$ ,  $\tilde{f}(0.1) = 1$ , the error is  $|0.82644628 - 1| = 0.1735537190$ .

For  $n = 1$ ,  $\tilde{f}(0.1) = 0.8$ , the error is  $|0.82644628 - 0.8| = 0.02644628$ .

For  $n = 2$ ,  $\tilde{f}(0.1) = 0.83$ , the error is  $|0.82644628 - 0.83| = 0.0035537190$ . Because this is within 1% then only three terms are needed. Therefore

$$f(x) \approx 1 - 2x + 3x^2$$

### 3 Problem 3

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Expand  $f(x) = \tan(x^2)$  to order  $x^6$  using (a) direct Taylor expansion. (b) The Taylor series for  $\sin(x)$  and  $\cos x$  with appropriate substitution.

Solution

#### 3.1 Part a

Using Taylor series

$$f(x) \approx \sum_{n=0}^{\infty} \frac{x^n}{n!} f^{(n)}(0)$$

Where  $f(x) = \tan(x^2)$  and the expansion is around  $x = 0$ . The Taylor series for  $f(u) = \tan(u)$  is found instead of  $\tan(x^2)$ , and then at the end  $u$  is replaced by  $x^2$ . This is called the substitution method. This simplifies the derivations. Therefore  $f(0) = 0$ . The first derivative is

$$\begin{aligned} f'(u) &= \frac{d}{du} \tan(u) \\ &= \frac{d}{du} \left( \frac{\sin u}{\cos u} \right) \\ &= \frac{\cos^2 u + \sin^2 u}{\cos^2 u} \\ &= \frac{1}{\cos^2 u} \end{aligned}$$

At  $u = 0$  this gives  $f'(0) = 1$ .

The next derivative using the above result gives

$$\begin{aligned} f''(u) &= \frac{d}{du} \left( \frac{1}{\cos^2 u} \right) \\ &= \frac{2 \cos u \sin u}{\cos^4 u} \\ &= \frac{2 \sin u}{\cos^3 u} \end{aligned}$$

At  $u = 0$  this gives  $f^{(2)}(0) = 0$ . The next derivative gives

$$\begin{aligned}
 f^{(3)}(u) &= 2 \frac{d}{du} \left( \frac{\sin u}{\cos^3 u} \right) \\
 &= 2 \frac{\cos u \cos^3 u - \sin u (3 \cos^2 u (-\sin u))}{\cos^6 u} \\
 &= 2 \frac{\cos^4 u + 3 \sin^2 u \cos^2 u}{\cos^6 u} \\
 &= \frac{2 \cos^4 u}{\cos^6 u} + \frac{6 \sin^2 u \cos^2 u}{\cos^6 u} \\
 &= \frac{2}{\cos^2 u} + \frac{6 \sin^2 u}{\cos^4 u} \\
 &= \frac{2}{\cos^2 u} + \frac{6(1 - \cos^2 u)}{\cos^4 u} \\
 &= \frac{2}{\cos^2 u} + \frac{6}{\cos^4 u} - \frac{6}{\cos^2 u} \\
 &= -\frac{4}{\cos^2 u} + \frac{6}{\cos^4 u}
 \end{aligned}$$

At  $u = 0$  this gives  $f^{(3)}(0) = -\frac{4}{1} + \frac{6}{1} = 2$ . Since the problem is asking for order  $x^6$  the process stops here, as this is the same as order  $u^3$  when  $u$  is replaced by  $x^2$ .

Therefore the Taylor series for  $\tan(u)$  is (for up to  $n = 3$ )

$$\begin{aligned}
 f(u) &\approx f(0) + uf'(0) + \frac{u^2}{2!}f^{(2)}(0) + \frac{u^3}{3!}f^{(3)}(0) + \dots \\
 &\approx 0 + u + 0 + 2\frac{u^3}{3!} \\
 &\approx u + \frac{1}{3}u^3
 \end{aligned}$$

Replacing  $u = x^2$ , gives the Taylor series for  $\tan(x^2)$  for up to  $x^6$  term as

$$\tan(x^2) \approx x^2 + \frac{1}{3}x^6$$

### 3.2 Part b

To obtain the above result using the Taylor series for  $\sin(x^2)$ ,  $\cos(x^2)$ , the Taylor series for  $\sin(x^2)$  and  $\cos(x^2)$  is found, and long division is applied using the definition of  $\tan(x^2) = \frac{\sin(x^2)}{\cos(x^2)}$ . Terms with order larger than  $x^6$  are ignored. The Taylor series for  $\sin(x)$  is

$$\sin(x) \approx x - \frac{x^3}{3!} + \frac{x^5}{5!} - \dots$$

Using the substitution method, the Taylor series for  $\sin(x^2)$  becomes

$$\begin{aligned}\sin(x^2) &\approx x^2 - \frac{x^6}{3!} + \frac{x^{10}}{5!} - \dots \\ &\approx x^2 - \frac{x^6}{6} + \frac{x^{10}}{120} - \dots\end{aligned}\tag{1}$$

