### 4.3.24(a)

(a)  $v(y) = u(e^y)$  solves a constant coefficient second-order ordinary differential equation with a double root r, and hence  $v(y) = c_1 e^{ry} + c_2 y e^{ry}$ . Therefore,

$$u(x) = c_1 |x|^r + c_2 |x|^r \log |x|.$$

# 4.3.25(c)(d)

(c) 
$$u(x,y) = \frac{1}{8}r^4\cos 4\theta + 2r^2\cos 2\theta + 6 = \frac{1}{8}x^4 - \frac{3}{4}x^2y^2 + \frac{1}{8}y^4 + 2x^2 - 2y^2 + 6$$

(d) 
$$u(x,y) = r \cos \theta = x$$
.

## 4.3.33

$$u(r,\theta) = \frac{a_0}{2} \, \frac{\log r}{\log 2} + \sum_{n=1}^{\infty} \, \frac{r^n - r^{-n}}{2^n - 2^{-n}} \, (a_n \cos n\theta + b_n \sin n\theta),$$

where  $a_n, b_n$  are the usual Fourier coefficients of  $h(\theta)$ .

## 4.3.38

First, if  $C = \frac{1}{\pi} \int_{-\pi}^{\pi} |h(\theta)| d\theta$ , then the Fourier coefficients are bounded by

$$\left| \begin{array}{l} \left| \, a_n \, \right| \leq \frac{1}{\pi} \int_{-\pi}^{\pi} \left| \, h(\theta) \cos n \, \theta \, \right| d\theta \\ \left| \, b_n \, \right| \leq \frac{1}{\pi} \int_{-\pi}^{\pi} \left| \, h(\theta) \sin n \, \theta \, \right| d\theta \end{array} \right. \right\} \; \leq \frac{1}{\pi} \int_{-\pi}^{\pi} \left| \, h(\theta) \, \right| d\theta = C.$$

Thus, the summands in (4.115) are bounded by

$$|a_n r^n \cos n\theta + b_n r^n \sin n\theta| \le r^n (|a_n| + |b_n|) \le 2Cr_{\star}^n$$

According to the Weierstrass M test, since the geometric series  $\sum_{n=1}^{\infty} 2Cr_{\star}^{n} < \infty$  converges, the series (4.115) converges uniformly. Q.E.D.

#### 4.3.42

Given such a curve, let  $\delta>0$  be the minimum distance between C and the boundary  $\partial\Omega$ , which is positive since C is assumed to lie in the interior of  $\Omega$  and both curves are compact (closed and bounded). Let  $(x_i,y_i)\in C,\ i=0,\ldots,n$ , be a finite sequence of points on the curve with  $(x_n,y_n)=(x,y)$  and such that the distance from  $(x_i,y_i)$  to  $(x_{i+1},y_{i+1})$  is  $\leq \frac{1}{2}\delta$ , which implies that the disk centered at  $(x_i,y_i)$  whose boundary circle passes through  $(x_{i+1},y_{i+1})$  is contained in  $\Omega$ . Using the preceding argument, a straightforward induction then shows that

$$M^{\star} = u(x_0,y_0) = u(x_1,y_1) = u(x_2,y_2) = \ \cdots \ = u(x_n,y_n) = u(x,y),$$
 as desired.   
 
$$Q.E.D.$$

#### 4.3.46

The rescaled function  $\widehat{u}(x,y) = u(Rx,Ry)$  satisfies the boundary value problem (4.101) on the unit disk, and hence by (4.126)

$$u(r,\theta) = \widehat{u}(r/R,\theta) = \frac{1}{2\pi} \int_{-\pi}^{\pi} h(\phi) \; \frac{1 - r^2/R^2}{1 + r^2/R^2 - 2(r/R)\cos(\theta - \phi)} \; d\phi.$$

# 4.4.4(a)(c)

- (a) Elliptic when  $x \neq 0$ ; parabolic when x = 0.
- (c) Parabolic when  $x + t \neq 0$ ; degenerate when t = -x.

## 4.4.11

By the chain rule,

$$\frac{\partial u}{\partial y} = i \frac{\partial u}{\partial t}, \quad \frac{\partial^2 u}{\partial y^2} = i^2 \frac{\partial^2 u}{\partial t^2} = -\frac{\partial^2 u}{\partial t^2}, \quad \text{and hence} \quad \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} = \frac{\partial^2 u}{\partial x^2} - \frac{\partial^2 u}{\partial t^2}.$$

Thus, this complex change of variables maps the elliptic Laplace equation to the hyperbolic wave equation, and the type is not preserved.

#### 4.4.16

False. The equation is elliptic and so has no real characteristics.