The Taylor series for  $\cos(x)$  is

$$\cos(x) \approx 1 - \frac{x^2}{2!} + \frac{x^4}{4!} - \dots$$

Using the substitution method, the Taylor series for  $\cos(x^2)$  becomes

$$\begin{aligned}\cos(x^2) &\approx 1 - \frac{x^4}{2!} + \frac{x^8}{4!} - \dots \\ &\approx 1 - \frac{x^4}{2} + \frac{x^8}{24} - \dots\end{aligned}\tag{2}$$

Since  $\tan(x^2) = \frac{\sin(x^2)}{\cos(x^2)}$  then the Taylor series for  $\tan(x^2)$  is

$$\tan(x^2) \approx \frac{x^2 - \frac{x^6}{3!} + \frac{x^{10}}{5!} - \dots}{1 - \frac{x^4}{2!} + \frac{x^8}{4!} - \dots}$$

Performing long division and stopping when the remainder has powers larger than  $x^6$  gives

$$\tan(x^2) \approx x^2 + \frac{1}{3}x^6 + \dots$$

Which is same result as part(a).

$$\begin{array}{r}
 \phantom{1 - \frac{x^4}{2} + \frac{x^8}{24} - \dots} \\
 \hline
 \phantom{1 - \frac{x^4}{2} + \frac{x^8}{24} - \dots} \quad X^2 + \frac{X^6}{3} \\
 \hline
 \phantom{1 - \frac{x^4}{2} + \frac{x^8}{24} - \dots} \quad X^2 - \frac{X^6}{6} + \frac{X^{10}}{120} + \dots \\
 \phantom{1 - \frac{x^4}{2} + \frac{x^8}{24} - \dots} \quad X^2 - \frac{X^6}{2} + \frac{X^{10}}{24} - \dots \\
 \hline
 \phantom{1 - \frac{x^4}{2} + \frac{x^8}{24} - \dots} \quad \frac{X^6}{3} - \frac{X^{10}}{120} + \dots \\
 \phantom{1 - \frac{x^4}{2} + \frac{x^8}{24} - \dots} \quad \frac{X^6}{3} - \frac{X^{10}}{6} + \dots \\
 \hline
 \phantom{1 - \frac{x^4}{2} + \frac{x^8}{24} - \dots} \quad \frac{31}{120} X^{10} + \dots \\
 \phantom{1 - \frac{x^4}{2} + \frac{x^8}{24} - \dots} \quad \text{ignore. stop here.}
 \end{array}$$

hence 
$$\frac{X^2 - \frac{X^6}{6} + \frac{X^{10}}{120} + \dots}{1 - \frac{X^4}{2} + \frac{X^8}{24} - \dots} = X^2 + \frac{X^6}{3} + \text{higher order}$$

Figure 1: Polynomials long division

## 4 Problem 4

---

A particle of mass  $m$  moves along the  $+x$  axis (i.e.  $x > 0$ ) with potential energy

$$V(x) = \frac{a}{2x^2} - \frac{b}{x}$$

Where  $a$  and  $b$  are positive parameters. (a) Find the equilibrium position  $x_0$ . (b) Show that the particle executes harmonic oscillations near  $x = x_0$ . (c) Find the angular frequency of oscillations.

Solution

### 4.1 Part a

Equilibrium position is where the slope of the potential energy is zero. This position  $x_0$  is found by solving for  $x$  from

$$\frac{dV}{dx} = 0$$

But

$$\begin{aligned} \frac{dV}{dx} &= \frac{a}{2}(-2x^{-3}) - b(-x^{-2}) \\ &= \frac{-a}{x^3} + \frac{b}{x^2} \\ &= \frac{-a + bx}{x^3} \end{aligned}$$

Hence

$$\begin{aligned} \frac{-a + bx}{x^3} &= 0 \\ bx &= a \end{aligned}$$

Therefore

$$x_0 = \frac{a}{b}$$

### 4.2 Part b

Approximating  $V(x)$  around  $x_0$  using Taylor series gives

$$V(x) \approx V(x_0) + (x - x_0)V'(x_0) + \frac{(x - x_0)^2}{2!}V''(x_0) + \dots$$

But  $\frac{dV}{dx}$  evaluated at  $x_0$  is zero, since this is the equilibrium point. The above simplifies to

$$V(x) \approx V(x_0) + \frac{(x - x_0)^2}{2!}V''(x_0) + \dots \quad (\text{A})$$



Higher terms are ignored, because  $(x - x_0)$  is assumed small and mass remain close to  $x_0$ .  
But

$$V(x_0) = \frac{a}{2x_0^2} - \frac{b}{x_0}$$

And since  $x_0 = \frac{a}{b}$  from part (a), the above simplifies to

$$\begin{aligned} V(x_0) &= \frac{a}{2\left(\frac{a}{b}\right)^2} - \frac{b}{\left(\frac{a}{b}\right)} \\ &= \frac{ab^2}{2a^2} - \frac{b^2}{a} \\ &= \frac{b^2}{2a} - \frac{b^2}{a} \\ &= -\frac{1}{2} \frac{b^2}{a} \end{aligned} \tag{A1}$$

And

$$\begin{aligned} \frac{d^2V}{dx^2} &= \frac{d}{dx} \left( \frac{-a}{x^3} + \frac{b}{x^2} \right) \\ &= \frac{3a}{x^4} - \frac{b}{x^3} \end{aligned}$$

At  $x = x_0$  the above becomes

$$\begin{aligned} V''(x_0) &= \frac{3a}{\left(\frac{a}{b}\right)^4} - \frac{b}{\left(\frac{a}{b}\right)^3} \\ &= \frac{b^4}{a^3} \end{aligned} \tag{A2}$$

Using (A1,A2) into A gives

$$\begin{aligned} V(x) &\approx -\frac{1}{2} \frac{b^2}{a} + \frac{(x - x_0)^2}{2!} \frac{b^4}{a^3} + \dots \\ &\approx -\frac{1}{2} \frac{b^2}{a} + \frac{\left(x - \frac{a}{b}\right)^2}{2} \frac{b^4}{a^3} + \dots \\ &\approx -\frac{1}{2} \frac{b^2}{a} + \frac{1}{2} \left( x^2 + \frac{a^2}{b^2} - 2x \frac{a}{b} \right) \frac{b^4}{a^3} + \dots \\ &\approx -\frac{1}{2} \frac{b^2}{a} + \frac{1}{2a} b^2 + \frac{1}{2a^3} b^4 x^2 - \frac{1}{a^2} b^3 x + \dots \\ &\approx \frac{b^4}{2a^3} x^2 - \frac{b^3}{a^2} x + \dots \end{aligned}$$

Therefore near  $x_0$  the potential energy is approximated as

$$V(x) \approx \frac{b^4}{2a^3}x^2 - \frac{b^3}{a^2}x \quad (1)$$

The force on the mass is given by

$$F = -\frac{dV}{dx}$$

Using  $V(x)$  in (1) the force becomes

$$F = -\frac{b^4}{a^3}x - \frac{b^3}{a^2}$$

But  $F = m\frac{d^2x}{dt^2}$ . Hence we obtain the equation of motion as

$$\begin{aligned} m\frac{d^2x}{dt^2} &= F \\ &= -\frac{b^4}{a^3}x - \frac{b^3}{a^2} \end{aligned}$$

Therefore

$$\begin{aligned} m\frac{d^2x(t)}{dt^2} + \frac{b^4}{a^3}x(t) &= -\frac{b^3}{a^2} \\ \frac{d^2x(t)}{dt^2} + \left(\frac{b^4}{ma^3}\right)x(t) &= -\frac{b^3}{ma^2} \end{aligned} \quad (B)$$

Let

$$\frac{b^4}{ma^3} = \omega^2$$

The equation of motion (B) becomes

$$\frac{d^2x(t)}{dt^2} + \omega^2x(t) = -\frac{b^3}{ma^2}$$

But this is standard second order ode whose solution is

$$x(t) = A \cos(\omega t) + B \sin(\omega t) + x_p(t)$$

Where  $x_p(t)$  is the particular solution due to the forcing function  $-\frac{b^3}{ma^2}$  and  $A, B$  are constants of integrations found from initial conditions. Since the forcing function is just constant, and not function function of time, the above becomes

$$\begin{aligned} x(t) &= A \cos(\omega t) + B \sin(\omega t) + F_p \\ &= A \cos(\omega t + \phi) + F_p \end{aligned}$$

Therefore the motion is simple harmonic motion since  $\cos(\omega t + \phi)$  is harmonic. The forcing function  $F_p$  has no effect on the nature of the harmonic motion, other than adding an extra constant displacement shift to  $x(t)$  for all time. Since there is no damping, the particle will continue this motion forever.

The following is a plot of the solution for 10 seconds using arbitrary values for  $a, b, m$  and with initial conditions  $x(0) = 1, x'(0) = 0$ . The solution shows the motion is harmonic as expected.

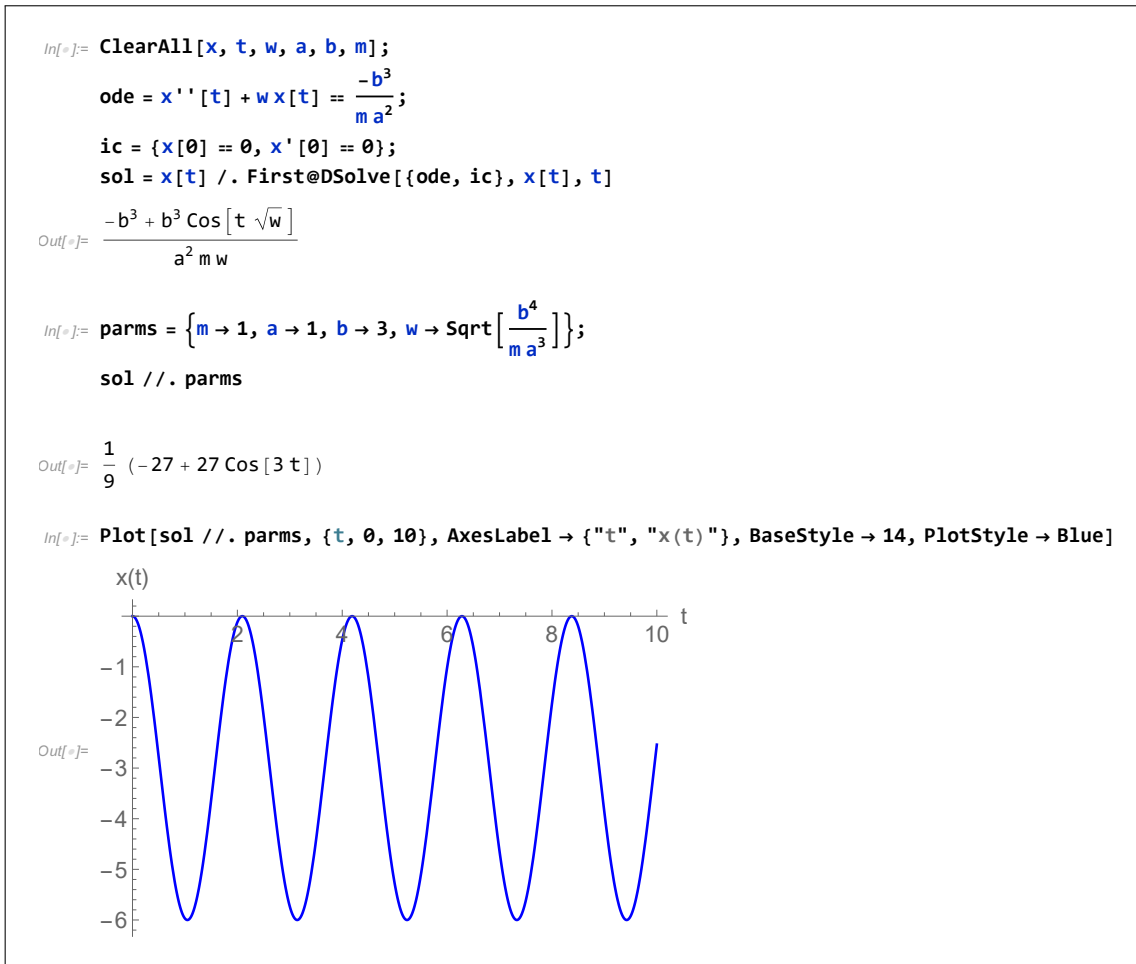


Figure 2: Plot of solution

### 4.3 Part c

The angular frequency of oscillation is

$$\omega = \sqrt{\frac{b^4}{m a^3}}$$

In radians per second. The quantity  $\frac{b^4}{a^3}$  can be called the stiffness  $k$  (Newton per meter).

Hence  $\omega = \sqrt{\frac{k}{m}}$ .

#### 4.4 Appendix

An easier way to do part b, is to keep  $(x - x_0)$  intact and replace this with  $y$  at the end. Like this

Using (A1,A2) into A gives

$$V(x) \approx -\frac{1}{2} \frac{b^2}{a} + \frac{(x - x_0)^2}{2!} \frac{b^4}{a^3} + \dots$$

The force on the mass is given by

$$\begin{aligned} F &= -\frac{dV}{dx} \\ &= -(x - x_0) \frac{b^4}{a^3} \end{aligned}$$

But  $F = m \frac{d^2x}{dt^2}$ . Hence we obtain the equation of motion as

$$\begin{aligned} m \frac{d^2x}{dt^2} &= F \\ &= -(x - x_0) \frac{b^4}{a^3} \end{aligned}$$

Now let  $y = x - x_0$ . the above becomes

$$\begin{aligned} m \frac{d^2y}{dt^2} &= -y \frac{b^4}{a^3} \\ m \frac{d^2y}{dt^2} + y \frac{b^4}{a^3} &= 0 \\ \frac{d^2y}{dt^2} + y \frac{b^4}{ma^3} &= 0 \end{aligned}$$

Which is SHM. Using this method, it is faster to show.